# 27. On the Existence of the Poles of the Scattering Matrix for Several Convex Bodies 

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1. Introduction. Let $\mathcal{O}$ be an open bounded set in $\boldsymbol{R}^{3}$ with smooth boundary $\Gamma$. We set

$$
\Omega=\boldsymbol{R}^{3}-\overline{\mathcal{O}},
$$

and suppose that $\Omega$ is connected. Consider the following acoustic problem

$$
\begin{cases}\square u(x, t)=\frac{\partial^{2} u}{\partial t^{2}}-\Delta u=0 & \text { in } \Omega \times(-\infty, \infty)  \tag{1.1}\\ B u(x, t)=0 & \text { on } \Gamma \times(-\infty, \infty) \\ u(x, 0)=f_{1}(x) & \\ \frac{\partial u}{\partial t}(x, 0)=f_{2}(x) & \end{cases}
$$

where $\Delta=\sum_{j=1}^{3} \partial^{2} / \partial x_{j}^{2}$. As boundary operator $B$ we shall consider the following two operators,

$$
B_{D}=1 \quad \text { (Dirichlet condition) }
$$

and

$$
B_{N}=\sum_{j=1}^{3} n_{j}(x) \partial / \partial x_{j} \quad \text { (Neumann condition) }
$$

where $n(x)=\left(n_{1}(x), n_{2}(x), n_{3}(x)\right)$ denotes the unit outer normal of $\Gamma$ at $x$.
Denote by $\mathcal{S}_{\dagger}(z), \dagger=D, N$, the scattering matrix for the scatterer $\mathcal{O}$ under the boundary condition $B_{+} u=0$ (for the definition, see [6]). It is well known that $\mathcal{S}_{\mathrm{f}}(z)$ is an $\mathcal{L}\left(L^{2}\left(S^{2}\right)\right.$ )-valued meromorphic function in the whole complex domain $C$.

As to the modified Lax and Phillips conjecture, ${ }^{11}$ that is, when $\mathcal{O}$ is trapping, there exists $\alpha>0$ such that a slub domain $\{z ; \operatorname{Im} z<\alpha\}$ contains an infinite number of poles of the scattering matrix, we have a few examples. Especially for the Dirichlet boundary condition an obstacle consisting of two disjoint convex bodies is the only example ( $[2,3]$ ). The purpose of this note is to study the modified Lax and Phillips conjecture in the case that $\mathcal{O}$ consists of several disjoint strictly convex bodies. Our theorem gives a sufficient condition for the existence of such $\alpha$, which is stated by means of an analytic function defined by using purely geometric informations of $\Omega$.

This work was done during my stay at Massachusetts Institute of Technology. I would like to express my sincere gratitude to Professor Melrose for the invitation and stimulating conversations.

[^0]2. Statement of theorems. Let $\mathcal{O}_{j}, j=1,2, \cdots, j$ be open bounded sets in $\boldsymbol{R}^{3}$ with smooth boundary $\Gamma_{j}$. We assume the following:
(H.1) Every $\mathcal{O}_{j}$ is strictly convex, that is, the Gaussian curvature of $\Gamma_{j}$ is positive everywhere.
(H.2) For all $\left\{j_{1}, j_{2}, j_{3}\right\} \in\{1,2, \cdots, J\}^{3}$ such that $j_{l} \neq j_{h}$ if $l \neq h$, the convex hull of $\overline{\mathcal{O}}_{j_{1}}$ and $\overline{\mathcal{O}}_{j_{2}}$ has no intersection with $\overline{\mathcal{O}}_{j_{3}}$.

We set

$$
\begin{equation*}
\mathcal{O}=\bigcup_{j=1}^{J} \mathcal{O}_{j} . \tag{2.1}
\end{equation*}
$$

Let $\gamma$ be a periodic ray in $\Omega$. We shall use the following notations:
$d_{r}$ : the length of $\gamma$,
$T_{r}$ : the primitive period of $\gamma$,
$i_{r}$ : the number of the reflecting points of $\gamma$,
$P_{r}$ : the Poincare map of $\gamma$.
Concerning the periodic rays in $\Omega$, we have
$\#\left\{\gamma:\right.$ periodic ray in $\Omega$ such that $\left.d_{r}<r\right\} \leq e^{a_{0} r}$,
$\left|I-P_{r}\right| \geq e^{2 a_{1} d_{r}}$,

$$
\begin{equation*}
\left|I-P_{r}\right| \geq e^{2 a_{1} a_{r}}, \tag{2.2}
\end{equation*}
$$

where $a_{0}$ and $a_{1}$ are positive constants determined by $\mathcal{O}$, and we denote by $|A|$ the determinant of matrix $A$.

Define functions $F_{+}(\mu), \dagger=D, N$, by

$$
\begin{equation*}
F_{\dagger}(\mu)=\sum_{r}(-1)^{a_{+} i_{r}} T_{r}\left|I-P_{r}\right|^{-1 / 2} e^{-\mu d_{r}}, \quad a_{D}=1, \quad a_{N}=0 \tag{2.4}
\end{equation*}
$$

where the summation is taken over all the periodic rays in $\Omega$. Note that it follows from (2.2) and (2.3) that $F_{D}$ and $F_{N}$ are holomorphic in $\{\mu: \operatorname{Re} \mu$ $\left.>a_{0}-a_{1}\right\}$.

Theorem 1. Let $\mathcal{O}$ be an obstacle given by (2.1) satisfying (H.1) and (H.2). If $F_{\dagger}, \dagger=D$ or $N$, cannot be prolonged analytically to an entire function, there exists $\alpha>0$ such that the scattering matrix $\mathcal{S}_{\dagger}(z)$ has infinitely many poles in $\{z ; \operatorname{Im} z<\alpha\}$.

Theorem 2. Let $\mathcal{O}_{1}, \mathcal{O}_{2}$ and $\widetilde{\mathcal{O}}_{3}$ be open sets in $\boldsymbol{R}^{3}$ satisfying (H.1) and (H.2). If $\mathcal{O}_{3} \subset \tilde{\mathcal{O}}_{3}$, and the principal curvatures of $\Gamma_{3}=\partial \mathcal{O}_{3}$ are greater than $\kappa$ everywhere of $\Gamma_{3}$, then $F_{D}$ for $\mathcal{O}=\bigcup_{j=1}^{3} \mathcal{O}_{j}$ cannot be prolonged analytically to an entire function. Here $\kappa$ is a positive constant depending on $\mathcal{O}_{1}, \mathcal{O}_{2}$ and $\widetilde{\mathcal{O}}_{3}$.

Remark. It is easy to show that $F_{N}$ has a singularities on the real axis. Thus in the case of the Neumann condition, the modified Lax and Phillips conjecture holds for $\mathcal{O}$ satisfying (H.1) and (H.2) ([5]).
3. On the proofs of theorems. In order to prove Theorem 1 we shall use the trace formula due to Bardos, Guillot and Ralston [1], and follow the argument in [5]. In the proof of the main estimate of the trace the following lemma is crucial.

Lemma 3. Let $\rho \in C_{0}^{\infty}(-2,2)$ such that $\rho \geq 0$ for all $t$ and $\rho(t)=1$ for $t \in[-1,1]$. Suppose that $F_{\dagger}$ cannot be prolonged analytically to an entire function. Then there exists a positive constant $\alpha_{0}$ such that for any large $\beta>0$ we can find sequences $\left\{l_{q}\right\}_{q=1}^{\infty}$ and $\left\{m_{q}\right\}_{q=1}^{\infty}$ with the following properties:
(i)

$$
l_{q} \longrightarrow \infty \quad \text { as } q \longrightarrow \infty .
$$

(ii)

$$
e^{\beta l_{q}} \leq m_{q} \leq e^{2 \beta l_{q}}
$$

(iii) $\left|\left\langle\boldsymbol{\rho}_{q}, \hat{F}_{+}\right\rangle_{\mathscr{( \boldsymbol { R } _ { + } ) \times \mathscr { Q } ^ { \prime } ( \boldsymbol { R } _ { + } )}}\right| \geq e^{a_{1} l_{q}} \quad$ for all $q$,
where, $\hat{F}_{+}$is a distribution in $(0, \infty)$ given by

$$
\hat{F}_{\dagger}(t)=\sum_{r}(-1)^{a+i_{r}} T_{r}\left|I-P_{r}\right|^{-1 / 2} \delta\left(t-d_{r}\right)
$$

and $\rho_{q}(t)=\rho\left(m_{q}\left(t-l_{q}\right)\right)$.
In order to show Theorem 2 we shall make a rearrangement of the summation in (2.4), and use the results in [3, 4] on asymptotic behavior of phase functions and periodic rays in $\Omega$.

## References

[1] C. Bardos, J. C. Guillot, and J. Ralston: La relation de Poisson pour l'équation des ondes dans un ouvert non borné. Application à la théorie de la diffusion. Comm. Partial Diff. Equ., 7, 905-958 (1982).
[2] C. Gérard: Asymptotique des poles de la matrice de scattering pour deux obstacles strictement convexes. Univ. Paris-Sud (preprint).
[3] M. Ikawa: On the poles of the scattering matrix for two strictly convex obstacles. J. Math. Kyoto Univ., 23, 127-194 (1983).
[4] -: Decay of solutions of the wave equation in the exterior of several convex bodies (to appear in Ann. Inst. Fourier).
[5] --: On the poles of the scattering matrix for several strictly convex bodies (to appear).
[6] P. D. Lax and R. S. Phillips: Scattering Theory. Academic Press (1967).
[7] R. Melrose: Singularities and energy decay in acoustical scattering. Duke Math. J., 46, 43-59 (1979).


[^0]:    1) The original one is given in [6, page 158], but $\mathcal{O}$ considered in [4] is a counter example.
