

### 113. The Initial Value Problem for Linear Partial Differential Equations with Variable Coefficients. III

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In the present paper we consider mixed problems of linear parabolic equations with boundary conditions formulated by J. L. Lions [4] such that, using his notations,  $V$  is independent of the time variables, but  $N_t$  depends on them.

Our methods (§2), are also applicable to mixed problem of linear equations of many other types with above-mentioned boundary conditions, with which it seems interesting to me to compare Kato's methods [2].

As an illustration of our considerations we consider in §3 the Fokker-Planck's equations formulated by K. Yosida [7].

Only a sketch of this proof will be given, however, the details with further investigations will be published elsewhere.

1. Preliminary. Let  $\mathcal{Q}$  be a domain of the Euclidean space. Let  $((u, v))_t$  be real bilinear forms defined on a real separable Hilbert space  $V$  with following conditions, where  $\mathfrak{D}(\mathcal{Q}) \subset V \subset L_2(\mathcal{Q})$  and the injections  $\mathfrak{D}(\mathcal{Q}) \rightarrow V$ ,  $V \rightarrow L_2(\mathcal{Q})$  are both continuous: there are positive constants  $a, b, c$  such that for any  $t (-\infty < t < \infty)$

$$(I) \quad \begin{aligned} ((u, u))_t &\geq a \|u\|_V^2 \\ b \|u\|_V \|v\|_V &\geq |((u, v))_t|, \end{aligned}$$

(II) for fixed  $u, v \in V$ ,

$$c |t - t'| \|u\|_V \|v\|_V \geq |((u, v))_t - ((u, v))_{t'}|.$$

Furthermore let  $\bar{A}_t$  be an operator in  $L_2(\mathcal{Q})$  into itself such that  $f \in D(\bar{A}_t)$  if and only if  $((f, v))_t = (\bar{A}_t f, v)_{L_2(\mathcal{Q})}$  for every  $v \in V$ , where  $(\bar{A}_t f, v)_{L_2(\mathcal{Q})} = A_t f(v)$  for the distribution  $A_t f$  defined by the relation:  $((f, v))_t = A_t f(v)$  for every  $v \in \mathfrak{D}(\mathcal{Q})$ . Then  $\bar{A}_t$  is a densely defined, closed operator in  $L_2(\mathcal{Q})$  into itself whose adjoint coincides with operator  $\bar{A}_t^*$  defined as above from  $((u, v))_t^* = ((v, u))_t$  [3, 4]. Let  $G_t^*$  be the Green operators with respect to the form  $((u, v))_t^*$ . Then from (II) we see the following

**Lemma 1.** For any  $u \in L_2(\mathcal{Q})$ ,  $G_t^* u$  is differentiable in  $V$  (a.e.t) and  $\frac{d}{dt} G_t^* u$  is measurable with respect to  $V$  such that

$$\left\| \frac{d}{dt} G_t^* u \right\|_V \leq c \|u\|_V \quad (\text{a.e.t}).$$

**Definition.** Let  $E$  be a real separable Hilbert space. Then we denote by  $\mathfrak{Q}^n(E)$  the completion of the real linear space  $\mathfrak{D}_t(E)$  with

the following norm:

$$\int_{-\infty}^{\infty} ((q^n u(t), u(t)))_E dt = \int_{-\infty}^{\infty} (1 + |\xi|^2)^n \|\mathfrak{F}(u)(\xi)\|_E^2 d\xi,$$

where  $q = \alpha - \frac{d^2}{dt^2}$  ( $\alpha > 0$ ),  $\mathfrak{F}(u)$  is the Fourier transform of  $u$  and where  $\|\mathfrak{F}(u)(\xi)\|_E$  is the complex Hilbert space-norm extended from the real norm of  $E$ .

**Lemma 2.** *For any integers  $s, m, q^s$  is an isomorphism from  $\mathfrak{Q}^{m+s}(E)$  onto  $\mathfrak{Q}^{m-s}(E)$ . Furthermore  $\frac{d}{dt}$  is a continuous operator from  $\mathfrak{Q}^m(E)$  into  $\mathfrak{Q}^{m-1}(E)$ .*

2. Parabolic equations. Let  $((u, v))_t$  be a real bilinear form defined on  $V$  with the following properties: there are positive constants  $\alpha$  such that the bilinear form

$$((u, v))_t + \alpha(u, v)_{L_2(\Omega)}$$

satisfies the conditions (I), (II) and in §1. Then for sufficiently large  $\alpha$  and  $\beta, q = \beta - \frac{d}{dt^2},$

$$\int_{-\infty}^{\infty} ((qu, u))_t dt \geq \gamma \int_{-\infty}^{\infty} (qu, u)_{L_2(\Omega)} dt \quad (\gamma > 0)$$

where  $u = u(t) \in \mathfrak{D}_t(V)$ . Therefore setting

$$((u, v))_t = ((u, v))_t + \alpha(u, v)_{L_2(\Omega)}$$

we see the following

**Lemma 3.** *For any  $u(t) \in q^{-1}G_t^*(\mathfrak{D}((-\infty, \infty) \times \Omega))$  and for sufficiently large  $\beta$ , and for some  $\gamma(\beta) > 0$*

$$\int_{-\infty}^{\infty} \int_{\Omega} \left(-\frac{d}{dt} + \bar{A}_t^*\right) q u(t) \cdot u(t) dx dt \geq \gamma(\beta) \int_{-\infty}^{\infty} \int_{\Omega} qu \cdot u dx dt.$$

From Lemma 3 it follows that for any  $u(t) \in q^{-1}G_t^*(\mathfrak{D}(-\infty, \infty) \times \Omega)$

$$\left\| q^{-1} \left(-\frac{d}{dt} + \bar{A}_t^*\right) q u \right\|_{\mathfrak{Q}^1(L_2(\Omega))} \geq \gamma' \|u\|_{\mathfrak{Q}^1(L_2(\Omega))} \quad (\gamma' > 0).$$

Therefore by a limit process we see the following

**Theorem 1.** *For any  $g(t) \in \mathfrak{Q}^1(L_2(\Omega))$  there is a solution  $f(t) \in \mathfrak{Q}^1(L_2(\Omega)) \cap \mathfrak{Q}^0(V)$  such that*

$$f(t) \in D(\bar{A}_t) \quad (a.e.t)$$

$$\left(\frac{d}{dt} + \bar{A}_t\right) f(t) = g(t) \quad (a.e.t).$$

Furthermore such a solution  $f(t)$  satisfies the following inequality:

$$\left\| \left(\frac{d}{dt} + \bar{A}_t\right) f \right\|_{\mathfrak{Q}^0(L_2(\Omega))} \geq \gamma'' \|f\|_{\mathfrak{Q}^0(L_2(\Omega))} \quad (\gamma'' > 0),$$

which implies the uniqueness of solution in Theorem 1. Moreover, since  $\mathfrak{Q}^1(L_2(\Omega)) \subset C_t(L_2(\Omega))$ , we see that for such a solution  $f(t)$ , and for  $b, a$  ( $b > a$ )

$$\|f(b)\|_{L_2(\Omega)}^2 \leq 2 \left\{ \int_a^b \|g(t)\|_{L_2(\Omega)}^2 dt \right\}^{\frac{1}{2}} \left\{ \int_a^b \|f(t)\|_{L_2(\Omega)}^2 dt \right\}^{\frac{1}{2}} + \|f(a)\|_{L_2(\Omega)}^2.$$

Thus from the above inequalities and Theorem 1, by a limit process analogous as one used in my former paper [6] we see the following

**Theorem 2.** *For any  $g \in \mathfrak{L}^1(L_2(\Omega))$  such that  $g(t)=0 \ t < a$  and for any  $b > a$ , there is a unique solution  $f(t) \in \mathfrak{L}^1(L_2(\Omega))[a, b] \cap \mathfrak{L}^0(V)[a, b] \cap C_t(L_2(\Omega))[a, b]$  such that*

$$\begin{aligned} f(t) &\in D(\bar{A}_t) \quad (a.e.t) \\ \left(\frac{d}{dt} + \bar{A}_t\right)f(t) &= g(t) \quad (a.e.t) \text{ in } L_2(\Omega) \\ f(a) &= 0 \text{ in } L_2(\Omega). \end{aligned}$$

Here we remark that if  $((u, v))_t$  satisfy furthermore a condition with respect to perturbations (see § 3), then Theorem 2\* can be strengthened.

3. Example (Fokker-Planck's equations). For the sake of simplicity let  $\Omega$  be a bounded domain with sufficiently smooth boundary  $S$  in the Euclidean space  $R^N (N \geq 2)$ . Let  $A_t$  be the following differential operator: for any sufficiently smooth function  $f$ ,

$$A_t f = \frac{\partial^2}{\partial x^i \partial x^j} (b^{ij}(t, x) f(x)) + \frac{\partial}{\partial x^i} (-a^i(t, x) f(x)) \text{ on } \Omega$$

$f \in D(A_t)$  if and only if

$$b^{ij}(t, x) \frac{\partial f}{\partial x^j} \pi^i(x) + \left(\frac{\partial b^{ij}}{\partial x^j}(t, x) - a^i(t, x)\right) \pi^i(x) f(x) (= B(f, 1)) = 0 \text{ on } S,$$

where  $b^{ij}(t, x)$ ,  $a^i(t, x)$  are sufficiently smooth real functions defined on  $[0, T] \times \bar{\Omega}$ ,  $\pi^i(x) = \cos(n(x), x^i)$  on  $S$ .

Furthermore we assume that

$$b^{ij}(t, x) \xi_i \xi_j > 0$$

for any real  $\xi_i : \sqrt{\sum \xi_i^2} \neq 0$  and for  $(t, x) \in [0, T] \times \bar{\Omega}$ . Then we see that for any  $f \in D(A_t) \cap C^2(\bar{\Omega})$ ,  $v \in C^1(\bar{\Omega})$ ,

$$\begin{aligned} (A_t f, v)_{L_2(\Omega)} &= - \left( b^{ij}(t, x) \frac{\partial}{\partial x^j} f(x), \frac{\partial}{\partial x^i} v(x) \right)_{L_2(\Omega)} \\ &\quad + \left( a^i(t, x) f(x), \frac{\partial}{\partial x^i} v(x) \right)_{L_2(\Omega)} \\ &\quad - \left( \frac{\partial b^{ij}}{\partial x^j}(t, x) f, \frac{\partial}{\partial x^i} v(x) \right)_{L_2(\Omega)} \\ &\quad + (B(f, 1), v)_{L_2(\Omega)}. \end{aligned}$$

Let  $((u, v))_t$  be the following:

$$\begin{aligned} &\left( b^{ij}(t, x) \frac{\partial}{\partial x^j} u(x), \frac{\partial}{\partial x^i} v(x) \right)_{L_2(\Omega)} \\ &\quad + \left( u(x), \left( -a^i(t, x) + \frac{\partial b^{ij}}{\partial x^j}(t, x) \right) \cdot \frac{\partial}{\partial x^i} v \right)_{L_2(\Omega)}. \end{aligned}$$

Then, using some extending  $b^{ij}$ ,  $a^i$ , we see that the bilinear form

$((u, v))_t$ , satisfies the condition in §2, in fact, that in the end of §2, with  $V = H^1(\bar{\Omega})$ .

Furthermore we see the following

**Lemma 4.** *Let  $u(t) \in \mathcal{V}^1(L_2(\Omega))$  be a solution such that*

$$\left(\frac{\partial}{\partial t} - A_t\right)u(t) = 0 \quad (\text{a.e.t.})$$

$$u(t) \in D(\bar{A}_t) \quad (\text{a.e.t.}),$$

then  $u(s) \geq 0$  whenever  $u(t) \geq 0$  for some  $t (s > t)$ .

For, let  $h(t, x)$  be the following function such that  $h(t, x) = 1, -1$  and  $0$  when  $u(t, x) > 0, < 0$  and  $= 0$  respectively.

Then for  $b > a$

$$\begin{aligned} \int_a^b \left\| \left( \frac{\partial}{\partial t} + A \right) u \right\|_{L_1(\Omega)} dt &\geq \int_a^b \int_{\Omega} h(t, x) \cdot \frac{\partial}{\partial t} u(t, x) dx dt \\ &\quad - \int_a^b \int_{\Omega} h(t, x) \cdot A(t, x) u(t, x) dx dt. \end{aligned}$$

Furthermore by the regularity of solutions of elliptic equations on the boundary [1, 5] and by Yosida's lemma [8] we see that

$$\int_{\Omega} h(t, x) \cdot A(t, x) f(t, x) dx \leq 0 \quad (\text{a.e.t.})$$

Therefore  $0 \geq \int_{\Omega} \int_a^b h(t, x) \cdot \frac{\partial}{\partial t} u(t, x) dt dx = \|f(b)\|_{L_1(\Omega)} - \|f(a)\|_{L_1(\Omega)}$ .

Furthermore from the type of  $A$  we see that

$$\int_{\Omega} f(b) dx = \int_{\Omega} f(a) dx.$$

Thus we see that the mapping  $u(t) \rightarrow u(s) (s > t)$  is norm preserving with respect to  $L_1(\Omega)$  and  $u(s) \geq 0$ , whenever  $u(t) \geq 0$ .

Now let  $T_{s,t}(u) = u(s) (s > t)$  when  $u(s)$  is a solution of our equation such that  $u(t) = u$  and  $u(s) \in \mathcal{V}^1(L_2(\Omega))(t, T)$ . Then from Lemma 4 and Theorem 2, we see that these  $T_{s,t}$  are extended over  $L_1(\Omega)$  into itself and that these extensions are transition operators on  $L_1(\Omega)$ . Therefore by the hypoellipticity of parabolic equations [6] and the regularity of the solutions of elliptic equations on the boundary [1, 5] we see the following

**Theorem 3.** *The diffusion problem with the differential operators  $A_t$  as generating operators can be solved, i.e. there is a transition density  $p(t, s, x, y)$  such that for  $s > t$  and for  $f \in L_1(\Omega)$  setting*

$$T_{s,t}f(y) = \int_{\Omega} p(t, s, x, y) f(x) dx,$$

$T_{s,t}$  is a transition operator [ $0 < t < s < T$ ] and such that for any  $\tilde{f} \in D(\bar{A}_t)$

$$\left(\frac{\partial}{\partial t} - A_s\right)T_{s,t}\tilde{f} = 0 \quad \text{on } [0, T] \times \Omega$$

$$\begin{aligned} T_{t,t} \tilde{f} &= \tilde{f} && \text{on } L_2(\mathcal{Q}) \\ B_s(T_{s,t} \tilde{f}, 1) &= 0 && \text{(a.e. } x \in S) \text{ for all } s \in [0, T]. \end{aligned}$$

Finally we remark that our consideration can be applicable to other diffusion problems.

### References

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