

### 13. Spectral Analysis for the Casimir Operator on the Quantum Group $SU_q(1, 1)$

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In this letter we will determine the spectrum of the Casimir operator for zonal spherical functions on the quantum group  $SU_q(1, 1)$ , and show the Plancherel formula and the expansion theorem for them.

§ 1.  $\mathcal{U}_q(su(1, 1))$  ( $0 < q < 1$ ) be the real form of the universal quantum enveloping algebra [3] with the  $*$ -structure,  $k^* = k$ ,  $e^* = -f$ ,  $f^* = -e$  [1]. In [1] we have classified irreducible unitary representations of  $\mathcal{U}_q(su(1, 1))$  as follows: We set a left  $\mathcal{U}_q(su(1, 1))$ -module  $V_l = \bigoplus_{i \in I_l} \mathcal{C}\xi_i$  by

$$(1.1) \quad \begin{aligned} k \cdot \xi_j &= q^{-j} \xi_j, \\ e \cdot \xi_j &= q^{1/2-l} \frac{1 - q^{2(l-j+1)}}{1 - q^2} \xi_{j-1}, \\ f \cdot \xi_j &= q^{1/2-l} \frac{1 - q^{2(l+j+1)}}{1 - q^2} \xi_{j+1}. \end{aligned}$$

Here the complex spin  $l$  and the indices set  $I_l$  are listed in the table below ( $q = e^{-h}$ ,  $h > 0$ ).

(1.2)

$l$	$I_l$
$l = -1/2$	$I_l = \{1/2, 3/2, \dots\}, \{-1/2, -3/2, \dots\}$
$-1/2 < l < 0$	$I_l = \mathbb{Z}$
$l \in 1/2\mathbb{N}$	$I_l = \{l+1, l+2, \dots\}, \{-l-1, -l-2, \dots\}$
$l = -1/2 + i\theta$	$(0 \leq \theta \leq \pi/2h) I_l = \mathbb{Z}, (0 < \theta \leq \pi/2h) I_l = \mathbb{Z} + 1/2$
$l = -1/2 + i(\pi/2h + it) \quad (0 < t)$	$I_l = \mathbb{Z}, \mathbb{Z} + 1/2$

Let  $w_{ij}^{(l)}$  be the matrix elements on  $SU_q(1, 1)$  corresponding to each representations. They are represented in terms of the basic hypergeometric functions

$${}_2\varphi_1 \left( \begin{matrix} a, b \\ c \end{matrix}; q, z \right) = \sum_{n=0}^{\infty} \frac{(a; q)_n (b; q)_n}{(q; q)_n (c; q)_n} z^n,$$

where  $(a; q)_n = \prod_{r=0}^{n-1} (1 - aq^r)$  ( $1 \leq n \leq \infty$ ). For example, in the case of  $i+j \leq 0, j \leq i$ .

$$(1.3) \quad w_{ij}^{(l)} = q^{(j-i)(l+j)} \frac{(q^{2(l-i+1)}; q^2)_{i-j}}{(q^2; q^2)_{i-j}} x^{-i-j} v^{i-j} {}_2\varphi_1 \left( \begin{matrix} q^{2(l-j+1)}, q^{-2(l+j)} \\ q^{2(l-j+1)} \end{matrix}; q^2, q^{2\zeta} \right),$$

where  $x, u, v$  and  $y$  are the coordinate elements on  $SU_q(1, 1)$  and  $\zeta = -q^{-1}uv$ . These matrix elements satisfy the eigen-equation

$$(1.4) \quad \pi_i(C)w_{ij}^{(l)} = [l+1/2]^2 w_{ij}^{(l)}$$

where  $\pi_i$  is the left invariant differential representation of  $\mathcal{U}_q(sl(2))$  on  $\mathcal{A}$  (the dual space of  $\mathcal{U}_q(sl(2))$ ), and

$$(1.5) \quad C = \frac{ql^2 + q^{-1}k^{-2} - 2}{(q - q^{-1})^2} + fe$$

is the Casimir element of  $\mathcal{U}_q(sl(2))$  [3], and  $[a] = (q^a - q^{-a}) / (q - q^{-1})$ .

Equation (1.4) is, in fact, a  $q$ -difference equation of the second order. In particular, the zonal spherical function

$$(1.6) \quad w_{00}^{(l)} = {}_2\phi_1 \left( \begin{matrix} q^{2l+2}, q^{-2l} \\ q^2 \end{matrix}; q^2, q^2\zeta \right)$$

satisfies a  $q$ -analogue of the Legendre equation

$$(1.7) \quad qD_{q^2}\{z(z+1)D_{q^2}T_{q^2}^{-1}\varphi(z)\} + [1/2]^2\varphi(z) = [l+1/2]^2\varphi(z)$$

where  $z = -\zeta$ , and

$$D_{q^2}\varphi(z) = \frac{\varphi(z) - \varphi(q^2z)}{(1 - q^2)z}, \quad T_{q^2}\varphi(z) = \varphi(q^2z).$$

§ 2. We will consider the spectral theory for the difference equation which arises from the equation (1.7).

For a solution  $\varphi(z)$  to (1.7), we set  $\varphi(n) = \varphi(q^{2n})$ . Then we see that  $\varphi = (\varphi(n))_{n \in \mathbb{Z}}$  solves the following difference equation

$$(2.1) \quad (q + q^{1-2n})\varphi(n-1) - 2(q^2 + q^{1-2n})\varphi(n) + (q^3 + q^{1-2n})\varphi(n+1) = \lambda\varphi(n),$$

where  $n$  runs over  $\mathbb{Z}$ , and

$$(2.2) \quad \lambda = (1 - q^2)^2 [l + 1/2]^2.$$

Taking into account the Haar measure on  $SU_q(1, 1)$  (cf. [2]), we introduce a Hilbert space

$$l_{00}^2 = \{\varphi = (\varphi(n))_{n \in \mathbb{Z}} \mid \|\varphi\|_{00} < +\infty\}$$

with an inner product

$$(2.3) \quad (\varphi, \psi) = (1 - q^2) \sum_{n=-\infty}^{+\infty} \varphi(n) \overline{\psi(n)} q^{2n}.$$

It is easy to see that the equation (2.1) is formally self-adjoint in  $l_{00}^2$ . Moreover we impose on (2.1) the following boundary condition;

$$(2.4) \quad \lim_{n \rightarrow +\infty} \{\varphi(n) - \varphi(n-1)\} = 0.$$

**Theorem 1.** Equation (2.1) with the boundary condition (2.4) constitutes a self-adjoint boundary value problem in  $l_{00}^2$ .

**Theorem 2.** The Green kernel  $G(n, m; \lambda)$  for the above boundary value problem is given as follows:

$$(2.5) \quad G(n, m; \lambda) = \varphi_{+\infty}(n; \lambda) \cdot \varphi_{-\infty}(m; \lambda) \quad (m \leq n)$$

where, setting  $l = -1/2 + i\theta$ ,

$$(2.6) \quad \varphi_{+\infty}(n; \lambda) = {}_2\phi_1 \left( \begin{matrix} q^{1+2i\theta}, q^{1-2i\theta} \\ q^2 \end{matrix}; q^2, -q^{2n+2} \right)$$

and  $\varphi_{-\infty}(n; \lambda) = \gamma(\theta)\phi(q^{2n}; \theta)$  with

$$(2.7) \quad \phi(z; \theta) = (-q^{1-2i\theta}; q^2)_{\infty} (q^{4-4i\theta}; q^4)_{\infty} (q^2z)^{-1/2+i\theta}$$

$$(2.8) \quad \gamma(\theta) = \frac{\times {}_2\varphi_1\left(\begin{matrix} q^{1-2i\theta}, q^{1-2i\theta} \\ q^{2-4i\theta} \end{matrix}; q^2, -z^{-1}\right) - 2(q^4; q^4)_\infty q^{-2i\theta}}{(1-q^2)(q^2; q^4)_\infty (-q^{1+2i\theta}; q^2)_\infty (q^{1-2i\theta}; q^2)_\infty (-q^{1-2i\theta}; q^2)_\infty}.$$

Investigating the singularities of the Green function (cf. [4]), we can determine the spectrum and the spectral measure of this boundary value problem, and finally establish the eigen-function expansion theorem.

**Theorem 3.** *Let  $\varphi(\theta) = (\varphi_{+\infty}(n; \lambda))_{n \in \mathbb{Z}}$  and  $\theta_k = \pi/2h + (2k+1)i/2$ . Then, for any  $f \in l^2_{00}$ , we have*

$$(2.9) \quad \|f\|_{00}^2 = \int_0^{2\pi/h} d\theta c(\theta) |(f, \varphi(\theta))_{00}|^2 + \sum_{k=0}^{\infty} c_k |(f, \varphi(\theta_k))_{00}|^2,$$

and

$$(2.10) \quad f = \int_0^{\pi/2h} d\theta c(\theta) \varphi(\theta) \cdot (f, \varphi(\theta))_{00} + \sum_{k=0}^{\infty} c_k \varphi(\theta_k) (f, \varphi(\theta_k))_{00},$$

where

$$(2.11) \quad c(\theta) = \frac{4q^2(q^4; q^4)_\infty (q^{4i\theta}; q^4)_\infty (q^{-4i\theta}; q^4)_\infty}{\pi(1-q^2)(q^2; q^4)_\infty (q^{2+4i\theta}; q^4)_\infty (q^{2-4i\theta}; q^4)_\infty},$$

$$(2.12) \quad c_k = [2k+1].$$

The formula (2.9) is viewed as the Plancherel formula for zonal spherical functions on the quantum group  $SU_q(1, 1)$ , and, in the classical case, (2.10) is known as the Fok-Mehler formula.

Extending the theory developed here to the whole of functions on  $SU_q(1, 1)$ , one can determine the spectrum of the Casimir operator and establish the Plancherel formula in the  $L^2$ -space on  $SU_q(1, 1)$  [5]. In particular, one sees that the principal series of unitary representations are parametrized as follows:

$$(2.13) \quad \begin{aligned} l &= -1/2 + i\theta \quad (0 \leq \theta \leq \pi/2h), \text{ the principal continuous series;} \\ l &\in \frac{1}{2}N, \text{ the first discrete series;} \\ l &= -1/2 + i\theta_k \quad (k \in N), \text{ the second discrete series.} \end{aligned}$$

After this work was completed, the author learned the paper [6] of L. L. Vaksman and L. I. Korodskii. They also obtained the formula (2.9).

### References

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