# Branching laws for square integrable representations 

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#### Abstract

In this note, we study square integrable representations of a real reductive Lie group with admissible restriction to some reductive subgroup. We give a simple condition which insures admissibility of the restriction, and which allows to compute the branching numbers in a simple explicit manner by means of partition functions, generalizing the multiplicity formulas due to Kostant-Heckman and Hecht-Schmid. We consider also the semi-classical analogue of these results for coadjoint orbits.


Key words: Discrete series; square integrable representations; branching laws; multiplicity formulas.

1. Introduction. Let $G$ be a connected reductive Lie group, and $H$ a closed connected reductive subgroup of $G$. Recall that a square integrable representation (also called a discrete series) of $G$ is an irreducible unitary representation of $G$ which can be realized as a subspace of the left regular representation in $L^{2}(G)$. We consider a square integrable representation $\pi$ of $G$ which has an admissible restriction to $H$. This means that we may express the restriction $\pi_{\left.\right|_{H}}$ as a Hilbert sum

$$
\pi_{\left.\right|_{H}}=\sum_{i \in \mathbf{N}} n_{i} \sigma_{i}
$$

(here, the $\sigma_{i}$ are classes of irreducible unitary representations of $H$ ) with $0 \leq n_{i}<\infty$. The $\sigma_{i}$ for which $n_{i}>0$ are known to be square integrable [11], 8.7. In this note we present some results on the multiplicity (also called branching number)

$$
m\left(\pi, \sigma_{i}\right):=n_{i}
$$

of a square integrable representation $\sigma_{i}$ of $H$ in $\pi$. Proofs will appear elsewhere.

A theorem of Harish-Chandra insures that a connected reductive Lie group admits square integrable representations if and only if it has a compact Cartan subgroup. Henceforth, we assume that $G$ and $H$ have a compact Cartan subgroup. We fix a compact Cartan subgroup $T$ of $G$, and a maximal compact subgroup $K$ of $G$ containing $T$ so that:

[^0]$L:=H \cap K$ is a maximal compact subgroup of $H$,
$U:=H \cap T$ is a compact Cartan subgroup of $H$. Let $W_{K}$ denote the Weyl group of $K$. We use a similar notation for other compact connected groups. We denote by $V^{*}$ the dual of a vector space $V$. We denote by $\mathfrak{g}_{\mathbf{R}}$ the Lie algebra of $G$ and $\mathfrak{g}$ the complexified Lie algebra. We use the same system of notations for any Lie group. We denote by $\Phi(\mathfrak{g}, \mathfrak{t}) \subset$ $i \mathrm{t}_{\mathbf{R}}^{*} \subset \mathrm{t}^{*}$ the set of roots. If $\Psi \subset \Phi(\mathfrak{g}, \mathrm{t})$ is a positive system of roots, we consider the corresponding sets of positive compact roots $\Psi_{c}:=\Psi \cap \Phi(\mathfrak{f}, \mathrm{t})$, and of noncompact roots $\Psi_{n}:=\Psi \backslash \Psi_{c}$. We denote as usual by $\rho, \rho_{c}$ and $\rho_{n}$ the corresponding half-sums.

Harish-Chandra parameterized the set of equivalence classes of square integrable representations by means of the set of $\mathfrak{g}$-regular $\lambda \in i \mathrm{t}_{\mathbf{R}}^{*}$ such that $e^{\lambda+\rho}$ is a character of $T$. We call such a $\lambda$ a $G$-Harish-Chandra parameter, and denote by $\pi_{\lambda}$ the corresponding equivalent class of square integrable representation. Moreover, $\pi_{\lambda}=\pi_{\lambda^{\prime}}$ if and only if $W_{K} \lambda^{\prime}=W_{K} \lambda$.

In the special case of $K$, we denote by $\tau_{\xi}$ the irreducible representation of $K$ associated with the $K$-Harish-Chandra parameter $\xi \in i \mathrm{t}_{\mathbf{R}}^{*}$. It is the irreducible representation of $K$ with infinitesimal character defined by $\xi$.

Similarly, we denote by $\sigma_{\mu}$ the square integrable representation of $H$ associated with a $H$-HarishChandra parameter $\mu \in i \mathfrak{u}_{\mathbf{R}}^{*}$, and by $\kappa_{\nu}$ the irreducible representation of $L$ associated with the $L$-Harish-Chandra parameter $\nu \in i \mathfrak{u}_{\mathbf{R}}^{*}$.

Let $\lambda$ be a $G$-Harish-Chandra parameter. This defines a positive system of roots

$$
\Psi^{\lambda}=\left\{\alpha \in \Phi(\mathfrak{g}, \mathfrak{t}): \lambda\left(h_{\alpha}\right)>0\right\}
$$

where $h_{\alpha} \in \mathrm{t}$ is the coroot. We recall the formula of Blattner (proven by Hecht-Schmid [4]) for the multiplicity $m\left(\pi_{\lambda}, \tau_{\xi}\right)$ of $\tau_{\xi}$ in $\pi_{\lambda}$. We denote by $p_{\Psi_{n}^{\lambda}}$ the partition function associated with $\Psi_{n}^{\lambda}$. Assuming that $\xi$ is in the positive Weyl Chamber for $\Psi_{c}^{\lambda}$, we have:

$$
m\left(\pi_{\lambda}, \tau_{\xi}\right)=\sum_{w \in W_{K}} \epsilon(w) p_{\Psi_{n}^{\lambda}}\left(w \xi-\lambda-\rho_{n}\right)
$$

There are similar formulas in terms of partition functions, due to Heckman and Kostant [5], for the restriction of an irreducible representation of $K$ to $L$.

We prove that for a square integrable representation $\pi_{\lambda}$, its restriction $\pi_{\lambda_{\mid H}}$ to $H$ is admissible if and only if its restriction $\pi_{\lambda_{L}}$ to $L$ is admissible. This allows a strategy to compute multiplicities in $\pi_{\lambda_{I_{H}}}$ in case of admissibility. First, we compute multiplicities in $\pi_{\lambda_{L}}$, combining the formulas of Blattner and Heckman-Kostant, and deduce from it the multiplicities in $\pi_{\lambda_{I_{H}}}$. We explain the latter step in theorem 5 . However, we obtain nice formulas for the multiplicities in $\pi_{\lambda_{L}}$ only with a supplementary assumption which we call condition ( $C$ ).

In section 2 we introduce this condition $(C)$ for $\Psi$ and $L$. Under this condition, we prove that $\pi_{\lambda_{L_{L}}}$ is admissible. In sections 3 and 4 we give formulas for the multiplicities in $\pi_{\lambda_{L}}$ and $\pi_{\lambda_{\left.\right|_{H}}}$, in term of a linear combination of partition functions. Condition $(C)$ is only a sufficient condition of admissibility, but it allows a tremendous simplification in the computation of multiplicities. However, for $(G, H)$ a symmetric pair we obtain the equivalence: $\pi_{\lambda_{l_{H}}}$ is admissible if and only if condition (C) holds for $\Psi$ and $L$. An interesting fact is the role of a particular invariant connected subgroup $Z_{1} K_{1} \subset K$ associated to $\Psi$. We compare our results with previous works, notably by Gross-Wallach [3] on small representations, and Kobayashi [9] on symmetric spaces.

In section 5, we give semi-classical analogues of these results. We consider the coadjoint orbit $\Omega:=$ $G \lambda \subset i \mathfrak{g}_{\mathbf{R}}^{*}$. The admissibility of $\pi_{\lambda_{H}}$ is equivalent to the fact that the restriction $q_{\mathfrak{h}}$ from $\Omega$ to $i \mathfrak{h}_{\mathbf{R}}^{*}$ is proper, so we may consider the measure on $i \mathfrak{h}_{\mathbf{R}}^{*}$ which is the push-forward of the Liouville measure $\beta_{\Omega}$. Under condition $(C)$, we give a formula for this measure in terms of convolution product of Heaviside functions, generalizing formulas given in [5] for the pair $(K, L)$ and in [2] for the pair $(G, K)$.

We say that two square integrable representa-
tions $\pi_{\lambda}$ and $\pi_{\lambda^{\prime}}$ of $G$ are in the same family if $\Psi^{\lambda}$ and $\Psi^{\lambda^{\prime}}$ are $W_{K}$-conjugate. There is a similar definition for the square integrable representations of $H$. We conjecture that, for a square integrable representation $\pi$ of $G$ such that $\pi_{\left.\right|_{H}}$ is admissible, all the representations $\sigma_{i}$ which occur in $\pi_{\left.\right|_{H}}$ belong to a unique family of square integrable representations of $H$. We prove that this is true under condition $(C)$ : this follows from our formulas for multiplicities and for Liouville measures, a convexity result of Weinstein [17], and deep results on partition functions, in particular those of Szenes and Vergne [15].

The problem of finding branching numbers for discrete series, or more generally for $A_{\mathfrak{q}}(\lambda)$ modules, has received much attention. Let us mention Schmid [14], Kobayashi [6, 8], Gross-Wallach [3]. For an update on the subject we refer to Kobayashi [12, 13] and references therein. There exists different types of formulas for branching numbers. In this note we consider only formulas in terms of partition functions. Moreover, we do not expect that our results extend to more singular $A_{\mathfrak{q}}(\lambda)$ modules. For instance, in this larger context, the conjecture on the unique family stated in the previous paragraph is not always true: see examples in [6].
2. Condition (C). We shall use multisets $S$ of vectors of $i \mathfrak{u}_{\mathbf{R}}^{*}$, that is non ordered lists of vectors, allowing repetitions. We say that a multiset $S$ is strict if it is contained in an open half-subspace of $i \mathfrak{u}_{\mathbf{R}}^{*}$; in particular all elements of $S$ are non zero. A submultiset $S^{\prime} \subset S$ is called a positive system for $S$ if $S^{\prime}$ is the intersection of $S$ with an open halfsubspace.

Let $q_{\mathfrak{u}}: \mathrm{t}^{*} \rightarrow \mathfrak{u}^{*}$ be the restriction map. Let $Z$ be the connected component of the centralizer of $\mathfrak{u}$ in $K$. We define $\Phi_{z}:=\Phi(\mathfrak{\jmath}, \mathrm{t})$. It is the set of $\alpha \in \Phi(\mathfrak{f}, \mathrm{t})$ such that $q_{\mathrm{u}}(\alpha)=0$.

For any space $V$ in which $\mathfrak{u}$ acts, we denote by $\Phi(V, \mathfrak{u}) \subset \mathfrak{u}^{*}$ the multiset of non zero roots. For instance, we have $\Phi(\mathfrak{f} / \mathfrak{l}, \mathfrak{u})=q_{\mathfrak{u}}\left(\Phi(\mathfrak{f}, \mathfrak{t}) \backslash \Phi_{\mathfrak{z}}\right) \backslash \Phi(\mathfrak{l}, \mathfrak{u})$.

Definition 1. Consider a positive system $\Psi \subset$ $\Phi(\mathfrak{g}, \mathfrak{t})$. Condition $(C)$ holds for $\Psi$ and $L$ if there exists a positive system $\Delta(\mathfrak{f} / \mathfrak{l}, \mathfrak{u}) \subset \Phi(\mathfrak{f} / \mathfrak{l}, \mathfrak{u})$ such that, for each $w \in W_{K}$, the multiset

$$
q_{\mathfrak{u}}\left(w \Psi_{n}\right) \cup \Delta(\mathfrak{f} / \mathfrak{l}, \mathfrak{u})
$$

is strict.
Theorem 1. Let $\lambda$ be a G-Harish-Chandra parameter. If condition $(C)$ holds for $\Psi^{\lambda}$ and L, then $\pi_{\lambda}$ restricted to $H$, as well to $L$, is admissible.

Remark 1. Condition ( $C$ ) implies in par-
ticular that $q_{\mathfrak{u}}(\alpha) \neq 0$ for all non compact roots of $\Phi(\mathfrak{g}, \mathfrak{t})$. This is in fact always satisfied when $\pi_{\lambda}$ restricted to $L$ is admissible. The rest of condition ( $C$ ) is only a sufficient condition of admissibility - see the examples given below after the definition of $K_{2}(\Psi)$. However it is a natural one to consider for the problem of expressing branching numbers in term of partition functions.

Hereinafter, for simplicity, we assume that $G$ is a simple group. We consider some examples. We denote by $Z_{K}$ the connected component of the center of $K$, and $\mathfrak{z}_{K}$ its Lie algebra. Recall that $\operatorname{dim}_{\mathfrak{z}_{K}} \leq 1$, and that if $\operatorname{dim}_{\mathfrak{3}_{K}}=1$-this is the Hermitian symmetric case - there are (up to $W_{K}$-conjugacy) exactly two positive systems $\Psi \subset \Phi(\mathfrak{g}, \mathrm{t})$ such that $\Psi_{n}$ is $W_{K}$ invariant. They are called holomorphic.

- We suppose that $\operatorname{dim} \mathfrak{z}_{K}=1$.

Let $U$ be a subtorus of $T$ and $\Psi^{\lambda}$ a holomorphic system. Then, condition (C) holds for $\Psi^{\lambda}$ and $U$ if and only if $\pi_{\lambda}$ restricted to $U$ is admissible. Moreover, condition (C) holds if $U$ is sufficiently close to $Z_{K}$ (see [16]).

For an arbitrary system $\Psi$ condition (C) holds for $\Psi$ and $Z_{K}$ if and only if $\Psi$ is a holomorphic system.

- Let $K_{s}$ denote the semisimple factor of $K$. Then, condition (C) holds for $\Psi$ and $K_{s}$ if and only if $\mathbf{R}^{+} \Psi_{n} \cap i 3_{K}^{*}=\{0\}$.
- The fact that condition (C) holds for $\Psi$ and $L$ implies that condition (C) holds for $\Psi$ and any subgroup between $L$ and $K$.
- Let $G$ be a noncompact real form for $G_{2}$, and $\Psi_{c}=\left\{\alpha_{1}, \alpha_{2}\right\}$, with $\alpha_{1}$ short.
Let $L=S U_{2}\left(\alpha_{2}\right), \Psi \supset \Psi_{c}$ be the positive system of roots for which $\alpha_{1}$ is simple. Then condition (C) holds and $\Psi$ and $L$. This is the quaternionic case, and in particular $\Psi$ is small in the sense of [3].
Let $L=S U_{2}\left(\alpha_{1}\right)$ and $\Psi \supset \Psi_{c}$ be the positive system of roots for which $\alpha_{2}$ is simple. Then condition (C) holds for $\Psi$ and $L$. In this case $\Psi$ is not small in the sense of [3].

In the next paragraph we analyze the constraints on $K$ when we assume that condition (C) holds for certain subgroup. To $\Psi$ we associate two ideals of $\mathfrak{f}$. We define

$$
\mathfrak{F}_{1}(\Psi)_{\mathbf{C}}:=\text { ideal spanned by } \sum_{\alpha, \beta \in \Psi_{n}} \mathbf{C} X_{\alpha+\beta} .
$$

Still assuming for simplicity that $\mathfrak{g}$ is simple, we define $\mathfrak{z}_{1}(\Psi)=\mathfrak{z}_{K}$ if $\mathbf{R}^{+} \Psi_{n} \cap i_{\mathfrak{b}}^{K} \neq\{0\}$, and $\mathfrak{z}_{1}=\{0\}$
otherwise. We denote by $Z_{1}(\Psi) \subset Z_{K}$ and $K_{1}(\Psi) \subset$ $K_{s}$ the corresponding invariant connected subgroups of $K$.

Proposition 1. (i) If condition (C) holds for $\Psi$ and $L$, then we have $K_{1}(\Psi) \subseteq L$. If we assume moreover that $\mathfrak{I}=\mathfrak{I} \cap \mathfrak{j} \oplus \mathfrak{I} \cap \mathfrak{f}_{s}$, then we have $Z_{1}(\Psi) K_{1}(\Psi) \subseteq L$.
(ii) When $K_{1}(\Psi) Z_{1}(\Psi)$ is included in $L$, then condition $(C)$ holds for $\Psi$ and $L$.

When $Z_{1}(\Psi) \neq 1$, condition (C) alone does not imply that $L$ contains $Z_{1}(\Psi)$. This can been seen on the holomorphic example with $L=U$ considered above.

When $\Psi$ is small in the sense of [3], Gross and Wallach define a specific invariant proper subgroup of $K$, to which the corresponding discrete series of $G$ restrict in an admissible manner. In this case, we check that the group $K_{1}(\Psi) Z_{1}(\Psi)$ coincides with the group of Gross and Wallach. However, we computed all the positive systems $\Psi$ for which the group $K_{1}(\Psi)$ is a proper subgroup of $Z_{K} K_{s}$, and (for classical $\mathfrak{g}$ ) those for which $Z_{1}(\Psi)$ is a proper subgroup of $Z_{K}$. Many of them are not small. We gave above an example for $G=G_{2}$.

Let $K_{2}(\Psi) \subset K_{s}$ denote the invariant connected subgroup complementary to $K_{1}(\Psi)$. We show that $\mathfrak{f}_{2}(\Psi)$ is the largest semisimple ideal of $\mathfrak{f}_{s}$ contained in the subalgebra generated by the root vectors corresponding to the simple compact roots of $\Psi$, and their opposite.

The above results imply that $\pi_{\lambda}$ has an admissible restriction to the group $K_{1}\left(\Psi^{\lambda}\right) Z_{1}\left(\Psi^{\lambda}\right)$. It is an interesting fact that it may have also an admissible restriction to $K_{2}\left(\Psi^{\lambda}\right)$. This provides examples where there is admissibility of restriction, but where condition (C) is not satisfied.

Suppose first that $K_{2}\left(\Psi^{\lambda}\right)=K_{s}$. Then $\Psi^{\lambda}$ is holomorphic and one can show that $\pi_{\lambda}$ restricted to $K_{s}$ is admissible if and only if $G / K$ is not a tube domain.

Suppose now that $K_{2}\left(\Psi^{\lambda}\right)$ is a proper subgroup of $K_{s}$. We show that $\pi_{\lambda}$ restricted to $K_{2}\left(\Psi^{\lambda}\right)$ (and also to $Z_{K} K_{2}\left(\Psi^{\lambda}\right)$ ) is admissible if and only if $\mathfrak{g}=\mathfrak{s p}(p, 1)$ and $\Psi$ is quaternionic (that is, $\mathfrak{1}_{1}\left(\Psi^{\lambda}\right)=$ $\mathfrak{s p}(1)$ and $\left.\mathfrak{f}_{2}\left(\Psi^{\lambda}\right)=\mathfrak{s p}(q)\right)$. Note that in the particular case of $\mathfrak{s p}(1,1) \simeq \mathfrak{s v}(4,1)$, all $\Psi$ are quaternionic.

Next we relate condition (C) to some of the theorems of Kobayashi on admissible restrictions. For the rest of this section we assume $(G, H)$ is a symmetric
pair. Let $\sigma$ denote the involution determined by the symmetric pair $(G, H)$. We may and do assume that $\sigma$ stabilizes $K$ and $T$. Thus $(K, L)$ is again a symmetric pair. We choose a $\sigma$-stable compact Cartan subgroup $B \subset K$ such that the space $\mathfrak{b}_{-}:=\{X \in$ $\mathfrak{b}: \sigma X=-X\}$ is of maximal dimension: it is a Cartan subspace for the symmetric pair $(K, L)$. We fix a system of positive roots $\Delta \subset \Phi(\mathfrak{f}, \mathfrak{b})$ such that $\left\{\alpha_{\mathrm{l}_{-}} \neq 0: \alpha \in \Delta\right\}$ is a system of positive roots for $\Phi\left(\mathfrak{f}, \mathfrak{b}_{-}\right)$.

Let $\tilde{\Psi} \subset \Phi(\mathfrak{g}, \mathfrak{b})$ be a positive system. We say that $\tilde{\Psi}$ satisfies Kobayashi's condition if we have:

$$
\tilde{\Psi} \supset \Delta \quad \text { and } \quad \mathbf{R}^{+} \tilde{\Psi}_{n} \cap \mathfrak{b}_{-}^{*}=\{0\}
$$

Consider a $G$-Harish-Chandra parameter $\lambda \in \mathfrak{b}^{*}$ such that the positive system of positive roots $\tilde{\Psi}^{\lambda} \subset$ $\Phi(\mathfrak{g}, \mathfrak{b})$ contains $\Delta$. The importance of Kobayashi condition comes from the following criterium of admissibility:

Theorem 2. (Kobayashi [7, 9, 10]) The restriction of $\pi_{\lambda}$ to $H$ is admissible if and only if $\tilde{\Psi}^{\lambda}$ satisfies Kobayashi's condition.

We prove:
Proposition 2. Assume $\quad \mathbf{R}^{+} \tilde{\Psi}_{n} \cap \mathfrak{b}_{-}^{*}=\{0\}$. Then $K_{1}(\tilde{\Psi})$ is a subgroup of $L$.

Condition (C) is stated by means of the $\sigma$-stable compact Cartan subgroup $T$, for which the dimension of $t_{-}$is minimal. Kobayashi's condition is expressed through the $\sigma$-stable compact Cartan subgroup $B$, for which the dimension of $\mathfrak{b}_{-}$is maximal. We relate $\mathfrak{b}$ and $\mathfrak{t}$ via a Cayley transform. Consider a positive system $\tilde{\Psi} \subset \Phi(\mathfrak{g}, \mathfrak{b})$ which contains contains $\Delta$, and let $\Psi \subset \Phi(\mathfrak{g}, \mathrm{t})$ be the positive system obtained from $\tilde{\Psi}$ by Cayley transform. Using proposition 2, we prove:

Proposition 3. Assume $(G, H)$ is a symmetric pair. Then condition $(C)$ for $\Psi$ and $L$ holds if and only if Kobayashi's condition holds for $\tilde{\Psi}$.

Corollary 1. Let $(G, H)$ be a symmetric pair.
(i) The restriction of $\pi_{\lambda}$ to $L$ and $H$ is admissible if and only if condition $(C)$ holds for $\Psi^{\lambda}$ and $L$.
(ii) The restriction of $\pi_{\lambda}$ to $L$ and $H$ is admissible if and only $Z_{1}\left(\Psi^{\lambda}\right) K_{1}\left(\Psi^{\lambda}\right) \subset L$.
3. Multiplicity formulas. To express multiplicities, it is convenient to use centered partition functions, that is partitions functions appropriately shifted. Let $S=\left\{\gamma_{1}, \cdots, \gamma_{q}\right\}$ be a strict multiset of elements of $i \mathfrak{u}_{\mathbf{R}}^{*}$. We denote by $p_{S}$ the corresponding
partition function, by $\rho_{S}$ the half-sum of elements of $S$, and by $q_{S}$ the partition function shifted by $\rho_{S}$. Thus, for $\nu \in i \mathfrak{u}_{\mathbf{R}}^{*}$, we have $q_{S}(\nu)=p_{S}\left(\nu-\rho_{S}\right)$.

We choose a positive system $\Psi_{\mathfrak{\jmath}} \subset \Phi_{\mathfrak{\jmath}}$, and denote by $\varpi_{3}$ the Weyl dimension polynomial for $Z$ : for $\gamma \in \mathfrak{t}^{*}$

$$
\varpi_{\mathfrak{z}}(\gamma)=\frac{\prod_{\alpha \in \Psi_{\mathbf{3}}} \gamma\left(h_{\alpha}\right)}{\prod_{\alpha \in \Psi_{\mathfrak{z}}} \rho_{\Psi_{\mathfrak{z}}}\left(h_{\alpha}\right)} .
$$

Let $\lambda$ be a $G$-Harish-Chandra parameter. Let $m\left(\pi_{\lambda}, \sigma_{\mu}\right)$ denote the multiplicity of $\sigma_{\mu}$ in $\pi_{\lambda_{\left.\right|_{H}}}$, for all $H$-Harish-Chandra parameters $\mu$, and $m\left(\pi_{\lambda}, \kappa_{\nu}\right)$ the multiplicity of $\kappa_{\nu}$ in $\pi_{\lambda_{L}}$, for all L-Harish-Chandra parameters $\nu$.

Assume that $\Psi^{\lambda}$ satisfies condition (C) for $L$, and choose $\Delta(\mathfrak{f} / \mathfrak{l}, \mathfrak{u})$ like in definition 1 . For each $w \in W_{K}$, we define $S_{w}^{L}:=q_{\mathfrak{u}}\left(w \Psi_{n}\right) \cup \Delta(\mathfrak{f} / \mathfrak{l}, \mathfrak{u})$ and $S_{w}^{H}:=S_{w}^{L} \backslash \Phi(\mathfrak{h} / \mathfrak{l}, \mathfrak{u})$. Remark that that for all $w \in$ $W_{K}, S_{w}^{L}$ is a positive system for $\Phi(\mathfrak{g} / \mathfrak{l}, \mathfrak{u})$, and $S_{w}^{H}$ a positive system for $\Phi(\mathfrak{g} / \mathfrak{h}, \mathfrak{u})$.

Theorem 3. Assume that condition ( $C$ ) holds for $\Psi^{\lambda}$ and $L$. Then we have:

$$
m\left(\pi_{\lambda}, \sigma_{\mu}\right)= \pm \sum_{w \in W_{Z} \backslash W_{K}} \epsilon(w) \varpi_{\mathfrak{\jmath}}(w \lambda) q_{S_{w}^{H}}\left(\mu-q_{\mathfrak{u}}(w \lambda)\right) .
$$

Note that when $H=L$ and $\mu=\nu$, the theorem gives the multiplicities $m\left(\pi_{\lambda}, \kappa_{\nu}\right)$.

We may rewrite theorem 3 using the decomposition $W_{K}=W_{1} \times W_{2}$, where $W_{i}=W_{K_{i}\left(\Psi^{\lambda}\right)}$. We do it in the simplest case, when $H$ and $G$ have the same rank. Assume that $T=U$. We may choose $\Delta(\mathfrak{f} / \mathfrak{l}, \mathfrak{t}) \subset \Psi^{\lambda}$. Define $S_{\lambda}^{H}:=\Psi^{\lambda} \backslash \Phi(\mathfrak{h}, \mathfrak{t})$. We obtain

$$
m\left(\pi_{\lambda}, \sigma_{\mu}\right)= \pm \sum_{s \in W_{1}, t \in W_{2}} \epsilon(s t) q_{S_{\lambda}^{H}}(s \mu-t \lambda)
$$

## 4. Partition functions via discrete Heavi-

 side measures. To make the comparison with the results of the next paragraph easier, we restate theorem 3 in terms of measures on $i \mathfrak{u}_{\mathbf{R}}^{*}$. For $\gamma \in i \mathfrak{u}_{\mathbf{R}}^{*}$ let $\delta_{\gamma}$ denote the Dirac delta function attached to $\gamma$. If $\gamma \neq 0$, we consider the measure$$
y_{\gamma}=\sum_{n \geq 0} \delta_{\frac{\gamma}{2}+n \gamma}=\delta_{\frac{\gamma}{2}}+\delta_{\frac{\gamma}{2}+\gamma}+\delta_{\frac{\gamma}{2}+2 \gamma}+\cdots
$$

Let $S=\left\{\gamma_{1}, \cdots, \gamma_{q}\right\}$ be a strict multiset of elements of $i u_{\mathbf{R}}^{*}$. We write $*$ for the convolution product of measures on $i \mathfrak{u}_{\mathbf{R}}^{*}$. We consider the measure

$$
y_{S}:=y_{\gamma_{1}} * \cdots * y_{\gamma_{q}}=\sum_{\mu \in i u_{\mathbf{R}}^{*}} q_{S}(\mu) \delta_{\mu}
$$

Let $\lambda$ be a $G$-Harish-Chandra parameter. Assume that $\Psi^{\lambda}$ satisfy condition (C) for $L$, and choose $\Delta(\mathfrak{f} / \mathrm{l}, \mathfrak{u})$ like in definition 1 . We use the notations of section 3. For $\nu \in i \mathfrak{u}_{\mathbf{R}}^{*}$, we define $m^{L}(\lambda, \nu)$ by the formula

$$
\sum_{\nu \in i u_{\mathbf{R}}^{*}} m^{L}(\lambda, \nu) \delta_{\nu}=\sum_{w \in W_{Z} \backslash W_{K}} \epsilon(w) \varpi_{3}(w \lambda) \delta_{q_{\mathrm{u}}(w \lambda)} * y_{S_{w}^{L}} .
$$

For $\mu \in i \mathfrak{u}_{\mathbf{R}}^{*}$, we define $m^{H}(\lambda, \mu)$ similarly.
The function $m^{H}(\lambda,$.$) and m^{L}(\lambda,$.$) are skew$ invariant under the action of $W_{L}$. The following is essentially a reformulation of theorem 3.

Theorem 4. Assume that condition ( $C$ ) holds for $\Psi^{\lambda}$ and $L$.
(i) $m^{L}(\lambda, \nu)=0$ if $\nu$ is not a L-Harish-Chandra parameter. If $\nu$ is a L-Harish-Chandra parameter, we have $m\left(\pi_{\lambda}, \kappa_{\nu}\right)=\left|m^{L}(\lambda, \nu)\right|$.
(ii) $m^{H}(\lambda, \mu)=0$ if $\mu$ is not a H-Harish-Chandra parameter. If $\mu$ is a $H$-Harish-Chandra parameter, we have $m\left(\pi_{\lambda}, \sigma_{\mu}\right)=\left|m^{H}(\lambda, \mu)\right|$.

Let us give an outline of the proof of theorem 4. We consider first some special connected subgroups $L^{\prime} \subset K$ such that condition (C) is satisfied for $\Psi^{\lambda}$ and $L^{\prime}$, and prove theorem 4 by applying it to the pairs $(G, K)$ and $\left(K, L^{\prime}\right)$ (where it is known, by [4] and [5]). Then, as in [5], we use a trick which allows to deduce the multiplicities for the pair $(G, H)$ from the multiplicities for the pair $\left(G, L^{\prime}\right)$, whenever $H$ contains $L^{\prime}$ and $L^{\prime}$ contains $U$, by using product of difference operators associated to the roots in some positive system $\Psi\left(\mathfrak{h} / \mathrm{l}^{\prime}, \mathfrak{u}\right) \subset \Phi\left(\mathfrak{h} / \mathrm{l}^{\prime}, \mathfrak{u}\right)$. In the particular case $L=L^{\prime}$, it follows from theorem 5 below. To state this theorem, we introduce some notations.

For $\gamma \in i \mathfrak{u}_{\mathbf{R}}^{*}$ and $S$ as above, let $d_{\gamma}:=\delta_{-\frac{\gamma}{2}}-\delta_{\frac{\gamma}{2}}$, and

$$
d_{S}:=d_{\gamma_{1}} * \cdots * d_{\gamma_{q}}
$$

The operation of convolution by $d_{\gamma}$ is a difference operator, and we have $d_{\gamma} * y_{\gamma}=\delta_{0}$ and $d_{S} * y_{S}=\delta_{0}$.

Theorem 5. Let $S_{\mathfrak{h} / \mathrm{I}} \subset \Phi(\mathfrak{h} / \mathfrak{l})$ be a positive system of non compact roots for $\mathfrak{h}$. Then

$$
d_{S_{\mathfrak{l}} / \mathrm{I}} * \sum_{\nu \in i u_{\mathbf{R}}^{*}} m^{L}(\lambda, \nu) \delta_{\nu}= \pm \sum_{\mu \in i u_{\mathbf{R}}^{*}} m^{H}(\lambda, \mu) \delta_{\mu}
$$

5. Push forward of Liouville Measure. We now study the semi-classical analogues of multiplicity functions, as in Heckman [5] and Duflo-HeckmanVergne [2]. Consider an element $\lambda \in i t^{*}$. As usual, we extend $\lambda$ to a linear functional on $\mathfrak{g}$ which van-
ishes on each root subspace. The coadjoint orbit $\Omega:=G \lambda \subset i \mathrm{~g}_{\mathbf{R}}^{*}$ is a symplectic manifold. Let $\beta_{\Omega}$ denote the Kostant-Kirillov-Liouville measure, normalized as in [1], Chapter VII.

From now on, we assume that $\lambda$ is $\mathfrak{g}$-regular. We denote by $q_{\mathfrak{h}}: \mathfrak{g}^{*} \rightarrow \mathfrak{h}^{*}$ the restriction map, and we define similarly $q_{\mathrm{I}}$.

Proposition 4. The following are equivalent.
(i) $q_{I}$ is proper on $\Omega$.
(ii) $q_{\mathfrak{\natural}}$ is proper on $\Omega$.

If moreover $\lambda$ is a G-Harish-Chandra parameter, it is also equivalent to
(iii) $\pi_{\lambda}$ has an admissible restriction to $H$.

From now on, we assume that $q_{\mathfrak{h}}$ is proper on $\Omega$. To describe $q_{\mathfrak{h}}(\Omega)$, it is convenient to choose a positive Weyl Chamber $C_{L} \subset i \mathfrak{u}_{\mathbf{R}}^{*}$. The set $q_{\mathfrak{h}}(\Omega) \subset i \mathfrak{h}_{\mathbf{R}}^{*}$ is $H$-invariant, and each $H$-orbit $\omega \subset q_{\mathfrak{h}}(\Omega)$ meets $C_{L}$ in exactly one point. According to Weinstein [17], $q_{\mathfrak{b}}(\Omega) \cap C_{L}$ is a convex polyhedron.

This allows to describe the push-forward $\left(q_{\mathfrak{h}}\right)_{*}\left(\beta_{\Omega}\right)$ by means of a measure $M^{H}$ on $C_{L}$, called the Duistermaat-Heckman measure. For $\mu \in i \mathfrak{u}^{*}$, let $\omega_{\mu}:=H \mu \subset i \mathfrak{h}^{*}$ be its coadjoint orbit, and $\beta_{\omega_{\mu}}$ the corresponding Liouville measure. Then $M^{H}$ is the positive measure on $q_{\mathfrak{h}}(\Omega) \cap C_{L}$ such that

$$
\left(q_{\mathfrak{h}}\right)_{*}\left(\beta_{\Omega}\right)=\int_{C_{L}} d M^{H}(\mu) \beta_{\omega_{\mu}}
$$

There is a canonical way to extend $M^{H}$ to a signed measure on $i u_{\mathbf{R}}^{*}$ which is $W_{L}$-skew invariant.

In analogy with section 4 , for $\gamma \in i \mathfrak{u}_{\mathbf{R}}^{*}$ with $\gamma \neq 0$, we define the Heaviside measure $Y_{\gamma}$ by $Y_{\gamma}(\varphi)=\int_{0}^{\infty} \varphi(t \gamma) d t$, and for a strict multiset $S=$ $\left\{\gamma_{1}, \cdots, \gamma_{q}\right\}$ of elements of $i \mathfrak{u}_{\mathbf{R}}^{*}$, the measure

$$
Y_{S}:=Y_{\gamma_{1}} * \cdots * Y_{\gamma_{q}} .
$$

We also define a measure $t_{\gamma}$ by the formula

$$
t_{\gamma}(\varphi):=\int_{-\frac{1}{2}}^{\frac{1}{2}} \varphi(t \gamma) d t
$$

and $t_{S}:=t_{\gamma_{1}} * \cdots * t_{\gamma_{q}}$. We have $t_{\gamma} * y_{\gamma}=Y_{\gamma}$ and $t_{S} * y_{S}=Y_{S}$.

Let $\lambda$ be a $G$-Harish-Chandra parameter. Assume that $\Psi^{\lambda}$ satisfy condition (C) for $L$, and choose $\Delta(\mathfrak{f} / \mathfrak{l}, \mathfrak{u})$ like in definition 1 . We use the notations of section 3. We obtain the following analogue of theorem 4.

Theorem 6. Assume that condition ( $C$ ) holds for $\Psi^{\lambda}$ and L. We have:

$$
M^{H}= \pm \sum_{w \in W_{Z} \backslash W_{K}} \epsilon(w) \varpi_{\mathfrak{\jmath}}(w \lambda) \delta_{q_{\mathrm{u}}(w \lambda)} * Y_{S_{w}^{H}}
$$

Let $S$ be a positive system for $\Phi(\mathfrak{g} / \mathfrak{h}, \mathfrak{u})$. Since $t_{\gamma}=t_{-\gamma}$, the measure $t_{S}$ does not depend on the choice of $S$, and will be denoted by $r_{\mathfrak{g} / \mathfrak{b}}$. Comparing theorems 4 and 6 , we obtain.

## Proposition 5.

$$
r_{\mathfrak{g} / \mathfrak{h}} * \sum_{\mu \in i \mathfrak{u}_{\mathbf{R}}^{*}} m^{H}(\lambda, \mu) \delta_{\mu}=M^{H}
$$

Theorems about partitions functions due to [15] and others allow to invert this formula.

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