

54. Evans-Selberg's Theorem on Abstract Riemann Surfaces with Positive Boundaries. II

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Our $N_{V_{m\langle p \rangle}}(p, q)$ is increasing with respect to m . We define the value of $N(z, q)$ at a minimal point p by $\lim_{m \rightarrow M'} N_{V_{m\langle p \rangle}}(p, q)$ denoted by $N(p, q)$. If p or q belongs to R , this definition is equivalent to that defined before.

If $V_m(p)$ is not regular, we define $N_{V_{m\langle p \rangle}}(p, q)$ by $\lim_{m' \rightarrow m} N_{V_{m'\langle p \rangle}}(p, q)$, where $m' < m$ and $V_{m'}(p)$ is regular. In the case when $V_m(p)$ is regular, it is proved that $\lim_{m' \rightarrow m} N_{V_{m'\langle p \rangle}}(p, q) = N_{V_{m\langle p \rangle}}(p, q)$, hence we can define $N_{V_{m\langle p \rangle}}(p, q)$ for every $m < \sup_{z \in R} N(z, p) = M'$. As in case of a Riemann surface with a null-boundary, we can prove the following

Theorem 10. 1) $N(z, q)$ ($q \in \bar{R}$) is δ -lower semicontinuous in $R + B_1$.

2) $N(z, q)$ is superharmonic in weak sense at every point of $R + B_1$.

3) If p and q are in $R + B_1$, then $N(p, q) = N(q, p)$.

Till now $N(z, q)$ ($q \in \bar{R}$) is defined only on $R + B_1$. Next we define $N(z, q)$ at points belonging to B_0 . If $p \in B_0$, $N(z, p) = \int_{B_1} N(z, p_\alpha) d\mu(p_\alpha)$ ($p_\alpha \in B_1$) by Theorem 8. Although the uniqueness of this mass distribution is not proved by the present author, the value of $N(z, q)$ in $R + B_1$ is uniquely determined. On the other hand, by 3), for $q \in B_1$, $N(p_\alpha, q) = N(q, p_\alpha)$. Hence it is quite natural to define the value of $N(z, q)$ at $p \in B_0$ by $\int N(p_\alpha, q) d\mu(p_\alpha)$. Evidently by 3), in such definition, we have $N(q, p) = N(p, q)$, where the term of the right hand side does not depend on a particular distribution but on the behaviour of $N(z, q)$, because $N(p, q) = \lim_{m \rightarrow M'} N_{V_{m\langle p \rangle}}(p, q)$ and $N_{V_{m\langle p \rangle}}(p, q)$ is defined by the value of $N(z, q)$ on $\partial V_m(p)$. As for the behaviour of $N(z, q)$ ($q \in \bar{R}$), we have the following

Theorem 11. 1) If $q \in R + B_1$, then $N(p, q) = N(q, p)$ for $p \in \bar{R}$.

2) If $q \in \bar{R}$ and $p \in R + B_1$, then $N(p, q) = \int N(p, q_\alpha) d\mu(q_\alpha)$, where $N(z, q) = \int N(z, q_\alpha) d\mu(q_\alpha)$.

3) $N(z, q)$ ($q \in \bar{R}$) is δ -lower semicontinuous in \bar{R} .

1') For every p and q belonging to \bar{R} , $N(p, q) = N(q, p)$.

7. Potentials on \bar{R} . In the sequel, we shall study the mass distributions on \bar{R} . We have seen that $N(z, p)$ has the essential properties of logarithmic potential: lower semicontinuity in \bar{R} , symmetricity and superharmonicity in $R+B_1$. But there exists the fatal difference between our case and space, that is, the real mass distribution can be defined only on $R+B_1$, i.e. the distribution on B_0 is superficial and it can be replaced, by Theorem 8, by that on B_1 where $N(z, p)$ is superharmonic. Therefore only subsets $R+B_1$ of \bar{R} can be a kernel of mass distribution. Hence it is easy to construct the potential theory on \bar{R} .

The energy integral $I(\mu)$ of a mass distribution μ on a δ -closed subset of $R+B_1$ defined as in space

$$I(\mu) = \iint N(q, p) d\mu(p) d\mu(q)$$

and the $\overset{*}{\text{cap}}$ acity is defined as usual. In § 1, we defined capacity of F , we must study the relation between two capacities. At first, we have, if $\text{Cap}(F) > 0$, $\overset{*}{\text{Cap}}(F) > 0$. Now we have the following

Theorem 12. Let F be a δ -closed subset of $R+B_1$ of $\overset{*}{\text{cap}}$ acity positive. Then there exists a unit mass distribution on F whose energy is minimal and whose potential $U(z)$ has the following properties:

- 1) $U(z)$ is a constant C on the kernel of this distribution.
- 2) $U(z) = C$ on E except possibly a set of $\overset{*}{\text{cap}}$ acity zero.
- 3) $U(z) = U_F(z)$.
- 4) $U(z) = C\omega_F(z)$, where $\omega_F(z)$ is the equilibrium potential of F .

By 2) of this theorem and by 2) of Theorem 5, we have the following

Corollary. $\text{Cap}(F) = \overset{*}{\text{Cap}}(F)$.

Transfinite diameter. Since $N(z, p)$ ($p \in \bar{R}$) is δ -lower semicontinuous in \bar{R} , the transfinite diameter of a δ -closed subset A of \bar{R} is defined as follows:

$$\frac{1}{D_A} = \lim_{n \rightarrow \infty} \left(\min \left(\frac{1}{nC_2} \left(\sum_{\substack{i, j=1 \\ p_i, p_j \in A}}^{n, n} (p_i, p_j) \right) \right) \right).$$

Then as in the case of R^* with a null-boundary, we have the following

Theorem 13. If $D_A = 0$ for a δ -closed subset A of \bar{R} , then there exists a superharmonic function $U(z)$ in \bar{R} such that $U(z) = 0$ on ∂R_0 , $\int_0 \frac{\partial U(z)}{\partial n} ds = 2\pi$ and $U(z) = \infty$ at every point of A .

For a δ -closed subset A of $R+B_1$, it can be proved as in space

$D_A = \frac{1}{I(\mu)}$, where $I(\mu)$ is the energy of the equilibrium potential of

A. Hence we have the following

Theorem 14 (Extension of Evans-Selberg's theorem). Let A be a δ -closed subset of $R+B_1$, of capacity zero. Then there exists a unit mass distribution on A whose potential satisfies the following properties:

- 1) $U(z) = 0$ on ∂R_0 .
- 2) $U(z) = \infty$ at every point of A .
- 3) $U(z) = U_A(z)$.
- 4) $\int_{\partial R_0} \frac{\partial U(z)}{\partial n} ds = \int_{C_r} \frac{\partial U(z)}{\partial n} ds$, for the niveau curve C_r of $U(z)$

with $r \in E$, where E is a set in the interval $[0, \infty]$ such that $\text{mes } E = 0$.

In general cases we can not omit the condition that A is a subset of $R+B_1$. The reason is as follows: there may exist a set B_0 which is an F_σ and of capacity zero and any mass can not be distributed on B_0 , in other words, B_0 has behaviour like an empty set in space for mass distribution though B_0 is not empty.

The value of $U(z)$ at a point $p \in B_0$ is given as follows: since $N(z, p) = \int_{B_1} N(z, p_\alpha) d\mu(p_\alpha) (p_\alpha \in B_1)$, $U(p) = \int_{B_1} U(p_\alpha) d\mu(p_\alpha)$. Therefore $U(z)$ may be infinite at larger set A' than A .