

64. On the Isomonodromic Deformation of a Linear Ordinary Differential Equation of the Third Order

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§ 1. Introduction. Consider a third order linear ordinary differential equation of Fuchsian type

$$(1.1) \quad \frac{d^3y}{dx^3} + p_1 \frac{d^2y}{dx^2} + p_2 \frac{dy}{dx} + p_3 y = 0$$

with the following Riemannian scheme :

$$(1.2) \quad \left\{ \begin{array}{ccccc} x=0 & x=1 & x=t & x=\lambda_j & x=\infty \\ \alpha_0 & \alpha_1 & \beta & \gamma_j & \alpha_\infty \\ \alpha_0 + \kappa_0 & \alpha_1 + \kappa_1 & \beta + \theta & \gamma_j + 2 & \alpha_\infty + \kappa_\infty \\ \alpha_0 + \kappa'_0 & \alpha_1 + \kappa'_1 & \beta + \theta' & \gamma_j + 3 & \alpha_\infty + \kappa'_\infty \end{array} \right\} \\ (j=1, 2, 3, 4)$$

and we suppose that *the singularities* $x=\lambda_j$ ($j=1, 2, 3, 4$) *are non-logarithmic ones and the characteristic exponents at each singular point do not differ by integer.*

The purpose of this paper is to derive a system of isomonodromic deformation equations of (1.1) regarding t as deformation parameter.

§ 2. Hamiltonian system attached to (1.1). The coefficients $p_j(x)$ ($j=1, 2, 3$) of the equation (1.1) are given by

$$\begin{aligned} p_1(x) &= \frac{a_0^1}{x} + \frac{a_1^1}{x-1} + \frac{b^1}{x-t} + \sum_{k=1}^4 \frac{c_k^1}{x-\lambda_k}, \\ p_2(x) &= \frac{a_0^2}{x^2} + \frac{a_1^2}{(x-1)^2} + \frac{b^2}{(x-t)^2} + \sum_{k=1}^4 \frac{c_k^2}{(x-\lambda_k)^2} \\ &\quad + \frac{a_\infty^2}{x(x-1)} - \frac{t(t-1)H}{x(x-1)(x-t)} + \sum_{k=1}^4 \frac{\lambda_k(\lambda_k-1)\mu_k}{x(x-1)(x-\lambda_k)}, \\ p_3(x) &= \frac{a_0^3}{x^3} + \frac{a_1^3}{(x-1)^3} + \frac{b^3}{(x-t)^3} + \sum_{k=1}^4 \frac{c_k^3}{(x-\lambda_k)^3} \\ &\quad + \frac{1}{T(x)} \left[a_\infty^3 + \eta_0 \frac{t}{x} - \eta_1 \frac{t-1}{x-1} + \eta_t \frac{t(t-1)}{x-t} \right. \\ &\quad \left. + \sum_{k=1}^4 \left\{ \frac{T(\lambda_k)}{(x-\lambda_k)^2} + \frac{T'(\lambda_k)}{x-\lambda_k} \right\} \xi_k + \sum_{k=1}^4 \frac{\zeta_k}{x-\lambda_k} \right], \end{aligned}$$

where

$$T(x) = x(x-1)(x-t)$$

and a_i^j, b^j, c_k^j ($i=1, 2, 3; k=1, 2, 3, 4; j=0, 1, \infty$) are constants de-

terminated by the characteristic exponents.

We see from the assumption that $x = \lambda_k$ are non-logarithmic singularities that η_d, ξ_j, ζ_j ($d=0, 1, t; j=1, 2, 3, 4$) and H are determined as rational functions of t, λ_k, μ_k ($k=1, 2, 3, 4$). Using H thus determined, we obtain the following theorem.

Theorem 1. *The isomonodromic deformation of (1.1) is governed by the Hamiltonian system*

$$H(\alpha_0, \alpha_1, \beta, \gamma_k) \begin{cases} \frac{d\lambda_j}{dt} = \frac{\partial H}{\partial \mu_j} \\ \frac{d\mu_j}{dt} = -\frac{\partial H}{\partial \lambda_j} \end{cases} \quad (j=1, 2, 3, 4)$$

with the Hamiltonian

$$H = -\operatorname{Res}_{x=t} p_2(x)$$

and the canonical variables conjugate to λ_j :

$$\mu_j = \operatorname{Res}_{x=\lambda_j} p_2(x) \quad (j=1, 2, 3, 4).$$

The explicit form of Hamiltonian function H for the system $H(0, 0, 0, 0)$ is given in the last section.

§ 3. Canonical transformation. Consider a transformation of (1.1) of the form

$$(3.1) \quad y = \Phi(x)z.$$

Putting in (3.1)

$$\Phi(x) = x^{\alpha_0}(x-1)^{\alpha_1}(x-t)^\beta \prod_{k=1}^4 (x-\lambda_k)^{\gamma_k},$$

we obtain by this transformation the linear equation

$$(3.2) \quad \frac{d^3z}{dx^3} + q_1 \frac{d^2z}{dx^2} + q_2 \frac{dz}{dx} + q_3 z = 0$$

having the Riemannian scheme (1.2) with

$$\alpha_0 = \alpha_1 = \beta = \gamma_k = 0 \quad (k=1, 2, 3, 4).$$

A simple computation shows that the coefficients $q_i(x)$ ($i=1, 2, 3$) of (3.2) are related to those of (1.1) as

$$(3.3) \quad \begin{aligned} q_1(x) &= p_1(x) + 3\Phi'\Phi^{-1}, \\ q_2(x) &= p_2(x) + [2p_1(x)\Phi' + 3\Phi'']\Phi^{-1}, \\ q_3(x) &= p_3(x) + [p_2(x)\Phi' + p_1(x)\Phi'' + \Phi''']\Phi^{-1}. \end{aligned}$$

Set

$$K = -\operatorname{Res}_{x=t} q_2(x), \quad \nu_k = \operatorname{Res}_{x=\lambda_k} q_2(x),$$

then the relation (3.3) reads

$$(3.4) \quad K = H - \operatorname{Res}_{x=t} (2p_1(x)\Phi' + 3\Phi'')\Phi^{-1},$$

$$(3.5) \quad \nu_k = \mu_k + \operatorname{Res}_{x=\lambda_k} (2p_1(x)\Phi' + 3\Phi'')\Phi^{-1}.$$

Then we can prove

Theorem 2. *The change of variables (3.4), (3.5):*

$$(\lambda_k, \mu_k, H) \longrightarrow (\lambda_k, \nu_k, K)$$

defines a canonical transformation, which takes the Hamiltonian system $H(\alpha_0, \alpha_1, \beta, \gamma_k)$ to $H(0, 0, 0, 0)$.

Remark 1. The transformation in the above theorem is invertible. Hence the Hamiltonian systems $H(\alpha_0, \alpha_1, \beta, \gamma_k)$ are transformed to each other by the canonical transformation.

Remark 2. The transformation (3.1) with

$$\Phi(x) = \exp\left(-\frac{1}{3} \int^x p_1(x) dx\right)$$

takes (1.1) into the linear equation (3.2) with $q_1(x) \equiv 0$. The linear equation of this form is called of *SL-type*.

§ 4. Hamiltonian. We will give the explicit form of the Hamiltonian function for the system $H(0, 0, 0, 0)$. Suppose that $\alpha_0 = \alpha_1 = \beta = \gamma_k = 0$ for the equation (1.1). Then the condition that $x = \lambda_j$ ($j = 1, 2, 3, 4$) are non-logarithmic singular points reads as

$$(4.1) \quad \xi_j = -2(\mu_j + E_j)$$

$$(4.2) \quad \zeta_j = \frac{1}{2} \xi_j(\mu_j + \xi_j), \quad (j = 1, 2, 3, 4)$$

$$(4.3) \quad \xi_j(\xi_j + F_j) + 2G_j = 0,$$

where E_j, F_j and G_j are constant terms in the Laurent series expansion of $p_1(x), p_2(x)$ and $p_3(x)$ at $x = \lambda_j$ respectively.

Solving (4.3) with respect to H , we arrive at the

Proposition. *The Hamiltonian function H for the system $H(0, 0, 0, 0)$ is given by*

$$\left[t(t-1) \sum_{k=1}^4 (\mu_k + E_k) \frac{T(\lambda_k)}{A'(\lambda_k)} \right] H = \sum_{j=1}^4 A_j \frac{T(\lambda_j)}{A'(\lambda_j)},$$

where

$$\begin{aligned} A_j &= (\mu_j + E_j) T(\lambda_j) \left[(\mu_j + 2E_j) \left(\mu_j + E_j - \frac{T'(\lambda_j)}{T^2(\lambda_j)} \right) \right. \\ &\quad \left. - \left(\frac{1}{\lambda_j} + \frac{1}{\lambda_j - 1} \right) \mu_j + \sum_{\substack{k=1 \\ k \neq j}}^4 \frac{\lambda_k(\lambda_k - 1)\mu_k}{\lambda_j(\lambda_j - 1)(\lambda_j - \lambda_k)} + U_j \right] \\ &\quad + \sum_{\substack{k=1 \\ k \neq j}}^4 \frac{\mu_k + E_k}{\lambda_j - \lambda_k} \left[\mu_k + 2E_k - 2T'(\lambda_k) - \frac{2T(\lambda_k)}{\lambda_j - \lambda_k} \right] + \alpha_\infty^3, \\ U_j &= \frac{\alpha_0^2}{\lambda_j^2} + \frac{\alpha_1^2}{(\lambda_j - 1)^2} + \frac{b^2}{(\lambda_j - t)^2} + \sum_{\substack{k=1 \\ k \neq j}}^4 \frac{2}{(\lambda_j - \lambda_k)^2} + \frac{\alpha_\infty^2}{\lambda_j(\lambda_j - 1)} + \frac{T'(\lambda_j)}{T(\lambda_j)} \end{aligned}$$

and

$$A(x) = \prod_{j=1}^4 (x - \lambda_j).$$

Remark 3. ξ_j, ζ_j ($j = 1, 2, 3, 4$) are determined by (4.1), (4.2) as rational functions of λ_k, μ_k ($k = 1, 2, 3, 4$) and η_d ($d = 0, 1, t$) by (4.3).

References

- [1] Fuchs, R.: Über lineare homogene Differentialgleichungen zweiter Ordnung mit drei im Endlichen gelegene wesentlich singularen Stellen. *Math. Ann.*, **63**, 301 (1907).
- [2] Garnier, R.: Sur des équations différentielles du troisième ordre dont l'intégrale générale est uniforme et sur une classe d'équations nouvelles d'ordre supérieur dont l'intégrale générale a ses points critiques fixes. *Ann. Scient. Éc. Norm. Sup.*, (3) **29**, 96 (1912).
- [3] Okamoto, K.: Sur le problème de Fuchs sur un tore, II. *J. Fac. Sci. Univ. Tokyo, Sec. IA*, **24**, 357 (1977).
- [4] Jimbo, M., Miwa, T., and Ueno, K.: Monodromy preserving deformation of linear ordinary differential equations with rational coefficients, I. *Physica*, **2D**, 306 (1981).