114. Dual Pairs on Spinors

Cases of (C_m, C_n) and $(C_m^{(1)}, C_n^{(1)})$

By Kohji HASEGAWA

Department of Mathematics, Nagoya University

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§0. Introduction. Weyl's reciprocity theorem says that the symmetric group \mathfrak{S}_m and the general linear group $\operatorname{GL}(n, \mathbb{C})$ are mutually commutant (i.e. $(\mathfrak{S}_m, \operatorname{GL}(n, \mathbb{C}))$ forms a *dual pair* [3]) on the tensor space $(\mathbb{C}^n)^{\otimes m}$. The purpose of this paper is to give the spinor analogues of this theorem : we claim $(\mathfrak{sp}(2m), \mathfrak{sp}(2n))$ forms a dual pair on the $\mathfrak{o}(4mn)$ -module $\wedge (\mathbb{C}^{2mn})$, and describe its irreducible decomposition as a $\mathfrak{sp}(2m) \oplus \mathfrak{sp}(2n) - \operatorname{module}$ (Theorem A). The affine Lie algebra pair $(\mathbb{C}_m^{(1)}, \mathbb{C}_n^{(1)})$ also forms a dual pair on $\wedge (\hat{W}_{4mn})$ (Theorem B). As corollaries we deduce new dualities for branching rules. Details appear in our forthcoming paper [2], where we also construct various dual pairs for all classical Lie algebras, and for their affinizations. Our method is similar to that of [3], which deals with dual pairs on the Shale-Weil modules.

§1. Finite dimensional case. 1.1. After [1] we review the spinor representation of the orthogonal Lie algebra $\mathfrak{o}(2l) = \left\{ X \in \mathfrak{gl}(2l) \middle| {}^{t}X \begin{bmatrix} 0 & 1_{l} \\ 1_{l} & 0 \end{bmatrix} + \begin{bmatrix} 0 & 1_{l} \\ 1_{l} & 0 \end{bmatrix} X = 0 \right\}$. Let $\mathcal{C}(W_{2l})$ be the Clifford algebra over the vector space $W_{2l} := V^{l} \oplus V_{l} \simeq C^{2l}$, where $V^{l} := \bigoplus_{j=1}^{l} C\psi^{j}$ and $V_{l} := \bigoplus_{j=1}^{l} C\psi_{j}$, with a symmetric bilinear form (,) defined by

 $(\psi^i, \psi_i) = \delta^i_i$ and $(\psi^i, \psi^j) = 0 = (\psi_i, \psi_i)$ for $1 \le i, j \le l$.

As a C-algebra $\mathcal{C}(W_{2l}) \simeq \operatorname{Mat}(2^{l}, \mathbb{C})$. Its irreducible representation is realized on the exterior algebra $\wedge(V^{l})$, with defining 1 the vacuum vector and V^{l} (resp. V_{l}) the creation (resp. annihilation) operators. Write [a, b] for ab-ba, and the spinor representation s is defined by

$$s: \mathfrak{o}(2l) \ni \begin{bmatrix} E^{i}{}_{j} & 0\\ 0 & -E^{j}{}_{i} \end{bmatrix} \longmapsto \frac{1}{2} [\psi^{i}, \psi_{j}] \in \mathcal{C}(W_{2l}) \simeq \operatorname{End} \wedge (V^{l}),$$

$$\begin{bmatrix} 0 & E^{i}{}_{j} - E^{j}{}_{i}\\ 0 & 0 \end{bmatrix} \longmapsto \frac{1}{2} [\psi^{i}, \psi^{j}], \begin{bmatrix} 0 & 0\\ E^{i}{}_{j} - E^{j}{}_{i} & 0 \end{bmatrix} \longmapsto \frac{1}{2} [\psi_{i}, \psi_{j}] \quad (1 \leq i, j \leq l).$$
1.2. Now we deal with the dual pair ($\mathfrak{sp}(2m)$, $\mathfrak{sp}(2n)$). Recall that
$$\mathfrak{sp}(2n) := \left\{ X \in \mathfrak{gl}(2n) \mid {}^{t}X \begin{bmatrix} 1_{n} \\ -1_{n} \end{bmatrix} + \begin{bmatrix} 1_{n} \\ -1_{n} \end{bmatrix} X = 0 \right\}$$

$$= \left\{ \begin{bmatrix} A & B \\ C & -{}^{t}A \end{bmatrix} \middle| \substack{A, B, C \in \mathfrak{gl}(n) \\ {}^{t}B = B, {}^{t}C = C } \right\},$$

and let l=2mn. Then there exist two Lie algebra monomorphisms $R:\mathfrak{Sp}(2n) \to \mathfrak{o}(2l)$ and $L:\mathfrak{Sp}(2m) \to \mathfrak{o}(2l)$ so that $R(\mathfrak{Sp}(2n))' = L(\mathfrak{Sp}(2m))$ and $L(\mathfrak{Sp}(2m))'$

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= $R(\mathfrak{sp}(2n))$, where $A' := \{x \in \mathfrak{o}(2l) | [x, A] = 0\}$. Such R (resp. L) arises from the right (resp. left) action of $\mathfrak{sp}(2n)$ (resp. $\mathfrak{sp}(2m)$) on $Mat(2m \times 2n, C)$.

Considering
$$s \circ R$$
 and $s \circ L$, we get
Proposition 1. The map $\mathfrak{Sp}(2n) \rightarrow \mathcal{C}(W_{4mn}) \simeq \operatorname{End} \wedge (V^{2mn})$ defined by
 $\mathfrak{Sp}(2n) \ni \begin{bmatrix} E^{i}{}_{j} & 0\\ 0 & -E^{j}{}_{i} \end{bmatrix} \longrightarrow \frac{1}{2} \sum_{p=1}^{m} ([\psi^{i,p}, \psi_{j,p}] - [\psi^{-j,p}, \psi_{-i,p}]) \in \operatorname{End} \wedge (V^{2mn})$
 $\begin{bmatrix} 0 & E^{i}{}_{j} + E^{j}{}_{i} \\ 0 & 0 \end{bmatrix} \longrightarrow \frac{1}{2} \sum_{p=1}^{m} ([\psi^{i,p}, \psi_{-j,p}] + [\psi^{j,p}, \psi_{-i,p}])$
 $\begin{bmatrix} 0 & 0\\ E^{i}{}_{j} + E^{j}{}_{i} \end{bmatrix} \longrightarrow \frac{1}{2} \sum_{p=1}^{m} ([\psi^{-i,p}, \psi_{j,p}] + [\psi^{-j,p}, \psi_{i,p}])$ $(1 \le i, j \le n)$

is a Lie algebra monomorphim, and so is the map

$$\begin{split} \mathfrak{sp}(2m) \ni \begin{bmatrix} E^{p}_{q} & 0 \\ 0 & -E^{q}_{p} \end{bmatrix} & \longmapsto \frac{1}{2} \sum_{j=1}^{n} \left([\psi^{j,p}, \psi_{j,q}] + [\psi^{-j,p}, \psi_{-j,q}] \right) \in \mathrm{End} \wedge (V^{2mn}) \\ \begin{bmatrix} 0 & E^{p}_{q} + E^{q}_{p} \\ 0 & 0 \end{bmatrix} & \longmapsto \frac{1}{2} \sum_{j=1}^{n} \left([\psi^{j,p}, \psi^{-j,q}] + [\psi^{j,q}, \psi^{-j,p}] \right) \\ \begin{bmatrix} 0 & 0 \\ E^{p}_{q} + E^{q}_{p} & 0 \end{bmatrix} & \longmapsto \frac{1}{2} \sum_{j=1}^{n} \left([\psi_{j,p}, \psi_{-j,q}] + [\psi_{j,q}, \psi_{-j,p}] \right) & (1 \le p, q \le m) \end{split}$$

Here we use symbols $\psi^{\pm k,p}$ (resp. $\psi_{\pm k,p}$) $(1 \le p \le m, 1 \le k \le n)$ for the basis of V^i (resp. V_i), instead of ψ^j (resp. ψ_j) $(1 \le j \le l)$ in 1.1. Note that $\wedge (V^i) \simeq \wedge (V^{2n})^{\otimes m} \simeq \wedge (C^{2n})^{\otimes m}$ as $\mathfrak{Sp}(2n)$ -modules.

Let $\mathfrak{h} := \{h \in \mathfrak{Sp}(2m) \mid h \text{ is diagonal}\}, \mathfrak{n}_+ := \left\{ \begin{bmatrix} A & B \\ 0 & -{}^tA \end{bmatrix} \middle| A \text{ is strictly upper triangular, } {}^tB = B \right\}$ and $\mathfrak{n}_- := {}^t\mathfrak{n}_+$, then we fix a triangular decomposition $\mathfrak{Sp}(2n) = \mathfrak{n}_+ \oplus \mathfrak{h} \oplus \mathfrak{n}_-$. By $L^{\mathfrak{g}}(\lambda)$ we denote the irreducible g-module with highest weight λ . A sequence $Y = (y_1, \dots, y_m) \in \mathbb{Z}^m$ with $n \ge y_1 \ge \dots \ge y_m \ge 0$ is called a Young diagram contained in the $m \times n$ rectangle, and the set of all such sequences is denoted by \mathfrak{R}_{mn} . For $Y = (y_i) \in \mathfrak{R}_{mn}$ we set $|Y| := \Sigma_i y_i$, and define $Y^{\dagger} \in \mathfrak{R}_{nm}$ by taking the complement of Y in \mathfrak{R}_{mn} and transposing it: for example,

$$\mathbf{R}_{mn} \ni Y = \begin{bmatrix} y_1 \\ \vdots \\ y_m \end{bmatrix} = \begin{bmatrix} 3 \\ 3 \\ 2 \\ 0 \\ 0 \end{bmatrix} = \mathbf{m} \begin{bmatrix} \mathbf{m} \\ \mathbf{m} \\ \mathbf{m} \end{bmatrix} \longmapsto Y^{\dagger} = \begin{bmatrix} y_1^{\dagger} \\ \vdots \\ y_n^{\dagger} \end{bmatrix} := \mathbf{n} \begin{bmatrix} \mathbf{m} \\ \mathbf{m} \\ \mathbf{m} \\ \mathbf{m} \end{bmatrix} = \begin{bmatrix} 5 \\ 3 \\ 2 \\ 2 \end{bmatrix} \in \mathbf{R}_{nm}$$

$$(m = 5, n = 4).$$

Put $Y\left(\begin{bmatrix}E^{j_{j}}&\\&-E^{j_{j}}\end{bmatrix}\right):=y_{j}$ for $Y=(y_{1},\dots,y_{m})\in \mathbb{R}_{mn}$, and we identify a diagram with a dominant integral weight of $\mathfrak{Sp}(2m)$. Then our main result is

Theorem A. As a $\mathfrak{sp}(2m) \oplus \mathfrak{sp}(2n)$ -module,

 $\wedge (V^{2mn}) \simeq \bigoplus_{Y \in \mathbf{R}_{mn}} L^{\mathfrak{sp}(2m)}(Y) \otimes L^{\mathfrak{sp}(2n)}(Y^{\dagger}).$

The highest weight vector of the Y-component with respect to $\mathfrak{sp}(2m)$ $\oplus \mathfrak{sp}(2n)$ is $\wedge_{\tilde{y}_{j=1}^{p}} \psi^{p,j} \cdot 1$, where $Y \in \mathbb{R}_{mn}$ is identified with (\tilde{y}_{j}^{p}) as follows: K. HASEGAWA

All irreducible highest weight module of $\mathfrak{Sp}(2n)$ appears as the irreducible component of $\wedge(V^{2mn})$, when m varies over $\mathbb{Z}_{>0}$.

§ 2. Affine version. 2.1. Let $\hat{g} := g \otimes C[t, t^{-1}] \oplus C^{\circ \hat{g}}$ be the non-twisted affine Lie algebra associated to a simple Lie algebra g (see [5]). We review Frenkel's spinor representation of $\mathfrak{o}^{\circ}(2l)$ [1]. Let $C(\hat{W}_{2l})$ be the Clifford algebra over the *C*-vector space $\hat{W}_{2l} := W_{2l} \otimes C[t, t^{-1}]$ with a symmetric form $(\psi(\mu), \psi'(\mu')) := \delta_{\mu+\mu',0}(\psi, \psi')$, where $\psi(\mu) := \psi \otimes t^{\mu}$. Put $\hat{W}_{2l}^{+} := (W_{2l} \otimes tC[t]) \oplus$ $(V_l \otimes 1)$ and $\hat{W}_{2l}^{-} := (W_{2l} \otimes t^{-1}C[t^{-1}]) \oplus (V^l \otimes 1)$, then $\wedge(\hat{W}_{2l})$ becomes an irreducible $C(\hat{W}_{2l})$ -module by defining 1 the vacuum, \hat{W}_{2l}^{-} (resp. \hat{W}_{2l}^{+}) the creation (resp. annihilation) operators. Define

$$:a(\mu)b(\nu)::=a(\mu)b(\nu)-(a(\mu),b(\nu))\varepsilon, \text{ where } \varepsilon:=\begin{cases}1 & \text{if } \mu > 0 > \nu,\\1/2 & \text{if } \mu = 0 = \nu,\\0 & \text{otherwise}\end{cases}$$

and the spinor representation \hat{s} is defined by

$$\begin{split} \hat{s} : \hat{\mathfrak{o}}(2l) &\ni c^{\mathfrak{o}^{*}(2l)} \longmapsto id \in \operatorname{End} \wedge (\hat{W}_{\overline{2l}}) \\ X(k) :&= X \otimes t^{k} \longmapsto \Sigma_{\mu \in \mathbf{Z}} : a(\mu)b(k-\mu) :, \end{split}$$

for $X \in \mathfrak{o}(2l)$ that satisfies s(X) = (1/2)[a, b] $(a, b \in W_{2l})$ and $k \in \mathbb{Z}$.

2.2. We proceed to the dual pair $(C_m^{(1)}, C_n^{(1)})$. Again taking l=2mn we get a Lie algebra monomorphism \hat{L} (resp. \hat{R}) by defining

$$\begin{split} \hat{L} : \mathfrak{sp}^{*}(2m) \ni A(k) \longmapsto L(A)(k) \in \mathfrak{o}^{*}(2l) \\ c^{\mathfrak{sp}^{*}(2m)} \longmapsto n \cdot c^{\mathfrak{o}^{*}(2l)} \\ \end{split} \qquad \begin{pmatrix} \operatorname{resp.} \hat{R} : \mathfrak{sp}^{*}(2n) \ni A(k) \\ \longmapsto R(A)(k) \in \mathfrak{o}^{*}(2l) \\ c^{\mathfrak{sp}^{*}(2n)} \longmapsto m \cdot c^{\mathfrak{o}^{*}(2l)} \\ \end{pmatrix}$$

Proposition 2. The map $\hat{s} \circ \hat{L}$: $\mathfrak{sp}^{(2m)} \to \operatorname{End} \wedge (\hat{W}_{4mn})$ (resp. $\hat{s} \circ \hat{R}$: $\mathfrak{sp}^{(2n)} \to \operatorname{End} \wedge (\hat{W}_{4mn})$) is a level *n* (resp. *m*) integrable representation of $\mathfrak{sp}^{(2m)}$ (resp. $\mathfrak{sp}^{(2m)}$), and $[\hat{s} \circ \hat{L}(\mathfrak{sp}^{(2m)}), \hat{s} \circ \hat{R}(\mathfrak{sp}^{(2n)})] = 0$.

Theorem B. As a $\mathfrak{sp}^{(2n)} \oplus \mathfrak{sp}^{(2n)} - module$,

 $\wedge(\hat{W}_{4mn})\simeq \bigoplus_{Y\in \mathbf{R}_{mn}} L^{\mathfrak{sp}^{(2m)}}(Y,n)\otimes L^{\mathfrak{sp}^{(2n)}}(Y^{\dagger},m).$

The highest weight vector of the Y-component is $\wedge_{\tilde{y}_{j}^{p}=1}\psi^{p,j}(0)\cdot 1$, where (\tilde{y}_{j}^{p}) is as in Theorem A. All level m irreducible integrable highest weight module of $C_{n}^{(1)}$ appears as the irreducible component of $\wedge(\hat{W}_{4mn})$.

Here we write $L^{\hat{\mathfrak{g}}}(Y, n)$ for $L^{\hat{\mathfrak{g}}}(\lambda)$, when the highest weight $\lambda \in (\mathfrak{h} \otimes 1 \oplus Cc^{\hat{\mathfrak{g}}})^*$ satisfies $\lambda(c^{\hat{\mathfrak{g}}}) = n$ and $\lambda|_{\mathfrak{g} \otimes 1} = Y \in \mathbb{R}_{mn} \longrightarrow \mathfrak{h}^*$ (see 1.2).

Theorems A and B are shown by the Weyl-Kac character formula and its application derived by Jimbo-Miwa [4].

§ 3. Dualities of branching rules. We derive two affine cases here. First, noting $\wedge(\hat{W}_{4(l+m)n}) \simeq \wedge(\hat{W}_{4ln}) \otimes \wedge(\hat{W}_{4mn})$ we deduce

Corollary 3. Define the coset Virasoro module $\mathcal{A}_{y,y'}^{Y}$ (resp. $\mathcal{B}_{Y}^{y,y'}$) by

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$$\begin{bmatrix} 0 & 0 \\ E^{p}_{q} + E^{q}_{p} & 0 \end{bmatrix} (i+kl) \longmapsto \sum_{j=1}^{l-i} \begin{bmatrix} 0 & 0 \\ E^{p}_{q} \otimes E^{j}_{l-i-j+1} + E^{q}_{p} \otimes E^{l-i-j+1}_{j-i-j+1} & 0 \end{bmatrix} (k)$$

$$+ \sum_{j=l+1-i}^{l} \begin{bmatrix} 0 & 0 \\ E^{p}_{q} \otimes E^{j}_{2l-i-j+1} + E^{q}_{p} \otimes E^{2l-i-j+1}_{j-i-j+1} & 0 \end{bmatrix} (k+1),$$

$$c^{\mathfrak{sp}(2m)} \longmapsto c^{\mathfrak{sp}(2lm)} \qquad (1 < i < l, \ k \in \mathbb{Z} \ and \ 1 < p, \ q < m).$$

(ii) Let l=2. Then $\iota(\mathfrak{sp}^{(2m)})=\{x\in\mathfrak{sp}^{(4m)}|\sigma(x)=x\}$, where σ is the order 2 diagram automorphism of $C_{2m}^{(1)}$ [4].

Corollary 4. Define A_Y^{v} (resp. B_y^{v}) by $L^{\mathfrak{sp}^{(2n)}}(y,m) \simeq \bigoplus_{Y \in \mathbf{R}_{n,im}} A_Y^{v} \otimes L^{\mathfrak{sp}^{(2n)}}(Y,lm)$

$$(resp. L^{\mathfrak{sp}(2lm)}(Y, n) \simeq \bigoplus_{y \in \mathbf{R}_{m,n}} B_y^Y \otimes L^{\mathfrak{sp}(2m)}(y, n)).$$

Then $A_{Y^{\dagger}}^{y^{\dagger}} \simeq B_{y}^{Y}$ as a Virasoro module, for $Y \in \mathbb{R}_{lm,n}$ and $y \in \mathbb{R}_{m,n}$.

This is our second duality, which is shown by using "principal picture" on $\wedge(\hat{W}_{4mn})$. The case n=1 of Cor. 3 appears in [6] and [7], and the case l=2 of Cor. 4 appears in [4].

References

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