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## Higher Order Estimates for the Yukawa<sub>2</sub> Quantum Field Theory\*

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Abstract. Higher order estimates of the form

$$\prod_{1}^{n} N_{\tau_i} \leq \operatorname{const}(H(g) + I)^n, \qquad \sum_{1}^{n} \tau_i < 1, \, \tau_i \geq 0$$

are proved for the Yukawa<sub>2</sub> models with and without SU<sub>3</sub> symmetry. We also prove norm convergence of  $\prod_1^n N_{\tau_1}^{i} \cdot R_{\kappa}^{n/2+\delta}$  as  $\kappa \to \infty$  where  $R_{\kappa} = (H(g, \kappa) + I)^{-1}$ .

## **Introduction and Results**

Higher order estimates, bounding powers of the fractional energy operator by powers of the Hamiltonian, have proved useful in studying the  $\mathscr{P}(\phi)_2$  model [1]. In this paper we obtain similar estimates for the Yukawa<sub>2</sub> model as well as for the Yukawa<sub>2</sub> model with internal SU<sub>3</sub> symmetry discussed in [2].

In the following we will use even, positive odd and negative odd values of  $\varepsilon$  to label bosons, fermions and anti-fermions respectively. Thus  $b(k, \varepsilon)$  denotes the annihilation operator for free particles of momentum k and type  $\varepsilon$ . The fractional energy operator is:

$$\begin{split} N_{\tau} &= \sum_{\varepsilon} N_{\tau}^{(\varepsilon)} = \sum_{\varepsilon} \int dk \; \mu(k, \varepsilon)^{\tau} \; b^{*}(k, \varepsilon) \; b(k, \varepsilon) \; , \\ \mu(k, \varepsilon) &= (k^{2} + m(\varepsilon)^{2})^{\frac{1}{2}} \; , \end{split}$$

where  $m(\varepsilon) = m$  for bosons,  $m(\varepsilon) = M$  for fermions. For convenience we define  $E(k) = \sqrt{k^2 + 1}$ . We will work with the dense domain  $\mathcal{D}$  of vectors in Fock space with finite numbers of particles and wave functions in Schwartz space.

Formally, the finite volume Hamiltonian H(g) has the form

$$H(g) = H_0 + H_I(g) + C(g)$$
  
=  $N_1 + \lambda \int dx g(x) : \overline{\psi} \psi \phi : -\frac{1}{2} \delta m^2 \int dx g^2(x) : \phi^2 : -E(g),$ 

where  $g \ge 0 \in C_0^{\infty}$  and  $\delta m^2$ , E(g) provide infinite renormalizations. To define the momentum cutoff Hamiltonian  $H(g, \kappa)$  we multiply the momentum space kernels  $w^c(k, p_1, p_2)$ ,  $w(k, p_1, p_2)$  of the interaction term  $H_I(g)$  by a general momentum cutoff function  $\chi_{\kappa}(k, p_1, p_2)$  in the sense of [3]. The renormalization constants

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are then defined as:

$$\begin{split} \delta m_{\kappa}^{2} &= -\frac{\lambda^{2}}{2\pi} \int dk \, \omega(k)^{-1} |\chi_{\kappa}(0, k/2, -k/2)|^{2} + \text{const} + o(1) \,, \\ E(g, \kappa) &= -\int dk \, dp_{1} \, dp_{2} |w^{c}(k, p_{1}, p_{2}) \, \chi_{\kappa}(k, p_{1}, p_{2})|^{2} \left(\mu(k) + \omega(p_{1}) + \omega(p_{2})\right)^{-1} \\ &+ \text{const} + o(1) \,, \end{split}$$

where  $\mu(k) \equiv \mu(k, 0)$ ,  $\omega(k) = \mu(k, 1)$ . The SU<sub>3</sub> Hamiltonians involve slight generalizations and are defined in [2]. Glimm and Jaffe [3] have shown that the Hamiltonians, with a suitable choice of the constant in  $E(q, \kappa)$ , define positive self-adjoint operators converging in the sense of resolvents to a positive selfadjoint operator H(q) which satisfies:

$$N_{\tau} \leq \operatorname{const}(H(g) + I), \quad \tau < 1.$$

Furthermore the operators  $H(g, \kappa)$  are essentially self-adjoint on  $\mathcal{D}$  and satisfy  $\kappa$ -dependent estimates of the form:

$$\prod_{i=1}^n N_{\tau_i} \leq C_{\kappa} (H(g,\kappa) + I)^n, \qquad \sum_{i=1}^n \tau_i < 1.$$

The proof of these estimates requires the essential self-adjointness of  $H(g, \kappa, \sigma)^n$ on  $\mathcal{D}$ , where  $H_I(g, \kappa, \sigma)$  and  $C(g, \kappa, \sigma)$  have momentum space kernels in Schwartz space converging to those of  $H_I(q,\kappa)$  and  $C(q,\kappa)$  as  $\sigma \to \infty$ , which follows by techniques similar to those of Jaffe, Lanford, and Wightman [6]. For notation and general techniques we refer to Glimm and Jaffe [3], Dimock [4], and McBryan [2, 5]. Our main results are:

**Theorem 1.** Provided  $\tau = \sum_{i=1}^{n} \tau_i < 1$ ,  $\tau_i \ge 0$ , then there is a constant, depending on n,  $\tau$ , g, such that n

$$\prod_{i=1}^{n} N_{\tau_i} \leq \operatorname{const}(H(g) + I)^n.$$
(1)

The restriction  $\tau < 1$  is indicated by perturbation theory and so we expect that (1) are the most general  $n^{th}$  order estimates. It is also useful to have estimates controlling the inequality (1) as the momentum cutoff  $\kappa$  is removed. With  $R_{\kappa} = (H(g, \kappa) + I)^{-1}, R \equiv R_{\infty} = (H(g) + I)^{-1}$ , we define  $\delta R^{\beta} = R_{\kappa_2}^{\beta} - R_{\kappa_1}^{\beta}$  where one of  $\kappa_1, \kappa_2$  may be infinite.

**Theorem 2.** Provided  $\tau = \sum_{i=1}^{n} \tau_i < 1$ ,  $\tau_i \ge 0$ , and  $\delta > 0$ , there is a constant and an  $\varepsilon > 0$  such that

- (i)  $\Pi_{i=1}^{n} N_{\tau_{i}}^{\frac{1}{2}} \cdot \delta R^{n/2}$  converges weakly to 0 as  $\kappa = \min(\kappa_{1}, \kappa_{2})$  tends to  $\infty$ . (ii)  $\|\Pi_{1}^{n} N_{\tau_{i}}^{\frac{1}{2}-\delta} \cdot \delta R^{n/2}\| \leq \operatorname{const} \kappa^{-\varepsilon}$ . (iii)  $\|\Pi_{1}^{n} N_{\tau_{i}}^{\frac{1}{2}} \cdot \delta R^{n/2+\delta}\| \leq \operatorname{const} \kappa^{-\varepsilon}$ .
- (ii)
- (iii)

Theorem 1 follows from two lemmas:

**Lemma 3.** For  $\psi \in D(\prod_{i=1}^{n} N_{\tau_{i}}^{\frac{1}{2}})$  and an arbitrary choice of  $\varepsilon_{i}, \tau_{i}$ :

$$\left\|\prod_{i=1}^{n} N_{\tau_{i}}^{(\varepsilon_{i})^{\frac{1}{2}}}\psi\right\|^{2} = \sum_{r=1}^{n} \int dk_{1} \dots dk_{r} \sum_{\substack{1=i_{1} < \cdots i_{r} \\ i_{1} \dots i_{r} < k_{1}, \dots, k_{r}}}^{r, \varepsilon} N_{i_{1} \dots i_{r}}^{r, \varepsilon}(k_{1}, \dots, k_{r}) \\ \cdot \|b(k_{r}, \varepsilon_{i_{r}}) \dots b(k_{1}, \varepsilon_{i_{1}})\psi\|^{2}, \qquad (2)$$

where  $P_{i_1...i_r}^{\tau, \epsilon}$  are homogeneous expressions of degree  $\sum_{i=1}^{n} \tau_i$  in  $\mu(k_i, \epsilon_i)$ . Explicitly:

$$P_{i_1\ldots i_r}^{\boldsymbol{\tau},\,\boldsymbol{\varepsilon}}(k_1,\ldots,\,k_r) = \prod_{t=1}^r \left[ \mu^{\boldsymbol{\tau}_{t_t}}(k_t,\varepsilon_{i_t}) \sum_{s=i_t+1}^{i_{t+1}-1} \cdot \left\{ \sum_{l=1}^t \,\delta_{\varepsilon_{i_l},\,\varepsilon_s} \mu^{\boldsymbol{\tau}_s}(k_l,\,\varepsilon_s) \right\} \right],$$

where for t = r we define  $i_{r+1} - 1 \equiv n$  and a vacuous product is always taken to be 1, *i.e.*,  $\prod_{s=3}^{2} \{ \} \equiv 1$ .

**Lemma 4.** For any choice of  $\varepsilon_i$ ,  $\sum_{i=1}^n \tau_i = \tau < 1$  and for any  $\theta$ , there is a constant independent of  $\kappa$  with:

$$\int dk_1 \dots dk_n E^{\tau_1}(k_1) \dots E^{\tau_n}(k_n) \| b(k_n, \varepsilon_n) \dots b(k_1, \varepsilon_1) R_{\kappa}^{n/2} \theta \|^2 \leq \text{const} \| \theta \|^2.$$
(3)

*Proof of Theorem 1.* From the form of  $P_{i_1...i_r}^{\tau, \varepsilon}$  we have

$$|P_{i_1\ldots i_r}^{\tau,\epsilon}(k_1,\ldots k_r)| \leq \operatorname{const} \prod_{t=1}^r \left[ E^{\tau_{i_t}}(k_t) \prod_{s=i_t+1}^{i_{t+1}-1} \left\{ \sum_{l=1}^t E^{\tau_s}(k_l) \right\} \right].$$

Inserting this in (2) and applying Lemma 4 with  $\psi = R_{\kappa}^{n/2} \theta$  we obtain:

$$\left\|\prod_{1}^{n} N_{\tau_{i}}^{(\varepsilon_{i})^{\frac{1}{2}}} R_{k}^{n/2} \theta\right\|^{2} \leq \operatorname{const} \|\theta\|^{2}$$

$$\tag{4}$$

with a constant independent of  $\kappa$ . Since  $R_{\kappa}$  converges in norm to R,  $\prod_{i=1}^{n} N_{\tau_{i}}^{(\varepsilon_{i})^{\frac{1}{2}}} R_{\kappa}^{n/2}$  converges weakly on  $D(\prod_{i=1}^{n} N_{\tau_{i}}^{(\varepsilon_{i})^{\frac{1}{2}}})$  to  $\prod_{i=1}^{n} N_{\tau_{i}}^{(\varepsilon_{i})^{\frac{1}{2}}} R^{n/2}$  and the uniform bounds (4) then apply also to the case  $\kappa = \infty$ . This completes the proof of Theorem 1 and of Theorem 2 (i).

Proof of Theorem 2. (ii) We use  $||A|| = ||A^*A||^{\frac{1}{2}}$ . Thus

$$\left\| \left(\prod_{1}^{n} N_{\tau_{i}}\right)^{\frac{1}{2}-\delta} \delta R^{n/2} \right\| = \left\| \delta R^{n/2} \left(\prod_{1}^{n} N_{\tau_{i}}\right)^{1-2\delta} \delta R^{n/2} \right\|^{\frac{1}{2}}$$

$$\leq \left\| \left(\prod_{1}^{n} N_{\tau_{i}}\right)^{\frac{1}{2}-2(1-\alpha)\delta} \delta R^{n/2} \right\|^{\frac{1}{2}} \left\| \left(\prod_{1}^{n} N_{\tau_{i}}\right)^{\frac{1}{2}-2\alpha\delta} \delta R^{n/2} \right\|^{\frac{1}{2}}$$

$$\leq \operatorname{const} \left\| \left(\prod_{1}^{n} N_{\tau_{i}}\right)^{\frac{1}{2}-2\alpha\delta} \delta R^{n/2} \right\|^{\frac{1}{2}}, \qquad (5)$$

where we have used Theorem 1 in the form

$$\left\| \left( \prod_{1}^{n} N_{\tau_{i}} \right)^{\frac{1}{2} - \delta'} \delta R^{n/2} \right\|^{\frac{1}{2}} \leq \left( \left\| \left( \prod_{1}^{n} N_{\tau_{i}} \right)^{\frac{1}{2}} R_{\kappa_{1}}^{n/2} \right\| + \left\| \left( \prod_{1}^{n} N_{\tau_{i}} \right)^{\frac{1}{2}} R_{\kappa_{2}}^{n/2} \right\| \right)^{\frac{1}{2}} \leq \text{const}.$$

The inequality (5) allows us to reduce the exponent  $\frac{1}{2} - \delta$  to  $\frac{1}{2} - 2\alpha\delta$ ,  $0 < \alpha \leq 1$ . By iterating *m* times and choosing *m*,  $\alpha$  carefully we reduce the exponent to  $\frac{1}{2} - (2\alpha)^m \delta$  with  $0 \leq \frac{1}{2} - (2\alpha)^m \delta \leq \frac{\gamma}{2n}$ ,  $\gamma < 1$ . Thus  $\left\| \left( \prod_{n=1}^{n} N \right)^{\frac{1}{2} - \delta} \delta R^{n/2} \right\| \leq \text{const} \left\| \left( \prod_{n=1}^{n} N \right)^{\frac{\gamma}{2n}} \delta R^{n/2} \right\|^{2^{-m}}$ 

$$\left\| \left( \prod_{1}^{n} N_{\tau_{1}} \right)^{\frac{1}{2}-\delta} \delta R^{n/2} \right\| \leq \operatorname{const} \left\| \left( \prod_{1}^{n} N_{\tau_{1}} \right)^{\frac{1}{2}n} \delta R^{n/2} \right\|^{\frac{1}{2}-m}$$
$$\leq \operatorname{const} \| N_{\tau}^{\frac{1}{2}} \delta R^{n/2} \|^{\frac{1}{2}-m}$$
$$\leq \operatorname{const} \kappa^{-\varepsilon} \quad \text{for some } \varepsilon > 0 \,,$$

where we have used the norm convergence of  $N_{\tau}^{\gamma/2} \delta R^{\frac{1}{2}}$ ,  $\gamma < 1$ , which we have proved in [2, Theorem 2.4.1]. There remains only the specification of *m* and  $\alpha$ . A suitable choice is

$$m > \log 2\delta / \log \left(1 - \frac{\gamma}{n}\right), \quad \alpha = \frac{1}{2} \left(\frac{1}{2\delta}\right)^{\overline{m+1}}, \quad \delta < \frac{1}{2}.$$

(iii) We use the inequality, valid for  $\tau \ge 0$  and  $\beta \ge 1$ 

Thus with 
$$0 < \delta < \frac{1}{2}$$
 and  $\beta > 1$ :  

$$\prod_{i=1}^{n} N_{\tau_{i}} = \prod_{i=1}^{n} N_{\tau_{i}}^{1-2\delta} \prod_{i=1}^{n} N_{\tau_{i}}^{2\delta} \le \prod_{i=1}^{n} N_{\tau_{i}}^{1-2\delta} \prod_{i=1}^{n} N_{\tau_{i}}^{2\beta\delta} \le \prod_{i=1}^{n} N_{\tau_{i}}^{1-2\delta} \cdot N_{\tau/\beta}^{2n\beta\delta}, \quad \tau = \sum_{i=1}^{n} \tau_{i}.$$

We now choose  $\beta$  so that  $\tau + \tau/\beta < 1$ , i.e.  $\beta > \max\left(1, \frac{\tau}{1-\tau}\right)$ , and we choose  $\delta > 0$  sufficiently small that  $2n\beta\delta = 1-2\delta$ , i.e.,  $\delta = (2(n\beta+1))^{-1}$ . Then defining  $\tau_{n+1} = \tau/\beta$  we have

$$\prod_{1}^{n} N_{\tau_{i}} \leq \prod_{1}^{n+1} N_{\tau_{i}}^{1-2\delta} \text{ and } \sum_{1}^{n+1} \tau_{i} < 1$$

and with  $R_{\kappa}(\zeta) = (H(g, \kappa) - \zeta)^{-1}, \zeta \leq -1$ , we obtain by (ii):

$$\left\|\prod_{1}^{n} N_{\tau_{i}}^{\frac{1}{2}} \,\delta R(\zeta)^{(n+1)/2}\right\| \leq \left\|\prod_{1}^{n+1} N_{\tau_{i}}^{\frac{1}{2}-\delta} \,\delta R(\zeta)^{(n+1)/2}\right\| \leq \operatorname{const} \kappa^{-\varepsilon'}.$$
(6)

However, by Theorem 1 we have

$$\left\|\prod_{1}^{n} N_{\tau_{i}}^{\frac{1}{2}} \delta R(\zeta)^{(n+1)/2}\right\| \leq \operatorname{const} |\zeta|^{-\frac{1}{2}}.$$
(7)

Combining (6) and (7) using  $||A|| = ||A||^{\nu} ||A||^{1-\nu}, 0 \le \nu \le 1$ , we get

$$\left\|\prod_{1}^{n} N_{\tau_{1}}^{\frac{1}{2}} \delta R(\zeta)^{(n+1)/2}\right\| \leq \operatorname{const} |\zeta|^{-\frac{1}{2}+\delta/2} \kappa^{-\varepsilon}, \quad \varepsilon = \varepsilon' \delta > 0,$$

where  $\delta > 0$  has no relation to the  $\delta$  used previously. Finally, using the identity, valid for  $0 < \delta < \frac{1}{2}$ :

$$R_{\kappa}(\zeta)^{n/2+\delta} = c(\delta) \int_{0}^{\infty} d\lambda \, \lambda^{-\frac{1}{2}-\delta} R_{\kappa}(\zeta-\lambda)^{(n+1)/2} ,$$

where

$$c(\delta) = \left[\int_{0}^{\infty} d\lambda \, \lambda^{-\frac{1}{2}-\delta} (1+\lambda)^{-(n+1)/2}\right]^{-1}$$

we obtain for  $\delta > 0$ :

$$\begin{aligned} \left\| \prod_{1}^{n} N_{\tau_{i}}^{\frac{1}{2}} \,\delta R(\zeta)^{n/2+\delta} \right\| &\leq c(\delta) \int_{0}^{\infty} d\lambda \,\lambda^{-\frac{1}{2}-\delta} \left\| \prod_{1}^{n} N_{\tau_{i}}^{\frac{1}{2}} \,\delta R(\zeta-\lambda)^{(n+1)/2} \right\| \\ &\leq c(\delta) \int_{0}^{\infty} d\lambda \,\lambda^{-\frac{1}{2}-\delta} \operatorname{const} |\zeta-\lambda|^{-\frac{1}{2}+\delta/2} \,\kappa^{-\varepsilon} \\ &\leq \operatorname{const} \kappa^{-\varepsilon} \end{aligned}$$

which completes the proof of Theorem 2.

Proof of Lemmas 3, 4. Lemma 3 follows easily by induction from

$$b(k, \varepsilon_1) N_{\tau}^{(\varepsilon_2)^{\frac{1}{2}}} = \left( N_{\tau}^{(\varepsilon_2)} + \delta_{\varepsilon_1, \varepsilon_2} \mu(k, \varepsilon_1)^{\tau} \right)^{\frac{1}{2}} b(k, \varepsilon_1)$$

and

$$\|(N_{\tau}^{(\varepsilon)}+a)^{\frac{1}{2}}\psi\|^{2} = \|N_{\tau}^{(\varepsilon)^{\frac{1}{2}}}\psi\|^{2} + a\|\psi\|^{2}$$

The proof of Lemma 4 depends on the renormalized pull-through formula [5]. This expansion takes the form:

$$b(k_{n},\varepsilon_{n})\dots b(k_{1},\varepsilon_{1}) R_{\kappa}^{\alpha}(\zeta) = \sum_{r=0}^{n} \sum_{i_{r}>\dots>i_{1}\geq1}^{n} (-)^{\delta_{i_{r}}^{r,\varepsilon_{\dots i_{1}}}} \cdot U_{r,\kappa}^{(\alpha)}(\zeta-\mu(k_{n},\varepsilon_{n})\dots-\mu(k_{\overline{i_{r}}},\overline{\varepsilon_{i_{r}}})\dots-\mu(k_{\overline{i_{1}}},\overline{\varepsilon_{i_{1}}})\dots-\mu(k_{1},\varepsilon_{1}); k_{i_{r}}\varepsilon_{i_{r}},\dots,k_{i_{1}}\varepsilon_{i_{1}})$$
(8)  
$$\cdot b(k_{n},\varepsilon_{n})\dots b(k_{\overline{i_{r}}},\overline{\varepsilon_{i_{r}}})\dots b(k_{\overline{i_{1}}},\overline{\varepsilon_{i_{1}}})\dots b(k_{1},\varepsilon_{1}),$$

where  $0 < \alpha \leq 1$  and a slash denotes absence of a term. The indices  $\delta_{i_r...i_1}^{r,\epsilon}$  are given by:

$$\delta_{i_r\ldots i_1}^{r,e} = \sum_{s=1}^r \left( \sum_{j=i_s+1}^{i_{s+1}-1} \varepsilon_j \right) \left( \sum_{l=1}^s \varepsilon_{i_l} \right),$$

while the  $U_{n,\kappa}^{(\alpha)}$  are defined inductively by:

$$U_{0,\kappa}^{(\alpha)}(\zeta) = R_{\kappa}(\zeta)^{\alpha},$$
$$U_{n,\kappa}^{(\alpha)}(\zeta; k_{n}\varepsilon_{n}, \dots, k_{1}\varepsilon_{1}) = b(k_{n}, \varepsilon_{n}) U_{n-1,\kappa}^{(\alpha)}(\zeta; k_{n-1}\varepsilon_{n-1}, \dots, k_{1}\varepsilon_{1})$$
$$- (-)^{\varepsilon_{n}\frac{n}{2}\varepsilon_{i}} U_{n-1,\kappa}(\zeta - \mu(k_{n}, \varepsilon_{n}); k_{n-1}\varepsilon_{n-1}, \dots, k_{1}\varepsilon_{1}) b(k_{n}, \varepsilon_{n}).$$

For a fuller treatment of these defining relations and of (8) see [5]. In [5] we have proved that:

$$\int dk_1 \dots dk_n E^{\tau_n}(k_n) \dots E^{\tau_1}(k_1) \| U_{n,\kappa}^{(\frac{1}{2})}(\zeta; k_n \varepsilon_n, \dots, k_1 \varepsilon_1) \|^2 \leq \text{const},$$
(9)

uniformly in  $\kappa$  provided  $\sum_{i=1}^{n} \tau_i < 1$ .

 $||b(k_n,$ 

Returning to the proof of Lemma 4, we pull all of the  $b(k_i, \varepsilon_i)$  through the first  $R_{\kappa}^{\frac{1}{2}}$  in (3) obtaining by (8):

$$\int dk_{1} \dots dk_{n} E^{\tau_{n}}(k_{n}) \dots E^{\tau_{1}}(k_{1}) \|b(k_{n}, \varepsilon_{n}) \dots b(k_{1}, \varepsilon_{1}) R^{n/2} \theta\|^{2} \leq 2^{n} \int dk_{1} \dots dk_{n} E^{\tau_{n}}(k_{n}) \dots E^{\tau_{1}}(k_{1}) \|R^{\frac{1}{2}}(1-\mu(k_{n}, \varepsilon_{n}) \dots -\mu(k_{1}, \varepsilon_{1})) \cdot b(k_{n}, \varepsilon_{n}) \dots b(k_{1}, \varepsilon_{1}) R^{(n-1)/2} \theta\|^{2} + 2^{n} \int dk_{1} \dots dk_{n} E^{\tau_{n}}(k_{n}) \dots E^{\tau_{1}}(k_{1}) \sum_{r=1}^{n} \sum_{i_{r} > \dots > i_{1} \ge 1}^{n}, . \\\|U_{r,\kappa}^{(\frac{1}{2})}(1-\mu(k_{n}, \varepsilon_{n}) \dots -\mu(k_{t_{r}} \varepsilon_{t_{r}}) \dots -\mu(k_{t_{1}} \varepsilon_{t_{1}}) \dots -\mu(k_{1}, \varepsilon_{1}); k_{i_{r}} \varepsilon_{i_{r}}, \dots, k_{i_{1}} \varepsilon_{i_{1}})\|^{2} \\\|b(k_{n}, \varepsilon_{n}) \dots b(k_{t_{r}, r} \varepsilon_{t_{r}}) \dots b(k_{1}, \varepsilon_{1}) R^{(n-1)/2} \theta\|^{2},$$

$$(10)$$

where we have supressed  $\kappa$  for convenience and used

$$\left\|\sum_{i=1}^{m} a_{i}\right\|^{2} \leq m \sum_{i=1}^{m} \|a_{i}\|^{2}.$$

The proof of Lemma 4 now follows by induction on *n*. We assume the result (3) for all possible choices of  $\tau_i$ , i = 1, ..., m with  $\sum_{i=1}^{m} \tau_i \leq \tau$  and for all  $m, 1 \leq m \leq (n-1)$ . The first term in (6) is

$$2^{n} \int dk_{n} \dots dk_{1} E^{\tau_{n}}(k_{n}) \dots E^{\tau_{1}}(k_{1}) \\ \cdot \|R^{\frac{1}{2}}(\zeta - \mu(k_{n},\varepsilon_{n})\dots - \mu(k_{1},\varepsilon_{1})) b(k_{n},\varepsilon_{n})\dots b(k_{1},\varepsilon_{1}) R^{(n-1)/2} \theta\|^{2} \\ = 2^{n} \int dk_{n-1} \dots dk_{1} E^{\tau_{n-1}}(k_{n-1})\dots E^{\tau_{1}}(k_{1}) \{ \int dk_{n} E^{\tau_{n}}(k_{n}) \|R^{\frac{1}{2}}(\zeta - \mu(k_{n},\varepsilon_{n})\dots \mu(k_{1},\varepsilon_{1})) \\ \cdot (N^{(\varepsilon_{n})}_{\tau_{n}} + \mu^{\tau_{n}}(k_{n},\varepsilon_{n}))^{\frac{1}{2}} b(k_{n},\varepsilon_{n}) N^{(\varepsilon_{n})-\frac{1}{2}}_{\tau_{n}} b(k_{n-1},\varepsilon_{n-1})\dots b(k_{1},\varepsilon_{1}) R^{(n-1)/2} \theta\|^{2} \} \\ \leq \text{const} \int dk_{n-1} \dots dk_{1} E^{\tau_{n-1}}(k_{n-1})\dots E^{\tau_{1}}(k_{1}) \|b(k_{n-1},\varepsilon_{n-1})\dots b(k_{1},\varepsilon_{1}) R^{(n-1)/2} \theta\|^{2} \\ \leq \text{const} \|\theta\|^{2}$$

by the induction hypothesis.

We have used the first order estimate in the form

$$\int dk_n E^{\tau_n}(k_n) \| R^{\frac{1}{2}}(\zeta - \mu(k_n, \varepsilon_n)) \left( N_{\tau_n}^{(\varepsilon_n)} + \mu^{\tau_n}(k_n, \varepsilon_n) \right)^{\frac{1}{2}} b(k_n, \varepsilon_n) N_{\tau_n}^{(\varepsilon_n) - \frac{1}{2}} \chi \|^2$$
  

$$\leq \operatorname{const} \int dk_n E^{\tau_n}(k_n) \| b(k_n, \varepsilon_n) N_{\tau_n}^{(\varepsilon_n) - \frac{1}{2}} \chi \|^2$$
  

$$\leq \operatorname{const} \| N_{\tau_n}^{(\varepsilon_n)^{\frac{1}{2}}} N_{\tau_n}^{(\varepsilon_n) - \frac{1}{2}} \chi \|^2 = \operatorname{const} \| \chi \|^2.$$

For the remaining terms  $(r \neq 0)$  in (10) we use the estimate (9) for  $U_{r,\kappa}^{(\frac{1}{2})}$ . Thus

$$2^{n} \int dk_{n} \dots dk_{1} E(k_{n})^{\iota_{n}} \dots E^{\iota_{1}}(k_{1}) \sum_{r=1}^{n} \sum_{i_{r} > \dots > i_{1} \ge 1}^{n},$$

$$\| U_{r,\kappa}^{(\frac{1}{2})}(1 - \mu(k_{n}, \varepsilon_{n}) \dots - \mu(k_{\tau_{r}} \varepsilon_{i_{r}}) \dots - \mu(k_{\tau_{1}} \varepsilon_{i_{1}}) \dots - \mu(k_{\tau_{1}} \varepsilon_{i_{1}}); k_{i_{r}} \varepsilon_{i_{r}}, \dots, k_{i_{1}} \varepsilon_{i_{1}}) \|^{2},$$

$$\| b(k_{n}, \varepsilon_{n}) \dots b(k_{\tau_{r}}, \varepsilon_{i_{r}}) \dots b(k_{\tau_{1}}, \varepsilon_{i_{1}}) \dots b(k_{1}, \varepsilon_{1}) R^{(n-1)/2} \theta \|^{2}$$

$$\leq \text{const} \int dk_{n} \dots dk_{i_{r}} \dots dk_{i_{1}} \dots dk_{i_{1}} \dots b(k_{1}, \varepsilon_{1}) R^{(n-r)/2} \theta \|^{2}$$

$$\leq \text{const} \| \theta \|_{2}^{2}$$

by the induction hypothesis.

Since the induction hypothesis is certainly valid for n = 1:

$$\int dk \ E^{\tau}(k) \ \|b(k,\varepsilon) \ R_{\kappa}^{\frac{1}{2}}\theta\|^{2} \leq \int dk \ \text{const} \ \mu(k,\varepsilon)^{\tau} \ \|b(k,\varepsilon) \ R_{\kappa}^{\frac{1}{2}}\theta\|^{2}$$
$$\leq \text{const} \ \|N_{\tau}^{(\varepsilon)\frac{1}{2}} \ R_{\kappa}^{\frac{1}{2}}\theta\|^{2}$$
$$\leq \text{const} \ \|\theta\|^{2}$$

by the first order estimate, Lemma 4 follows by induction.

## References

- 1. Rosen, L.: The  $(\phi^{2n})_2$  quantum field theory: higher order estimates, Commun. Pure Appl. Math. 24, 417–457 (1971)
- McByan, O.: Vector currents in the Yukawa<sub>2</sub> quantum field theory. University of Toronto (preprint)
- 3. Glimm, J., Jaffe, A.: Quantum field theory models. In: Dewitt, C., Stora, A. (Eds.): Statistical mechanics and quantum field theory. Gordon and Breach 1971

- 4. Dimock, J.: Estimates, renormalized currents, and field equations for the Yukawa<sub>2</sub> field theory. Ann. Phys. **72**, 177–242 (1972)
- 5. McBryan, O.: The *n*-th order renormalized pull-through formula for the Yukawa<sub>2</sub> quantum field theory. Rockefeller University (preprint)
- Jaffe, A., Lanford, O., Wightman, A.: A general class of cut-off model field theories. Commun. math. Phys. 15, 47–68 (1969)

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