

97. On the Generalized Loxodromes in the Conformally Connected Manifold.

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In a previous paper¹⁾, we have defined a conformal arc length σ on the curves in the conformally connected manifold and have proved the Frenet fomulae

$$(1) \quad \left\{ \begin{array}{l} \frac{d}{d\sigma} A_{(0)} = A_{(1)}, \\ \frac{d}{d\sigma} A_{(1)} = \lambda A_{(0)} + A_{(2)}, \\ \frac{d}{d\sigma} A_{(2)} = \lambda A_{(1)} + A_{(3)}, \\ \frac{d}{d\sigma} A_{(3)} = A_{(0)} + \lambda^4 A_{(4)}, \\ \frac{d}{d\sigma} A_{(4)} = -\lambda^4 A_{(3)} + \lambda^5 A_{(5)}, \\ \dots\dots\dots \\ \dots\dots\dots \\ \frac{d}{d\sigma} A_{(\infty)} = -\lambda^{\infty} A_{(n)}. \end{array} \right.$$

The conformal arc length does not exist for the generalized circles.

The most simple curves having the conformal arc length are the ones for which we have

$$(2) \quad \lambda = \lambda^4 = \lambda^5 = \dots = \lambda^n = \lambda^{\infty} = 0.$$

We shall, first, consider the properties of these curves.

§ 1. Substituting the equations (2) in the Frenet formulae (1), we have

$$(3) \quad \frac{d}{d\sigma} A_{(0)} = A_{(1)}, \quad \frac{d}{d\sigma} A_{(1)} = A_{(2)}, \quad \frac{d}{d\sigma} A_{(2)} = A_{(3)}, \quad \frac{d}{d\sigma} A_{(3)} = A_{(0)},$$

which shows that, if we develop the curve on the tangent conformal space at $A_{(0)}$, we shall obtain a curve lying always on a two-dimensional sphere determined by $A_{(0)}, A_{(1)}, A_{(2)}$ and $A_{(3)}$.

Considering always the development of the curve on the tangent conformal space at a fixed point of the conformally connected manifold, we can treat such curves as if they were on a two-dimensional flat conformal space.

1) K. Yano and Y. Mutô: On the conformal arc length, Proc. **17** (1941), 318-322.

Now, putting

$$(4) \quad \begin{cases} P = A_{(0)} + A_{(1)} + A_{(2)} + A_{(3)}, \\ Q = A_{(0)} - A_{(1)} + A_{(2)} - A_{(3)}, \end{cases}$$

we have

$$(5) \quad \frac{d}{d\sigma} P = P \quad \text{and} \quad \frac{d}{d\sigma} Q = -Q,$$

consequently, we can see that P and Q are two fixed points.

The circle passing through the three points $A_{(0)}$, P and Q being $A_{(1)} - A_{(3)}$, the angle θ between the curve and $A_{(1)} - A_{(3)}$ is given by

$$\sin \theta = \frac{A_{(1)}(A_{(1)} - A_{(3)})}{\sqrt{A_{(1)}A_{(1)}} \sqrt{(A_{(1)} - A_{(3)})(A_{(1)} - A_{(3)})}} = \frac{1}{\sqrt{2}}$$

from which we find $\theta = \frac{\pi}{4}$. Thus we have the

Theorem I. The curve whose all conformal curvatures $\lambda, \lambda^4, \dots, \lambda^n, \lambda^\infty$ vanish cuts the circles passing through two fixed points always by the fixed angle $\frac{\pi}{4}$, consequently, it is a loxodrome.

The four points $A_{(0)}$, $A_{(2)}$, P and Q lying on the circle $A_{(1)} - A_{(3)}$, the points $A_{(0)}$ and $A_{(2)}$ are harmonically separated by the points P and Q on this circle.

We shall now find the differential equations of these generalized loxodromes.

Differentiating the equation

$$(6) \quad A_{(0)} = \left(\frac{d\sigma}{ds} \right) A_0,$$

and making use of the formulae

$$\begin{cases} dA_0 = & du^i A_i, \\ dA_j = \Pi_{jk}^0 du^k A_0 + \Pi_{jk}^i du^k A_i + \Pi_{jk}^\infty du^k A_\infty, \\ dA_\infty = & \Pi_{\infty k}^i du^k A_i, \end{cases}$$

which define the conformal connection, we find successively

$$(7) \quad A_{(1)} = \frac{d^2\sigma}{d\sigma^2} A_0 + \frac{du^i}{ds} A_i,$$

$$(8) \quad A_{(2)} = \left[\frac{\frac{d^3\sigma}{ds^3}}{\left(\frac{d\sigma}{ds}\right)^2} - \frac{\left(\frac{d^2\sigma}{ds^2}\right)^2}{\left(\frac{d\sigma}{ds}\right)^3} + \frac{\Pi_{jk}^0 \frac{du^j}{ds} \frac{du^k}{ds}}{\frac{d\sigma}{ds}} \right] A_0 + \left[\frac{\frac{\partial^2 u^i}{\partial s^2}}{\frac{d\sigma}{ds}} + \frac{\frac{d^2\sigma}{ds^2} \frac{du^i}{ds}}{\left(\frac{d\sigma}{ds}\right)^2} \right] A_i + \frac{1}{\frac{d\sigma}{ds}} A_\infty,$$

$$(9) \quad A_{(3)} = \frac{1}{\left(\frac{d\sigma}{ds}\right)^2} \left[g_{jk} \frac{\partial^3 u^j}{\partial s^3} \frac{\partial^2 u^k}{\partial s^2} + \Pi_{jk}^0 \frac{\partial^2 u^j}{\partial s^2} \frac{du^k}{ds} \right] A_0 + \frac{1}{\left(\frac{d\sigma}{ds}\right)^2} \left[\frac{\partial^3 u^i}{\partial s^3} + \frac{du^i}{ds} g_{jk} \frac{\partial^2 u^j}{\partial s^2} \frac{\partial^2 u^k}{\partial s^2} - \frac{du^i}{ds} \Pi_{jk}^0 \frac{du^j}{ds} \frac{du^k}{ds} + \Pi_{\infty k}^i \frac{du^k}{ds} \right] A_i^*),$$

where we have denoted by $\partial/\partial s$ the covariant differentiation along the curve with respect to the Christoffel symbols $\Pi_{jk}^i = \{i_{jk}\}$.

Putting

$$(10) \quad V^i = \frac{\partial^3 u^i}{\partial s^3} + \frac{du^i}{ds} g_{jk} \frac{\partial^2 u^j}{\partial s^2} \frac{\partial^2 u^k}{\partial s^2} - \frac{du^i}{ds} \Pi_{jk}^0 \frac{du^j}{ds} \frac{du^k}{ds} + \Pi_{\infty k}^i \frac{du^k}{ds},$$

we have, from (9),

$$(11) \quad A_{(3)} = \frac{1}{\left(\frac{d\sigma}{ds}\right)^2} \left[g_{jk} V^j \frac{\partial^2 u^k}{\partial s^2} A_0 + V^i A_i \right].$$

Differentiating once more the equation (11) along the curve, we have, in consequence of (3),

$$(12) \quad \frac{d}{d\sigma} A_{(3)} = A_{(0)} = \left[-\frac{2 \frac{d^2\sigma}{ds^2}}{\left(\frac{d\sigma}{ds}\right)^4} g_{jk} V^j \frac{\partial^2 u^k}{\partial s^2} + \frac{1}{\left(\frac{d\sigma}{ds}\right)^3} \left(g_{jk} \frac{\partial V^j}{\partial s} \frac{\partial^2 u^k}{\partial s^2} + g_{jk} V^j \frac{\partial^3 u^k}{\partial s^3} + \Pi_{jk}^0 V^j \frac{du^k}{ds} \right) \right] A_0 + \left[-\frac{2 \frac{d^2\sigma}{ds^2}}{\left(\frac{d\sigma}{ds}\right)^4} V^i + \frac{1}{\left(\frac{d\sigma}{ds}\right)^3} \left(g_{jk} V^j \frac{\partial^2 u^k}{\partial s^2} \frac{du^i}{ds} + \frac{\partial V^i}{\partial s} \right) \right] A_i,$$

*) We have used the relations $0 = \lambda = -\{t, \sigma\} = -[\{t, s\} - \{\sigma, s\}] \left(\frac{ds}{d\sigma}\right)^2$ and $\{\sigma, s\} = \{t, s\} = \frac{1}{2} g_{jk} \frac{\partial^2 u^j}{\partial s^2} \frac{\partial^2 u^k}{\partial s^2} - \Pi_{jk}^0 \frac{du^j}{ds} \frac{du^k}{ds}$ for the reduction.

from which we obtain

$$(13) \quad g_{jk} \frac{\delta V^j}{\delta s} \frac{\delta^2 u^k}{\delta s^2} + g_{jk} V^j \frac{\delta^3 u^k}{\delta s^3} + H_{jk}^0 V^j \frac{du^k}{ds} - \frac{2 \frac{d^2 \sigma}{ds^2}}{\frac{d\sigma}{ds}} g_{jk} V^j \frac{\delta^2 u^k}{\delta s^2} = \left(\frac{d\sigma}{ds} \right)^4$$

and

$$(14) \quad \frac{\delta V^i}{\delta s} + g_{jk} V^j \frac{\delta^2 u^k}{\delta s^2} \frac{du^i}{ds} = \frac{2 \frac{d^2 \sigma}{ds^2}}{\frac{d\sigma}{ds}} V^i.$$

But, as we can see from (10), we have

$$g_{jk} V^j \frac{du^k}{ds} = 0$$

which gives, on covariant differentiation,

$$(15) \quad g_{jk} V^j \frac{\delta^2 u^k}{\delta s^2} = -g_{jk} \frac{\delta V^j}{\delta s} \frac{du^k}{ds}.$$

Substituting these equations in (14), we find

$$(16) \quad \frac{\delta V^i}{\delta s} - \frac{du^i}{ds} g_{jk} \frac{\delta V^j}{\delta s} \frac{du^k}{ds} = \frac{2 \frac{d^2 \sigma}{ds^2}}{\frac{d\sigma}{ds}} V^i$$

or

$$(17) \quad \frac{Dv^i}{Ds} = 0$$

where D/Ds is a conformal covariant derivative defined in a previous paper¹⁾, and

$$(18) \quad v^i = \frac{V^i}{\left(\frac{d\sigma}{ds} \right)^2}$$

σ being defined by

$$(19) \quad \left(\frac{d\sigma}{ds} \right)^4 = g_{jk} V^j V^k.$$

If the equations (16) hold good, the left-hand member of (13) will be written as

$$\frac{2 \frac{d^2 \sigma}{ds^2}}{\frac{d\sigma}{ds}} g_{jk} V^j \frac{\delta^2 u^k}{\delta s^2} + g_{jk} V^j V^k - \frac{2 \frac{d^2 \sigma}{ds^2}}{\frac{d\sigma}{ds}} g_{jk} V^j \frac{\delta^2 u^k}{\delta s^2}$$

then, in consequence of (19), the equations (13) are automatically satisfied.

1) K. Yano: Sur la connexion de Weyl-Hlavatý et ses applications à la géométrie conforme, Proc. Physico-Math. Soc. Japan. Vol. 22 (1940), 595-621.

So that, (17) are the differential equations of the generalized loxodromes to be obtained.

§ 2. Next, we shall consider the curves for which we have

$$(20) \quad \lambda = \text{const.} \quad \overset{4}{\lambda} = \overset{5}{\lambda} = \dots = \overset{n}{\lambda} = \overset{\infty}{\lambda} = 0.$$

For these curves, we have, from (1),

$$(21) \quad \frac{d}{d\sigma} \underset{(0)}{A} = \underset{(1)}{A}, \quad \frac{d}{d\sigma} \underset{(1)}{A} = \lambda \underset{(0)}{A} + \underset{(2)}{A}, \quad \frac{d}{d\sigma} \underset{(2)}{A} = \lambda \underset{(1)}{A} + \underset{(3)}{A}, \quad \frac{d}{d\sigma} \underset{(3)}{A} = \underset{(0)}{A},$$

consequently, we can, as in § 1, treat the curves as if they were in a two-dimensional conformal space. In order that the circle

$$(22) \quad \alpha_0 \underset{(0)}{A} + \alpha_1 \underset{(1)}{A} + \alpha_2 \underset{(2)}{A} + \alpha_3 \underset{(3)}{A}$$

where $\alpha_0, \alpha_1, \alpha_2$ and α_3 are constants, be a fixed circle, the constants $\alpha_0, \alpha_1, \alpha_2$ and α_3 must satisfy the equations

$$(23) \quad \alpha_1 \lambda + \alpha_3 = \rho \alpha_0, \quad \alpha_2 \lambda + \alpha_0 = \rho \alpha_1, \quad \alpha_1 = \rho \alpha_2, \quad \alpha_2 = \rho \alpha_3$$

where ρ is a proportional constant.

Eliminating $\alpha_0, \alpha_1, \alpha_2$ and α_3 from these equations, we have

$$(24) \quad \rho^4 - 2\lambda\rho^2 - 1 = 0$$

which gives two real values of ρ

$$(25) \quad \begin{cases} \rho = +\sqrt{\lambda + \sqrt{\lambda^2 + 1}}, \\ \rho' = -\sqrt{\lambda + \sqrt{\lambda^2 + 1}}. \end{cases}$$

Corresponding to these values of ρ , we have

$$\alpha_0 : \alpha_1 : \alpha_2 : \alpha_3 = \sqrt{\lambda^2 + 1} \cdot \rho : \rho^2 : \rho : 1,$$

and

$$\alpha_0 : \alpha_1 : \alpha_2 : \alpha_3 = \sqrt{\lambda^2 + 1} \cdot \rho' : \rho'^2 : \rho' : 1 = -\sqrt{\lambda^2 + 1} \cdot \rho : \rho^2 : -\rho : 1.$$

Consequently, the two circles defined by

$$(26) \quad \begin{cases} P = +\sqrt{\lambda^2 + 1} \underset{(0)}{\rho A} + \underset{(1)}{\rho^2 A} + \underset{(2)}{\rho A} + \underset{(3)}{A}, \\ Q = -\sqrt{\lambda^2 + 1} \underset{(0)}{\rho A} + \underset{(1)}{\rho^2 A} - \underset{(2)}{\rho A} + \underset{(3)}{A}, \end{cases}$$

are both fixed circles. Moreover, we can easily verify that

$$(27) \quad P \cdot P = 0 \quad \text{and} \quad Q \cdot Q = 0,$$

that is to say that the P and Q are both point-circles.

The circle passing through the three points $\underset{(0)}{A}, P$ and Q being $\underset{(1)}{A} - \underset{(3)}{\rho^2 A}$, the angle θ ($0 \leq \theta \leq \pi$) at $\underset{(0)}{A}$ between the curve and this circle is given by

$$(28) \quad \sin \theta = \frac{A(A - \rho^2 A)}{\sqrt{AA} \sqrt{(A - \rho^2 A)(A - \rho^2 A)}} = \frac{1}{\sqrt{\rho^4 + 1}}$$

$$= \frac{1}{\sqrt{2(\lambda^2 + 1) + 2\lambda\sqrt{\lambda^2 + 1}}}.$$

Putting

$$(29) \quad \lambda = \tan \varphi \quad \left(-\frac{\pi}{2} \leq \varphi \leq \frac{\pi}{2} \right)$$

we have, from (28),

$$\sin \theta = \frac{1}{\sqrt{2 \sec^2 \varphi + 2 \tan \varphi \sec \varphi}} = \frac{1}{\sqrt{2}} \left(\cos \frac{\varphi}{2} - \sin \frac{\varphi}{2} \right) = \sin \left(\frac{\pi}{4} - \frac{\varphi}{2} \right)$$

hence,

$$(29) \quad \theta = \frac{\pi}{4} - \frac{\varphi}{2}.$$

Then, we have the

Theorem II. The curve whose conformal curvature $\lambda = \text{const.}$ and $\lambda = \lambda = \dots = \lambda = 0$, cuts all the circles passing through two fixed points by the fixed angle $\frac{\pi}{4} - \frac{\varphi}{2}$ where $\tan \varphi = \lambda$, consequently it is a loxodrome.

The four points A, A, P and Q lying on the circle $A - \rho^2 A$, the points A and A are harmonically separated by the points P and Q on this circle.