# Finite element approximations of harmonic differentials on a Riemann surface 

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## Introduction

In the present paper we aim to establish a method of finite element approximations on a Riemann surface. Our method matches the abstract definition of Riemann surfaces, and also will offer a new technique of high practical use in numerical calculation not only for the case of Riemann surfaces but also for the case of plane domains (cf. Mizumoto and Hara [18]). It is characteristic of our method that by adopting an ordinary finite element approximation on every parametric disk, the approximations of high precision are obtained.

Let $\bar{\Omega}$ be a compact bordered or closed Riemann surface. We choose a fixed finite collection $\Phi=\left\{z=\varphi_{j}(p), p \in U_{j} ; j=1, \cdots, m\right\}$ of local parameters $z$ $=\varphi_{j}(p)$ and parametric disks $U_{j}$ so that $\bar{\Omega} \subset \bigcup_{j=1}^{m} U_{j}$. 11 is devoted to construction of a triangulation $K$ of $\bar{\Omega}$ with width $h$ associated to $\Phi$ (cf. §1.2), a normal subdivision of $K$ (cf. §1.3), and a naturalized triangulation $K^{\prime}$ associated to $K$ (cf. §1.4). The triangulation $K$ of $\bar{\Omega}$ is constructed as the sum of subtriangulations $K_{j}(j=1, \cdots, m)$ in such a way that $\left|K_{j}\right| \subset U_{j}$, each 2-simplex $s$ of $K$ belongs to one and only one $K_{j}$, each $s \in K_{j}$ is natural (see $\S 1.2$ ) at most except for the case when it has a common side with another $s^{\prime} \in K_{k}(k \neq j)$, and the diameter of $\varphi_{j}(s)$ is at most $h$ for each $s \in K_{j}(j=1, \cdots, m)$. Let $K_{j}^{\prime}(j=1, \cdots$, $m$ ) be triangulations consisting of all 2 -simplices of $K_{j}$ which are not minor or major, and all naturalized simplices of $K_{j}$ (see $\S 1.4$ ). Then the triangulation $K^{\prime}$ is defined as the sum of $K_{j}^{\prime}(j=1, \cdots, m)$.

In §2, we introduce and investigate two spaces $\Lambda=\Lambda(K)$ and $\Lambda^{\prime}=\Lambda^{\prime}\left(K^{\prime}\right)$ of differentials. The space $\Lambda$ consists of locally exact differentials $\sigma_{h}$ such that for each 2-simplex $s \in K_{j}(j=1, \cdots, m)$ the coefficients of $\sigma_{h}$ are constant on $\varphi_{j}(s)$ except that $\sigma_{h}$ is modified on all lunes of minor or major simplices (see $\S 1.4$ and §2.1). To each $\sigma_{h} \in \Lambda$, we associate a differential $\sigma_{h}^{\prime}=F\left(\sigma_{h}\right)$ on $K^{\prime}$ whose coefficients are constant on $\varphi_{j}(s)$ for each 2-simplex $s \in K_{j}^{\prime}(j=1, \cdots, m)$ and which is equal to $\sigma_{h}$ on $\bar{\Omega}$ except for all lunes of $K$ (cf. §2.2). The space $\Lambda^{\prime}$ consists of all $\sigma_{h}^{\prime}=F\left(\sigma_{h}\right)\left(\sigma_{h} \in \Lambda\right)$. We shall investigate estimates of differences of

[^0]Dirichlet norms $\left\|\sigma_{h}\right\|_{\Omega}^{2}$ and $\left\|\sigma_{h}^{\prime}\right\|_{K^{\prime}}^{2}$ (see Lemma 2.2).
Let $\omega$ be a harmonic differential on $\Omega$ which satisfies some period conditions and some boundary conditions (see §3.1). The finite element approximations $\psi_{h}$ and $\omega_{h}^{\prime}$ of $\omega$ are defined in the spaces $\Lambda$ and $\Lambda^{\prime}$ respectively (cf. $\S 3.2$ and $\S 3.3$ resp.). The differential $\omega_{h}^{\prime}$ can be numerically calculated. $\S 3$ is devoted to error estimates of $\psi_{h}$ and $\omega_{h}$ for $\omega$, where $\omega_{h}=F^{-1}\left(\omega_{h}^{\prime}\right)$. We shall make use of Bramble and Zlámal's lemma (see Lemma 3.5). In Theorems 3.1 and 3.2, we obtain error estimates:

$$
\left\|\psi_{h}-\omega\right\|^{2} \leqq C h^{2} \quad \text { and } \quad\left\|\omega_{h}-\omega\right\|^{2} \leqq C^{\prime} h^{2}
$$

where $C$ and $C^{\prime}$ are constants which depend only on the differential $\omega$ and the smallest value of interior angles of triangles $\varphi_{j}(s)$ for all $s \in K_{j}^{\prime} \quad(j=1, \cdots, m)$. Further, in Theorem 3.2, we obtain an estimate for $\|\omega\|^{2}$ :

$$
\|\omega\|^{2} \leqq\left\|\omega_{h}^{\prime}\right\|^{2}+\varepsilon\left(\omega_{h}^{\prime}\right)
$$

in a special case (see $\S 3.2$ ), where $\varepsilon\left(\omega_{h}^{\prime}\right)$ is a quantity of $O\left(h^{2}\right)$ which can be numerically calculated.

Finally, in $\S 4$ we apply our results to numerical calculation of periodicity moduli of closed and compact bordered Riemann surfaces, and we shall show that calculation results for some concrete Riemann surfaces of genus one are fairly good. With respect to the problems of this type, there have been some investigations by means of finite-difference method (Gaier [11], [12], Mizumoto [13], [14], [15], Opfer [20], [21]).

Our treatment at critical points of a Riemann surface is closely related to that at boundary singularities on a plane (cf. Akin [2], Babuška [3], Babuška and Rosenzweig [4], Babuška, Szabo and Katz [5], Barnhill and Whiteman [6], Blackburn [7], Craig, Zhu and Zienkiewicz [10], Mizumoto and Hara [18], Opfer and Puri [22], Rivara [23], Schatz and Wahlbin [24], [25], Thatcher [28], Tsamasphyros [29], Whiteman and Akin [30], Yserentant [31]).

The results in the present paper (Theorems 3.1 and 3.2) may be generalized to the case of harmonic differentials on a higher dimensional Riemannian manifold.

## §1. Triangulation

1. Collection $\Phi$ of local parameters Let $\Omega$ be a closed Riemann surface or a subdomain of a Riemann surface $W$ whose closure $\bar{\Omega}$ is a compact bordered subregion of $W$. In the latter case, we assume that the boundary $\partial \Omega$ consists of a finite number of analytic arcs meeting at vertices $p_{k}(k=1, \cdots, v)$, and there exist parametric disks $V_{k}(k=1, \cdots, v)$ with the centers $p_{k}$ and local parameters $z$
$=\psi_{k}(p)$ by which $V_{k} \cap \bar{\Omega}$ are mapped onto sectors $\left\{|z| \leqq r_{k}\right\} \cap\left\{0 \leqq \arg z \leqq \beta_{k}\right\} \quad(0$ $\left.<\beta_{k} \leqq 2 \pi, \beta_{k} \neq \pi\right)$. For conformity, if $\Omega$ is a closed Riemann surface, then we interpret that $\Omega=W$.

By $\Phi=\left\{z=\varphi_{j}(p), U_{j} ; j=1, \cdots, m\right\}$ we denote a finite collection of local parameters $z=\varphi_{j}(p)(j=1, \cdots, m)$ and parametric disks $U_{j}(j=1, \cdots, m)$ on $W$ which satisfies the following conditions (i)~(iv):
(i) Each $U_{j} \quad(j=1, \cdots, m)$ is a parametric disk and by the mapping $z$ $=\varphi_{j}(p), U_{j}$ is mapped onto a disk $|z|<\rho_{j}$. Furthermore, each vertex $p_{k}(k$ $=1, \cdots, v)$ is the center of some $U_{j}$.
(ii) $\bar{\Omega}$ is covered by $\left\{U_{j}\right\}_{j=1}^{m}$.
(iii) If $U_{j} \cap U_{k} \neq \varnothing$, then there exists a constant $L(>1)$ such that for the mapping $\zeta=f(z) \equiv \varphi_{k} \circ \varphi_{j}^{-1}(z), 1 / L<\left|f^{\prime}(z)\right|<L$ on $\varphi_{j}\left(\mathrm{U}_{j} \cap \mathrm{U}_{k}\right)$.
(iv) If $U_{j} \cap \partial \Omega \neq \varnothing$, then $\varphi_{j}\left(U_{j} \cap \Omega\right)$ is a half disk $\left\{|z|<\rho_{j}\right\} \cap\{\operatorname{Im} z>0\}$ or a sector $\left\{|z|<\rho_{j}\right\} \cap\left\{0<\arg z<\alpha_{j}\right\} \quad\left(0<\alpha_{j} \leqq 2 \pi, \alpha_{j} \neq \pi\right)$.

In the latter case of (iv), by the mapping $\zeta=\left(\varphi_{j}(p)\right)^{\pi / \alpha_{j}}, \mathrm{U}_{j} \cap \Omega$ is mapped onto a half disk $\left\{|\zeta|<\rho_{j}^{\pi / \alpha}\right\}_{j} \cap\{\operatorname{Im} \zeta>0\}$. In this case we define anew $z=\varphi_{j}(p)$ and $\rho_{j}$ by $\zeta=\left(\varphi_{j}(p)\right)^{\pi / \alpha_{j}}$ and $\rho_{j}^{\pi / \alpha_{j}}$ respectively. Then, the local parameter $z$ $=\varphi_{j}(p)$ is no longer conformal at the center of $U_{j}$.
2. Triangulation $K$ associated to $\Phi$ For the collection $\Phi$ of local parameters and parametric disks defined in §1.1, and for a sufficiently small positive number $h$, we construct a triangulation $K=K^{h}$ of $\bar{\Omega}$ which satisfies the following conditions (i) $\sim(\mathrm{v})$. This is called a triangulation of $\bar{\Omega}$ with width $h$ associated to $\Phi$.
(i) Each point at which $\partial \Omega$ is not analytic is a carrier of some 0 -simplex of $K$.
(ii) $K$ is the sum of subtriangulations $K_{1}, \cdots, K_{m}$ of $K$ such that each 2simplex of $K$ belongs to one and only one $K_{j}(j=1, \cdots, m)$, and the carrier $|s|$ of each 2-simplex $s$ of $K_{j}$ is contained in $U_{j}$.

If a 1 -simplex $e \in K_{j}$ does not belong to another $K_{k}(k \neq j)$, or a 1-simplex $e$ belongs to $K_{j} \cap K_{k}(j \neq k)$ and the mapping $\varphi_{k}{ }^{\circ} \varphi_{j}^{-1}$ is an affine transformation, then $e$ is said to be linear. If each edge of a 2 -simplex $s \in K_{j}$ is linear and $\varphi_{j}(s)$ is an ordinary triangle, then $s$ is called a natural simplex.
(iii) Each 2-simplex $s \in K_{j}$ which has not a common edge with any 2simplex of another $K_{k}(k \neq j)$, is a natural simplex.

A 2 -simplex of $K_{k}$ which has a common edge with a 2 -simplex $s \in K_{j}$ $(j \neq k)$, is said to be an adjoint (simplex) of $s$ and is denoted by $s^{\prime}$.
(iv) For each pair of a 2 -simplex $s \in K_{j}$ and its adjoint $s^{\prime} \in K_{k}$ with a common edge $e$, either one of the following three cases (a), (b), (c) occurs.
(a) Both $s$ and $s^{\prime}$ are natural simplices.
(b) $\varphi_{j}(s)$ is a curvilinear triangle such that $\varphi_{j}(e)$ is a strictly concave arc w.r.t. $\varphi_{j}(s), \varphi_{k}\left(s^{\prime}\right)$ is an ordinary triangle, and all edges of $s$ and $s^{\prime}$ except for $e$ are linear (cf. Fig. 1). Then $s$ is called a minor simplex. The case where $s^{\prime}$ is a minor simplex and $s$ is its adjoint may also occur.


Fig. 1
(c) $\varphi_{j}(s)$ is a curvilinear triangle such that $\varphi_{j}(e)$ is a strictly convex arc w.r.t. $\varphi_{j}(s), \varphi_{k}\left(s^{\prime}\right)$ is an ordinary triangle, and all edges of $s$ and $s^{\prime}$ except for $e$ are linear (cf. Fig. 2). Then $s$ is called a major simplex. The case where $s^{\prime}$ is a major simplex and $s$ is its adjoint may also occur.


Fig. 2
If $s$ is a minor or major simplex of $K_{j}$, then it is assumed that $\left|s^{\prime}\right| \subset U_{j}$ for its adjoint $s^{\prime}$.
(v) For each 2-simplex $s \in K_{j}(j=1, \cdots, m), d\left(\varphi_{j}(s)\right) \leqq h$, where throughout the present paper we denote the diameter of a region $G$ by $d(G)$.

Next, we assume that for the fixed $\Phi$ the class of the triangulations $K=K^{h}$ satisfies the following conditions ( $\mathrm{i}^{\prime}$ ) and (ii'):
(i') For each $j=1, \cdots, m$ the union of carriers of all minor and major simplices of $K_{j}$, and all their adjoints is contained in a closed subset $R_{j}$ of $U_{j} \cap \bar{\Omega}$ which is independent of the individual triangulation $K$.
(ii') The number $N$ of minor and major simplices of $K$ satisfies the inequality:

$$
\begin{equation*}
N \leqq M \cdot \frac{1}{h} \tag{1.1}
\end{equation*}
$$

where $M$ is a constant which is independent of the individual triangulation $K$.
3. Normal subdivision of triangulation $K$ For a triangulation $K=K^{h}$ of $\bar{\Omega}$ with width $h$ associated to $\Phi$ we can construct a subdivision $K^{1}=K^{1, h / 2}$, called the normal subdivision of $K=K^{h}$ by the following procedure:
(i) $K^{1}$ is the sum of the subtriangulations $K_{1}^{1}, \cdots, K_{m}^{1}$ which are the subdivisions of $K_{1}, \cdots, K_{m}$ respectively which are defined in the following (ii), (iii).
(ii) If $s \in K_{j}$ is a 2 -simplex which is not minor or major, then $s$ is subdivided to four 2 -simplices $s_{1}, s_{2}, s_{3}$ and $s_{4}$ of $K_{j}^{1}$ so that $\varphi_{j}\left(s_{1}\right), \varphi_{j}\left(s_{2}\right)$, $\varphi_{j}\left(s_{3}\right)$ and $\varphi_{j}\left(s_{4}\right)$ are mutually congruent ordinary triangles as in Fig. 3.


Fig. 3
(iii) Let $s \in K_{j}$ and $s^{\prime} \in K_{k}$ be a minor (or major) simplex and its adjoint, and let $e_{1}, e_{2}$ and $e_{3}$ be edges of $s$ such that $e_{1}$ is the common edge of $s$ and $s^{\prime}$. We subdivide the edges $e_{1}, e_{2}$ and $e_{3}$ to two edges $e_{11}$ and $e_{12}, e_{21}$ and $e_{22}$, and $e_{31}$ and $e_{32}$ respectively so that $\varphi_{k}\left(e_{11}\right)$ and $\varphi_{k}\left(e_{12}\right), \varphi_{j}\left(e_{21}\right)$ and $\varphi_{j}\left(e_{22}\right)$, and


Fig. 4
$\varphi_{j}\left(e_{31}\right)$ and $\varphi_{j}\left(e_{32}\right)$ have the same length respectively. Then we subdivide the simplex $s$ to two minor (or major resp.) simplices $s_{1}$ and $s_{2}$ of $K_{j}^{1}$ and, two natural simplices $s_{3}$ and $s_{4}$ of $K_{j}^{1}$ so that $e_{11}, e_{12}, e_{21}, e_{22}, e_{31}$ and $e_{32}$ are edges of $s_{1}, s_{2}$ and $s_{3}$ (cf. Fig. 4). Here we note that such a subdivision is always possible if $h$ is sufficiently small.

We can easily see that the normal subdivision $K^{1}=\sum_{j=1}^{m} K_{j}^{1}$ is a triangulation of $\bar{\Omega}$ with width $h / 2+O\left(h^{2}\right)$ associated to $\Phi$ (cf. (1.10)).
4. Naturalized triangulation For each minor (or major) simplex $s \in K_{j}$ we
 $\subset|s|$ resp.) and $\left.\varphi_{j}(\varphi) s\right)$ is the ordinary triangle which has two common sides with $\varphi_{j}(s)$. Further we define a 2 -simplex $b \ell=b \ell(s)(\# \ell=\# \ell(s)$ resp.) with two edges whose carrier is the closed region $\overline{|4 s|-|s|}(\overline{|s|-|4 s|}$ resp.). b $\ell(s)$ (\# $\ell(s)$ resp.) is called the deficient (excessive resp.) lune of $s$.

Each triple of a minor (or major) simplex $s \in K_{j}$, its adjoint $s^{\prime} \in K_{k}$ and its deficient lune $b \ell$ (excessive lune $\# \ell$ resp.) is denoted by $\left(s, s^{\prime}, b \ell\right)\left(\left(s, s^{\prime}, \# \ell\right)\right.$ resp.), and is called a triple for a minor (major resp.) simplex $s$ or for a deficient (excessive resp.) lune $b \ell$ (\# $\#$ resp.) (cf. Fig. 5), where it is always assumed that $|b \ell|$ $\subset\left|s^{\prime}\right|$ for each ( $s, s^{\prime}, b \ell$ ).


Fig. 5

For simplicity of notation, we also denote $b \ell=b \ell(s)$ or $\# \ell=\# \ell(s)$ by $\ell$ $=\ell(s)$. If a minor or major simplex $s$ is in $K_{j}$, then we say that $\ell=\ell(s)$ is a lune of $K_{j}$ and write $l \in K_{j}$.

Now we shall define the naturalized triangulation $K^{\prime}$ associated to $K$.
First, $K_{j}^{\prime}(j=1, \cdots, m)$ are defined as triangulations such that the collection of all 2-simplices of $K_{j}^{\prime}$ consists of all 2-simplices of $K_{j}$ which are not minor or major, and of all naturalized simplices of minor and major ones of $K_{j}$. Then the triangulation $K^{\prime}$ is defined as the sum of $K_{j}^{\prime}(j=1, \cdots, m)$. We should note that $K^{\prime}$ is no longer a triangulation of $\bar{\Omega}$, and also is not an ordinary triangulation.
5. Parametrization of lunar domains Let $\left(s, s^{\prime}, \ell\right)$ be an arbitrary triple for a deficient or excessive lune $\ell$, and let $e_{1}$ and $e_{2}$ be two edges of $\ell$ such that $e_{1} \subset \partial s$. Further, let

$$
\begin{equation*}
z^{\prime}=(1-t) z_{1}+t z_{2} \quad(0 \leqq t \leqq 1) \tag{1.2}
\end{equation*}
$$

and

$$
\begin{equation*}
\zeta^{\prime \prime}=(1-t) \zeta_{1}+t \zeta_{2} \quad(0 \leqq t \leqq 1) \tag{1.3}
\end{equation*}
$$

be parameter representations of the oriented segments $\varphi_{j}\left(-e_{2}\right)$ and $\varphi_{k}\left(e_{1}\right)$ respectively. The representation (1.3) induces a parameter representation of the curve $\varphi_{j}\left(e_{1}\right)$ :

$$
\begin{equation*}
z^{\prime \prime}=g\left((1-t) \zeta_{1}+t \zeta_{2}\right) \quad(0 \leqq t \leqq 1) \tag{1.4}
\end{equation*}
$$

where $z=g(\zeta) \equiv \varphi_{j} \varphi_{k}^{-1}(\zeta)$. By (1.2) and (1.4) we obtain a parameter representation of the lunar domain $\varphi_{j}(\ell)$ :

$$
\begin{align*}
z= & z(t, \tau) \equiv(1-\tau) z^{\prime}+\tau z^{\prime \prime}  \tag{1.5}\\
& =(1-\tau)\left((1-t) z_{1}+t z_{2}\right)+\tau g\left((1-t) \zeta_{1}+t \zeta_{2}\right) \\
& (0 \leqq t \leqq 1,0 \leqq \tau \leqq 1) .
\end{align*}
$$

## 6. Area of lune

Lemma 1.1. Let $\left(s, s^{\prime}, \ell\right)$ be a triple for an arbitrary deficient or excessive lune $\ell$. Then, the estimate

$$
\begin{equation*}
A\left(\varphi_{j}(\ell)\right) \leqq \frac{h_{1}^{3}}{8}\left(\left|\frac{g^{\prime \prime}\left(\zeta_{1}\right)}{g^{\prime}\left(\zeta_{1}\right)^{2}}\right|+O\left(h_{1}\right)\right) \tag{1.6}
\end{equation*}
$$

holds, where throughout the present paper we denote the area of a region $G$ by $A(G), z=g(\zeta) \equiv \varphi_{j}{ }^{\circ} \varphi_{k}^{-1}(\zeta), h_{1}=d\left(\varphi_{j}(\ell)\right)$ and $\zeta_{1}$ is one of the vertices of the lunar domain $\varphi_{k}(\ell)$.

Proof. Here we shall preserve the notations in §1.5. By Taylor's expansion we have

$$
\begin{equation*}
z^{\prime \prime}-z_{1}=g^{\prime}\left(\zeta_{1}\right)\left(\zeta_{2}-\zeta_{1}\right) t+\frac{1}{2} g^{\prime \prime}\left(\zeta_{1}\right)\left(\zeta_{2}-\zeta_{1}\right)^{2} t^{2}+\cdots \tag{1.7}
\end{equation*}
$$

for the point $z^{\prime \prime}$ of (1.4) on $\varphi_{j}\left(e_{1}\right)$, and

$$
\begin{align*}
z^{\prime}-z_{1} & =t\left(z_{2}-z_{1}\right)  \tag{1.8}\\
& =g^{\prime}\left(\zeta_{1}\right)\left(\zeta_{2}-\zeta_{1}\right) t+\frac{1}{2} g^{\prime \prime}\left(\zeta_{1}\right)\left(\zeta_{2}-\zeta_{1}\right)^{2} t+\cdots
\end{align*}
$$

for the point $z^{\prime}$ of (1.2) on $\varphi_{j}\left(-e_{2}\right)$, where we assume that the triangulation $K$ is so chosen that $\varphi_{k}\left(e_{1}\right)$ is contained in a disk $V$ centered at $\zeta_{1}$ such that $\varphi_{k}^{-1}(V)$ $\subset U_{j} \cap U_{k}$. By (1.7) and (1.8) we find that the equality

$$
\begin{equation*}
z^{\prime \prime}-z^{\prime}=\left(\zeta_{2}-\zeta_{1}\right)^{2} \cdot \frac{t(t-1)}{2} \cdot g^{\prime \prime}\left(\zeta_{1}\right)+O\left(\left(\zeta_{2}-\zeta_{1}\right)^{3}\right) \tag{1.9}
\end{equation*}
$$

holds for the point $z^{\prime}$ of (1.2) on $\varphi_{j}\left(-e_{2}\right)$ and the point $z^{\prime \prime}$ of (1.4) on $\varphi_{j}\left(e_{1}\right)$ with common $t$.

Since $\left|\zeta_{2}-\zeta_{1}\right| \leqq h_{1}\left(1 /\left|g^{\prime}\left(\zeta_{1}\right)\right|+O\left(h_{1}\right)\right)$, the equality (1.9) implies

$$
\begin{equation*}
\left|z^{\prime \prime}-z^{\prime}\right| \leqq \frac{h_{1}^{2}}{8}\left(\left|\frac{g^{\prime \prime}\left(\zeta_{1}\right)}{g^{\prime}\left(\zeta_{1}\right)^{2}}\right|+O\left(h_{1}\right)\right) . \tag{1.10}
\end{equation*}
$$

Therefore we obtain the estimates

$$
\begin{aligned}
A\left(\varphi_{j}(\ell)\right) & \leqq\left|z_{2}-z_{1}\right| \cdot \max _{0 \leq t \leq 1}\left|z^{\prime}-z^{\prime \prime}\right| \\
& \leqq \frac{h_{1}^{3}}{8}\left(\left|\frac{g^{\prime \prime}\left(\zeta_{1}\right)}{g^{\prime}\left(\zeta_{1}\right)^{2}}\right|+O\left(h_{1}\right)\right) .
\end{aligned}
$$

## §2. Spaces of differentials

1. Subspace $\Lambda$ of $\Gamma_{c}$ Let $\Gamma_{c}^{0}=\Gamma_{c}^{0}(\bar{\Omega})$ be the set of all locally exact differentials $\sigma$ in the class $C^{0}$ on $\bar{\Omega}$ with the finite Dirichlet norm

$$
\|\sigma\|^{2}=\|\sigma\|_{\Omega}^{2}=\int_{\Omega} \sigma * \sigma<\infty
$$

where by $* \sigma$ we denote the conjugate differential of $\sigma$. Let $\Gamma_{c}=\Gamma_{c}(\bar{\Omega})$ be the completion of $\Gamma_{c}^{0}$. We should note that in Chapter 5 of Ahlfors and Sario [1], $\Gamma_{c}$ is defined as the completion of $\Gamma_{c}^{1} \equiv \Gamma_{c}^{0} \cap C^{1}$.

We define a subspace $\Lambda=\Lambda(K)$ of $\Gamma_{c}$ as the space of differentials $\sigma_{h}$ which satisfy the following conditions (i) $\sim$ (iv):
(i) $\sigma_{h} \in \Gamma_{c}$.
(ii) If $s \in K_{j}(j=1, \cdots, m)$ is a natural simplex, then

$$
\sigma_{h}=a_{0} d x+b_{0} d y \quad \text { on } \quad \varphi_{j}(s) \quad(z=x+i y),
$$

where $a_{0}$ and $b_{0}$ are constants.
(iii) Let ( $s, s^{\prime}$, b $\ell$ ) be a triple for a minor simplex $s$, and let $e_{1}$ and $e_{2}$ be two edges of $b \ell$ such that $-e_{1} \subset \partial s$. Then

$$
\sigma_{h}=a_{0} d x+b_{0} d y \quad \text { on } \quad \varphi_{j}(s)
$$

$$
\sigma_{h}=\alpha_{0} d \xi+\beta_{0} d \eta \quad \text { on } \quad \varphi_{k}\left(s^{\prime}\right)-\varphi_{k}(b \ell)
$$

and $\sigma_{h}$ is a harmonic differential in $b \ell$ which satisfies the boundary conditions

$$
\sigma_{h}=a_{0} d x+b_{0} d y \quad \text { along } \quad \varphi_{j}\left(e_{1}\right)
$$

and

$$
\sigma_{h}=\left(\alpha_{0} \frac{\partial \xi}{\partial x}+\beta_{0} \frac{\partial \eta}{\partial x}\right) d x+\left(\alpha_{0} \frac{\partial \xi}{\partial y}+\beta_{0} \frac{\partial \eta}{\partial y}\right) d y \quad \text { along } \quad \varphi_{j}\left(e_{2}\right)
$$

where $a_{0}, b_{0}, \alpha_{0}$ and $\beta_{0}$ are constants, and

$$
\zeta=f(z) \equiv \varphi_{k} \circ \varphi_{j}^{-1}(z) \quad(z=x+i y, \zeta=\xi+i \eta)
$$

(iv) Let $\left(s, s^{\prime}, \# \ell\right)$ be a triple for a major simplex $s$, and let $e_{1}$ and $e_{2}$ be two edges of $\# \ell$ such that $e_{1} \subset \partial$ s. Then

$$
\begin{array}{lll}
\sigma_{h}=a_{0} d x+b_{0} d y & \text { on } & \varphi_{j}(\text { ( } s), \\
\sigma_{h}=\alpha_{0} d \xi+\beta_{0} d \eta & \text { on } & \varphi_{k}\left(s^{\prime}\right),
\end{array}
$$

and $\sigma_{h}$ is a harmonic differential in $\# \ell$ which satisfies the boundary conditions

$$
\sigma_{h}=a_{0} d x+b_{0} d y \quad \text { along } \quad \varphi_{j}\left(e_{2}\right)
$$

and

$$
\sigma_{h}=\left(\alpha_{0} \frac{\partial \xi}{\partial x}+\beta_{0} \frac{\partial \eta}{\partial x}\right) d x+\left(\alpha_{0} \frac{\partial \xi}{\partial y}+\beta_{0} \frac{\partial \eta}{\partial y}\right) d y \quad \text { along } \quad \varphi_{j}\left(e_{1}\right)
$$

where $a_{0}, b_{0}, \alpha_{0}$ and $\beta_{0}$ are constants, and $\zeta=\xi+i \eta$ is as in (iii).
We note that $\sigma_{h} \in \Lambda$ is generally discontinuous on each edge of 2-simplices of $K$.
2. Space $\Lambda^{\prime}$ Let $K^{\prime}$ be the naturalized triangulation associated to $K$. For each differential $\sigma_{h} \in \Lambda$, we define the differential $\sigma_{h}^{\prime}$ on $K^{\prime}$ associated to $\sigma_{h}$ as the differential $\sigma_{h}^{\prime}$ which satisfies the following conditions (i) $\sim$ (iv):
(i) For each 2 -simplex $s \in K_{j}^{\prime}(j=1, \cdots, m)$

$$
\sigma_{h}^{\prime}=a_{0} d x+b_{0} d y \quad \text { on } \quad \varphi_{j}(s)
$$

where $a_{0}$ and $b_{0}$ are constants.
(ii) If $s \in K$ is a natural simplex, then

$$
\sigma_{h}^{\prime}=\sigma_{h} \quad \text { on } \quad|s| .
$$

(iii) If ( $s, s^{\prime}, b \ell$ ) is a triple for a minor simplex $s$, then

$$
\sigma_{h}^{\prime}=\sigma_{h} \quad \text { on } \quad|s| \cup\left|s^{\prime}\right|-|b \ell| .
$$

(iv) If $\left(s, s^{\prime}, \# \ell\right)$ is a triple for a major simplex $s$, then

$$
\sigma_{h}^{\prime}=\sigma_{h} \quad \text { on } \quad|4 s| \cup\left|s^{\prime}\right| .
$$

We should note that the differential $\sigma_{h}^{\prime}$ is defined just twice on each deficient lune $b \ell$, while it is never defined on any excessive lune \# $\ell$. In the former case, for each triple ( $s, s^{\prime}$, b $\ell$ ) we shall denote the differential $\sigma_{h}^{\prime}$ on $\hbar_{s} \in K_{j}^{\prime}$ and $s^{\prime} \in K_{k}^{\prime}$ by $\sigma_{h, \text {,ss }}^{\prime}$ and $\sigma_{h, s^{\prime}}^{\prime}$ respectively.

The space of all differentials $\sigma_{h}^{\prime}$ associated to $\sigma_{h} \in \Lambda$ is denoted by $\Lambda^{\prime}$ $=\Lambda^{\prime}\left(K^{\prime}\right)$. Let $\sigma_{h}^{\prime}$ and $\chi_{h}^{\prime}$ be two differentials of $\Lambda^{\prime}$. Then the inner product ( $\sigma_{h}^{\prime}$, $\chi_{h}^{\prime}$ ) of $\sigma_{h}^{\prime}$ and $\chi_{h}^{\prime}$ is defined by

$$
\begin{aligned}
\left(\sigma_{h}^{\prime}, \chi_{h}^{\prime}\right) & =\left(\sigma_{h}^{\prime}, \chi_{h}^{\prime}\right)_{\mathbf{K}^{\prime}} \\
& =\sum_{s \in K^{\prime}} \int_{|s|} \sigma_{h}^{\prime} * \chi_{h}^{\prime},
\end{aligned}
$$

and the norm $\left\|\sigma_{h}^{\prime}\right\|$ of $\sigma_{h}^{\prime}$ is defined by

$$
\left\|\sigma_{h}^{\prime}\right\|=\left\|\sigma_{h}^{\prime}\right\|_{K^{\prime}}=\sqrt{\left(\sigma_{h}^{\prime}, \sigma_{h}^{\prime}\right)_{K^{\prime}}} .^{1)}
$$

We see that $\sigma_{h}^{\prime}=F\left(\sigma_{h}\right)$ defines a one-to-one mapping of $\Lambda$ onto $\Lambda^{\prime}$.
3. Finite element interpolations Let $\sigma$ be an element of $\Gamma_{c}$. We define the finite element interpolation $\hat{\sigma}$ of $\sigma$ in the space $\Lambda$ as the differential uniquely determined by the following conditions (i) and (ii):
(i) $\hat{\sigma} \in \Lambda$;
(ii) For each 1 -simplex $e \in K$,

$$
\int_{e} \hat{\sigma}=\int_{e} \sigma
$$

## 4. Harmonic differentials on a lune

Lemma 2.1. Let $\ell=\ell(s)$ be a deficient or excessive lune of $K_{j}$, let $e_{1}$ and $e_{2}$ be two edges of $\ell$, and let $\sigma_{1}$ and $\sigma_{2}$ be exact differentials in the class $C^{0}$ on $\ell$ which satisfy the condition

[^1]$$
\int_{e_{1}} \sigma_{1}=-\int_{e_{2}} \sigma_{2}
$$

Further, let $\chi$ be the differential harmonic in $\ell$ and continuous on $\ell$ which satisfies the boundary conditions

$$
\chi=\sigma_{i} \quad \text { along } \quad e_{i} \quad(i=1,2)
$$

Then the inequalities

$$
\begin{align*}
\|\chi\|_{\ell}^{2} & \leqq \iint_{\varphi_{j}(\ell)} \max \left\{\left(a_{1}^{2}+{b_{1}}^{2}\right),\left({a_{2}}^{2}+{b_{2}}^{2}\right)\right\} d x d y  \tag{2.1}\\
& \leqq\left\|\sigma_{1}\right\|_{\ell}^{2}+\left\|\sigma_{2}\right\|_{\ell}^{2}
\end{align*}
$$

hold, where

$$
\begin{gathered}
\|\chi\|_{\ell}^{2}=\int_{|\ell|} \chi * \chi, \quad \text { etc., and } \\
\sigma_{1}=a_{1} d x+b_{1} d y \text { and } \sigma_{2}=a_{2} d x+b_{2} d y \text { on } \varphi_{j}(\ell)
\end{gathered}
$$

Proof. By making use of the parameter representation (1.5) of the lunar domain $\varphi_{j}(\ell)$, we define a differential $\sigma$ on $\ell$ by

$$
\sigma^{\circ} \varphi_{j}^{-1}(z)=(1-\tau) \sigma_{1}^{\circ} \varphi_{j}^{-1}(z)+\tau \sigma_{2} \circ \varphi_{j}^{-1}(z) \quad\left(z=z(t, \tau) \in \varphi_{j}(\ell)\right)
$$

We note that $\sigma$ satisfies the same boundary conditions as $\chi$ on $\partial \ell$. Since $\chi$ is harmonic in $\ell$, the inequality

$$
\begin{equation*}
\|\chi\|_{\ell}^{2} \leqq\|\sigma\|_{\ell}^{2} \tag{2.2}
\end{equation*}
$$

holds. Further, the inequalities

$$
\begin{align*}
\|\sigma\|_{\ell}^{2} & \leqq \iint_{\varphi_{j}(\ell)}\left((1-\tau) \sqrt{{a_{1}^{2}+b_{1}^{2}}^{2}}+\tau \sqrt{{a_{2}^{2}}^{2}+b_{2}^{2}}\right)^{2} d x d y  \tag{2.3}\\
& \leqq \iint_{\varphi_{j}(\ell)} \max \left\{\left(a_{1}^{2}+{b_{1}^{2}}^{2}\right), \quad\left(a_{2}^{2}+{b_{2}}^{2}\right)\right\} d x d y
\end{align*}
$$

hold. The inequalities (2.2) and (2.3) imply the inequality (2.1).

## 5. Difference of norms of $\sigma_{h}$ and $\sigma_{h}^{\prime}$

Lemma 2.2. Let $\sigma_{h}$ be an arbitrary differential of the space $\Lambda$ and let $\sigma_{h}^{\prime}$ $=F\left(\sigma_{h}\right)$.
(i) The inequalities

$$
\begin{align*}
\left\|\sigma_{h}\right\|^{2} & \leqq\left\|\sigma_{h}^{\prime}\right\|^{2}+\sum_{\sharp \ell \in K}\left\|\sigma_{h}\right\|_{\sharp \ell}^{2}  \tag{2.4}\\
& \leqq\left\|\sigma_{h}^{\prime}\right\|^{2}+\sum_{j=1}^{m} \sum_{\forall \ell \in K_{j}} A\left(\varphi_{j}(\# \ell)\right) \cdot\left(\frac{1}{\lambda} \int_{e_{2}} \sigma_{h}^{\prime}\right)^{2} \cdot(1+\kappa h)
\end{align*}
$$

hold, where $e_{2}$ is the edge of $\# \ell$ such that $\varphi_{j}\left(e_{2}\right)$ is a segment, $\lambda$ is the length of $\varphi_{j}\left(e_{2}\right)$ and $\kappa$ is a constant which depends only on the transformations $f(z)$ $=\varphi_{k}{ }^{\circ} \varphi_{j}^{-1}(z)$.
(ii)

$$
\begin{align*}
\left\|\sigma_{h}^{\prime}\right\|^{2} & \leqq\left\|\sigma_{h}\right\|^{2}+\sum_{b \ell \in K}\left(\left\|\sigma_{h, s s}^{\prime}\right\|_{b \ell}^{2}+\left\|\sigma_{h, s^{\prime}}^{\prime}\right\|_{b \ell}^{2}\right)  \tag{2.5}\\
& =\left\|\sigma_{h}\right\|^{2}+\sum_{j=1}^{m} \sum_{b \ell \in K_{j}}\left\{A\left(\varphi_{j}(b \ell)\right) \cdot\left(a_{0}^{2}+b_{0}^{2}\right)+A\left(\varphi_{k}(b \ell)\right) \cdot\left(\alpha_{0}^{2}+\beta_{0}^{2}\right)\right\},
\end{align*}
$$

where for each triple ( $s, s^{\prime}$, b $\ell$ ) the notations in (iii) of §2.1 are preserved.
Proof. (i) By Lemma 2.1 we see that for each triple $\left(s, s^{\prime}, b \ell\right)$

$$
\begin{equation*}
\left\|\sigma_{h}\right\|_{b \ell}^{2} \leqq\left\|\sigma_{h, b s}^{\prime}\right\|_{b \ell}^{2}+\left\|\sigma_{h, s^{\prime}}^{\prime}\right\|_{b \ell}^{2} . \tag{2.6}
\end{equation*}
$$

Hence the first inequality of (2.4) is obtained.
Let ( $s, s^{\prime}, \# \ell$ ) be a triple for an excessive lune $\# \ell$. We preserve the notations in (iv) of $\S 2.1$. We shall prove the inequality

$$
\begin{equation*}
\left\|\sigma_{h}\right\|_{\sharp \ell}^{2} \leqq A\left(\varphi_{j}(\# \ell)\right) \cdot\left(\frac{1}{\lambda} \int_{e_{2}} \sigma_{h}^{\prime}\right)^{2} \cdot(1+\kappa h), \tag{2.7}
\end{equation*}
$$

from which the second inequality of (2.4) follows.
By $\gamma$ and $\delta$ we denote the arguments of the oriented segments $\varphi_{j}\left(-e_{2}\right)$ and $\varphi_{k}\left(e_{1}\right)$ respectively. By making use of the parameter representation (1.5) of the lunar domain $\varphi_{j}(\# \ell)$, we define a differential $\sigma$ on $\# \ell$ by

$$
\begin{align*}
\sigma= & a d x+b d y  \tag{2.8}\\
\equiv & (1-\tau)\left(a_{0} \cos \gamma+b_{0} \sin \gamma\right) \cdot((\cos \gamma) d x+(\sin \gamma) d y) \\
& +\tau\left(\alpha_{0} \cos \delta+\beta_{0} \sin \delta\right) \cdot \\
& \cdot\left((\cos \delta)\left(\frac{\partial \xi}{\partial x} d x+\frac{\partial \xi}{\partial y} d y\right)+(\sin \delta)\left(\frac{\partial \eta}{\partial x} d x+\frac{\partial \eta}{\partial y} d y\right)\right) \\
& \quad\left(z=z(t, \tau) \in \varphi_{j}(\# \ell)\right) .
\end{align*}
$$

We note that $\sigma$ satisfies the same boundary conditions as $\sigma_{h}$ on $\partial(\# \ell)$. Hence

$$
\begin{equation*}
\left\|\sigma_{h}\right\|_{\sharp \ell}^{2} \leqq\|\sigma\|_{\sharp \ell}^{2} \leqq A\left(\varphi_{j}(\# \ell)\right) \max _{\varphi_{j}(\notin)}\left(a^{2}+b^{2}\right), \tag{2.9}
\end{equation*}
$$

since $\sigma_{h}$ is harmonic in \# .
From the equation (2.8) it follows that

$$
\begin{align*}
& \max _{\varphi_{j}(t \ell)}\left(a^{2}+b^{2}\right) \leqq \max \left\{\left(a_{0} \cos \gamma+b_{0} \sin \gamma\right)^{2}\right.  \tag{2.10}\\
& \\
& \left.\quad\left(\alpha_{0} \cos \delta+\beta_{0} \sin \delta\right)^{2} \max _{\varphi_{j}(t \ell)}\left|f^{\prime}(z)\right|^{2}\right\}
\end{align*}
$$

Further we note that

$$
\begin{equation*}
a_{0} \cos \gamma+b_{0} \sin \gamma=\frac{1}{\lambda} \int_{-e_{2}} \sigma_{h}^{\prime} \tag{2.11}
\end{equation*}
$$

and

$$
\begin{equation*}
\alpha_{0} \cos \delta+\beta_{0} \sin \delta=\frac{1}{\mu} \int_{e_{1}} \sigma_{h}^{\prime}=\frac{1}{\mu} \int_{-e_{2}} \sigma_{h}^{\prime} \tag{2.12}
\end{equation*}
$$

where

$$
\begin{equation*}
\lambda=\int_{\varphi_{j}\left(e_{2}\right)}|d z| \quad \text { and } \quad \mu=\int_{\varphi_{j}\left(e_{1}\right)}\left|f^{\prime}(z) d z\right| . \tag{2.13}
\end{equation*}
$$

By making use of the power series expansion of $f^{\prime}$ around a vertex $z_{1}$ of the lunar domain $\varphi_{j}(\# \ell)$, we see that

$$
\begin{equation*}
\max _{\varphi_{j}(\nexists)}\left|f^{\prime}(z)\right|^{2} \leqq\left|f^{\prime}\left(z_{1}\right)\right|^{2}\left(1+\kappa_{1} h\right) \tag{2.14}
\end{equation*}
$$

and

$$
\begin{equation*}
\mu \geqq\left(\left|f^{\prime}\left(z_{1}\right)\right|-\kappa_{2} h\right) \int_{\varphi_{j}\left(e_{2}\right)}|d z|=\lambda\left(\left|f^{\prime}\left(z_{1}\right)\right|-k_{2} h\right) \tag{2.15}
\end{equation*}
$$

with constants $\kappa_{1}, \kappa_{2}>0$ depending only on $f$. Then the estimate (2.7) follows from (2.9)~(2.15).
(ii) The inequality (2.5) is obvious from the definition of $\sigma_{h}^{\prime}$.

## §3. Finite element approximations

1. Formulation of problems Let $\left\{C_{1}, C_{2}, C_{3}\right\}$ be a partition to three parts of the boundary $\partial \Omega$ such that each $C_{j}(j=1,2,3)$ is a sum of boundary components of $\partial \Omega$, and let $\gamma_{k}(k=1, \cdots, \kappa)$ be the boundary components of $C_{2}$.

Let $\Theta$ be a differential in $\Gamma_{c}$ which satisfies the following conditions (i), (ii)
and (iii):
(i) If $U_{j \cap} \cap C_{1} \neq \varnothing$, then $\Theta^{\circ} \varphi_{j}^{-1}$ is harmonic on a neighborhood of $\varphi_{j}\left(U_{j} \cap C_{1}\right) ;$
(ii) $\Theta=0$ along $C_{2}$;
(iii) $\Theta$ is exact on a neighborhood of each boundary component of $C_{3}$, where the conditions (i), (ii) and (iii) may be ignored if $\partial \Omega=\varnothing$.
By $\Gamma_{\theta}$ we denote the subspace of $\Gamma_{c}$ consisting of all differentials $\sigma$ for which there exists a function $v$ on $\bar{\Omega}$ such that

$$
\begin{aligned}
d v & =\Theta-\sigma & & \text { on } \bar{\Omega}, \\
v & =0 & & \text { on } C_{1}, \\
v & =\text { const. } & & \text { on } \quad \gamma_{k} \quad(k=1, \cdots, \kappa) .
\end{aligned}
$$

By $\omega$ we denote the harmonic differential in $\Gamma_{\theta}$ uniquely determined by the conditions

$$
\begin{equation*}
\int_{\gamma_{k}} * \omega=0 \quad(k=1, \cdots, \kappa) \tag{3.1}
\end{equation*}
$$

and

$$
\begin{equation*}
* \omega=0 \quad \text { along } \quad C_{3} \tag{3.2}
\end{equation*}
$$

The differential $\omega$ can be constructed by the following procedure. Let $\chi$ be the harmonic component of $\Theta$ in the orthogonal decomposition of $\Gamma_{c}$ (cf. Chapter 5 of Ahlfors and Sario [1]), and let $u$ be the solution of the boundary value problem:

$$
\begin{aligned}
& u \text { is a harmonic function on } \Omega, \\
& u=0 \quad \text { on } C_{1}, \\
& u=\text { const. } \quad \text { on } \gamma_{k}, \\
& \int_{\gamma_{k}} * d u=\int_{\gamma_{k}} * \chi \quad(k=1, \cdots, \kappa)
\end{aligned}
$$

and

$$
* d u=* \chi \quad \text { along } \quad C_{3} .
$$

Then, $\omega=\chi-d u$. We note that the differential $\omega^{\circ} \varphi_{j}^{-1}(j=1, \cdots, m)$ is harmonic on $\varphi_{j}\left(U_{j} \cap \bar{\Omega}\right) .{ }^{1)}$

[^2]Lemma 3.1. The harmonic differential $\omega$ satisfies the minimal property

$$
\begin{equation*}
\|\omega\|=\min _{\sigma \in \Gamma_{\boldsymbol{\theta}}}\|\sigma\| . \tag{3.3}
\end{equation*}
$$

In the equality (3.3), the minimum of the right hand side is attained if and only if $\sigma=\omega$.

Proof. For each $\sigma \in \Gamma_{\boldsymbol{\theta}}$ there exists a function $v$ such that

$$
\begin{cases}d v=\sigma-\omega, &  \tag{3.4}\\ v=0 & \text { on } \quad C_{1} \\ v=\text { const. } & \text { on } \quad \gamma_{k} \quad(k=1, \cdots, \kappa)\end{cases}
$$

From (3.1), (3.2) and (3.4) it follows that

$$
\begin{align*}
(\sigma-\omega, \omega) & =\int_{\partial \Omega} v * \omega  \tag{3.5}\\
& =\int_{C_{1}} v * \omega+\sum_{k=1}^{\kappa} \int_{v_{k}} v * \omega+\int_{C_{3}} v * \omega=0
\end{align*}
$$

where

$$
(\sigma, \tau)=(\sigma, \tau)_{\Omega}=\int_{\Omega} \sigma * \tau
$$

The equality (3.5) implies that

$$
\|\sigma\|^{2}=\|\omega\|^{2}+\|\sigma-\omega\|^{2} \geqq\|\omega\|^{2} .
$$

In the last inequality, the equality holds if and only if $\sigma=\omega$.
The unique harmonic differential $\omega$ in $\Gamma_{\boldsymbol{\theta}}$ is called the harmonic solution in $\Gamma_{\theta}$.

Our aim is to obtain finite element approximations of $\omega$ in the spaces $\Lambda$ and $\Lambda^{\prime}$, and error estimates between them and $\omega$.
2. Finite element approximation $\psi_{h}$ in $\Lambda$ Let $\hat{\Theta}$ be the finite element interpolation of $\Theta$ in the space $\Lambda$. By $\Lambda_{\boldsymbol{\theta}}$ we denote the subspace of $\Lambda$ consisting of all differentials $\sigma_{h} \in \Lambda$ for which there exists a function $v$ on $\bar{\Omega}$ such that

$$
\begin{aligned}
d v & =\hat{\Theta}-\sigma_{h}, & & \\
v & =0 & & \text { on } \quad C_{1}, \\
v & =\text { const. } & & \text { on } \quad \gamma_{k} \quad(k=1, \cdots, \kappa) .
\end{aligned}
$$

By $\psi_{h}$ we denote the differential of $\Lambda_{\boldsymbol{\theta}}$ such that

$$
\begin{equation*}
\left\|\psi_{h}\right\|=\min _{\sigma_{h} \in \Lambda_{\boldsymbol{\theta}}}\left\|\sigma_{h}\right\| . \tag{3.6}
\end{equation*}
$$

We call $\psi_{h}$ the finite element approximation of $\omega$ in the space $\Lambda$.
Next, we consider the special case where the differential $\Theta$ satisfies the condition:

$$
\Theta=0 \quad \text { along } \quad C_{1}
$$

We denote such a differential $\Theta$ by $\Theta_{0}$. Since $\Lambda_{\boldsymbol{\Theta}_{0}} \subset \Gamma_{\boldsymbol{\theta}_{0}}$, we see that

$$
\begin{equation*}
\|\omega\| \leqq\left\|\psi_{h}\right\| . \tag{3.7}
\end{equation*}
$$

Lemma 3.2. (i) In the case of general $\Theta$, the equality

$$
\begin{equation*}
\left\|\psi_{h}-\omega\right\|=\min _{\sigma_{h} \in \Lambda_{\theta}}\left\|\sigma_{h}-\omega\right\| \tag{3.8}
\end{equation*}
$$

holds, where the minimum is attained if and only if $\sigma_{h}=\psi_{h}$.
(ii) In the case of $\Theta=\Theta_{0}$, the equality

$$
\begin{equation*}
\left\|\psi_{h}-\omega\right\|^{2}=\left\|\psi_{h}\right\|^{2}-\|\omega\|^{2} \tag{3.9}
\end{equation*}
$$

holds.
Proof. (i) First, by a method similar to (3.5), it is shown that

$$
\begin{equation*}
\left(\omega, \sigma_{h}-\psi_{h}\right)=0 \quad \text { for each } \quad \sigma_{h} \in \Lambda_{\Theta} \tag{3.10}
\end{equation*}
$$

By (3.6), standard arguments imply that

$$
\begin{equation*}
\left(\psi_{h}, \sigma_{h}-\psi_{h}\right)=0 \quad \text { for } \quad \text { each } \quad \sigma_{h} \in \Lambda_{\boldsymbol{\Theta}} \tag{3.11}
\end{equation*}
$$

From (3.10) and (3.11), it follows that

$$
\left\|\omega-\sigma_{h}\right\|^{2}=\left\|\omega-\psi_{h}\right\|^{2}+\left\|\sigma_{h}-\psi_{h}\right\|^{2} \geqq\left\|\omega-\psi_{h}\right\|^{2} .
$$

In the last inequality, the equality holds if and only if $\sigma_{h}=\psi_{h}$.
(ii) Since $\Lambda_{\boldsymbol{\theta}_{0}} \subset \Gamma_{\boldsymbol{\theta}_{0}}$, both $\psi_{h}$ and $\omega$ are elements of $\Gamma_{\boldsymbol{\theta}_{0}}$. Hence, by (3.5) $\left(\omega, \psi_{h}-\omega\right)=0$ and thus

$$
\left\|\psi_{h}-\omega\right\|^{2}=\left\|\psi_{h}\right\|^{2}-\|\omega\|^{2} .
$$

From (3.11) the following lemma immediately follows.
Lemma 3.3. In the case of general $\Theta$, the equality

$$
\begin{equation*}
\left\|\sigma_{h}-\psi_{h}\right\|^{2}=\left\|\sigma_{h}\right\|^{2}-\left\|\psi_{h}\right\|^{2} \tag{3.12}
\end{equation*}
$$

holds for each $\sigma_{h} \in \Lambda_{\boldsymbol{\theta}}$.
3. Finite element approximation $\omega_{h}^{\prime}$ in $\Lambda^{\prime}$ Let $\Lambda_{\theta}^{\prime}=\left\{\sigma_{h}^{\prime} \mid \sigma_{h}^{\prime}=F\left(\sigma_{h}\right), \sigma_{h} \in \Lambda_{\Theta}\right\}$. By $\omega_{h}^{\prime}$ we denote the differential of $\Lambda_{\boldsymbol{\theta}}^{\prime}$ such that

$$
\begin{equation*}
\left\|\omega_{h}^{\prime}\right\|=\min _{\sigma_{h} \in \Lambda_{\theta}}\left\|\sigma_{h}^{\prime}\right\| \tag{3.13}
\end{equation*}
$$

We call $\omega_{h}^{\prime}$ the finite element approximation of $\omega$ in the space $\Lambda^{\prime}$.
Lemma 3.4. The equality

$$
\begin{equation*}
\left\|\sigma_{h}^{\prime}-\omega_{h}^{\prime}\right\|^{2}=\left\|\sigma_{h}^{\prime}\right\|^{2}-\left\|\omega_{h}^{\prime}\right\|^{2} \tag{3.14}
\end{equation*}
$$

holds for each $\sigma_{h}^{\prime} \in \Lambda_{\boldsymbol{\theta}}^{\prime}$.
Proof. By (3.13), standard arguments imply that

$$
\begin{equation*}
\left(\omega_{h}^{\prime}, \sigma_{h}^{\prime}-\omega_{h}^{\prime}\right)=0 \quad \text { for each } \sigma_{h}^{\prime} \in \Lambda_{\theta}^{\prime} \tag{3.15}
\end{equation*}
$$

This implies (3.14).
4. Lemma of Bramble and Zlámal The following lemma is due to J. H. Bramble and M. Zlámal (cf. [9]).

Lemma 3.5. Let $\Delta$ be a closed triangle on the $z$-plane $(z=x+i y)$ with $d(\Delta)$ $\leqq h$, let $v$ be a function of the class $C^{2}$ defined on $\Delta$ such that $v=0$ at each vertex of $\Delta$. Then, the inequality

$$
\begin{align*}
& \iint_{\Delta}\left(\left(\frac{\partial v}{\partial x}\right)^{2}+\left(\frac{\partial v}{\partial y}\right)^{2}\right) d x d y  \tag{3.16}\\
& \quad \leqq \frac{B}{\sin ^{2} \theta} h^{2} \iint_{\Delta}\left(\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2}+2\left(\frac{\partial^{2} v}{\partial x \partial y}\right)^{2}+\left(\frac{\partial^{2} v}{\partial y^{2}}\right)^{2}\right) d x d y
\end{align*}
$$

holds, where $B$ is an absolute constant and $\theta$ is the smallest interior angle of the triangle $\Delta$.

## 5. Pointwise estimate

Lemma 3.6. Let $\Delta$ be a closed curvilinear triangle on the $z$-plane $(z=x+i y)$ with $d(\Delta) \leqq h$ which is the image of some 2-simplex $s \in K_{j}(j=1, \cdots, m)$ by $z$ $=\varphi_{j}(p)$, and let $v$ be a function of the class $C^{2}$ defined on $\Delta$ such that $v=0$ at each vertex of $\Delta$. Then,

$$
\begin{aligned}
& \left|\frac{\partial v}{\partial x}\right|, \quad\left|\frac{\partial v}{\partial y}\right| \\
& \quad \leqq h \cdot \frac{4}{\sin \theta} \max _{z \in A}\left(\left|\frac{\partial^{2} v}{\partial x^{2}}\right|+2\left|\frac{\partial^{2} v}{\partial x \partial y}\right|+\left|\frac{\partial^{2} v}{\partial y^{2}}\right|\right)(1+\kappa h)
\end{aligned}
$$

on $\Delta$, where $\theta$ is the smallest interior angle of the ordinary triangle which has common vertices with $\Delta$, and $\kappa$ is a constant which depends only on $f(z)$ $=\varphi_{k}{ }^{\circ} \varphi_{j}^{-1}(z)$.

Proof. (Cf. Theorem 3.1 of Strang and Fix [27].) Let $z_{0}=x_{0}+i y_{0}$ be a fixed point and $z=x+i y$ an arbitrary point in $\Delta$, and let $k=x-x_{0}$ and $l=y$ $-y_{0}$. Here we choose the point $z_{0}$ so that for each $z \in \Delta$ the segment between $z_{0}$ and $z$ is contained in $\Delta$.

By Taylor's theorem we have that

$$
v(z)=P(z)+r(z),
$$

where

$$
\begin{align*}
& P(z)=v\left(z_{0}\right)+\left(k \frac{\partial}{\partial x}+l \frac{\partial}{\partial y}\right) v\left(z_{0}\right), \\
& r(z)=\frac{1}{2!}\left(k \frac{\partial}{\partial x}+l \frac{\partial}{\partial y}\right)^{2} v\left(z^{\prime}\right) \tag{3.17}
\end{align*}
$$

with some point $z^{\prime}$ on the segment between $z_{0}$ and $z$. First, from (3.17) the estimate

$$
\begin{equation*}
|r(z)| \leqq \frac{h^{2}}{2} \max _{z \in \Delta}\left(\left|\frac{\partial^{2} v}{\partial x^{2}}\right|+2\left|\frac{\partial^{2} v}{\partial x \partial y}\right|+\left|\frac{\partial^{2} v}{\partial y^{2}}\right|\right) \quad(z \in \Delta) \tag{3.18}
\end{equation*}
$$

immediately follows. Let $z_{j}(j=1,2,3)$ be the vertices of $\Delta$. Then, by the assumption of the lemma

$$
\begin{equation*}
v\left(z_{j}\right)=P\left(z_{j}\right)+r\left(z_{j}\right)=0 \quad(j=1,2,3) \tag{3.19}
\end{equation*}
$$

Since $P(z)$ is a linear function of $x$ and $y$, by (3.19) we have the expression

$$
\begin{equation*}
P(z)=-r\left(z_{1}\right) \phi_{1}(z)-r\left(z_{2}\right) \phi_{2}(z)-r\left(z_{3}\right) \phi_{3}(z), \tag{3.20}
\end{equation*}
$$

where $\phi_{j}(j=1,2,3)$ are linear functions of $x$ and $y$ such that

$$
\phi_{j}\left(z_{k}\right)=\delta_{j k} \quad(j, k=1,2,3)
$$

with Kronecker's symbol $\delta_{j k}$. (3.18) and (3.20) imply the estimate

$$
\begin{align*}
\left|\frac{\partial P}{\partial x}\right| & \leqq\left|r\left(z_{1}\right)\right|\left|\frac{\partial \phi_{1}}{\partial x}\right|+\left|r\left(z_{2}\right)\right|\left|\frac{\partial \phi_{2}}{\partial x}\right|+\left|r\left(z_{3}\right)\right|\left|\frac{\partial \phi_{3}}{\partial x}\right|  \tag{3.21}\\
& \leqq \frac{3}{2} h^{2} \max _{z \in A}\left(\left|\frac{\partial^{2} v}{\partial x^{2}}\right|+2\left|\frac{\partial^{2} v}{\partial x \partial y}\right|+\left|\frac{\partial^{2} v}{\partial y^{2}}\right|\right) \cdot \max _{1 \leqq j \leq 3}\left|\frac{\partial \phi_{j}}{\partial x}\right| .
\end{align*}
$$

Here we can easily verify that

$$
\begin{equation*}
\left|\frac{\partial \phi_{j}}{\partial x}\right| \leqq \frac{1}{h_{1}} \cdot \frac{2}{\sin \theta} \quad(j=1,2,3) \tag{3.22}
\end{equation*}
$$

where $h_{1}$ is the diameter of the ordinary triangle which has common vertices with $\Delta$. From (3.21) and (3.22) it follows that

$$
\begin{equation*}
\left|\frac{\partial P}{\partial x}\right| \leqq 3 h \cdot \frac{1}{\sin \theta} \max _{z \in \mathcal{A}}\left(\left|\frac{\partial^{2} v}{\partial x^{2}}\right|+2\left|\frac{\partial^{2} v}{\partial x \partial y}\right|+\left|\frac{\partial^{2} v}{\partial y^{2}}\right|\right)(1+\kappa h) . \tag{3.23}
\end{equation*}
$$

By Taylor's theorem we have that

$$
\frac{\partial v(z)}{\partial x}=\frac{\partial v\left(z_{0}\right)}{\partial x}+\left(k \frac{\partial}{\partial x}+l \frac{\partial}{\partial y}\right) \frac{\partial}{\partial x} v\left(z^{\prime \prime}\right)
$$

with some point $z^{\prime \prime}$ on the segment between $z_{0}$ and $z$. Since $\partial v\left(z_{0}\right) / \partial x$ $=\partial P\left(z_{0}\right) / \partial x$ and

$$
\left|\left(k \frac{\partial}{\partial x}+l \frac{\partial}{\partial y}\right) \frac{\partial}{\partial x} v\left(z^{\prime \prime}\right)\right| \leqq h \max _{z \in \Lambda}\left(\left|\frac{\partial^{2} v}{\partial x^{2}}\right|+\left|\frac{\partial^{2} v}{\partial x \partial y}\right|\right)
$$

by (3.23) we obtain the estimate

$$
\left|\frac{\partial v(z)}{\partial x}\right| \leqq \frac{4 h}{\sin \theta} \max _{z \in A}\left(\left|\frac{\partial^{2} v}{\partial x^{2}}\right|+2\left|\frac{\partial^{2} v}{\partial x \partial y}\right|+\left|\frac{\partial^{2} v}{\partial y^{2}}\right|\right)(1+\kappa h) .
$$

Analogously the estimate for $|\partial v / \partial y|$ is obtained.

## 6. Approximation by $\psi_{h}$

Theorem 3.1. Let $\omega$ be the harmonic solution in $\Gamma_{\theta}$ defined in $\S 3.1$ and let $\psi_{h}$ be the finite element approximation of $\omega$ in the space $\Lambda$. Then,

$$
\begin{align*}
& \left\|\psi_{h}-\omega\right\|^{2}  \tag{3.24}\\
& \leqq \frac{h^{2}}{\sin ^{2} \theta}\left(B \sum_{j=1}^{m} \iint_{\varphi_{j}\left(\mathbf{K}_{j}^{\prime}\right)}\left(\left(\frac{\partial a}{\partial x}\right)^{2}+\left(\frac{\partial a}{\partial y}\right)^{2}+\left(\frac{\partial b}{\partial x}\right)^{2}+\left(\frac{\partial b}{\partial y}\right)^{2}\right) d x d y\right. \\
& \left.+C h^{2} \sum_{j=1}^{m} \max _{\varphi_{j}\left(R_{j}\right)}\left(\left(\frac{\partial a}{\partial x}\right)^{2}+\left(\frac{\partial a}{\partial y}\right)^{2}+\left(\frac{\partial b}{\partial x}\right)^{2}+\left(\frac{\partial b}{\partial y}\right)^{2}\right)\right),
\end{align*}
$$

where $B$ and $C$ are constants independent of the triangulation $K$ and the differential $\Theta, \theta$ is the smallest value of interior angles of all triangles $\varphi_{j}(s)$ $\left(s \in K_{j}^{\prime} ; j=1, \cdots, m\right)$,

$$
\omega=a d x+b d y \quad \text { on } \quad \varphi_{j}\left(U_{j} \cap \bar{\Omega}\right) \quad(j=1, \cdots, m),
$$

by $\varphi_{j}\left(K_{j}^{\prime}\right)$ we denote the image set by $\varphi_{j}$ of the carrier of $K_{j}^{\prime}$, and $R_{j}(j=1, \cdots, m)$
are the closed subsets of $U_{j} \cap \bar{\Omega}$ defined in (i') of $\S 1.2$.
Proof. First, by (i) of Lemma 3.2,

$$
\begin{equation*}
\left\|\psi_{h}-\omega\right\| \leqq\|\hat{\omega}-\omega\| \tag{3.25}
\end{equation*}
$$

Hence it is sufficient to estimate $\|\hat{\omega}-\omega\|$.
We have

$$
\begin{equation*}
\|\hat{\omega}-\omega\|_{\Omega}^{2}=\sum_{j=1}^{m} \sum_{s \in K_{j}}\|\hat{\omega}-\omega\|_{s}^{2} . \tag{3.26}
\end{equation*}
$$

Here we note that $\omega^{\circ} \varphi_{j}^{-1}(j=1, \cdots, m)$ is of the class $C^{1}$ on $\varphi_{j}\left(U_{j} \cap \bar{\Omega}\right)$. Then, by Lemma 3.5,

$$
\begin{align*}
& \|\hat{\omega}-\omega\|_{s}^{2}  \tag{3.27}\\
& \leqq \frac{B}{\sin ^{2} \theta} h^{2} \iint_{\varphi_{j}(s)}\left(\left(\frac{\partial a}{\partial x}\right)^{2}+\left(\frac{\partial a}{\partial y}\right)^{2}+\left(\frac{\partial b}{\partial x}\right)^{2}+\left(\frac{\partial b}{\partial y}\right)^{2}\right) d x d y
\end{align*}
$$

for each natural simplex $s$ of $K_{j}$. For simplicity, we denote the right hand side of (3.27) by $I\left[\varphi_{j}(s)\right]$.

For a triple ( $s, s^{\prime}, \ell$ ) fo. a minor simplex $s$, we denote the differential $\hat{\omega}^{\prime}$ on $\forall s \in K_{j}^{\prime}$ and $s^{\prime} \in K_{k}^{\prime}$ by $\hat{\omega}_{\natural s}^{\prime}$ and $\hat{\omega}_{s^{\prime}}^{\prime}$ respectively. Then, by Lemma 2.1

$$
\begin{equation*}
\|\hat{\omega}-\omega\|_{\ell}^{2} \leqq\left\|\hat{\omega}_{\sharp s}^{\prime}-\omega\right\|_{\ell}^{2}+\left\|\hat{\omega}_{s^{\prime}}^{\prime}-\omega\right\|_{\ell}^{2} . \tag{3.28}
\end{equation*}
$$

This inequality and Lemma 3.5 imply that

$$
\begin{align*}
\|\hat{\omega}-\omega\|_{s+s^{\prime}}^{2} & \leqq \hat{\omega}_{\mathrm{\natural} s}^{\prime}-\omega\left\|_{\mathfrak{G} s}^{2}+\right\| \hat{\omega}_{s^{\prime}}^{\prime}-\omega \|_{s^{\prime}}^{2}  \tag{3.29}\\
& \leqq I\left[\varphi_{j}(\emptyset s)\right]+I\left[\varphi_{k}\left(s^{\prime}\right)\right] .
\end{align*}
$$

Let ( $s, s^{\prime}, \ell$ ) be a triple for a major simplex $s$. Then, by Lemma 3.5

$$
\begin{equation*}
\|\hat{\omega}-\omega\|_{s}^{2} \leqq I\left[\varphi_{j}(\text { Һs })\right]+\|\hat{\omega}-\omega\|_{\ell}^{2} \tag{3.30}
\end{equation*}
$$

and

$$
\begin{equation*}
\|\hat{\omega}-\omega\|_{s^{\prime}}^{2} \leqq I\left[\varphi_{k}\left(s^{\prime}\right)\right] . \tag{3.31}
\end{equation*}
$$

Let

$$
\begin{array}{ll}
\hat{\omega}=a_{0} d x+b_{0} d y & \text { on } \varphi_{j}(\text { Ł } s), \quad \text { and } \\
\hat{\omega}=\alpha_{0} d \xi+\beta_{0} d \eta & \text { on } \quad \varphi_{k}\left(s^{\prime}\right),
\end{array}
$$

where $a_{0}, b_{0}, \alpha_{0}$ and $\beta_{0}$ are constants. Then we define differentials $\hat{\omega}_{s}$ and $\hat{\omega}_{s^{\prime}+\ell}$ on $s$ and $s^{\prime}+\ell$ respectively by

$$
\hat{\omega}_{s}=a_{0} d x+b_{0} d y \quad \text { on } \quad \varphi_{j}(s), \text { and }
$$

$$
\hat{\omega}_{s^{\prime}+\ell}=\alpha_{0} d \xi+\beta_{0} d \eta \quad \text { on } \quad \varphi_{k}\left(s^{\prime}+\ell\right) .
$$

Then, by Lemma 2.1

$$
\begin{equation*}
\|\hat{\omega}-\omega\|_{\ell}^{2} \leqq\left\|\hat{\omega}_{s}-\omega\right\|_{\ell}^{2}+\left\|\hat{\omega}_{s^{\prime}+\ell}-\omega\right\|_{\ell}^{2} . \tag{3.32}
\end{equation*}
$$

Further, by Lemma 3.6

$$
\begin{align*}
& \left\|\hat{\omega}_{s}-\omega\right\|_{\ell}^{2}  \tag{3.33}\\
& \quad \leqq A\left(\varphi_{j}(\ell)\right) \cdot \frac{32 h^{2}}{\sin ^{2} \theta} \cdot \max _{\boldsymbol{\varphi}_{j^{(s)}}}\left(\left|\frac{\partial a}{\partial x}\right|+\left|\frac{\partial a}{\partial y}\right|+\left|\frac{\partial b}{\partial x}\right|+\left|\frac{\partial \dot{b}}{\partial y}\right|\right)^{2}(1+\kappa h)^{2}
\end{align*}
$$

and

$$
\begin{align*}
& \left\|\hat{\omega}_{s^{\prime}+\ell}-\omega\right\|_{\ell}^{2}  \tag{3.34}\\
& \quad \leqq A\left(\varphi_{k}(\ell)\right) \cdot \frac{32 h^{2}}{\sin ^{2} \theta} \cdot \max _{\varphi_{k}\left(s^{\prime}+\ell\right)}\left(\left|\frac{\partial \alpha}{\partial \xi}\right|+\left|\frac{\partial \alpha}{\partial \eta}\right|+\left|\frac{\partial \beta}{\partial \xi}\right|+\left|\frac{\partial \beta}{\partial \eta}\right|\right)^{2}(1+\kappa h)^{2},
\end{align*}
$$

where $\omega=a d x+b d y$ on $\varphi_{j}(s)$ and $\omega=\alpha d \xi+\beta d \eta$ on $\varphi_{k}\left(s^{\prime}+\ell\right)$.
By (3.25)~(3.34), Lemma 1.1 and (1.1), the estimate (3.24) is obtained.

## 7. Approximation by $\omega_{h}^{\prime}$

Theorem 3.2. (i) Let $\omega$ be the harmonic solution in $\Gamma_{\boldsymbol{\theta}}$ defined in §3.1, let $\omega_{h}^{\prime}$ be the finite element approximation of $\omega$ in the space $\Lambda^{\prime}$ and let $\omega_{h}=F^{-1}\left(\omega_{h}^{\prime}\right)$. Then

$$
\begin{align*}
& \left\|\omega_{h}-\omega\right\|^{2}  \tag{3.35}\\
& \left.\begin{array}{l}
\leqq \\
\sin ^{2} \theta \\
\left.\quad+B^{\prime} h^{2} \sum_{j=1}^{m} \sum_{j=1}^{m} \max _{\varphi_{j}\left(R_{j}\right)}\left(\left(\frac{\partial a}{\partial x}\right)^{2}+\left(\frac{\partial a}{\partial y}\right)^{2}+\left(\frac{\partial b}{\partial x}\right)^{2}+\left(\frac{\partial b}{\partial y}\right)^{2}\right)\right) \\
\quad+C^{\prime} h^{2} \sum_{j=1}^{m} \max _{\varphi_{j}\left(R_{j}\right)}\left(a^{2}+b^{2}\right)
\end{array}, \$\left(\frac{\partial a}{\partial x}\right)^{2}+\left(\frac{\partial b}{\partial x}\right)^{2}+\left(\frac{\partial b}{\partial y}\right)^{2}\right) d x d y \\
& \quad
\end{align*}
$$

where $A^{\prime}, B^{\prime}$ and $C^{\prime}$ are constants independent of the triangulation $K$ and the differential $\Theta$, and other notations are the same as in Theorem 3.1.
(ii) Let $\Theta_{0}$ be the differential defined in $\S 3.2$, let $\omega$ be the harmonic solution in $\Gamma_{\Theta_{0}}$ and let $\omega_{h}^{\prime}$ be the finite element approximation of $\omega$ in the space $\Lambda^{\prime}$. Then the estimate

$$
\begin{equation*}
\|\omega\|^{2} \leqq\left\|\omega_{h}^{\prime}\right\|^{2}+\varepsilon\left(\omega_{h}^{\prime}\right) \tag{3.36}
\end{equation*}
$$

holds with

$$
\begin{equation*}
\varepsilon\left(\omega_{h}^{\prime}\right) \equiv \sum_{j=1}^{m} \sum_{\# \ell \in K_{j}} A\left(\varphi_{j}(\# \ell)\right) \cdot\left(\frac{1}{\lambda} \int_{e_{2}} \omega_{h}^{\prime}\right)^{2} \cdot \max \left\{1,\left(\frac{\lambda}{\mu}\right)^{2} \max _{\varphi_{j}(\notin \ell)}\left|f^{\prime}(z)\right|^{2}\right\} \tag{3.37}
\end{equation*}
$$

where $e_{1}$ and $e_{2}$ are the edges of $\# \ell$ such that $\varphi_{j}\left(e_{2}\right)$ is a straight segment, $\lambda$ and $\mu$ are the lengths of the segments $\varphi_{j}\left(e_{2}\right)$ and $\varphi_{k}\left(e_{1}\right)$ resp., and $f(z) \equiv \varphi_{k}{ }^{\circ} \varphi_{j}^{-1}(z)$.

Proof. (i) First, note that

$$
\begin{equation*}
\left\|\omega_{h}-\omega\right\|^{2} \leqq 2\left\|\psi_{h}-\omega\right\|^{2}+2\left\|\omega_{h}-\psi_{h}\right\|^{2} \tag{3.38}
\end{equation*}
$$

From Lemmas 2.1, 2.2 and 3.3, and (3.13), it follows that

$$
\begin{align*}
&\left\|\omega_{h}-\psi_{h}\right\|^{2}=\left\|\omega_{h}\right\|^{2}-\left\|\psi_{h}\right\|^{2}  \tag{3.39}\\
& \leqq\left\|\omega_{h}^{\prime}\right\|^{2}-\left\|\psi_{h}\right\|^{2}+\sum_{\# \ell \in K}\left\|\omega_{h}\right\|_{\# \ell}^{2} \\
& \leqq\left\|\psi_{h}^{\prime}\right\|^{2}-\left\|\psi_{h}\right\|^{2}+\sum_{\# \ell \in K}\left\|\omega_{h}\right\|_{\# \ell}^{2} \\
& \leqq \sum_{j=1}^{m} \sum_{b \ell \in K_{j}}\left(A\left(\varphi_{j}(b \ell)\right) \cdot\left(a_{0}^{\prime 2}+b_{0}^{\prime 2}\right)+A\left(\varphi_{k}(b \ell)\right) \cdot\left(\alpha_{0}^{\prime 2}+{\beta_{0}^{\prime}}^{2}\right)\right) \\
& \quad+\sum_{j=1}^{m} \sum_{\# \ell \in K_{j}}\left(A\left(\varphi_{j}(\# \ell)\right) \cdot\left(a_{0}^{2}+b_{0}^{2}\right)+A\left(\varphi_{k}(\# \ell)\right) \cdot\left(\alpha_{0}{ }^{2}+\beta_{0}^{2}\right)\right),
\end{align*}
$$

where for each triple ( $s, s^{\prime}, b \ell$ ) for $b \ell \in K_{j}$

$$
\begin{array}{ll}
\psi_{h}^{\prime}=a_{0}^{\prime} d x+b_{0}^{\prime} d y & \text { on } \quad \varphi_{j}\left(\varphi_{s}\right) \quad \text { and } \\
\psi_{h}^{\prime}=\alpha_{0}^{\prime} d \xi+\beta_{0}^{\prime} d \eta & \text { on } \quad \varphi_{k}\left(\mathrm{~s}^{\prime}\right)
\end{array}
$$

and for each triple $\left(s, s^{\prime}, \# \ell\right)$ for $\# \ell \in K_{j}$

$$
\begin{array}{ll}
\omega_{h}=a_{0} d x+b_{0} d y & \text { on } \varphi_{j}(\text { (थs }) \quad \text { and } \\
\omega_{h}=\alpha_{0} d \xi+\beta_{0} d \eta & \text { on } \varphi_{k}\left(s^{\prime}\right)
\end{array}
$$

with constants $a_{0}^{\prime}, b_{0}^{\prime}, \alpha_{0}^{\prime}, \beta_{0}^{\prime}, a_{0}, b_{0}, \alpha_{0}$ and $\beta_{0}$.
In the inequality (3.39), we have

$$
\begin{align*}
& A\left(\varphi_{j}(b \ell)\right) \cdot\left(a_{0}^{\prime 2}+b_{0}^{\prime 2}\right)  \tag{3.40}\\
& \quad=\frac{A\left(\varphi_{j}(b \ell)\right)}{A\left(\varphi_{j}(s)\right)}\left\|\psi_{h}\right\|_{s}^{2} \\
& \quad \leqq 2 \frac{A\left(\varphi_{j}(b \ell)\right)}{A\left(\varphi_{j}(s)\right)}\left(\left\|\psi_{h}-\omega\right\|_{s}^{2}+\|\omega\|_{s}^{2}\right) \\
& \quad \leqq 2 \frac{A\left(\varphi_{j}(b \ell)\right)}{A\left(\varphi_{j}(s)\right)}\left\|\psi_{h}-\omega\right\|_{s}^{2}+2 A\left(\varphi_{j}(b \ell)\right) \cdot \max _{\varphi_{j}(s)}\left(a^{2}+b^{2}\right) .
\end{align*}
$$

Since we can easily verify that

$$
A\left(\varphi_{j}(\text { Ч } s)\right)>\frac{h_{1}{ }^{2}}{4} \sin \theta \quad\left(h_{1}=d\left(\varphi_{j}(\text { 七s })\right)\right),
$$

by Lemma 1.1 we have

$$
\begin{align*}
\frac{A\left(\varphi_{j}(b \ell)\right)}{A\left(\varphi_{j}(s)\right)} & =\frac{A\left(\varphi_{j}(b \ell)\right)}{A\left(\varphi_{j}(\llcorner s))-A\left(\varphi_{j}(b \ell)\right)\right.}  \tag{3.41}\\
& \leqq \frac{h}{2 \sin \theta}\left(\left|\frac{g^{\prime \prime}\left(\zeta_{1}\right)}{g^{\prime}\left(\zeta_{1}\right)^{2}}\right|+O(h)\right)
\end{align*}
$$

with the notations in Lemma 1.1. (3.40) and (3.14) imply

$$
\begin{align*}
& \sum_{j=1}^{m} \quad \sum_{b \ell \in K_{j}} A\left(\varphi_{j}(b \ell)\right) \cdot\left(a_{0}^{\prime 2}+b_{0}^{\prime 2}\right)  \tag{3.42}\\
& \quad \leqq \frac{C h}{\sin \theta} \sum_{j=1}^{m} \sum_{b \ell \in K_{j}}\left\|\psi_{h}-\omega\right\|_{s}^{2}+2 \sum_{j=1}^{m} \sum_{b \ell \in K_{j}} A\left(\varphi_{j}(b \ell)\right) \max _{\varphi_{j}(s)}\left(a^{2}+b^{2}\right)
\end{align*}
$$

where $C$ is a constant depending only on the transformations of local parameters. Since similar estimates for other terms of the right hand side of (3.39) are obtained, from (3.39) it follows that

$$
\begin{align*}
& \left\|\omega_{h}-\psi_{h}\right\|^{2}  \tag{3.43}\\
& \leqq \\
& \leqq \frac{C h}{\sin \theta}\left\|\omega_{h}-\omega\right\|^{2}+\frac{C h}{\sin \theta}\left\|\psi_{h}-\omega\right\|^{2} \\
& \quad+2 \sum_{j=1}^{m} \sum_{\ell \in K_{j}}\left(A\left(\varphi_{j}(\ell)\right) \max _{\varphi_{j}(s)}\left(a^{2}+b^{2}\right)+A\left(\varphi_{k}(\ell)\right) \max _{\varphi_{k^{(s)}}}\left(\alpha^{2}+\beta^{2}\right)\right),
\end{align*}
$$

where for each triple $\left(s, s^{\prime}, \ell\right)$ for $\ell \in K_{j}$

$$
\begin{array}{ll}
\omega=a d x+b d y & \text { on } \quad \varphi_{j}(s), \quad \text { and } \\
\omega=\alpha d \xi+\beta d \eta & \text { on } \quad \varphi_{k}\left(s^{\prime}\right)
\end{array}
$$

(3.38), (3.43), Theorem 3.1, Lemma 1.1 and (1.1) imply the estimate (3.35).
(ii) (3.7) and Lemma 3.3 and the proof of Lemma 2.2 (i) imply the inequalities

$$
\begin{aligned}
&\|\omega\|^{2} \leqq\left\|\psi_{h}\right\|^{2} \leqq\left\|\omega_{h}\right\|^{2} \\
& \leqq\left\|\omega_{h}^{\prime}\right\|^{2}+\sum_{j=1}^{m} \sum_{\# \ell \in K_{j}} A\left(\varphi_{j}(\# \ell)\right)\left(\frac{1}{\lambda} \int_{e_{2}} \omega_{h}^{\prime}\right)^{2} \cdot \\
& \cdot \max \left\{1,\left(\frac{\lambda}{\mu}\right)^{2} \max _{\varphi_{j}(\# \ell)}\left|f^{\prime}(z)\right|^{2}\right\} .
\end{aligned}
$$

From Lemma 2.2, (3.41), the proof of Theorem 3.1 and Theorem 3.2, the following corollary follows.

Corollary 3.1. Let $\omega$ and $\omega_{h}^{\prime}$ be the same as in Theorem 3.2, $\hat{\omega}$ be the finite element interpolation of $\omega$ in the space $\Lambda$, and $\hat{\omega}^{\prime}=F(\hat{\omega})$. Then, the estimate

$$
\begin{equation*}
\left\|\omega_{h}^{\prime}-\hat{\omega}^{\prime}\right\| \leqq A^{\prime \prime} h \tag{3.44}
\end{equation*}
$$

holds, where $A^{\prime \prime}$ is a constant dependent only on $\omega$ and $\theta$ in Theorem 3.1.

## § 4. Applications

1. Periodicity moduli of Riemann surfaces Let $\bar{\Omega}$ be a closed or compact bordered Riemann surface of genus 1 with no or one boundary component. Let $\{A, B\}$ be a canonical homology basis of $\bar{\Omega}$ such that $A \times B=1$. Then there exists a unique system of harmonic differentials $\{\phi, \varrho, \chi, \tau\}$ on $\Omega$ satisfying the period and boundary conditions:

$$
\begin{array}{ll}
\int_{B} \phi=\int_{B} \chi=1 . & \int_{A} \phi=\int_{A} \chi=0, \\
\int_{A} \varrho=\int_{A} \tau=-1, & \int_{B} \varrho=\int_{B} \tau=0, \\
\phi=\varrho=* \chi=* \tau=0 & \text { along } \partial \Omega \tag{4.3}
\end{array}
$$

and

$$
\begin{equation*}
\int_{\partial \Omega} * \phi=\int_{\partial \Omega} * \varrho=\int_{\partial \Omega} \chi=\int_{\partial \Omega} \tau=0, \tag{4.4}
\end{equation*}
$$

where the conditions (4.3) and (4.4) may be ignored if $\partial \Omega=\varnothing$. If $\partial \Omega=\varnothing$, then $\phi$ $=\chi$ and $\varrho=\tau$.

We can easily see that

$$
\left\{\begin{array}{l}
\|\phi\|^{2}=\int_{A} * \phi, \quad\|\varrho\|^{2}=\int_{B} * \varrho, \quad \text { and }  \tag{4.5}\\
(\phi, \varrho)=\int_{B} * \phi=\int_{A} * \varrho=0
\end{array}\right.
$$

We call

$$
p_{1}=\int_{A} * \phi \quad \text { and } \quad p_{2}=\int_{B} * \varrho
$$

periodicity moduli of $\dot{\Omega}$ with respect to $A$ and $B$ respectively, which are the
quantities determining the conformal structure of $\Omega$. By (4.1) $\sim(4.5)$ we see that

$$
\tau=-\frac{* \phi}{\|\phi\|^{2}} \quad \text { and } \quad \chi=\frac{* \varrho}{\|\varrho\|^{2}}
$$

These relations imply that

$$
\begin{equation*}
p_{1}=\|\boldsymbol{\phi}\|^{2}=\frac{1}{\|\tau\|^{2}} \quad \text { and } \quad p_{2}=\|\varrho\|^{2}=\frac{1}{\|\chi\|^{2}} \tag{4.6}
\end{equation*}
$$

If $\partial \Omega=\varnothing$, then

$$
\begin{equation*}
p_{1}=\|\phi\|^{2}=\frac{1}{\|\varrho\|^{2}}=\frac{1}{p_{2}} . \tag{4.7}
\end{equation*}
$$

By making use of a relation analogous to (4.7) for the modulus of quadrilaterals on the complex plane, Gaier [11] presented a method to obtain upper and lower bounds for the modulus by the finite difference approximation.
2. Calculation of periodicity moduli Let $\left\{\Theta_{1}, \Theta_{2}, \Theta_{3}, \Theta_{4}\right\}$ be a system of differentials in $\Gamma_{c}(\bar{\Omega})$ satisfying the period and boundary conditions:

$$
\begin{array}{ll}
\int_{B} \Theta_{1}=\int_{B} \Theta_{3}=1, & \int_{A} \Theta_{1}=\int_{A} \Theta_{3}=0 \\
\int_{A} \Theta_{2}=\int_{A} \Theta_{4}=-1, & \int_{B} \Theta_{2}=\int_{B} \Theta_{4}=0, \\
\Theta_{1}=\Theta_{2}=0 & \text { along } \quad \partial \Omega,
\end{array}
$$

and $\Theta_{3}$ and $\Theta_{4}$ are exact on a neighborhood of $\partial \Omega$. Here we interpret that $\partial \Omega$ $=C_{2}$ for $\Theta_{1}$ and $\Theta_{2}$, and $\partial \Omega=C_{3}$ for $\Theta_{3}$ and $\Theta_{4}$ in the notations in §3.1. We note that $\Theta_{1}, \Theta_{2}, \Theta_{3}$ and $\Theta_{4}$ satisfy the conditions for the differential $\Theta_{0}$ in $\S 3.2$. Then we can easily see that $\phi, \varrho, \chi$ and $\tau$ are the harmonic solutions in $\Gamma_{\theta_{1}}, \Gamma_{\theta_{2}}$, $\Gamma_{\theta_{3}}$ and $\Gamma_{\theta_{4}}$, respectively. Let $\phi_{h}^{\prime}, \varrho_{h}^{\prime}, \chi_{h}^{\prime}$ and $\tau_{h}^{\prime}$ be the finite element approximations of $\phi, \varrho, \chi$ and $\tau$ in the space $\Lambda^{\prime}$ respectively. Then by (ii) of Theorem 3.2 and (4.6), we obtain upper and lower bounds for $p_{1}$ and $p_{2}$ :

$$
\begin{equation*}
\frac{1}{\left\|\tau_{h}^{\prime}\right\|^{2}+\varepsilon\left(\tau_{h}^{\prime}\right)} \leqq p_{1} \leqq\left\|\phi_{h}^{\prime}\right\|^{2}+\varepsilon\left(\phi_{h}^{\prime}\right) \tag{4.8}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{1}{\left\|\chi_{h}^{\prime}\right\|^{2}+\varepsilon\left(\chi_{h}^{\prime}\right)} \leqq p_{2} \leqq\left\|\varrho_{h}^{\prime}\right\|^{2}+\varepsilon\left(\varrho_{h}^{\prime}\right) \tag{4.9}
\end{equation*}
$$

If $\partial \Omega=\varnothing$, then $\phi=\chi$ and $\varrho=\tau$, and thus (4.8) and (4.9) imply the inequalities

$$
\frac{1}{\left\|\varrho_{h}^{\prime}\right\|^{2}+\varepsilon\left(\varrho_{h}^{\prime}\right)} \leqq p_{1}=\frac{1}{p_{2}} \leqq\left\|\phi_{h}^{\prime}\right\|^{2}+\varepsilon\left(\phi_{h}^{\prime}\right) .
$$

3. Numerical example 1 (the case of a closed Riemann surface) Let $\bar{\Omega}$ be the two-sheeted covering surface with four branch points $z=-3,-1,1,3$ over the extended $z$-plane. Then $\Omega$ is a closed Riemann surface of genus one. A canonical homology basis $\{A, B\}$ of $\Omega$ is chosen as in Fig. 6. We aim to obtain good upper and lower approximate values of the periodicity moduli $p_{1}$ and $p_{2}$ of $\Omega$ with respect to $A$ and $B$ respectively.


Fig. 6

First, we construct a triangulation of the closed region:

$$
\bar{D}=\{z| | z \mid \leqq \sqrt{3}, \operatorname{Re} z \geqq 0, \operatorname{Im} z \geqq 0\}
$$

as in Fig. 7. The closed regions $G_{2}$ and $G_{3}$ are mapped onto the regions $G_{2}^{*}$ and $G_{3}^{*}$ resp. by the local parameters $\zeta=\varphi_{2}(z)=a \sqrt{z-1}$ and $w=\varphi_{3}(z)=b \log z \quad(a$ $=2(\sqrt{3}-1)^{1 / 2}$ and $\left.b=\sqrt{3}\right)$ respectively, where $a$ and $b$ are so determined that $|d \zeta / d z|=1$ and $|d w / d z|=1$ on $|z-1|=\sqrt{3}-1$ and $|z|=\sqrt{3}$ respectively. We construct ordinary triangulations $K_{2}^{*}$ and $K_{3}^{*}$ of $G_{2}^{*}$ and $G_{3}^{*}$ as in Fig. 7 respectively. By $K_{2}$ and $K_{3}$ we denote the image triangulations of $K_{2}^{*}$ and $K_{3}^{*}$ by the mappings $\varphi_{2}^{-1}$ and $\varphi_{3}^{-1}$ respectively. The triangulation $K_{1}$ of the region $G_{1}=\overline{D-\left(G_{2} \cup G_{3}\right)}$ in Fig. 7 is so constructed that each 2-simplex $s$ of $K_{1}$ is natural, minor or major according as $|s| \cap\left|K_{2}+K_{3}\right|=\varnothing,|s| \cap\left|K_{2}\right| \neq \varnothing$, or $|s| \cap\left|K_{3}\right| \neq \varnothing$, where if some intersection is a point then it is interpreted to be vacuous, and the local parameter $\varphi_{1}(z)$ of $K_{1}$ is the identity mapping $\varphi_{1}(z) \equiv z$.

A triangulation $L_{1}$ of the region $\bar{D}_{1}=\{z| | z \mid \geqq \sqrt{3}, \operatorname{Re} z \geqq 0, \operatorname{Im} z \geqq 0\}$ is defined by the reflection of the triangulation $L \equiv K_{1}+K_{2}+K_{3}$ with respect to


Fig. 7


Fig. 8
the circle $|z|=\sqrt{3}$ (cf. Fig. 8). Next we define a triangulation $L_{2}$ of the fourth quadrant by the reflection of the triangulation $L+L_{1}$ with respect to the real axis and then a triangulation $L_{3}$ of the left half-plane by the reflection of $L+L_{1}$ $+L_{2}$ with respect to the imaginary axis. Consequently, a triangulation $L_{4}$ of the extended $z$-plane is defined by $L_{4}=L+L_{1}+L_{2}+L_{3}$. Then, a triangulation $K$ of the covering surface $\Omega$ is so constructed that the projection $T$ of $K$ onto the extended $z$-plane is the triangulation $L_{4} . W e$ see that the triangulation $K$ conforms to the definition in $\S 1.2$. We denote the parts of $T^{-1}(\bar{D})$ and $T^{-1}(L)$ on the upper sheet of $\Omega$ by $\bar{D}$ and $L$ again respectively.

Let $\phi=\chi$ and $\varrho=\tau$ be the differentials on the present $\Omega$ defined in $\S 4.1$, and let $\phi_{h}^{\prime}$ and $\varrho_{h}^{\prime}$ be the finite element approximations of $\phi$ and $\varrho$ respectively in the space $\Lambda^{\prime}\left(K^{\prime}\right)$, where $K^{\prime}$ is the naturalized triangulation associated to the present K.

Let $\Lambda(L)$ be the space of differentials on $\bar{D}$ which are the restrictions of those in $\Lambda(K)$ to $\bar{D}$. Let $\Lambda_{\phi}(L)$ be the subspace of $\Lambda(L)$ which consists of the differentials $\sigma_{h}$ in $\Lambda(L)$ satisfying the conditions:

$$
\begin{array}{lll}
\sigma_{h}=0 & \text { along } & c_{0}=\{z \mid 0 \leqq \operatorname{Im} z \leqq \sqrt{3}, \operatorname{Re} z=0\}, \\
\sigma_{h}=0 & \text { along } & c_{1}=\{z \mid 1 \leqq \operatorname{Re} z \leqq \sqrt{3}, \operatorname{Im} z=0\}
\end{array}
$$

and

$$
\int_{B \cap \bar{D}} \sigma_{h}=\frac{1}{4},
$$

and let $\Lambda_{\phi}^{\prime}\left(L^{\prime}\right)=\left\{\sigma_{h}^{\prime}=F\left(\sigma_{h}\right), \sigma_{h} \in \Lambda_{\phi}(L)\right\}$. Further, let $\Lambda_{e}(L)$ be the subspace of $\Lambda(L)$ which consists of the differentials $\sigma_{h}$ in $\Lambda(L)$ satisfying the conditions:

$$
\begin{array}{lll}
\sigma_{h}=0 & \text { along } & c_{0}^{*}=\{z \mid 0 \leqq \operatorname{Re} z \leqq 1, \operatorname{Im} z=0\}, \\
\sigma_{h}=0 & \text { along } & c_{1}^{*}=\left\{z| | z \mid=\sqrt{3}, 0 \leqq \arg z \leqq \frac{\pi}{2}\right\}
\end{array}
$$

and

$$
\int_{A \cap \bar{D}} \sigma_{h}=-\frac{1}{4},
$$

and let $\Lambda_{\rho}^{\prime}\left(L^{\prime}\right)=\left\{\sigma_{h}^{\prime}=F\left(\sigma_{h}\right), \sigma_{h} \in \Lambda_{\varrho}(L)\right\}$. By $\phi_{h, L}^{\prime}$ and $\varrho_{h, L}^{\prime}$ we denote the
Table 1. Periodicity moduli $p_{1}$ of closed Riemann surface

| Exact value | $p_{1}=\int_{A} * \phi=0.781701$ |  |  |
| :---: | :---: | :---: | :---: |
| Finite element approximations | Original triangulation ( $h=0.213758$ ) |  |  |
|  | Upper bound | $\begin{aligned} & \left\\|\phi_{n}^{\prime}\right\\|^{2}+\varepsilon\left(\phi_{n}^{\prime}\right) \\ & =0.7822184+0.429347 \times 10^{-3} \\ & =0.782613 \quad(0.000912) \end{aligned}$ | $\begin{aligned} & \left\\|\phi_{n}^{\prime}-\phi^{\prime}\right\\| \\ & =3.76256 \times 10^{-3} \end{aligned}$ |
|  | Lower bound | $\begin{aligned} & \frac{1}{\left\\|e_{h}^{\prime}\right\\|^{2}+\varepsilon\left(\varrho_{h}^{\prime}\right)}{ }_{1} \\ &= \frac{1.280878+0.150405 \times 10^{-5}}{} \\ &=0.780714 \quad(-0.000987) \end{aligned}$ | $\begin{aligned} & \left\\|\varrho_{h}^{\prime}-\hat{\varrho}^{\prime}\right\\| \\ & =6.14254 \times 10^{-3} \end{aligned}$ |
|  |  | Normal subdivision ( $h$ |  |
|  | Upper bound | $\begin{aligned} & \left\\|\phi^{\prime}\right\\|^{2}+\varepsilon\left(\phi_{h}^{\prime}\right) \\ & =0.781968+0.107413 \times 10^{-3} \\ & =0.782075 \quad(0.000374) \end{aligned}$ | $\begin{aligned} & \quad\left\\|\phi_{n}^{\prime}-\phi^{\prime}\right\\| \\ & = \\ & =1.12050 \times 10^{-3} \end{aligned}$ |
|  | Lower bound | $\begin{aligned} & \frac{1}{\left\\|\varrho_{h}^{\prime}\right\\|^{2}+\varepsilon\left(\varrho_{h}^{\prime}\right)} \\ & =\frac{1}{1.279506+0.381486 \times 10^{-6}} \\ & =0.781551 \quad(-0.000150) \end{aligned}$ | $\begin{aligned} & \left\\|\varrho_{h}^{\prime}-\varrho^{\prime}\right\\| \\ & =1.83821 \times 10^{-3} \end{aligned}$ |

( ): Deviation from exact value.


Fig. 9
differentials in $\Lambda_{\phi}^{\prime}\left(L^{\prime}\right)$ and $\Lambda_{e}^{\prime}\left(L^{\prime}\right)$ respectively which minimize norms $\left\|\sigma_{h}^{\prime}\right\|_{L^{\prime}}$ in $\Lambda_{\phi}^{\prime}\left(L^{\prime}\right)$ and $\Lambda_{\varrho}^{\prime}\left(L^{\prime}\right)$ respectively. Then, by making use of the symmetricity of $K^{\prime}$, the period and boundary conditions of $\phi_{h}^{\prime}, \varrho_{h}^{\prime}, \phi_{h, L}^{\prime}$ and $\varrho_{h, L}^{\prime}$, and their minimality w.r.t. norm, we can verify that $\phi_{h, L}^{\prime}$ and $\varrho_{h, L}^{\prime}$ are the restrictions of $\phi_{h}^{\prime}$ and $\varrho_{h}^{\prime}$ to $L^{\prime}$ respectively, and $\left\|\phi_{h}^{\prime}\right\|_{K^{\prime}}^{2}=16\left\|\phi_{h, L}^{\prime}\right\|_{L^{\prime}}^{2}$ and $\left\|\varrho_{h}^{\prime}\right\|_{K^{\prime}}^{2}=16\left\|\varrho_{h, L}^{\prime}\right\|_{L^{\prime}}^{2}$. Consequently, to attain our aim it is sufficient to make numerical calculations of $\phi_{h, L}^{\prime}$ and $\varrho_{h, L}^{\prime}$ (cf. Mizumoto and Hara [16], [17] for the calculation method).

We should note that the symmetricity of $\phi$ and $\varrho$ on $\Omega$ has not been used and thus our method does not reject an application to the differentials which do not have symmetricity on $\Omega$.

Table 1 shows the exact value of the periodicity moduli $p_{1}$ which can be calculated by making use of a complete elliptic integral, and the values of our finite element approximations. Furthermore, computation results for the normal subdivision $K^{1}$ (see Fig. 9) of the present $K$ are shown. It can be said that the both of upper and lower bounds of $p_{1}$ are close to the exact value.
4. Numerical example 2 (the case of a compact bordered Riemann surface) Let $\bar{\Omega}$ be a two-sheeted compact bordered covering surface with three branch points $z=-1,1,3$ over the ellipse:

$$
E=\left\{z=x+i y \left\lvert\, \frac{x^{2}}{16}+\frac{y}{15} \leqq 1\right.\right\}
$$

Then $\bar{\Omega}$ is a compact bordered Riemann surface of genus one with one boundary component $C$. A canonical homology basis $\{A, B\}$ of $\bar{\Omega}$ is chosen as in Fig. 10. We aim to obtain good upper and lower approximate values of the


Fig. 10

Fig. 11
periodicity moduli $p_{1}$ and $p_{2}$ of $\bar{\Omega}$ with respect to $A$ and $B$ respectively.
First, we construct a triangulation of the upper half ellipse $\bar{D}=E \bigcap\{z \mid \operatorname{Im} z$ $\geqq 0\}$ as in Fig. 11. The closed regions $G_{2}, G_{3}, G_{4}$ and $G_{5}$ are mapped onto the regions $G_{2}^{*}, \mathrm{G}_{3}^{*}, G_{4}^{*}$ and $G_{5}^{*}$ resp. by the local parameters $\zeta=\varphi_{2}(z)=a \sqrt{z+1}, \zeta$ $=\varphi_{3}(z)=a \sqrt{z-1}, \zeta=\varphi_{4}(z)=b \sqrt{z-3}$ and $w=\varphi_{5}(z)=\cosh ^{-1} z \quad\left(a=2 / 5^{1 / 4}\right.$ and $b=2 / 85^{1 / 4}$ ) respectively, where $a$ and $b$ are so determined that $|d \zeta / d z|$ are equal to $|d w / d z|$ at $z=z_{0}+i \quad\left(z_{0}=-1,1\right.$ or 3$)$. We construct ordinary triangulations $K_{2}^{*}, K_{3}^{*}, K_{4}^{*}$ and $K_{5}^{*}$ of $G_{2}^{*}, G_{3}^{*}, G_{4}^{*}$ and $G_{5}^{*}$ as in Fig. 11 respectively. By $K_{2}, K_{3}, K_{4}$ and $K_{5}$ we denote the image triangulations of $K_{2}^{*}$, $K_{3}^{*}, K_{4}^{*}$ and $K_{5}^{*}$ by the mappings $\varphi_{2}^{-1}, \varphi_{3}^{-1}, \varphi_{4}^{-1}$ and $\varphi_{5}^{-1}$ respectively. The triangulation $K_{1}$ of the region $G_{1}=\overline{\Omega-\left(G_{2} \cup G_{3} \cup G_{4} \cup G_{5}\right)}$ in Fig. 11 is so constructed that each 2-simplex $s$ of $K_{1}$ is natural, minor or major according as $|s| \cap\left|K_{2}+K_{3}+K_{4}+K_{5}\right|=\varnothing,|s| \cap\left|K_{2}+K_{3}+K_{4}\right| \neq \varnothing$, or $|s| \cap\left|K_{5}\right| \neq \varnothing$, with the convention as in the previous section, and the local parameter of $K_{1}$ is $\varphi_{1}(z) \equiv z$.

A triangulation $L_{1}$ of the lower half ellipse $\bar{D}_{1}=E \cap\{z \mid \operatorname{Im} z \leqq 0\}$ is defined by the reflection of the triangulation $L \equiv K_{1}+K_{2}+K_{3}+K_{4}+K_{5}$ with respect to the real axis and a triangulation $L_{2}$ of $E$ is defined by $L_{2}=L+L_{1}$. Then, a triangulation $K$ of the covering surface $\bar{\Omega}$ is so constructed that the projection $T$ of $K$ onto the $z$-plane is the triangulation $L_{2}$. We see that the triangulation $K$ conforms to the definition in §1.2. We denote the parts of $T^{-1}(\bar{D})$ and $T^{-1}(L)$ on the upper sheet of $\bar{\Omega}$ by $\bar{D}$ and $L$ again respectively.

Let $\phi, \varrho, \chi$ and $\tau$ be the differentials on the present $\bar{\Omega}$ defined in $\S 4.1$, and let $\phi_{h}^{\prime}, \varrho_{h}^{\prime}, \chi_{h}^{\prime}$ and $\tau_{h}^{\prime}$ be the finite element approximations of $\phi, \varrho, \chi$ and $\tau$ respectively in the space $\Lambda^{\prime}\left(K^{\prime}\right)$, where $K^{\prime}$ is the naturalized triangulation associated to the present $K$.

Let $\Lambda(L)$ be the space of differentials on $\bar{D}$ which are the restrictions of those in $\Lambda(K)$ to $\bar{D}$. Let $\Lambda_{\phi}(L), \Lambda_{\ell}(L), \Lambda_{x}(L)$ and $\Lambda_{\tau}(L)$ be the subspaces of $\Lambda(L)$ which consist of the differentials $\sigma_{h 1}, \sigma_{h 2}, \sigma_{h 3}$ and $\sigma_{h 4}$ in $\Lambda(L)$ respectively satisfying the conditions:

$$
\begin{array}{lll}
\sigma_{h 1}=\sigma_{h 3}=0 & \text { along } & c_{0}=\{z \mid \quad 3 \leqq \operatorname{Re} z \leqq 4, \operatorname{Im} z=0\}, \\
\sigma_{h 1}=\sigma_{h 3}=0 & \text { along } & c_{1}=\{z \mid-1 \leqq \operatorname{Re} z \leqq 1, \operatorname{Im} z=0\}, \\
\sigma_{h 2}=\sigma_{h 4}=0 & \text { along } & c_{0}^{*}=\{z \mid \quad 1 \leqq \operatorname{Re} z \leqq 3, \operatorname{Im} z=0\}, \\
\sigma_{h 2}=\sigma_{h 4}=0 & \text { along } & c_{1}^{*}=\{z \mid-4 \leqq \operatorname{Re} z \leqq-1, \operatorname{Im} z=0\}, \\
\sigma_{h 1}=\sigma_{h 2}=0 & \text { along } \quad c=\left\{z=x+i y \mid x^{2} / 16+y^{2} / 15=1, y \leqq 0\right\}, \\
& \int_{B \cap \bar{D}} \sigma_{h 1}=\int_{B \cap \bar{D}} \sigma_{h 3}=\frac{1}{2}
\end{array}
$$

and

$$
\int_{A \cap \bar{D}} \sigma_{h 2}=\int_{A \cap \bar{D}} \sigma_{h 4}=-\frac{1}{2}
$$

Further, let $\Lambda_{\phi}^{\prime}\left(L^{\prime}\right)=\left\{\sigma_{h 1}^{\prime}\right\}, \Lambda_{e}^{\prime}\left(L^{\prime}\right)=\left\{\sigma_{h 2}^{\prime}\right\}, \Lambda_{x}^{\prime}\left(L^{\prime}\right)=\left\{\sigma_{h 3}^{\prime}\right\}$ and $\Lambda_{\tau}^{\prime}\left(L^{\prime}\right)=\left\{\sigma_{h 4}^{\prime}\right\}$, where $\sigma_{h j}^{\prime}=F\left(\sigma_{h j}\right)(j=1,2,3,4)$. By $\phi_{h, L}^{\prime}, \varrho_{h, L}^{\prime}, \chi_{h, L}^{\prime}$ and $\tau_{h, L}^{\prime}$ we denote the differentials of $\Lambda_{\phi}^{\prime}\left(L^{\prime}\right), \Lambda_{\theta}^{\prime}\left(L^{\prime}\right), \Lambda_{x}^{\prime}\left(L^{\prime}\right)$ and $\Lambda_{\tau}^{\prime}\left(L^{\prime}\right)$ respectively which minimize norms in $\Lambda_{\phi}^{\prime}\left(L^{\prime}\right), \Lambda_{\rho}^{\prime}\left(L^{\prime}\right), \Lambda_{x}^{\prime}\left(L^{\prime}\right)$ and $\Lambda_{\tau}^{\prime}\left(L^{\prime}\right)$ respectively. Then, by making use of the symmetricity of $K^{\prime}$, the period and boundary conditions of $\phi_{h}^{\prime}, \varrho_{h}^{\prime}, \chi_{h}^{\prime}$, $\tau_{h}^{\prime}$, $\phi_{h, L}^{\prime}, \varrho_{h, L}^{\prime}, \chi_{h, L}^{\prime}$ and $\tau_{h, L}^{\prime}$, and their minimality w.r.t. norm, we can verify that $\phi_{h, L}^{\prime}$, $\varrho_{h, L}^{\prime}, \chi_{h, L}^{\prime}$ and $\tau_{h, L}^{\prime}$ are the restrictions of $\phi_{h}^{\prime}, \varrho_{h}^{\prime}, \chi_{h}^{\prime}$ and $\tau_{h}^{\prime}$ to $L^{\prime}$ respectively, and $\left\|\phi_{h}^{\prime}\right\|_{K^{\prime}}^{2}=4\left\|\phi_{h, L}^{\prime}\right\|_{L^{\prime}}^{2}, \quad\left\|\varrho_{h}^{\prime}\right\|_{K^{\prime}}^{2}=4\left\|\varrho_{h, L}^{\prime}\right\|_{L^{\prime}}^{2}, \quad\left\|\chi_{h}^{\prime}\right\|_{K^{\prime}}^{2}=4\left\|\chi_{h, L}^{\prime}\right\|_{L^{\prime}}^{2} \quad$ and $\left\|\tau_{h}^{\prime}\right\|_{K^{\prime}}^{2}=4\left\|\tau_{h, L}^{\prime}\right\|_{L^{\prime}}^{2}$. Consequently, to attain our aim it is sufficient to make numerical calculations of $\phi_{h, L}^{\prime}, \varrho_{h, L}^{\prime}, \chi_{h, L}^{\prime}$ and $\tau_{h, L}^{\prime}$.

The exact values of the periodicity moduli $p_{1}$ and $p_{2}$ can be calculated by the following procedure.

Let $C_{0}$ and $C_{1}$ be the boundary parts of the upper half ellipse domain $D$ defined by

$$
C_{0}=\{z \mid 3 \leqq \operatorname{Re} z \leqq 4, \operatorname{Im} z=0\} \cup\left\{z=x+i y \left\lvert\, \frac{x^{2}}{16}+\frac{y^{2}}{15}=1\right., y \geqq 0\right\}
$$

and

$$
C_{1}=\{z \mid-1 \leqq \operatorname{Re} z \leqq 1, \operatorname{Im} z=0\} .
$$

Let $\Delta$ be the rectangular domain

$$
\Delta=\{W \mid 0<\operatorname{Re} W<1,0<\operatorname{Im} W<\tau\}
$$

and let $\Gamma_{0}$ and $\Gamma_{1}$ be the boundary parts of $\Delta$ defined by

$$
\Gamma_{0}=\{W \mid 0 \leqq \operatorname{Im} W \leqq \tau, \operatorname{Re} W=0\}
$$

and

$$
\Gamma_{1}=\{W \mid 0 \leqq \operatorname{Im} W \leqq \tau, \operatorname{Re} W=1\}
$$

If $D$ is conformally mapped onto $\Delta$ so that $C_{0}$ and $C_{1}$ are mapped onto $\Gamma_{0}$ and $\Gamma_{1}$ respectively, then the periodicity moduli $p_{1}$ is equal to $\tau$. The conformal map $W=f(z): D \rightarrow \Delta$ is constructed by the composition of the following mappings:

$$
\begin{equation*}
w=\frac{2}{\cosh ^{-1} 4} \cdot \cosh ^{-1} z-1 \tag{i}
\end{equation*}
$$

(ii) $\zeta=\operatorname{sn}(K(k) \cdot w)$, where $\frac{K^{\prime}(k)}{K(k)}=\frac{2 \pi}{\cosh ^{-1} 4}$;
(iii) $\frac{Z-Z_{1}}{Z-Z_{2}} \cdot \frac{Z_{3}-Z_{2}}{Z_{3}-Z_{1}}=\frac{\zeta-\zeta_{1}}{\zeta-\zeta_{2}} \cdot \frac{\zeta_{3}-\zeta_{2}}{\zeta_{3}-\zeta_{1}}$,
where $\zeta_{j}=\operatorname{sn}\left(K(k) \cdot w_{j}\right)(j=1,2,3,4)$ with $w_{1}=-1+i\left(2 \pi / \cosh ^{-1} 4\right), w_{2}=-1$, $w_{3}=2 \cosh ^{-1} 3 / \cosh ^{-1} 4-1, w_{4}=1+i\left(2 \pi / \cosh ^{-1} 4\right)$, and $Z_{1}=-1 / \kappa, Z_{2}=-1$, $Z_{3}=1, Z_{4}=1 / \kappa$ with $\kappa=(\sqrt{1 / c}-\sqrt{1 / c-1})^{2}, \quad c=\left(\left(\zeta_{4}-\zeta_{1}\right) /\left(\zeta_{4}-\zeta_{2}\right)\right) \cdot\left(\left(\zeta_{3}\right.\right.$ $\left.\left.-\zeta_{2}\right) /\left(\zeta_{3}-\zeta_{1}\right)\right)$;
(iv) $W=-\frac{1}{2}\left(\frac{1}{K(\kappa)} \int_{0}^{Z} \frac{d Z}{\sqrt{\left(1-Z^{2}\right)\left(1-\kappa^{2} Z^{2}\right)}}-\left(1+i \frac{K^{\prime}(\kappa)}{K(\kappa)}\right)\right)$.

Then we see that

$$
p_{1}=\tau=\frac{K^{\prime}(\kappa)}{2 K(\kappa)}
$$

Next, let $C_{0}^{\prime}$ and $C_{1}^{\prime}$ be the boundary parts of $D$ given by

$$
C_{0}^{\prime}=\{z \mid 1 \leqq \operatorname{Re} z \leqq 3, \operatorname{Im} z=0\}
$$

and
Table 2. Periodicity moduli $p_{1}$ of compact bordered Riemann surface

| Exact value | $p_{1}=\int_{A} * \phi=1.539330$ |  |  |
| :---: | :---: | :---: | :---: |
| Finite element approximations | Original triangulation ( $h=0.138840$ ) |  |  |
|  | Upper bound | $\begin{aligned} & \left\\|\phi_{n}^{\prime}\right\\|^{2}+\varepsilon\left(\phi_{h}^{\prime}\right) \\ = & 1.540588+0.572262 \times 10^{-4} \\ = & 1.540645 \quad(0.00132) \end{aligned}$ | $\begin{aligned} & \left\\|\phi_{h}^{\prime}-\hat{\phi}^{\prime}\right\\| \\ & =1.15335 \times 10^{-2} \end{aligned}$ |
|  | Lower bound | $\begin{aligned} & \frac{1}{\left\\|\tau_{h}^{\prime}\right\\|^{2}+\varepsilon\left(\tau_{h}^{\prime}\right)} \\ & =\frac{1}{0.649700+0.225117 \times 10^{-3}} \\ & =1.538639 \quad(-0.00069) \end{aligned}$ | $\begin{aligned} & \left\\|\tau_{h}^{\prime}-\hat{\tau}^{\prime}\right\\| \\ & =3.74131 \times 10^{-3} \end{aligned}$ |
|  |  | Normal subdivision ( $h$ |  |
|  | Upper bound | $\begin{aligned} & \left\\|\phi_{h}^{\prime}\right\\|^{2}+\varepsilon\left(\phi_{h}^{\prime}\right) \\ & =1.539652+0.142916 \times 10^{-4} \\ & =1.539666 \quad(0.00034) \end{aligned}$ | $\begin{aligned} & \left\\|\phi_{h}^{\prime}-\hat{\phi}^{\prime}\right\\| \\ & =5.89447 \times 10^{-3} \end{aligned}$ |
|  | Lower bound | $\begin{aligned} & \frac{1}{\left\\|\tau_{h}^{\prime}\right\\|^{2}+\varepsilon\left(\tau_{h}^{\prime}\right)} \\ &= \frac{1}{0.649652+0.558093 \times 10^{-4}} \\ &=1.539153 \quad(-0.00018) \end{aligned}$ | $\begin{aligned} & \left\\|\tau_{h}^{\prime}-\hat{\tau}^{\prime}\right\\| \\ & =1.09209 \times 10^{-3} \end{aligned}$ |

( ): Deviation from exact value.

Table 3. Periodicity moduli $p_{2}$ of compact bordered Riemann surface

| Exact value | $p_{2}=\int_{B} * \varrho=1.839350$ |  |  |
| :---: | :---: | :---: | :---: |
| Finite element approximations | Original triangulation ( $h=0.138840$ ) |  |  |
|  | Upper bound | $\begin{aligned} & \left\\|\varrho_{\varrho^{\prime}}\right\\|^{2}+\varepsilon\left(\varrho_{h}^{\prime}\right) \\ & =1.841976+0.351532 \times 10^{-3} \\ & =1.842328 \quad(0.00298) \end{aligned}$ | $\begin{aligned} & \left\\|\varrho_{h}^{\prime}-\hat{\varrho}^{\prime}\right\\| \\ & =7.65797 \times 10^{-3} \end{aligned}$ |
|  | Lower bound | $\begin{aligned} & \frac{1}{\left\\|\chi_{h}^{\prime}\right\\|^{2}+\varepsilon\left(\chi_{h}^{\prime}\right)} \\ & =\frac{1}{0.5445888+0.145580 \times 10^{-3}} \\ & =1.835760 \quad(-0.00359) \end{aligned}$ | $\begin{aligned} & \left\\|\chi_{h}^{\prime}-\hat{\chi}^{\prime}\right\\| \\ & =5.22574 \times 10^{-3} \end{aligned}$ |
|  | Normal subdivision ( $h=0.069420$ ) |  |  |
|  | Upper bound | $\begin{aligned} & \left\\|\varrho_{\\|}^{\prime}\right\\|^{2}+\varepsilon\left(\varrho_{h}^{\prime}\right) \\ = & 1.840016+0.875764 \times 10^{-4} \\ = & 1.840104 \quad(0.00075) \end{aligned}$ | $\begin{aligned} & \left\\|\varrho_{h}^{\prime}-\hat{\varrho^{\prime} \\|}\right\\| \\ & =2.28613 \times 10^{-3} \end{aligned}$ |
|  | Lower bound | $\begin{aligned} & \frac{1}{\left\\|\chi_{h}^{\prime}\right\\|^{2}+\varepsilon\left(\chi_{h}^{\prime}\right)} \\ & =\frac{1}{0.543904+0.361871 \times 10^{-4}} \\ & =1.838437 \quad(-0.00091) \end{aligned}$ | $\begin{aligned} & \left\\|\chi_{h}^{\prime}-\hat{\chi}^{\prime}\right\\| \\ & =1.73332 \times 10^{-3} \end{aligned}$ |

( ): Deviation from exact value.

$$
C_{1}^{\prime}=\left\{z=x+i y \left\lvert\, \frac{x^{2}}{16}+\frac{y^{2}}{15}=1\right., y \geqq 0\right\} \cup\{z \mid-4 \leqq \operatorname{Re} z \leqq-1, \operatorname{Im} z=0\}
$$

Let $\Delta, \Gamma_{0}$ and $\Gamma_{1}$ be as above. If the domain $D$ is conformally mapped onto the domain $\Delta$ so that $C_{0}^{\prime}$ and $C_{1}^{\prime}$ are mapped onto $\Gamma_{0}$ and $\Gamma_{1}$ respectively, then the periodicity moduli $p_{2}$ is equal to $\tau$. The conformal map $W=f(p): D \rightarrow \Delta$ is constructed similarly to the case of periodicity moduli $p_{1}$.

Tables 2 and 3 show the exact values of the periodicity moduli $p_{1}$ and $p_{2}$, and the values of our finite element approximations. Furthermore, computation results for the normal subdivision $K^{1}$ of the present $K$ are shown. It can be said that the both of upper and lower bounds of $p_{1}$ and $p_{2}$ are close to the exact values.

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[^1]:    1) We shall use the common notations (, ) and |||| for both inner products and both norms of differentials of the spaces $\Lambda$ and $\Lambda^{\prime}$.
[^2]:    1) It is sufficient for our purpose that $\omega^{\circ} \varphi_{j}^{-1}$ is of the class $C^{1}$ on $\varphi_{j}\left(U_{j} \cap \bar{\Omega}\right)$ and hence we can weaken the assumption (i) for $\Theta$.
