

# ON THE SOLUTION OF ORDINARY DIFFERENTIAL EQUATIONS OF THE FIRST ORDER INVARIANT UNDER CONTACT TRANSFORMATIONS

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In the classical Lie theory it is shown how to construct a differential equation invariant under a given group, and how to solve an equation when a group leaving the equation invariant is known. However, little is said about the problem of determining the group for a given differential equation, which is by far the most interesting problem.

In the present paper, necessary and sufficient conditions for the existence of an infinitesimal contact transformation leaving a given equation invariant are determined along with the general form of the characteristic function of the group. It will also be shown how to reduce, by a proper change of variables, the infinitesimal contact transformation to a point transformation. This enables one to solve the transformed differential equation by Lie's methods. Passing back to the original variables, a new differential equation is obtained which combined with the original equation gives its solution in parametric form.

Let

$$Bf = \xi \frac{\partial f}{\partial x} + \eta \frac{\partial f}{\partial y} + \pi \frac{\partial f}{\partial p}$$

be the symbol of the infinitesimal contact transformation leaving invariant the differential equation  $u = F(v)$ , with  $u = u(x, y, p)$ ,  $v = v(x, y, p)$ ,  $p = dy/dx$ , and  $F$  such that the equation  $G(x, y, p) = u - F(v) = 0$  satisfies the various conditions for the existence of solutions (but otherwise arbitrary). Throughout this paper we shall assume that:

(A) Both  $u$  and  $v$  have first derivatives with respect to  $x$ ,  $y$  and  $p$ , at least in some region  $R$  of the  $(x, y, p)$ -space.

(B) The Jacobians

$$J_1 = \frac{\partial(u, v)}{\partial(y, p)}, \quad J_2 = \frac{\partial(u, v)}{\partial(p, x)}, \quad J_3 = \frac{\partial(u, v)}{\partial(x, y)}$$

have in  $R$  derivatives of the first and second orders, while  $J_1$  and  $J_2$  have also derivatives of the third order with respect to  $x$ ,  $y$  and  $p$ ,

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as are involved in the discussion.

(C) The functions  $u$  and  $v$  are not in involution, that is,

$$[uv] = \begin{vmatrix} u_p & u_x + pu_y \\ v_p & v_x + pv_y \end{vmatrix} = J_2 - pJ_1 \neq 0.$$

Since  $u$  and  $v$  are to be invariants under  $Bf$  they will satisfy the partial differential equations

$$\begin{aligned} \xi \frac{\partial u}{\partial x} + \eta \frac{\partial u}{\partial y} + \pi \frac{\partial u}{\partial p} &= 0, \\ \xi \frac{\partial v}{\partial x} + \eta \frac{\partial v}{\partial y} + \pi \frac{\partial v}{\partial p} &= 0, \end{aligned}$$

from which we obtain

$$\frac{\xi}{\partial(u, v)/\partial(y, p)} = \frac{\eta}{\partial(u, v)/\partial(p, x)} = \frac{\pi}{\partial(u, v)/\partial(x, y)} = \sigma,$$

$\sigma = \sigma(x, y, p)$  being the common ratio. This can be written

$$(1) \quad \xi = \sigma J_1, \quad \eta = \sigma J_2, \quad \pi = \sigma J_3.$$

If  $W$  is the so-called characteristic function of the infinitesimal contact transformation, we have also

$$(2) \quad W = p\xi - \eta = \sigma(pJ_1 - J_2).$$

Now to find  $\sigma$  we recall that<sup>1</sup>

$$(3) \quad \xi = \frac{\partial W}{\partial p}, \quad \pi = -\frac{\partial W}{\partial x} - p \frac{\partial W}{\partial y}.$$

As a consequence of (1), (2) and (3) we obtain the system of equations

$$\begin{aligned} (pJ_1 - J_2) \frac{\partial \sigma}{\partial p} + \left( p \frac{\partial J_1}{\partial p} - \frac{\partial J_2}{\partial p} \right) \sigma &= 0, \\ (4) \quad (pJ_1 - J_2) \frac{\partial \sigma}{\partial x} + p(pJ_1 - J_2) \frac{\partial \sigma}{\partial y} \\ &+ \left[ \left( p \frac{\partial J_1}{\partial x} - \frac{\partial J_2}{\partial x} \right) + p \left( p \frac{\partial J_1}{\partial y} - \frac{\partial J_2}{\partial y} \right) + J_3 \right] \sigma = 0. \end{aligned}$$

This system may be written in the homogeneous form

<sup>1</sup> See Cohen, *An introduction to the Lie theory of one-parameter groups*, p. 186.

$$(5) \quad \begin{aligned} A_1 f &= \frac{\partial f}{\partial p} + M_1 \frac{\partial f}{\partial \sigma} = 0, \\ A_2 f &= \frac{\partial f}{\partial x} + p \frac{\partial f}{\partial y} + M_2 \frac{\partial f}{\partial \sigma} = 0, \end{aligned}$$

in which

$$(6) \quad M_1 = -\sigma \frac{p(\partial J_1 / \partial p) - (\partial J_2 / \partial p)}{pJ_1 - J_2},$$

$$(7) \quad M_2 = -\sigma \frac{(p(\partial J_1 / \partial x) - \partial J_2 / \partial x) + p(p(\partial J_1 / \partial y) - \partial J_2 / \partial y) + J_3}{pJ_1 - J_2}.$$

Adjoining to the system (5) the equations

$$(8) \quad A_3 f = (A_1 A_2) f = \frac{\partial f}{\partial y} + (A_1 M_2 - A_2 M_1) \frac{\partial f}{\partial \sigma} = \frac{\partial f}{\partial y} + M_3 \frac{\partial f}{\partial \sigma} = 0,$$

$$(9) \quad A_4 f = (A_1 A_3) f = (A_1 M_3 - A_3 M_1) \frac{\partial f}{\partial \sigma} = 0,$$

$$(10) \quad A_5 f = (A_2 A_3) f = (A_2 M_3 - A_3 M_2) \frac{\partial f}{\partial \sigma} = 0,$$

we see that the equations

$$(11) \quad A_1 M_3 - A_3 M_1 = 0, \quad A_2 M_3 - A_3 M_2 = 0,$$

are necessary and sufficient conditions in order that the system (4) have a solution. The system (5)–(8) implies the Jacobian complete system

$$(12) \quad \begin{aligned} K_1 f &= \frac{\partial f}{\partial x} + (M_2 - pM_3) \frac{\partial f}{\partial \sigma} = 0, \\ K_2 f &= \frac{\partial f}{\partial y} + M_3 \frac{\partial f}{\partial \sigma} = 0, \\ K_3 f &= \frac{\partial f}{\partial p} + M_1 \frac{\partial f}{\partial \sigma} = 0. \end{aligned}$$

Either we may solve (12) or the equivalent total differential equation

$$(13) \quad (M_2 - pM_3)dx + M_3 dy + M_1 dp - d\sigma = 0.$$

If  $f = \psi(x, y, p, \sigma)$  is the solution of (12), then

$$(14) \quad \psi(x, y, p, \sigma) = c$$

will be the solution of (13), and conversely. Equation (14) determines  $\sigma$  in terms of  $x, y,$  and  $p$ . Since  $\sigma$  enters as a factor in  $M_1$  and  $M_2$ , it is also a factor of  $M_3$ .<sup>2</sup> Hence, equation (13) can be written

$$d\sigma/\sigma = d\omega(x, y, p),$$

and so  $\sigma$  has the form

$$(15) \quad \sigma = ke^{\omega(x,y,p)}.$$

Several special formulas for  $\sigma$  may be found. For instance, if

$$M_1 = \phi_1(p)\sigma, \quad M_2 = \phi_2(x)\sigma,$$

then  $M_3=0$ , and equation (13) reduces to

$$\phi_2(x)\sigma dx + \phi_1(p)\sigma dp - d\sigma = 0,$$

from which we obtain

$$\sigma = k \exp \left( \int \phi_1(p)dp + \int \phi_2(x)dx \right).$$

Therefore, the characteristic function takes the form

$$(16) \quad W = k(pJ_1 - J_2) \exp \left( \int \phi_1(p)dp + \int \phi_2(x)dx \right)$$

by virtue of (2). This special case will be of use in some examples to be considered later.

We summarize our results in the following theorem.

**THEOREM.** *The characteristic function  $W$  of the infinitesimal contact transformation leaving invariant a given differential equation  $u = F(v)$  can be found by the formula*

$$W = k(pJ_1 - J_2)e^{\omega(x,y,p)}$$

if, and only if, the equations

$$A_1M_3 - A_3M_1 = 0, \quad A_2M_3 - A_3M_2 = 0$$

are both satisfied for all values of  $x, y$  and  $p$ .

Now, to solve the differential equation  $u = F(v)$  invariant under the known group

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<sup>2</sup> If  $M_1 = \sigma N_1, M_2 = \sigma N_2$ , then  $M_3 = A_1M_2 - A_2M_1 = \sigma(\partial N_2/\partial p - \partial N_1/\partial x - p\partial N_1/\partial y)$ . This relation, together with (11), are the conditions in order that (13) be an exact differential when divided by  $\sigma$ .

$$(17) \quad Bf = W_p \frac{\partial f}{\partial x} + (pW_p - W) \frac{\partial f}{\partial y} - (W_x + pW_y) \frac{\partial f}{\partial p},$$

we consider two cases:

(A) Both  $\xi = W_p$  and  $\eta = pW_p - W$  are free of  $p$ . This case occurs when  $W$  is linear in  $p$ . Then  $Bf$  represents an extended point transformation and the equation may be solved by introducing canonical variables.

(B) Either  $\xi$  or  $\eta$ , or both, contain  $p$ . Then  $Bf$  represents a general contact transformation.

In this case we may show that by a suitable change of variables the transformation reduces to a point transformation.<sup>3</sup> To this aim, let us define a finite contact transformation

$$(18) \quad X = X(x, y, p), \quad Y = Y(x, y, p), \quad P = P(x, y, p)$$

in the following manner:  $X = u$ ,  $Y \neq X$  in involution with  $X$ , that is, such that  $[XY] = 0$ , or

$$(19) \quad X_p \frac{\partial Y}{\partial x} + pX_p \frac{\partial Y}{\partial y} - (X_x + pX_y) \frac{\partial Y}{\partial p} = 0,$$

and  $P$  by the equation  $P = Y_p/X_p$ . The symbol for the transformed group will be

$$(20) \quad \overline{B}f = \xi \frac{\partial f}{\partial X} + \eta \frac{\partial f}{\partial Y} + \pi \frac{\partial f}{\partial P} = BX \frac{\partial f}{\partial X} + BY \frac{\partial f}{\partial Y} + BP \frac{\partial f}{\partial P}.$$

But  $\xi = BX = Bu = 0$  since  $u$  is invariant under  $Bf$ . Since  $\xi = \overline{W}_P$  this implies that  $\overline{W}$  is free of  $P$ . Also, we find that  $\eta$  does not contain  $P$  because  $\eta = P\overline{W}_P - \overline{W} = -\overline{W}$ . Hence,  $\overline{B}f$  is an extension of the point transformation group

$$(21) \quad Uf = -\overline{W}(X, Y) \frac{\partial f}{\partial Y}.$$

This group can be reduced further by introducing the canonical variables

$$X^* = X, \quad Y^* = -\int \frac{\partial Y}{\overline{W}(X, Y)}.$$

Then the symbol of the infinitesimal transformation assumes the

<sup>3</sup> Cohen, loc. cit. p. 195, proves that the contact transformation reduces to a point transformation by assuming the corresponding differential equation solvable for  $p$  in the form  $p = \omega(x, y)$ .

simplest form

$$U^*f = \frac{\partial f}{\partial Y^*}.$$

The equation  $u = F(v)$  when written with the variables  $X, Y, P$  takes the form

$$(22) \quad \phi(X, Y, P) = 0.$$

This is also a differential equation, that is,  $P = dY/dX$ , since the relation  $dY - PdX = \lambda(dy - p dx)$  which holds for any contact transformation implies  $dY - PdX = 0$  whenever  $dy - p dx = 0$ . Since (22) will be invariant under (21) we are in position to solve (22), either directly or by introducing the canonical variables  $X^*, Y^*$  [this last step reduces the equation to the form  $dY^*/dX^* = G(X^*)$ ]. Let

$$(23) \quad \psi(X, Y, c) = 0$$

be the solution of (22). Passing back to the original variables we get a second differential equation

$$\Psi(x, y, p, c) = 0$$

which together with  $u = F(v)$  determines the integral curves of the latter in terms of the parameter  $p$ .

*Examples.* I. Consider the differential equation

$$(24) \quad p + y/p = F(x + 2p).$$

Here  $u = p + y/p$ ,  $v = x + 2p$ . Hence, it follows that  $J_1 = 2/p$ ,  $J_2 = 1 - y/p^2$ ,  $J_3 = -1/p$ ,  $pJ_1 - J_2 = 1 + y/p^2$ ,  $M_1 = 2\sigma/p$ ,  $M_2 = M_3 = 0$ .

Formula (16) can be applied with  $\phi_1(p) = 2/p$ ,  $\phi_2(x) = 0$ . Therefore, the characteristic function of the group is

$$W = k(1 + y/p^2)p^2 = k(p^2 + y).$$

Since a constant factor is irrelevant, we see that equation (24) is invariant under the infinitesimal contact transformation

$$Bf = 2p \frac{\partial f}{\partial x} + (p^2 - y) \frac{\partial f}{\partial y} - p \frac{\partial f}{\partial p}.$$

By taking  $X = v = x + 2p$  equation (19) reduces to

$$(25) \quad 2 \frac{\partial Y}{\partial x} + 2p \frac{\partial Y}{\partial y} - \frac{\partial Y}{\partial p} = 0.$$

The corresponding system of ordinary differential equations is

$$\frac{dx}{2} = \frac{dy}{2p} = \frac{dp}{-1} = \frac{dY}{0},$$

from which we obtain  $Y = p^2 + y$  as a particular integral of (25). Finally we have  $P = 2p/2 = p$ . Introducing the new variables in (24) we get  $dY/Y = dX/F(X)$ . Hence, we have

$$Y = ce^{G(X)}, \quad G(X) = \int \frac{dX}{F(X)}.$$

Passing back to the variables  $x, y, p$  we obtain

$$(26) \quad p^2 + y - ce^{G(x+2p)} = 0.$$

The system (24)–(26) furnishes the solution of the equation (24).

For instance, if  $F(x+2p) = \tan(x+2p)$ , equations (24) and (26) are respectively

$$p + y/p = \tan(x + 2p), \quad p^2 + y = c \sin(x + 2p).$$

Solving for  $x$  and  $y$  we find

$$\begin{aligned} x &= -2t + \arccos(t/c), \\ y &= -t^2 \pm (c^2 - t^2)^{1/2}, \end{aligned}$$

which are the parametric equations of the solution, where  $t = p$  is the parameter.

II. To apply the method to find the group leaving invariant some familiar types of ordinary differential equations, let us consider first the homogeneous equation

$$p = F(y/x).$$

We have  $u = p, v = y/x, J_1 = -1/x, J_2 = -y/x^2, J_3 = 0, pJ_1 - J_2 = (y - px)/x^2, M_1 = 0, M_2 = 2\sigma/x, M_3 = 0$ . By using formula (16) with  $\phi_1(p) = 0, \phi_2(x) = 2/x$ , we get (taking  $k = -1$ )

$$W = px - y.$$

Since  $W$  is linear in  $p$  we obtain the point transformation with symbol

$$Uf = x \frac{\partial f}{\partial x} + y \frac{\partial f}{\partial y},$$

which corresponds to the so-called homotetic transformation.

For the linear equation  $p + P(x)y = F(x)$  we have  $u = p + Py, v = x, J_1 = 0, J_2 = 1, J_3 = -P, pJ_1 - J_2 = -1, M_1 = 0, M_2 = -\sigma P, M_3 = 0$ .

By putting  $\phi_1(\rho) = 0$ ,  $\phi_2(x) = -P$ ,  $ke^e = 1$  in formula (16) we obtain

$$W = - \exp \left( - \int P dx \right).$$

Hence the symbol for the group has the form

$$Uf = \exp \left( - \int P dx \right) \frac{\partial f}{\partial y}.$$

III. Finally, we shall give a short table of some general types with the corresponding characteristic functions.<sup>4</sup>

Differential Equations	Characteristic Functions
$y = px + F[x\phi(\rho)]$	$kx\phi(\rho)$
$y = px + \rho F[y\phi(\rho)]$	$kyp\phi(\rho)$
$y + \phi(\rho) = \rho F \left[ x + \int \phi'(\rho) d\rho / \rho \right]$	$k[y + \phi(\rho)]$
$e^x \phi(x + y + \rho) = F[e^x(\rho + 1)]$	$k\phi(x + y + \rho)$
$\frac{y + x\phi(\rho)}{\rho + \phi(\rho)} = xF \left[ \log x + \int \frac{\phi'(\rho) d\rho}{\rho + \phi(\rho)} \right]$	$k[y + x\phi(\rho)].$

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