

## 28. A Kinematic Theory of Turbulence\*

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1. *Generalities.* In the theory of turbulence<sup>1)</sup> the deviation of the velocity from its mean may be considered as a system of random vectors  $u_\lambda(t, \mathfrak{x}, \omega)$ ,  $\lambda=1, 2, 3$ , where  $t(\in R)$  is the time parameter and  $\mathfrak{x}(\in R^3)$  denotes the position and  $\omega(\in(\Omega, P))$  is the elementary event. Then we have

$$(1) \quad \mathcal{E}_\omega(u_\lambda(t, \mathfrak{x}, \omega)) = 0.$$

When the system  $\{u_\lambda(t, \mathfrak{x}, \omega)\}$  is of Gaussian type<sup>2)</sup>, we say that the turbulence is of Gaussian type.

Now we define the moment tensor of the turbulence by

$$(2) \quad R_{\lambda\mu}(t, \mathfrak{x}; s, \mathfrak{y}) = \mathcal{E}_\omega\{u_\lambda(t, \mathfrak{x}, \omega)u_\mu(s, \mathfrak{y}, \omega)\}.$$

Then  $R_{\lambda\mu}(t, \mathfrak{x}; s, \mathfrak{y})$  is a positive-definite function of  $(\lambda, t, \mathfrak{x})$  and  $(\mu, s, \mathfrak{y})$  in the sense of Bochner, namely we have

$$(3) \quad R_{\lambda\mu}(t, \mathfrak{x}; s, \mathfrak{y}) = R_{\mu\lambda}(s, \mathfrak{y}; t, \mathfrak{x}) \quad \text{and}$$

$$(4) \quad \sum_{i,j} \xi_i \xi_j R_{\lambda_i \lambda_j}(t_i, \mathfrak{x}_i; t_j, \mathfrak{x}_j) \geq 0;$$

in fact (3) is evident by (2) and the left side of (4) is equal to  $\mathcal{E}_\omega\left\{\left(\sum_i \xi_i u_{\lambda_i}(t_i, \mathfrak{x}_i, \omega)\right)^2\right\}$ . Conversely the function  $R_{\lambda\mu}(t, \mathfrak{x}; s, \mathfrak{y})$  satisfying (3) and (4) may be considered as the moment tensor of a turbulence of Gaussian type<sup>3)</sup>.

A turbulence is defined as temporally homogeneous, if its moment tensor satisfies

$$(5) \quad R_{\lambda\mu}(t+\tau, \mathfrak{x}; s+\tau, \mathfrak{y}) = R_{\lambda\mu}(t, \mathfrak{x}; s, \mathfrak{y}).$$

It is defined as spatially homogeneous, if we have

$$(6) \quad R_{\lambda\mu}(t, \mathfrak{x}+\alpha; s, \mathfrak{y}+\alpha) = R_{\lambda\mu}(t, \mathfrak{x}; s, \mathfrak{y}).$$

We say that it is isotopic if we have always

$$(7) \quad \sum_{\lambda'\mu'} k_{\lambda'\lambda} k_{\mu'\mu} R_{\lambda'\mu'}(t, \mathfrak{x}; s, \mathfrak{x}+K(\mathfrak{y})+\mathfrak{x}) = R_{\lambda\mu}(t, \mathfrak{x}; s, \mathfrak{y})$$

for any orthogonal transformation  $K \equiv \{k_{\lambda\mu}; \lambda, \mu=1, 2, 3\}$ . We can easily prove by (3) that the isotropism implies the homogeneity.

\* The cost of this research has been defrayed from the Scientific Expenditure of the Department of Education.

1) H. P. Robertson: The invariant theory of isotropic turbulence, Proc. Camb. Phil. Soc. 36, 1940.

2) Cf. K. Itô: ガウス型確率変動系 = ツイテ (全國紙上數學談話會第 261 號).

3) See Theorem 3 in my above-cited note (2).

It seems to be an important and perhaps difficult problem to determine the canonical form of  $R_{\lambda\mu}(t, \mathfrak{x}; s, \mathfrak{y})$  which satisfies (3), (4), (5) and (7).

2. *The temporally homogeneous and isotropic turbulence at a point.* For the investigation of this subject we can consider  $u_i(t, \omega)$  and  $R_{\lambda\mu}(t, s)$  respectively instead of  $u_i(t, \mathfrak{x}, \omega)$  and  $R_{\lambda\mu}(t, \mathfrak{x}; s, \mathfrak{y})$ . By (3) and (4) we have

$$(3') \quad R_{\lambda\mu}(t, s) = R_{\mu\lambda}(s, t) \quad \text{and} \quad (4') \quad \sum_{ij} \xi_i \xi_j R_{\lambda_i \lambda_j}(t_i, t_j) \geq 0.$$

The conditions (5) and (7) may be written in the forms:

$$(5') \quad R_{\lambda\mu}(t+\tau, s+\tau) = R_{\lambda\mu}(t, s) \quad \text{and} \quad (6') \quad \sum_{\lambda'\mu'} k_{\lambda'\lambda} k_{\mu'\mu} R_{\lambda'\mu'}(t, s) = R_{\lambda\mu}(t, s).$$

Theorem 1. A necessary and sufficient condition that  $R_{\lambda\mu}(t, s)$  should be the moment tensor of a temporally homogeneous and isotropic turbulence at a point is that  $R_{\lambda\mu}(t, s)$  is expressible by the form:

$$(8) \quad R_{\lambda\mu}(t, s) = \delta_{\lambda\mu} \int_{-0}^{\infty} \cos(\xi(t-s)) F(d\xi),$$

where  $\delta_{\lambda\mu}$  is the Kronecker's delta and  $F$  is a measure distribution on  $[0, \infty)$  with the finite total measure.

Proof. Necessity. The isotropism (6') implies that  $R_{\lambda\mu}(t, s)$  is an invariant tensor. Therefore we obtain  $R_{\lambda\mu}(t, s) = \delta_{\lambda\mu} C(t, s)$ . From the temporal homogeneity (5') and the symmetric character (3') follows that  $C(t, s)$  is a function of  $|t-s|$  only, say  $C_1(|s-t|)$ . Now we see by (4') that  $C(|\tau|)$  is a positive-definite function of  $\tau$ . Making use of the Bochner's theorem we obtain (8). The sufficiency is evident.

According to this theorem  $u_i(t, \omega)$  and  $u_\mu(s, \omega)$  ( $\lambda \neq \mu$ ) are non-correlated in this turbulence. Therefore, if we assume further that the turbulence be of Gaussian type, the three stochastic processes ( $u_\lambda(t, \omega); -\infty < t < \infty$ ),  $\lambda=1, 2, 3$ , will become independent. Nevertheless each process is clearly a stationary process of Gaussian type with the correlation function  $\rho(\tau) \equiv \int_{-0}^{\infty} \cos \tau \xi F(d\xi) / F([0, \infty))$ . In this case the problem may be reduced to the investigation of such a process.

Next we mention a theorem concerning the ergodicity of this process, which includes the result<sup>1)</sup> obtained before by the author; the proof can be achieved by the same idea and so will be omitted.

Theorem 2. A necessity and sufficient condition that a normalized continuous (in mean) stationary process  $u(t, \omega)$  of Gaussian type should be ergodic in the strongly mixing type is that its correlation function  $\rho(\tau)$  satisfies

$$(9) \quad \lim_{\tau \rightarrow \infty} \rho(\tau) = 0.$$

The condition (9) means that the correlation coefficient tends to 0

1) See K. Itô: On the ergodicity of a certain stationary process, Proc. **20** (1944), 54-55.

as the time interval increases indefinitely. In practice we may assume that it is well satisfied. Then by Theorem 2 we can see that

$$(10) \quad P\left\{\omega; \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T u(t, \omega) u(t + \tau, \omega) dt = \rho(\tau)\right\} = 1.$$

This identity justifies the practical method in which we make use of the time-mean of  $u(t, \omega)u(t + \tau, \omega)$  at a certain (realized) value of  $\omega$  instead of its mathematical expectation  $\rho(\tau)$ . It is also the case with turbulence.

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