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# A Poisson Random Walk is Bernoulli\*

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Abstract. Particles distributed on the integers in Poisson distribution, each independently taking a random walk, form a stationary Markov chain. The canonical shift in this space is Bernoulli.

## 1. Introduction

Consider particles distributed on  $\mathbb{Z}$  with Poisson distribution, parameter 1. (That is for each  $n \in \mathbb{Z}$ , there is a nonnegative number  $\omega_n$  of particles on n. The values  $\{\omega_n\}_{n\in\mathbb{Z}}$  are independent identically distributed (i.i.d.) random variables, each with Poisson distribution parameter 1.) Let each particle take a random walk, described as follows. (Results can be shown to remain valid for any random walk but we choose this one for simplicity. A discussion of generalizing this result to arbitrary random walk will be made in a closing remark at the end of this paper.) We let each particle stand still, move one step forward, or one step backward, each with probability 1/3. They all move independently. Then the particles will remain Poisson distributed. By continuing this procedure we obtain a sequence of configurations  $\{X_i\}_{i\in\mathbb{N}}$ , where each  $X_i$  is a Poisson-distributed  $\Omega$ -valued random variable, where  $\Omega = \mathbb{N}^{\mathbb{Z}}$ ,  $\mathbb{N} = \{0, 1, 2, ...\}$ . Since  $\{X_i\}_{i \in \mathbb{N}}$  forms a stationary Markov chain, we can extend the process to a stationary Markov chain  $\{X_i\}_{i \in \mathbb{Z}}$ . This process will henceforth be called Poisson Random Walk, abbreviated P-walk. The shift  $\{X_i\} \rightarrow \{Y_i\}, Y_i = X_{i+1}$  for all *i*, is a stationary shift. Here it is shown to be a Bernoulli shift.

My result can easily be extended to a process where particles are labeled by their past and future, thereby strengthening a result of Goldstein and Lebowitz (Theorem 5.3, page 10 of [1]).

### 2. Probabilistic Statement of Theorem to be Proved

Let  $\{X_i\}_{i\in\mathbb{Z}}$  be as in the Introduction. The theorem proved by this paper is

**Theorem 1.**  $\{X_i\}_{i \in \mathbb{Z}}$  is Bernoulli.

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The purpose of this section is to reduce this theorem to a probabilistic statement.

 $\Omega = \mathbb{N}^{\mathbb{Z}}$ . For any  $\omega \in \Omega$  and any finite set  $S \subset \mathbb{Z}$ ; let  $\omega_s$  be the restriction of  $\omega$  to the set S. Thus, when we speak of the particles of  $(X_i)_S$  we mean the particles of  $X_i$  which lie on S.  $(X_i)_S$  refers to the function taking each  $n \in S$  to the number of particles which lie on the integer n at time i.

From a theorem of Ornstein, Theorem 4, p. 53 of [3], it follows that if the process  $\{(X_i)_S\}_{i\in\mathbb{Z}}$  is Bernoulli for each finite  $S \subset \mathbb{Z}$ , then the entire process  $\{X_i\}_{i\in\mathbb{Z}}$  is Bernoulli. By definition, a process is Bernoulli if it is isomorphic to an independent process. However, that condition is difficult to check so instead we use here a condition which is equivalent.

The condition we use is Shields' alteration of the very weak Bernoulli condition of Ornstein [2]. Proving that condition for the  $\{X_i\}_s$  process is easily seen to be equivalent to proving the following theorem.

**Theorem 2.** Given  $\varepsilon > 0$  and a finite set  $S \subset \mathbb{Z}$ , we can construct a sequence of jointly distributed random variables  $\{(X_i, Y_i)\}_{i \in \mathbb{Z}}$  such that

1)  $\{X_i\}_{i \leq 0}$  and  $\{Y_i\}_{i \in \mathbb{Z}}$  are independent.

2)  $\{X_i\}, \{Y_i\}$  are both *P*-walks.

3)  $P((X_n)_S \neq (Y_n)_S) < \varepsilon$  for all sufficiently large n.

Thus the remainder of this paper is a proof of Theorem 2.

#### 3. Proof of Theorem 2

First construct the  $\{Y_i\}$  process and construct  $\{X_i\}_{i \le 0}$  to be independent of the  $\{Y_i\}$  process. We allow the particles of  $X_0$  to take a random walk in order to form configurations  $X_1, X_2, \ldots$ . As the definition of *P*-walk implies, the particles of  $X_0$  walk independently of each other. However, we will allow their walk to depend upon the walk of the  $Y_0$  particles as follows.

A one-to-one correspondence between a certain subset of the  $Y_0$  particles and a certain subset of the  $X_0$  particles will be established. The uncorresponded  $X_0$  particles will walk independently of all  $Y_0$  particles and all other  $X_0$  particles. Each corresponded  $X_0$  particle will similarly walk independent of all other particles until it meets its corresponding  $Y_0$  particle (i.e., until the two particles are both on the same integer). This will eventually happen with probability 1 because the difference between the positions of the two particles takes a symmetric, aperiodic random walk with bounded step sizes; all such walks eventually reach zero. Once the corresponded  $X_0$  particle meets its corresponding  $Y_0$  particle it will continue to move exactly as does that  $Y_0$  particle so that the two particles will stay together from that time onward.

In order to complete the definition of  $\{X_i\}$  process, all that remains is to choose the subsets of  $X_0$  and  $Y_0$  which we will correspond and to exhibit that correspondence. Choose L > 0. We will later specify how L is chosen more precisely. Let

$$I_k = \{i \in \mathbb{Z} : kL + 1 \leq i \leq (k+1)L\}.$$

Our correspondence will consist of a correspondence, for each k, between the elements of a subset of the particles of  $(X_0)_{I_k}$  and those of a subset of the particles

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of  $(Y_0)_{I_k}$ . (Recall that  $(X_0)_{I_k}$  are those particles of  $X_0$  that lie in  $I_k$ .) Let N(X, k) (respectively, N(Y, k)) be the total number of particles of  $(X_0)_{I_k}$  (respectively,  $(Y_0)_{I_k}$ ). Let

$$N_{k} = \min \{ N(X, k), N(Y, k) \}.$$

The subset we choose from  $(X_0)_{I_k}$  will consist of  $N_k$  particles of  $(X_0)_{I_k}$  randomly chosen in a uniform manner, i.e., for any given set of  $N_k$  elements of  $X_0$ , we will choose that set with probability  $1 / {\binom{N(X,k)}{N_k}}$  where  ${\binom{N(X,k)}{N_k}}$  is the binomial coefficient

$$\frac{N(X,k)!}{N_k! (N(X,k) - N_k)!}.$$

In a similar uniform manner choose  $N_k$  particles from  $(Y_0)_{I_k}$ . Since both sets have  $N_k$  members, arbitrarily choose a bijection between them. This set of bijections, for all k, collectively forms one large bijection of a subset of  $X_0$  with a subset of  $Y_0$ . This completes our construction of a  $\{X_i, Y_i\}_{i \ge 0}$  process satisfying conditions 1 and 2.

We now check condition  $3.P((X_n)_S \neq (Y_n)_S) \leq P$  (at time *n*, *S* contains uncorresponded  $X_0$  particles) + *P* (at time *n*, *S* contains uncorresponded  $Y_0$  particles) + *P* (at time *n*, *S* contains corresponded  $X_0$  particles which have not yet met their corresponding  $Y_0$  particle) + *P* (at time *n*, *S* contains corresponded  $Y_0$  particles which have not yet met their corresponding  $X_0$  particle).

Let  $A, A_1, B, B_1$  be the four terms of the right-hand side of the above inequality in order. By symmetry,  $A = A_1, B = B_1$ . To complete the proof of Theorem 2 we need only show

1) By properly choosing L, we can insure that both  $A < \varepsilon/4$  and  $B < \varepsilon/4$  for sufficiently large n.

For  $z \in \mathbb{Z}$ , let  $c_z$  be the number of uncorresponded  $X_0$  particles which are on point z. Note that  $E(c_z)$  is independent of z. For  $z_1, z_2 \in \mathbb{Z}$ , let  $P_n(z_1, z_2)$  be the probability that a particle which is on  $z_1$  at time 0 will be at  $z_2$  at time n. Note that  $P_n(z_1, z_2)$  depends only on  $z_2 - z_1$ . Thus for any fixed  $z_0 \in \mathbb{Z}$ 

$$\sum_{z\in\mathbb{Z}}P_n(z,z_0)=\sum_{z\in\mathbb{Z}}P_n(z_0,z)=1.$$

2)  $A \leq E(\# \{\text{uncorresponded } X_0 \text{ particles which reach } S \text{ at time } n\})$ 

$$= \sum_{\substack{z_1 \in \mathbb{Z} \\ z_2 \in S}} E(c_{z_1}) P_n(z_1, z_2) = E(c_0) \sum_{\substack{z_1 \in \mathbb{Z} \\ z_2 \in S}} P_n(z_1, z_2)$$
  
$$= E(c_0) \sum_{z_2 \in S} (1) = \# SE(c_0) = \frac{1}{L} \# (S)E \sum_{i=1}^{L} (c_i)$$
  
$$= \frac{1}{L} \# SE(N(X, 0)) - \min(\{N(X, 0), N(Y, 0)\})$$

For any i.i.d. random variables X, Y with variance L,

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$$E(X - \min(X, Y)) = E(X - Y)^{+} = \frac{1}{2}E|X - Y| \le \frac{1}{2}\sqrt{E(X - Y)^{2}} = \sqrt{L/2}$$

Hence the right-hand side of 2 can be made arbitrarily small by choosing L sufficiently large.

Now suppose an  $X_0$  particle at x is corresponded to a  $Y_0$  particle at y. Let  $P_n(x, y, z)$  be the probability that the two particles do not meet by time n and that the  $X_0$  particle reaches z at time n. Note that  $P_n(x, y, z)$  depends only on z - x and y - x. Also recall that y - x lies somewhere between -L and L. Thus

$$P_n(x, y, z) \leq \sum_{i=-L}^{L} P_n(0, i, z - x).$$

For  $-L \leq i \leq L$ , let  $a_{n,i}$  be the probability that an  $X_0$  particle at 0 and a  $Y_0$  particle at *i* will not meet by time *n*. Then  $\lim_{n \to \infty} a_{n,i} = 0$ , and hence

$$\lim_{n\to\infty}\sum_{i=-L}^{L}a_{n,i}=0.$$

Note also that for any  $z \in \mathbb{Z}$ ,

$$\sum_{x \in \mathbb{Z}} P_n(0, i, z - x) = a_{n,i}$$

Recall that the expected number of  $X_0$  particles at any site is 1. Since the probability that a corresponded particle at x reaches z at time n without reaching its corresponding  $Y_0$  particle is bounded above by

$$\sum_{i=-L}^{L} P_n(0, i, z-x),$$

we get

$$B \leq \sum_{\substack{z \in S \\ x \in \mathbb{Z}}} \left( E \ \# \left\{ \text{corresponded } X_0 \text{ particles at } x \right\} \sum_{i=-L}^{L} P_n(0, i, z - x) \right)$$
$$\leq \sum_{\substack{z \in S \\ x \in \mathbb{Z}}} \left\{ \left( \sum_{i=-L}^{L} P_n(0, i, z - x) \right) \right\}$$
$$= \sum_{\substack{z \in S \\ z \in S}} \sum_{i=-L}^{L} \left( \sum_{\substack{x \in \mathbb{Z}}} P_n(0, i, z - x) \right) = \# S \sum_{i=-L}^{L} a_{n,i}$$

which goes to zero as n approaches  $\infty$ .

*Remark.* For simplicity I only mentioned the 1/3, 1/3, 1/3 random walk, but the proof goes through word for word for any aperiodic random walk whose step size has finite expectation. It is almost the same if we drop the aperiodic assumption. Here we carry out the same coupling, making sure we do not correspond two particles on non-communicating sites.

If the random walk has infinite expectation we use a coupling due to D. S. Ornstein. Couple corresponding particles so that they take the same step size if that size is larger than one large prechosen number N in absolute value, and

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take independent steps conditioned on these step sizes not exceeding N in absolute value. Then the difference between the two random walks is bounded and will eventually reach zero if periodicity is right and N is chosen large enough so that the difference is not deterministic.

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#### References

- 1. Goldstein, S. Lebowitz, J. L. : Commun. Math. Phys. 37, 1-18 (1974)
- 2. Shields, P.: Monatsh. Math. 84, 133-142 (1977)
- 3. Ornstein, D. S.: Ergodic Theory, Randomness, and Dynamical Systems. New Haven: Yale University Press 1974

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