

SELFADJOINT OPERATOR EXTENSIONS SATISFYING THE WEYL COMMUTATION RELATIONS

BY PALLE E. T. JØRGENSEN¹

ABSTRACT. Motivated by questions concerning uniqueness of unbounded derivations in commutative C^* -algebras, and related problems on singular perturbations, we define two mixed global and infinitesimal versions of the Weyl operator commutation relations (one degree of freedom and infinite multiplicity), a weak one and a strong one. We announce two structure theorems of a geometric nature which characterize the nonselfadjoint symmetric operators entering in the Weyl systems. Proofs are only indicated.

Our starting point is the following variant of the Stone-von Neumann Uniqueness Theorem [12], [4b]. Let (U, V) be a pair of unitary one-parameter groups (always assumed strongly continuous) of operators on a separable Hilbert space H , and suppose that the Weyl commutation relation

$$(1) \quad U(t)V(s) = V(s)U(t)e^{its} \quad (\text{for all } s, t \in \mathbf{R})$$

holds. Then it is possible to represent the system in the form $SU(t)S^{-1}f(x) = f(x + t)$, $SV(s)S^{-1}f(x) = e^{isx}f(x)$, where S is an isometry of a space $L^2(\mathbf{R}, N)$ of the norm-square integrable functions f , with values in a separable Hilbert space N , onto H ; the dimension of N being equal to the (uniform) multiplicity of the spectrum of U .²

Instead of (1) we consider the following *infinitesimal* Weyl relation with symmetric but generally nonselfadjoint generator. Let $\{U(t)\}_{t \in \mathbf{R}}$ be a unitary one-parameter group on H , and let Q be a symmetric operator with dense domain $\mathcal{D}(Q)$ in H . The corresponding relation

$$(2) \quad (U(t)Qf, g) = (U(t)f, Qg) + t(U(t)f, g) \quad \text{for all } f, g \in \mathcal{D}(Q)$$

is here called the infinitesimal Weyl relation for the triple (U, Q, H) . It is clearly equivalent to (1) if Q is essentially selfadjoint. But in scattering theory of singular perturbations, and in recent investigations of the author concerning uniqueness of unbounded derivations, the relation (2) for nonselfadjoint Q plays an interesting role. Simple examples show that the operator Q of a given system (U, Q, H) may

Received by the editors August 3, 1978.

AMS (MOS) subject classifications (1970). Primary 47B25, 81A20; Secondary 43A65, 46K99, 47B47.

Key words and phrases. Weyl relations, operator domains, unitary groups.

¹Research supported in part by NSF Grant MCS 77-02831 and sequel.

²We refer to [LP] as a general reference, containing in addition important applications.

©American Mathematical Society 1979
0002-9904/79/0000-0016/\$02.00

have nonzero defect indices [4a] which can be equal or unequal. If they are equal, it may or may not be possible to extend Q to a selfadjoint operator \tilde{Q} so that the one-parameter group $V(s) = e^{is\tilde{Q}}$ is part of a (global) Weyl system (U, V) , i.e. such that (1) holds. The triples (U, Q, H) which can be extended to global Weyl systems are called *extendable*.

In view of the negative examples alluded to above, it may be surprising that a certain canonical and *minimal extension* always exists.

THEOREM 1. *Let (U, Q, H) be an infinitesimal Weyl system. Then symmetric extensions Q_1 of Q exist such that*

$$(3) \quad \begin{aligned} &\mathcal{D}(Q_1) \text{ is invariant under } U(t) \text{ for all } t \in \mathbf{R}, \text{ and} \\ &Q_1 U(t)f = U(t)Q_1 f - tU(t)f \text{ for all } f \in \mathcal{D}(Q_1). \end{aligned}$$

There is a unique smallest symmetric extension Q_m satisfying (3), i.e. $Q_m \subseteq Q_1$ for all symmetric extensions Q_1 satisfying (3). The operator closure \bar{Q}_m satisfies (3) as well and is the unique smallest symmetric and closed extension of Q satisfying (3).

Here is an easy

COROLLARY 2. *If the system (U, Q, H) is extendable to a global Weyl system (U, V) with $V(s) = e^{is\tilde{Q}}$ then $Q_m \subseteq \tilde{Q}$.*

DEFINITION. *We say that the system (U, Q, H) has U -indices (p, q) if the minimal operator Q_m has defect indices (p, q) . It can readily be shown that the U -indices are equal whenever Q has equal defect indices, but in general the indices of Q may be infinite while those of Q_m are finite. Theorem 4 below is a converse to Corollary 2.*

THEOREM 3. *Let (U, Q, H) be an infinitesimal Weyl system with at least one finite U -index. Then there is a Hilbert space K containing H , and a global Weyl system (U, V) on H such that U reduces to U on H , H is semi-invariant [9] for V , and Q is contained in the infinitesimal generator for the contraction semigroup $s \rightarrow P_H V(s)|_H$.*

A structure theory for the associated contraction semigroups, due to P. Muhly, will appear in a joint article with the author [2b].

The next result concerns the relation (3) for extendable systems. Hence it is stated for operators in L^2 -spaces, and $U(t)$ is translation. Vanishing conditions of the Fourier transform on a set of measure zero clearly give rise to systems (3). The theorem is a partial converse to this statement.

THEOREM 4. *Let Q_1 be a symmetric and closed operator in $L^2(\mathbf{R})$ which is contained in the multiplication operator $\tilde{Q}h(x) = xh(x)$. Assume that $\mathcal{D}(Q_1)$ is translation invariant, and that $\mathcal{D}(\tilde{Q}Q_1)$ is a core for Q_1 . Then (3) holds and*

the closed set $\Lambda = \{\lambda \in \mathbf{R} | \hat{h}(\lambda) = 0 \text{ for } \forall h \in \mathcal{D}(Q_1)\}$ is of zero measure. Moreover $\mathcal{D}(Q_1) = \{h \in \mathcal{D}(\tilde{Q}) | \hat{h}(\lambda) = 0 \text{ for } \forall \lambda \in \Lambda\}$.

The core condition can be omitted in the following cases: (1) Λ is a Cantor set, or (2) Q_1 has finite defect. In general it can be slightly weakened, but not omitted. In case (2) the result generalizes to arbitrary multiplicity.

The proofs of results 1 through 3 involve general operator theory [1], including theorems of Phillips [5] and Naimark [3], while the proof of Theorem 4 is based on a function theoretic approach to the extension theory of [4a] made possible by the extendability assumption and the Stone-von Neumann Theorem. (The function theory is based on Wiener-Tauberian consideration.)

Generalizations to group representations [2a], [6], [7], [8], [11] and field theory [10] would be of interest. The possibility of removing the finiteness assumption in Theorem 3 is related via the proof to a well-known conjecture of Phillips [5]. We conjecture the conclusion of Theorem 4, also without the finiteness assumption on the U -index, but it seems hard to settle either of the conjectures. It appears difficult in the general case, for two different extensions to establish the existence of a wave operator similarity which commutes with U .

However the extendability question has an answer in full generality, i.e. no restriction on the indices. The proof uses arguments that naturally extend [4a].

THEOREM 5. *Let (U, Q, H) be an infinitesimal Weyl system, and let P_{\pm} denote the orthogonal projections onto the respective defect spaces \mathcal{D}_{\pm} for Q_m . The system is extendable if and only if there exists a partial isometry S of \mathcal{D}_{+} onto \mathcal{D}_{-} such that $P_{-}(2[S, U(t)] - i(I + S)U(t)(I + S))P_{+} = 0$ for all $t \in \mathbf{R}$. Here $[\cdot, \cdot]$ denotes the commutator bracket.*

If it is assumed in addition that the spectral measure dE_{λ} of U is absolutely continuous, then extendability is equivalent to the validity of the following identity

$$P_{-}(2[S, E_{\lambda}] - (I + S)D_{\lambda}(I + S))P_{+} = 0 \quad (\lambda \in \mathbf{R})$$

for some partial isometry S of \mathcal{D}_{+} onto \mathcal{D}_{-} . Here D_{λ} denotes the Radon-Nikodym derivative of dE_{λ} .

We finally point out that Theorem 5 has a complete generalization to the case when the W^* algebra generated by the spectral projections of U is replaced by an arbitrary noncommutative W^* -algebra, and when the map $U(t) \rightarrow it U(t)$ is replaced by a spatial derivation which is implemented by a symmetric nonself-adjoint operator. This answers a question raised in a recent article of the author.

The author has benefitted greatly from discussions with Professors P. Chernoff, W. Helton, G. Johnson, R. T. Moore, P. Muhly, and R. S. Phillips. I am pleased to record my gratitude.

REFERENCES

1. C. Foias and L. Gehér, *Über die Weylsche Vertauschungsrelationen*, Acta Sci. Math. 24 (1963), 97-102.

2. P. E. T. Jørgensen (a) (with R. T. Moore), *Commutation relations for operators, semigroups, and resolvents in mathematical physics and group representations* (preprint); (b) (with P. S. Muhly), *Selfadjoint extensions satisfying the Weyl operator commutation relations* (in preparation).
 3. M. A. Naïmark, *On commuting unitary operators in spaces with indefinite metric*, *Acta Sci. Math.* **24** (1963), 177–189.
 4. J. von Neumann, (a) *Allgemeine Eigenwerttheorie Hermitescher Funktionaloperatoren*, *Math. Ann.* **102** (1929–30) 49–131; (b) *Die Eindentigkeit der Schrödingerschen Operatoren*, *Math. Ann.* **104** (1931), 570–578.
 5. R. S. Phillips, *On dissipative operators*, *Lecture Series in Differential Equations*, vol. II, A. K. Aziz (Ed.), Van Nostrand Mathematical Studies no. 19, Van Nostrand, New York, 1969.
 6. N. S. Poulsen, *On C^∞ vectors and intertwining bilinear forms for representations of Lie groups*, *J. Functional Analysis*, (1972), 87–120.
 7. R. T. Powers, *On selfadjoint algebras of unbounded operators*, *Trans. Amer. Math. Soc.* **187** (1974), 261–293.
 8. R. R. Rao, *Unitary representations defined by boundary conditions—the case of $sl(2, \mathbb{R})$* , *Acta Math.* **139** (1977), 185–216.
 9. D. Sarason, *On spectral sets having connected complement*, *Acta Sci. Math.* **26** (1965), 289–299.
 10. I. E. Segal, *Nonlinear functions of weak processes. II*, *J. Functional Analysis* **6** (1970), 29–75.
 11. I. M. Singer, *Lie algebras of unbounded operators*, Thesis, Univ. of Chicago, 1950.
 12. M. H. Stone, *Linear transformations in Hilbert space. III. Operational methods and group theory*, *Proc. Nat. Acad. Sci. U.S.A.* **16** (1930), 172–175.
- [LP] P. D. Lax and R. S. Phillips, *Scattering theory*, Academic Press, New York, 1967.

DEPARTMENT OF MATHEMATICS, STANFORD UNIVERSITY, STANFORD,
CALIFORNIA 94305