

Kirillov-Kostant theory and Feynman path integrals on coadjoint orbits II

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Introduction

Let G be a connected and simply connected solvable Lie group. In this paper we construct irreducible unitary representations of G by using the Feynman path integrals on coadjoint orbits [1][13].

In §1, we compute the path integrals on $M = \mathbf{R}^n \times \mathbf{R}^n$. Let θ be a 1-form on M and H a C^∞ -function on M which satisfies certain conditions. The path integral $K_{\theta, H}(x'', x'; T)$ ($x', x'' \in \mathbf{R}^n$) computed by using the action $\int_0^T \gamma^* \theta - H(\gamma(t)) dt$ (where γ runs over a certain set of paths on M) can be written by the solution of differential equations defined by θ and H .

In §2, we investigate the path integrals on coadjoint orbits. Let \mathfrak{g} be the Lie algebra of G and \mathfrak{g}^* the dual space of \mathfrak{g} . Fix an element λ of \mathfrak{g}^* and choose a real polarization \mathfrak{p} . Following the Kirillov-Kostant theory [4][5][14], we construct an irreducible unitary representation $\pi_\lambda^\mathfrak{p}$ of G . We put $\theta_\lambda = \langle \lambda, g^{-1} dg \rangle$ and $H_Y = \langle \lambda, g^{-1} Yg \rangle$ for any $Y \in \mathfrak{g}$. We show that the integral operator of K_{θ_λ, H_Y} corresponds to $\pi_\lambda^\mathfrak{p}(\exp TY)$.

§1 Path integrals on $\mathbf{R}^n \times \mathbf{R}^n$

In this section, we shall compute the Feynman path integrals on $M = \mathbf{R}^n \times \mathbf{R}^n$. Let n_1, \dots, n_m be natural numbers such that $\sum_{i=1}^m n_i = n$. We put $U^i = \mathbf{R}^{n_i}$ and $V^i = \mathbf{R}^{n_i}$ for $i = 1 \dots m$. Let ${}^t(x, y) = (x^1 \dots x^m y^1 \dots y^m)$ be the normal coordinates on $M = U^1 \times \dots \times U^m \times V^1 \times \dots \times V^m$ where $x^i = (x^{i,1} \dots x^{i,n_i}) \in U^i$ and $y^i = (y^{i,1} \dots y^{i,n_i}) \in V^i$ for $i = 1 \dots m$. Let θ be a 1-form on M and H a C^∞ -function on M . Suppose that θ and H are expressed in the following forms respectively:

$$\theta = \sum_{i=1}^m {}^t y^i (dx^i + \sum_{j=1}^{i-1} f^{ij} dx^j) + {}^t a^i dx^i + {}^t b^i dy^i \quad (1.1)$$

where

$$\begin{aligned}
 f^{ij} &= f^{ij0} + \sum_{k=1}^{n_i} x^{i,k} f^{ijk}, \\
 f^{ijk} &\in C^\infty(U^1 \times \dots \times U^{i-1}, \mathfrak{M}(n_i, n_j, \mathbf{R})), \quad k = 0, 1, \dots, n_i, \\
 a^i &\in C^\infty(U^1 \times \dots \times U^m, U^i), \\
 b^i &\in C^\infty(M, U^i)
 \end{aligned}$$

and

$$H = \sum_{i=1}^m {}^t y^i (h^{i0} + \sum_{k=1}^{n_i} x^{ik} h^{ik}) + c \tag{1.2}$$

where

$$\begin{aligned}
 h^{ik} &\in C^\infty(U^1 \times \dots \times U^{i-1}, U^i), \quad k = 0, 1, \dots, n_i, \\
 c &\in C^\infty(U^1 \times \dots \times U^m, \mathbf{R}).
 \end{aligned}$$

For ${}^t(x, y) \in M$, we put $dx = dx^{1,1} \wedge \dots \wedge dx^{1,n_1} \wedge dx^{2,1} \wedge \dots \wedge dx^{m,n_m}$ and $d(y/2\pi) = d(y^{1,1}/2\pi) \wedge \dots \wedge d(y^{1,n_1}/2\pi) \wedge d(y^{2,1}/2\pi) \wedge \dots \wedge d(y^{m,n_m}/2\pi)$. Now for $x', x'' \in \mathbf{R}^n$ and $T \in \mathbf{R}$ define the path integral $K_{\theta, H}(x'', x'; T)$ by

$$\begin{aligned}
 &K_{\theta, H}(x'', x'; T) \\
 &= \lim_{N \rightarrow \infty} \int dx_1 \dots dx_{N-1} d \frac{y_1}{2\pi} \dots d \frac{y_N}{2\pi} \exp \left\{ \sqrt{-1} \int_0^T \gamma^* \theta - H(\gamma(t)) dt \right\} \tag{1.3}
 \end{aligned}$$

where

$$\gamma(t) = \left(x_{k-1} + \frac{x_k - x_{k-1}}{T/N} \left(t - \frac{k-1}{N} T \right), y_{k-1} \right) \in M \quad \text{for } t \in \left[\frac{k-1}{N} T, \frac{k}{N} T \right),$$

$x_N = x''$ and $x_0 = x'$.

For simplicity, we put

$$\begin{aligned}
 \kappa(x, \dot{x}) &= \sum_{j=1}^m {}^t a^j(x) \dot{x}^j + c(x), \\
 \tau^i(x, \dot{x}) &= \sum_{j=1}^{i-1} f^{ij0}(x^1, \dots, x^{i-1}) \dot{x}^j - h^i(x^1, \dots, x^{i-1}), \\
 v^i(x, \dot{x}) &= \sum_{j=1}^{i-1} (f^{ij1}(x^1, \dots, x^{i-1}) \dot{x}^j, \dots, f^{ijn_i}(x^1, \dots, x^{i-1}) \dot{x}^j) \\
 &\quad - (h^{i1}(x^1, \dots, x^{i-1}), \dots, h^{in_i}(x^1, \dots, x^{i-1}))
 \end{aligned}$$

and

$$\mu^i(x, \dot{x}) = \tau^i(x, \dot{x}) + v^i(x, \dot{x})x^i.$$

From the form of μ^i , for $x \in \mathbf{R}^n = U^1 \times \dots \times U^m$, the differential equation

$$\begin{cases} \dot{w} + \mu(w, \dot{w}) = 0, \\ w(0) = x \end{cases} \quad (1.4)$$

has a unique solution $w(x, t)$. As the next theorem shows, the path integral $K_{\theta, H}(x'', x'; T)$ is described by using the solution w .

THEOREM 1.

$$\begin{aligned} & K_{\theta, H}(x'', x'; T) \\ &= \delta(x'' - w(x', T)) \exp \left\{ \sqrt{-1} \int_0^T \kappa(w(x', t), \dot{w}(x', t)) dt \right\} \left| \frac{dw(x', T)}{dx'} \right|^{\frac{1}{2}} \end{aligned}$$

where $\frac{dw(x', T)}{dx'}$ denotes the Jacobian.

PROOF. By the definition of γ in (1.3), we can assume that $b^i = 0$ and by integrating with respect to y , we obtain

$$\begin{aligned} & K_{\theta, H}(x'', x'; T) \\ &= \lim_{N \rightarrow \infty} \int dx_1 \cdots dx_{N-1} \prod_{k=1}^N \delta \left(x_k - x_{k-1} + \int_{\frac{k-1}{N}T}^{\frac{k}{N}T} \mu(x(t), \dot{x}(t)) dt \right) \\ & \quad \times \exp \left\{ \sqrt{-1} \int_0^T \kappa(x(t), \dot{x}(t)) dt \right\} \end{aligned}$$

where

$$x(t) = x_{k-1} + \frac{x_k - x_{k-1}}{T/N} \left(t - \frac{k-1}{N} T \right) \quad \text{for } t \in \left[\frac{k-1}{N} T, \frac{k}{N} T \right],$$

$$x_N = x'' \text{ and } x_0 = x'.$$

In order to integrate with respect to x , we define $z_{N,k} \in \mathbf{R}^n$ by the equations

$$\begin{cases} z_{N,k} - z_{N,k-1} \\ = - \int_{\frac{k-1}{N}T}^{\frac{k}{N}T} \mu \left(z_{N,k-1} + \frac{z_{N,k} - z_{N,k-1}}{T/N} \left(t - \frac{k-1}{N} T \right), \frac{z_{N,k} - z_{N,k-1}}{T/N} \right) dt, \\ z_{N,0} = x'. \end{cases} \quad (1.5)$$

Now we suppose that (1.5) is well-defined when N is sufficiently large. We put

$$z_N(t) = z_{N,k-1} + \frac{z_{N,k} - z_{N,k-1}}{T/N} \left(t - \frac{k-1}{N} T \right) \quad \text{for } t \in \left[\frac{k-1}{N} T, \frac{k}{N} T \right].$$

Then we obtain

$$\begin{aligned} & K_{\theta,H}(x'', x'; T) \\ &= \lim_{N \rightarrow \infty} \delta(x'' - z_N(T)) \exp \left\{ \sqrt{-1} \int_0^T \kappa(z_N(t), \dot{z}_N(t)) dt \right\} \\ & \quad \times \prod_{k=1}^N \prod_{i=1}^m \left| I_{n_i} - \int_{\frac{k-1}{N}T}^{\frac{k}{N}T} v^i(z_N(t), \dot{z}_N(t)) \frac{\frac{k}{N}T - t}{T/N} dt \right|^{-1} \end{aligned}$$

Hence, to complete the proof, we have only to show that the following (i) and (ii) hold.

(i) $z_N^i(t)$ is well-defined when N is sufficiently large and $z_N^i(t)$ and $\dot{z}_N^i(t)$ converge to $w^i(x', t)$ and $\dot{w}^i(x', t)$ respectively uniformly on $[0, T]$ for $i = 1 \cdots m$.

(ii)

$$\lim_{N \rightarrow \infty} \prod_{k=1}^N \prod_{i=1}^m \left| I_{n_i} - \int_{\frac{k-1}{N}T}^{\frac{k}{N}T} v^i(z_N(t), \dot{z}_N(t)) \frac{\frac{k}{N}T - t}{T/N} dt \right|^{-1} = \left| \frac{dw(x', T)}{dx'} \right|^{\frac{1}{2}}.$$

Since $\mu^i(x, \dot{x})$ is independent of $x^{i+1}, \dots, x^m, \dot{x}^i, \dots, \dot{x}^m$, we can use the induction with respect to i to show (i). The following discussion shows the facts that $z_N^1(t)$ is well-defined when N is sufficiently large and $z_N^1(t)$ and $\dot{z}_N^1(t)$ converge to $w^1(x', t)$ and $\dot{w}^1(x', t)$ respectively uniformly on $[0, T]$.

Now for $i = 1 \cdots l - 1$ suppose that $z_N^i(t)$ is well-defined when N is sufficiently large and that $z_N^i(t)$ and $\dot{z}_N^i(t)$ converge to $w^i(x', t)$ and $\dot{w}^i(x', t)$ respectively uniformly on $[0, T]$.

For simplicity, we put $\alpha_N, \alpha, \beta_N$ and β as follows:

$$\begin{aligned} \alpha_N(t) &= v^1(z_N(t), \dot{z}_N(t)), \\ \alpha(t) &= v^1(w(x', t), \dot{w}(x', t)), \\ \beta_N(t) &= \tau^1(z_N(t), \dot{z}_N(t)) \end{aligned}$$

and

$$\beta(t) = \tau^1(w(x', t), \dot{w}(x', t)).$$

Now since $\dot{w}^1(x', t) = \alpha(t)w^1(x', t) + \beta(t)$, we obtain

$$w^1\left(x', \frac{k}{N} T\right) = w^1\left(x', \frac{k-1}{N} T\right) + \int_{\frac{k-1}{N}T}^{\frac{k}{N}T} \alpha(t)w(x', t) dt + \int_{\frac{k-1}{N}T}^{\frac{k}{N}T} \beta(t) dt.$$

We put

$$A_{N,k} = \left(I_{n_l} - \int_{\frac{k-1}{N}T}^{\frac{k}{N}T} \alpha_N(t) \frac{\frac{k}{N}T - t}{T/N} dt \right),$$

$$B_{N,k} = \left(I_{n_l} + \int_{\frac{k-1}{N}T}^{\frac{k}{N}T} \alpha_N(t) \frac{t - \frac{k-1}{N}T}{T/N} dt \right)$$

and

$$C_{N,k} = \int_{\frac{k-1}{N}T}^{\frac{k}{N}T} \beta_N(t) dt.$$

Then, from the definition of $z_{N,k}$, we have

$$A_{N,k} z_{N,k}^l = B_{N,k} z_{N,k-1}^l + C_{N,k}.$$

From the hypothesis of the induction, α_N and β_N converge to α and β respectively uniformly on $[0, T]$. When T/N is sufficiently small, $A_{N,k}$ has the inverse matrix. Hence when N is sufficiently large, $z_{N,k}$ is well-defined and

$$z_{N,k}^l = A_{N,k}^{-1} B_{N,k} z_{N,k-1}^l + A_{N,k}^{-1} C_{N,k}.$$

Now we can find a positive number Γ such that for any positive number ε there exists a number N_0 for which

$$\begin{aligned} & \left\| z_{N,k}^l - w^l \left(x', \frac{k}{N} T \right) \right\| \\ & \leq \left\| A_{N,k}^{-1} B_{N,k} \left(z_{N,k-1}^l - w^l \left(x', \frac{k-1}{N} T \right) \right) \right\| \\ & + \left\| A_{N,k}^{-1} B_{N,k} w^l \left(x', \frac{k-1}{N} T \right) - w^l \left(x', \frac{k-1}{N} T \right) - \int_{\frac{k-1}{N}T}^{\frac{k}{N}T} \alpha(t) w(x', t) dt \right\| \\ & + \left\| A_{N,k}^{-1} B_{N,k} - \int_{\frac{k-1}{N}T}^{\frac{k}{N}T} \beta(t) dt \right\| \\ & \leq \left(1 + \frac{\Gamma}{N} \right) \left\| \left(z_{N,k-1}^l - w^l \left(x', \frac{k-1}{N} T \right) \right) \right\| + \frac{1}{2} \frac{\varepsilon}{N} + \frac{1}{2} \frac{\varepsilon}{N} \end{aligned}$$

for any $N > N_0$. Using the above inequality repeatedly and using $z_{N,0}^l = w^l(x', 0)$, we obtain

$$\left\| z_{N,k}^l - w^l \left(x', \frac{k}{N} T \right) \right\| \leq \sum_{i=1}^l \left(1 + \frac{\Gamma}{N} \right)^i \frac{\varepsilon}{N} \leq (e^{(1+\frac{1}{N})\Gamma} - 1) \frac{\varepsilon}{\Gamma}.$$

Therefore

$$\lim_{N \rightarrow \infty} \max_{k=0 \cdots N} \left\| z_{N,k}^l - w^l \left(x', \frac{k}{N} T \right) \right\| = 0.$$

This shows that $z_N^l(t)$ converges to $w^l(t)$ uniformly on $[0, T]$. And from the following inequality

$$\begin{aligned} & \left\| \frac{z_{N,k}^l - z_{N,k-1}^l}{T/N} - \dot{w}^l \left(x', \frac{k-1}{N} T \right) \right\| \\ &= \left\| \frac{(A_{N,k} - I_{n_l}) z_{N,k-1}^l + B_{N,k}}{T/N} - \alpha \left(\frac{k-1}{N} T \right) w^l \left(x', \frac{k-1}{N} T \right) - \beta \left(\frac{k-1}{N} T \right) \right\| \\ &\leq \left\| \frac{A_{N,k} - I_{n_l}}{T/N} \left(z_{N,k-1}^l - w^l \left(x', \frac{k-1}{N} T \right) \right) \right\| \\ &+ \left\| \left(\frac{A_{N,k} - I_{n_l}}{T/N} - \alpha \left(\frac{k-1}{N} T \right) \right) w^l \left(x', \frac{k-1}{N} T \right) \right\| \\ &+ \left\| \frac{B_{n,k}}{T/N} - \beta \left(\frac{k-1}{N} T \right) \right\| \end{aligned}$$

we obtain

$$\lim_{N \rightarrow \infty} \max_{k=0 \cdots N} \left\| \frac{z_{N,k}^l - z_{N,k-1}^l}{T/N} - \dot{w}^l \left(x', \frac{k-1}{N} T \right) \right\| = 0.$$

This shows that $z_N^l(t)$ converges to $\dot{w}^l(t)$ uniformly on $[0, T]$. Thus we have proved (i).

Now we have

$$\lim_{N \rightarrow \infty} \max_{k=0 \cdots N} \left\| I_{n_l} - \int_{\frac{k-1}{N} T}^{\frac{k}{N} T} \alpha_N(t) \frac{\frac{k}{N} T - t}{T/N} dt - \exp \left(-\frac{1}{2} \int_{\frac{k-1}{N} T}^{\frac{k}{N} T} \alpha(t) dt \right) \right\| \frac{N}{T} = 0.$$

Hence we obtain

$$\begin{aligned} \lim_{N \rightarrow \infty} \prod_{k=1}^N \left| I_{n_l} - \int_{\frac{k-1}{N} T}^{\frac{k}{N} T} \alpha_N(t) \frac{\frac{k}{N} T - t}{T/N} dt \right|^{-1} &= \lim_{N \rightarrow \infty} \prod_{k=1}^N \left| \exp \left(-\frac{1}{2} \int_{\frac{k-1}{N} T}^{\frac{k}{N} T} \alpha(t) dt \right) \right|^{-1} \\ &= \exp \frac{1}{2} \int_0^T \alpha(t) dt \\ &= \left| \frac{dw^l(x', T)}{dx'^l} \right|^{\frac{1}{2}}. \end{aligned}$$

Since $dw^l(x', T)/dx'^i = 0$ for $i > l$, consequently we obtain (ii). \square

§2 Path integrals on coadjoint orbits

Let G be a connected and simply connected solvable Lie group, \mathfrak{g} the Lie algebra of G , \mathfrak{g}^* the dual space of \mathfrak{g} . For $\lambda \in \mathfrak{g}^*$, we put $G_\lambda = \{g \in G \mid \text{Ad}^*(g)\lambda = \lambda\}$. Then the coadjoint orbit O_λ is canonically identified with the homogeneous space G/G_λ . Let \mathfrak{g}_λ be the Lie algebra of G_λ and we consider a real polarization $\mathfrak{p}(\mathfrak{g}_\lambda \subset \mathfrak{p} \subset \mathfrak{g})$. We fix a Lie subgroup P of G the Lie algebra of which coincides with \mathfrak{p} . Here we suppose that $\text{Ad}^*(P)\lambda = \lambda + \mathfrak{p}^\perp$ where $\mathfrak{p}^\perp = \{\zeta \in \mathfrak{g}^* \mid \zeta(X) = 0 \text{ for any } X \in \mathfrak{p}\}$.

First we choose the coordinates on O_λ to define the path integrals on O_λ . We take a chain $\mathfrak{g}_1 \supset \mathfrak{g}_2 \supset \dots \supset \mathfrak{g}_m \supset \mathfrak{g}_{m+1}$ of ideals in \mathfrak{g} , beginning with $\mathfrak{g}_1 = \mathfrak{g}$, ending with $\mathfrak{g}_{m+1} = \{0\}$, such that the factor algebras $\mathfrak{g}_i/\mathfrak{g}_{i+1}$ ($i = 1, \dots, m$) are all abelian. We put $n = \dim G/P$ and $n_i = \dim \mathfrak{g}_i/(\mathfrak{g}_{i+1} + \mathfrak{g}_i \cap \mathfrak{p})$. Then we have $n = \sum_{i=1}^m n_i$. Now $\mathfrak{g}_i/(\mathfrak{g}_{i+1} + \mathfrak{g}_i \cap \mathfrak{p}) \simeq (\mathfrak{g}_i \cap \mathfrak{p})/(\mathfrak{g}_{i+1} \cap \mathfrak{p})$. Therefore we can take $X_{i,j} \in \mathfrak{g}_i$ ($i = 1 \dots m, j = 1 \dots n_i$) such that

$$\sum_{j=1}^{n_i} \mathbf{R}X_{i,j} \oplus (\mathfrak{g}_{i+1} + \mathfrak{g}_i \cap \mathfrak{p}) = \mathfrak{g}_i.$$

We define the mapping $\phi: \mathbf{R}^n = U^1 \times \dots \times U^m \rightarrow G$ by

$$\phi(x) = \exp(x^{1,1} X_{1,1}) \dots \exp(x^{1,n_1} X_{1,n_1}) \dots \exp(x^{m,n_m} X_{m,n_m})$$

where $x = {}^t(x^1 \dots x^m)$ and $x^i = {}^t(x^{i,1} \dots x^{i,n_i})$. Let G_i be the analytic subgroup of G corresponding to \mathfrak{g}_i . Then the analytic subgroup of G corresponding to $\mathfrak{g}_{i+1} + \mathfrak{g}_i \cap \mathfrak{p}$ can be written as $G_{i+1}(G_i \cap P)$. This shows that the mapping $\mathbf{R}^n \ni x \mapsto \phi(x)P \in G/P$ is an onto-diffeomorphism.

For $i = 1 \dots m$ and $j = 1 \dots n_i$, we take $\zeta_{ij} \in \mathfrak{p}^\perp$ such that $\zeta_{ij}(X_{k,l}) = \delta_{ik}\delta_{jl}$ and take an immersion map $\psi: \mathbf{R}^n = V^1 \times \dots \times V^m \hookrightarrow P$ such that

$$\text{Ad}^*(\psi(y))\lambda = \lambda + \sum_{i=1}^m \sum_{j=1}^{n_i} y^{i,j} \zeta_{ij}$$

where $y = {}^t(y^1 \dots y^m)$ and $y^i = {}^t(y^{i,1} \dots y^{i,n_i})$. Then the mapping $\mathbf{R}^n \ni y \mapsto \psi(y)G_\lambda \in P/G_\lambda$ and the mapping $M \ni {}^t(x, y) \mapsto \phi(x)\psi(y)G_\lambda \in G/G_\lambda$ are onto-diffeomorphisms.

We take a 1-form θ_λ on M and a C^∞ -function on M for $Y \in \mathfrak{g}$ as follows:

$$\theta_\lambda = \langle \lambda, g^{-1} dg \rangle,$$

$$H_Y = \langle \lambda, g^{-1} Yg \rangle$$

where $g = \phi(x)\psi(y)$ and $(x, y) \in M$. From the definition of ψ , we have

$$\begin{aligned} \theta_\lambda &= \sum_{\substack{i=1 \dots m \\ j=1 \dots n_i}} y^{i,j} \langle \zeta_{ij}, \phi^{-1}(x) d\phi(x) \rangle \\ &\quad + \langle \lambda, \phi^{-1}(x) d\phi(x) \rangle + \langle \lambda, \psi^{-1}(x) d\psi(x) \rangle \end{aligned}$$

and from the definition of ϕ , we have

$$\begin{aligned} &\langle \zeta_{ij}, \phi^{-1}(x) d\phi(x) \rangle \\ &= dx^{i,j} + \langle \zeta_{ij}, g^{-1} dg \rangle + \sum_{k=1}^{n_i} x^{i,k} \langle \zeta_{ij}, [X_{i,k}, g^{-1} dg] \rangle \end{aligned}$$

where

$$g = \exp(x^{1,1} X_{1,1}) \cdots \exp(x^{1,n_1} X_{1,n_1}) \cdots \exp(x^{i-1,n_{i-1}} X_{i-1,n_{i-1}}).$$

Hence θ_λ is expressed as in the form (1.1) and similarly we can see that H_Y is expressed as in the form (1.2).

Now following the Kirillov-Kostant theory, we construct a unitary representation of G . We assume that the Lie algebra homomorphism

$$\rho \ni X \longmapsto -\sqrt{-1} \langle \lambda, X \rangle \in \sqrt{-1} \mathbf{R}$$

lifts to a unitary character η_λ of P . We denote by η_ρ the character of P such that $|\Omega|^{\frac{1}{2}}$ is the line bundle associated with η_ρ , where $|\Omega|^{\frac{1}{2}}$ denotes the square root of absolute value of the volume bundle on $G/(G \cap P)$. We put $\xi_\lambda = \eta_\lambda \eta_\rho$.

Let L_{ξ_λ} denote the line bundle associated with ξ_λ over the homogeneous space G/P . Then the space $C^\infty(L_{\xi_\lambda})$ of all complex valued C^∞ -sections of L_{ξ_λ} can be identified with

$$\{f \in C^\infty(G); f(gp) = \xi_\lambda(p)^{-1} f(g) \ (g \in G, p \in P)\}.$$

For any $g \in G$ we define an operator $\pi_\lambda^g(g)$ on $C^\infty(L_{\xi_\lambda})$: For $f \in C^\infty(L_{\xi_\lambda})$

$$(\pi_\lambda^g(g)f)(x) = f(g^{-1}x) \quad (x \in G).$$

By using the diffeomorphism $\mathbf{R}^n \ni x \mapsto \phi(x) \in G/P$, we can regard $\pi_\lambda^g(g)$ as an operator on $C^\infty(\mathbf{R}^n)$. Then $\pi_\lambda^g(g)$ is an isometry on $L^2(\mathbf{R}^n)$ so that we obtain a unitary representation of G on $L^2(\mathbf{R}^n)$.

The next theorem shows that the integral operator whose kernel function is the path integral $K_{\theta_\lambda, H_Y}(x'', x'; T)$ coincides with the operator $\pi_\lambda^g(\exp TY)$.

THEOREM 2.

$$\int K_{\theta_\lambda, H_Y}(x'', x'; T) f(x') dx' = (\pi_\lambda^p(\exp TY)f)(x'')$$

where

$$f \in C_c^\infty(\mathbf{R}^n).$$

PROOF. First for $x \in \mathbf{R}^n$ and $t \in \mathbf{R}$, we define $v(x, t) \in \mathbf{R}^n$ and $p(x, t) \in P$ by

$$\exp(-tY)\phi(x) = \phi(v(x, t))p(x, t).$$

Then we have

$$(\pi_\lambda^p(\exp TY)f)(x) = \xi_\lambda^{-1}(p(x, T))f(v(x, T)) \left| \frac{dv(x, T)}{dx} \right|^{\frac{1}{2}}. \tag{2.1}$$

The differential equations (1.4) corresponding to θ_λ and H_Y is written as the system

$$\begin{cases} \left\langle \zeta_{ij}, \phi^{-1}(w) \frac{d\phi(w)}{dt} - \phi^{-1}(w) Y\phi(w) \right\rangle = 0, & i = 1 \cdots m, j = 1 \cdots n_i, \\ w(0) = x. \end{cases}$$

By the uniqueness of the solution $w(x, t)$, we have $w(x, t) = v(x, -t)$. Hence from Theorem 1, we get

$$\begin{aligned} & K_{\theta_\lambda, H_Y}(x'', x'; T) \\ &= \delta(x'' - v(x', -T)) \exp \left\{ \sqrt{-1} \int_0^T \left\langle \lambda, p(x', -t) \frac{dp^{-1}(x', -t)}{dt} \right\rangle dt \right\} \left| \frac{dv(x', -T)}{dx'} \right|^{\frac{1}{2}}. \end{aligned}$$

By using the second kind coordinates on P , it is easy to show that

$$\exp \left\{ \sqrt{-1} \int_0^T \left\langle \lambda, p(x', -t) \frac{dp^{-1}(x', -t)}{dt} \right\rangle dt \right\} = \xi_\lambda(p(x', -T)).$$

Hence we get

$$\begin{aligned} K_{\theta_\lambda, H_Y}(x'', x'; T) &= \delta(x'' - v(x', -T)) \xi_\lambda(p(x', -T)) \left| \frac{dv(x', -T)}{dx'} \right|^{\frac{1}{2}} \\ &= \delta(x' - v(x'', T)) \xi_\lambda^{-1}(p(x'', T)) \left| \frac{dv(x'', T)}{dx''} \right|^{\frac{1}{2}}. \end{aligned} \tag{2.2}$$

Now Theorem 2 follows from (2.1) and (2.2). □

References

- [1] A. Alekseev, L.D. Faddeev and S. Shatashvili, Quantization of symplectic orbits of compact Lie groups by means of the functional integral, *J. Geometry and Physics* **5** (1989), 391–406.
- [2] A. Alekseev and S. Shatashvili, Path integral quantization of the coadjoint orbits of the Virasoro group and $2d$ gravity (preprint LOMI-E-16–88).
- [3] A. Alekseev and S. Shatashvili, From Geometric Quantization to Conformal Field Theory, *Comm. Math. Phys.* **128** (1990), 197–212.
- [4] L. Auslander and B. Kostant, Quantization and representation of solvable Lie groups, *Bull. Amer. Math. Soc.* **73** (1967), 692–695.
- [5] L. Auslander and B. Kostant, Polarization and unitary representations of solvable Lie groups, *Invent. Math.* **14** (1971), 255–354.
- [6] C. Chevalley, On the topological structure of solvable groups, *Ann. of Math.* **42** (1941), 668–675.
- [7] L.J. Corwin and F.P. Greenleaf, “Representations of nilpotent Lie groups and their applications,” Cambridge Univ. Press, Cambridge, 1989.
- [8] L.D. Faddeev and A.A. Slavnov, “Gauge fields: Introduction to quantum theory,” Benjamin Inc., Massachusetts, 1980.
- [9] R.P. Feynman, Space-time approach to non-relativistic quantum mechanics, *Rev. Mod. Phys.* **20** (1948), 367–387.
- [10] R.P. Feynman and A.R. Hibbs, “Quantum mechanics and path integrals,” Mc-Graw Hill Inc., New York, 1965.
- [11] H. Fujiwara, Affine structures of some solvable Lie groups, *Mem. Fac. Sci. Kyusyu Univ.* **33** (1979), 343–353.
- [12] C. Garrod, Hamiltonian Path-Integral Methods, *Rev. Mod. Phys.* **38** (1966), 483–494.
- [13] T. Hashimoto, K. Ogura, K. Okamoto, R. Sawae and H. Yasunaga, Kirillov-Kostant theory and Feynman path integrals on coadjoint orbits I, (to appear).
- [14] A.A. Kirillov, “Elements of the theory of representations,” Springer-Verlag, Berlin, 1976.
- [15] B. Kostant, Quantization and unitary representations. Part I. Prequantization, in “Lect. Notes in Math. Vol. 170,” Springer-Verlag, Berlin-Heidelberg-New York, 1970, pp. 87–208.
- [16] G.D. Mostow, Factor spaces of solvable groups, *Ann. of Math.* **60** (1954), 1–27.
- [17] L. Pukanszky, “Leçons sur les représentations des groupes,” Dunod, Paris, 1967.
- [18] A.G. Reiman and M.A. Semenov-Tjan-Sanskii, Current algebras and nonlinear partial differential equations, *Soviet Math. Dokl.* **21** (1980), 630–634.
- [19] S.S. Schweber, On Feynman Quantization, *J. Math. Phys.* **3** (1962), 831–842.
- [20] E. Witten, Coadjoint orbits of the Virasoro group, *Comm. Math. Phys.* **114** (1988), 1–53.
- [21] N. Woodhouse, “Geometric quantization,” Oxford university press, Oxford, 1980.

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