ON LAMÉ'S SIX EQUATIONS CONNECTED WITH TRIPLY ORTHOGONAL SYSTEMS OF SURFACES.

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Lamé* has shown for a triply orthogonal system of surfaces given by the parameters ρ , ρ_1 , ρ_2 that if the square of the element of length is given by $ds^2 = H^2 d\rho + H_1^2 d\rho_1^2 + H_2^2 d\rho_2^2$, where H, H_1 , H_2 are certain functions of ρ , ρ_1 , ρ_2 , then H, H_1 , H_2 must satisfy the following system of equations:

$$\frac{\partial^2 H}{\partial \rho_1 \partial \rho_2} = \frac{1}{H_1} \frac{\partial H}{\partial \rho_1} \frac{\partial H_1}{\partial \rho_2} + \frac{1}{H_2} \frac{\partial H}{\partial \rho_2} \frac{\partial H_2}{\partial \rho_1}$$
(1)

and two others of the same type; (2), (3)

$$\frac{\partial}{\partial \rho_{1}} \left(\frac{1}{H_{1}} \frac{\partial H}{\partial \rho_{1}} \right) + \frac{\partial}{\partial \rho} \left(\frac{1}{H} \frac{\partial H_{1}}{\partial \rho} \right) + \frac{1}{H_{2}^{2}} \frac{\partial H}{\partial \rho_{2}} \frac{\partial H_{1}}{\partial \rho_{2}} = 0 \quad (4)$$

with two others of this type.

(5), (6)

Also if V is a function of x, y, z for which

$$\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0,$$

he has shown that

$$\frac{\partial}{\partial \rho} \left(\frac{H_1 H_2}{H} \frac{\partial V}{\partial \rho} \right) + \frac{\partial}{\partial \rho_1} \left(\frac{H_2 H}{H_1} \frac{\partial V}{\partial \rho_1} \right) + \frac{\partial}{\partial \rho_2} \left(\frac{H H_1}{H_2} \frac{\partial V}{\partial \rho_2} \right) = 0.$$

If the system of coordinates ρ , ρ_1 , ρ_2 is isothermal, this equation must be satisfied by $V=\rho$, or by $V=\rho_1$, or by $V=\rho_2$. Hence $H_1H_2/H=Q^2$, where Q is a function of ρ_1 and ρ_2 only. Similarly $H_2H/H_1=Q_1^2$, and $HH_1/H_2=Q_2^2$, where Q_i is a function not involving the variable ρ_i . Hence $H=Q_1Q_2$, $H_1=Q_2Q$, $H_2=QQ_1$. The six equations given above transform into six in the variables Q. Lamé† gives a solution of

^{*} Leçons sur les coordonées curvilignes (1859), pp. 76, 78.

[†] Loc. cit., p. 99.

these equations by first finding a solution of equations (1), (2), (3) and using this to solve the remainder. He makes the statement that his solution of the first three equations is the most general possible *; this statement is obviously inaccurate and it seems of interest to give a complete solution of the equations. The solution of the first three equations, or rather a comparison of two different solutions, leads to a curious result in the theory of elimination.

Equation (1) becomes

$$Q \frac{\partial Q_2}{\partial \rho_1} \frac{\partial Q_1}{\partial \rho_2} = Q_1 \frac{\partial Q_2}{\partial \rho_1} \frac{\partial Q}{\partial \rho_2} + Q_2 \frac{\partial Q_1}{\partial \rho_2} \frac{\partial Q}{\partial \rho_1}, \tag{7}$$

with similar expressions for (2) and (3). Multiply (7) by $\partial Q_1/\partial \rho$, and the transformed expression for (3) by $\partial Q_1/\partial \rho_2$, and add. The result is

$$Q_{1} \left[\frac{\partial Q_{2}}{\partial \rho} \frac{\partial Q}{\partial \rho_{1}} \frac{\partial Q_{1}}{\partial \rho_{2}} + \frac{\partial Q_{1}}{\partial \rho} \frac{\partial Q_{2}}{\partial \rho_{1}} \frac{\partial Q}{\partial \rho_{2}} \right] = 0.$$

Hence unless all the Q's vanish identically

$$\frac{\partial Q_2}{\partial \rho} \frac{\partial Q}{\partial \rho_1} \frac{\partial Q_1}{\partial \rho_2} + \frac{\partial Q_1}{\partial \rho} \frac{\partial Q_2}{\partial \rho_1} \frac{\partial Q}{\partial \rho_2} = 0. \tag{8}$$

Assume that none of the derivatives in (8) vanish identically and write

$$K = -\frac{\partial\,Q}{\partial\rho_1}\Big/\frac{\partial\,Q}{\partial\rho_2}, \quad K_1 = -\frac{\partial\,Q_1}{\partial\rho_2}\Big/\frac{\partial\,Q_1}{\partial\rho}\,, \quad K_2 = -\frac{\partial\,Q_2}{\partial\rho}\Big/\frac{\partial\,Q_2}{\partial\rho_1}.$$

Equation (8) becomes $KK_1K_2 = 1$, where K_i is a function not involving ρ_i . By taking logarithms and differentiating with respect to two of the variables ρ , it is easy to prove that the most general values for the K's are

$$K = \frac{a_2}{a_1}, \quad K_1 = \frac{a}{a_2}, \quad K_2 = \frac{a_1}{a},$$

where a is a function of ρ only and similarly a_1 and a_2 are functions of ρ_1 and ρ_2 alone respectively.

^{*} Loc. cit., p. 100, line 23.

If we put $\int ad\rho = \sigma$, with similar values for σ_1 and σ_2 , we deduce that Q is a function of $\sigma - \sigma_2$ only, Q_1 a function of $\sigma_2 - \sigma$ only, and Q_2 a function of $\sigma - \sigma_1$ only. Substitute these values in (7) and it readily reduces to

$$\frac{Q}{Q'} + \frac{Q_1}{Q_1'} + \frac{Q_2}{Q_2'} = 0.$$

Hence

$$\frac{Q}{Q'} = n(\sigma_1 - \sigma_2), \quad \frac{Q_1}{Q_1'} = n(\sigma_2 - \sigma), \quad \frac{Q_2}{Q_2'} = n(\sigma - \sigma_1),$$

where n is a constant, and therefore

$$Q = c(\sigma_1 - \sigma_2)^{1/n}, \quad Q_1 = c_1(\sigma_2 - \sigma)^{1/n}, \quad Q_2 = c_2(\sigma - \sigma_1)^{1/n}, \quad (A)$$

where c, c_1, c_2 are constants. This is the most general solution, provided none of the derivatives in (8) vanish. If, however, for example, $\partial Q_2/\partial \rho_1=0$, then either $\partial Q_2/\partial \rho=0$, or $\partial Q/\partial \rho_1=0$, or $\partial Q_1/\partial \rho_2=0$. Equation (7), however, shows that $\partial Q_2/\partial \rho_1=0$ implies either $\partial Q/\partial \rho_1=0$ or $\partial Q_1/\partial \rho_2=0$. $\partial Q_2/\partial \rho_1=0$, $\partial Q_1/\partial \rho_2=0$ lead to the solution

$$Q_1 = f(\rho), \quad Q_2 = cf(\rho), \quad Q = \phi(\rho_1, \rho_2),$$
 (B)

where f and ϕ are arbitrary functions of their arguments and c is a constant. $\partial Q_2/\partial \rho_1=0$, $\partial Q/\partial \rho_1=0$ lead to the solution

$$Q_2 = f(\rho), \quad Q = \phi(\rho_2), \quad Q_1 = F(Q, Q_2), \qquad (C)$$

where F is homogeneous and of unit degree in Q_1 , Q_2 . These types (A), (B), (C) are the only three types of solution.

We now proceed to the integration of the equations in a different manner. Write $\log Q = \lambda$, $\log Q_1 = \lambda_1$, $\log Q_2 = \lambda_2$. Equations (1), (2), (3) become

$$\frac{\partial \lambda_2}{\partial \rho_1} \frac{\partial \lambda_1}{\partial \rho_2} = \frac{\partial \lambda_2}{\partial \rho_1} \frac{\partial \lambda}{\partial \rho_2} + \frac{\partial \lambda}{\partial \rho_1} \frac{\partial \lambda_1}{\partial \rho_2}$$
(9)

and two similar equations.

Write $\lambda_1 - \lambda_2 = \omega$, $\lambda_2 - \lambda = \omega_1$, $\lambda - \lambda_1 = \omega_2$ and (9) becomes

$$\frac{\partial \omega_1}{\partial \rho_1} \frac{\partial \omega_2}{\partial \rho_2} - \frac{\partial \omega_1}{\partial \rho_2} \frac{\partial \omega_2}{\partial \rho_1} = 0,$$

or

$$J\begin{pmatrix} \omega_1 & \omega_2 \\ \rho_1 & \rho_2 \end{pmatrix} = 0.$$

Hence the equation implies the existence of a relation, $f(\omega_1, \omega_2, \rho) = 0$. Exactly similarly there are relations

$$f_1(\omega_2, \omega, \rho_1) = 0, \quad f_2(\omega, \omega_1, \rho_2) = 0.$$

From these equations, together with $\omega + \omega_1 + \omega_2 = 0$, it is easy to deduce that either ρ , ρ_1 , ρ_2 do not any of them occur explicitly in f, f_1 , f_2 or if, for example, ρ occurs in f, it is easy to prove that either ρ_1 is absent from f_1 , or ρ_2 from f_2 . The latter case is thus reduced to the former, for if one relation of the type $f(\omega, \omega_1) = 0$ exist, then two others of that type also exist in virtue of the relation $\omega + \omega_1 + \omega_2 = 0$.

Substituting for the ω 's in terms of the Q's, we immediately deduce that the solution is equivalent to the statement that a homogeneous relation $F(Q, Q_1, Q_2) = 0$ exists among the Q's.

Combining the two solutions we have the following theorem connected with the theory of elimination:

Let $F(Q, Q_1, Q_2) = 0$ be any homogeneous relation. It is possible to express Q as a function of two variables ρ_1 and ρ_2 , Q_1 as a function of two variables ρ_2 and ρ , and Q_2 as a function of the variables ρ and ρ_1 , in two cases only:

(A) If F is of the form $aQ^n + a_1Q_1^{n} + a_2Q_2^{n}$, where a, a_1, a_2 are constants.*

(C) If F is general, and e. g. Q is a function of ρ_1 only, and Q_1 a function of ρ only.

It is not difficult to complete the solution of the equations $(1), \dots, (6)$. It may readily be shown that for case (A) (4), (5), and (6) are not satisfied unless $n = \frac{1}{2}$, and then

$$\sigma = A \varphi \left(\frac{\rho + a}{\sqrt{c}}, g_2, g_3 \right) + B,$$

$$\sigma_{\scriptscriptstyle 1} = A \, \varphi \left(\frac{\rho_{\scriptscriptstyle 1} \, + \, a_{\scriptscriptstyle 1}}{\sqrt{c_{\scriptscriptstyle 1}}}, \, g_{\scriptscriptstyle 2}, \, g_{\scriptscriptstyle 3} \right) + \, B, \label{eq:sigma_1}$$

^{*} This includes (B).

$$\sigma_{2} = A \varphi \left(\frac{\rho_{2} + a_{2}}{\sqrt{c_{2}}}, g_{2}, g_{3} \right) + B,$$

where the a's, the g's, A and B are constants.

For case (c) the complete solution is

$$\begin{split} Q_1 &= a \ \text{cosec} \ (b\rho_2 + c), \quad Q_2 = a' \ \text{cosec} \ (b'\rho_1 + c'), \\ Q &= A \left[\text{cosec}^2 \ (b\rho_2 + c) - \text{cosec}^2 \ (b'\rho_1 + c') \right]^{\frac{1}{2}}, \end{split}$$

where A, a, a', b, b', c, c' are constants such that

$$a^2b'^2+a'^2b^2=0.$$
 In case (B)
$$Q_1=\frac{1}{a\rho+b},\quad Q_2=\frac{k}{a\rho+b},$$

where $a,\,b,\,k$ are constants, and Q is a function of ρ_1 and ρ_2 which satisfies the equation

$$\left(k^2 \frac{\partial^2}{\partial \rho_2^2} + \frac{\partial^2}{\partial \rho_1^2}\right) \log Q + a^2 Q^2 = 0.$$

Of these three types of solution, the first is the same as that given by Lamé.* He gives it in different form, and his method of obtaining it is different. He falls into the error of imagining that the most general solution of the first three of his equations corresponds to the case of $n = \frac{1}{2}$, and it happens that this error is largely corrected because the second three equations require this limitation in case (A); case (B), however, escapes his notice.

The surfaces corresponding to the three solutions are readily obtained. (A) gives a system of confocal quadrics, and (C) a system of confocal spheroids with their axial planes. (B) gives a system of concentric spheres, and the conical surfaces obtained by joining the common centre to any set of isothermal lines on one of them.

BRYN MAWR, October, 1905.

^{*} Loc. cit., p. 104.