A NOTE ON THE DISTRIBUTION OF HITTING TIMES¹

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We show that "hitting times" associated with stationary processes have an absolutely continuous distribution, except, possibly, for an atom at the origin. The density is then identified in several special cases.

- **0.** We show that "hitting times" associated with stationary processes have an absolutely continuous distribution except, possibly, for an atom at the origin. Several examples are given.
- 1. Let $R(R_+)$ denote the real line (half-line $[0, \infty)$) with Borel sets $\mathscr{B}(\mathscr{B}_+)$. We use m for Lebesgue measure. Let θ_t , $t \in R$, be a group of measurable, measure-preserving transformations of a probability space (Ω, \mathscr{F}, P) . In addition, assume $(t, \omega) \to \theta_t(\omega)$ is $\mathscr{B} \times \mathscr{F}$ measurable. We will consider Borel sets $M(\omega)$ indexed by Ω and satisfying
 - (i) $M(\theta_t \omega) = M(\omega) t$ for all $t \in R$, $\omega \in \Omega$ and
 - (ii) $\sup M(\omega) = +\infty$ a.s.

The hitting time of $M(\omega)$ is $\tau(\omega) = \inf\{t > 0 : t \in M(\omega)\}$. Note that $\tau(\omega) > t$ implies $\tau(\omega) = t + \tau(\theta_t \omega)$, the so-called "terminal" property of τ . Assume τ is measurable.

THEOREM. The distribution of τ is absolutely continuous except, possibly, for an atom at 0.

This theorem, as well as Section 2, is a direct consequence of the following equation:

The proof of (*) is elementary. For each $\omega \in \Omega$, the paths $\tau(\theta_t \omega)$ are right-continuous with left-hand limits on R_+ , and have "saw-tooth" appearance, with possibly infinitely many "teeth" in a finite interval. With $\omega \in \Omega$, b>0 and T>0 fixed, it is fairly apparent that $\{t\colon \tau(\theta_t\omega)\geq b\}$ meets [0,T] in at most finitely many disjoint (possibly degenerate) closed intervals, say J_1,J_2,\ldots,J_n , such that

(1)
$$m(J_k) = \int_b^\infty \nu(k, x, \omega) dx$$
, $k = 1, 2, \dots, n$
where $\nu(k, x, \omega)$ is the number of solutions (necessarily finite) of $\tau(\theta_t \omega) = x$ for

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 $t \in J_k$. Summing (1) over $k = 1, 2, \dots, n$ we find that

$$\int_0^T I_{(b,\infty)}(\tau(\theta_t \omega)) dt = \int_b^\infty \nu_T(x, \omega) dx.$$

Incorporating the case b=0 we obtain (*) for any $\Gamma \in \mathscr{B}_+$ of the form $\Gamma = [b, \infty)$, hence for any $\Gamma \in \mathscr{B}_+$. Taking expectations with T=1 in (*) gives the theorem:

$$P(\tau \in \Gamma) = \int_{\Gamma} E\nu_{1}(x) dx + \delta_{0}(\Gamma)P(\tau = 0), \qquad \Gamma \in \mathscr{B}_{+}.$$

(Notice that $P(\tau = 0) = Em(M \cap [0, 1])$ if $M(\omega)$ is right-closed a.s.)

2. By (*), the family of additive functionals

$$\beta_t(x, \omega) = (P(\tau = 0))^{-1} \int_0^t I_{\{0\}}(\tau(\theta_s \omega)) ds, \qquad x = 0$$

= $(E\nu_1(x))^{-1}\nu_t(x, \omega), \qquad x > 0$

is an "occupation-time density" for the process $\tau \circ \theta_t$. (By an additive functional we mean a right-continuous, non-decreasing process $\alpha_t(\omega)$ for which $\alpha_0 \equiv 0$ and $\alpha_{t+s} - \alpha_t = \alpha_s \circ \theta_t$ a.s. for each s, t.) Specifically

$$\int_{\Gamma} \beta_t(x, \, \omega) P(\tau \in dx) = \int_0^t I_{\Gamma}(\tau(\theta_s \, \omega)) \, ds$$

for all $t \in R_+$, $\omega \in \Omega$, $\Gamma \in \mathcal{B}_+$. A straightforward calculation (which we will omit) shows that the family of Palm measures (of $\beta_t(x)$)

$$P_{\beta(x)}(A) = E \int_0^1 I_A(\theta_t \omega) d\beta_t(x, \omega), \qquad x \in R_+$$

is a regular version of the family of conditional probabilities $P^x(\cdot) = P(\cdot \mid \tau = x)$, $x \in R_+$. Indeed, if \mathscr{F} is separable, $\hat{P}_{\beta(x)} = P^x$ a.e. (m).

3. In general, the density $E_{\nu_1}(x)$ of the positive part of τ does not admit a smooth version. For instance, let $h(x) \geq 0$, x > 0, be any non-increasing, integrable function and let $X_t(\omega)$ be the semilinear Markov process with characteristic $\{\alpha, h(x)\}$, $\alpha \geq 0$. When $h(0+) = \infty$, X_t has a unique, stationary initial distribution (see Horowitz [1]), under which the distribution of the first passage to 0 is

$$(\alpha + \int_0^\infty h(x) dx)^{-1} (\alpha \delta_0(dx) + h(x) dx)$$
.

4. When $M(\omega)$ is discrete, say countable and clustering only at $\pm \infty$, the additive functional $n_t(\omega) \equiv \lim_{x \downarrow 0} \nu_t(x, \omega)$ counts the number of points in $M(\omega) \cap (0, t]$ and is a stationary point process. When $En_1 = \alpha < \infty$, the (normalized) Palm measure \hat{P} of $n_t(\omega)$ is a probability and $\hat{E}\tau = 1/\alpha$. A now well-known inversion formula (see e.g. Ryll-Nardzewski [2]) is

(2)
$$E\xi = \alpha \hat{E} \int_0^{\tau(\omega)} \xi(\theta_s \omega) ds$$
, for any random variable ξ .

Denote by $\phi(\lambda)$ and $\hat{\phi}(\lambda)$ the Laplace transforms of τ under P and \hat{P} respectively. From (2), $\phi(\lambda) = \alpha(1 - \hat{\phi}(\lambda))/\lambda$ and consequently $\alpha(1 - \hat{P}(\tau \le t))$ is the density of τ under P.

5. As a final example, let N_t be a Poisson process with unit intensity and

 $X_t = N_{t+1} - N_t$. This is a strictly stationary (non-Markov) process; we will compute the distribution of $\tau(\omega) = \inf\{t > 0 : X_t(\omega) = 0\}$.

Let T_1, T_2, \cdots be the times between the jumps of N_t . Clearly $P(\tau = 0) = P(T_1 > 1) = e^{-1}$ and for t > 0,

$$P(\tau \leq t) = P(T_1 > 1) + \sum_{n=1}^{\infty} P(T_1 \leq 1, \dots, T_n \leq 1, T_{n+1} > 1, \sum_{n=1}^{\infty} T_n \leq t)$$

= $e^{-1}(1 + \sum_{n=1}^{\infty} u_n(t))$

where $u_n(t) = \int_{D_n(t)} \cdots \int_{D_n(t)} (\exp(-\sum_{i=1}^n x_i) \prod_{i=1}^n dx_i)$ and $D_n(t)$ is the region $\{(x_1, x_2, \dots, x_n) : 0 \le x_i \le 1, \sum_{i=1}^n x_i \le t\}$. Observe that $u_1(t) = 1 - \exp[-\min(1, t)]$ and that

$$u_{n+1}(t) = \int_0^{\min(1,t)} e^{-x} u_n(t-x) dx$$
, $n = 1, 2, \dots$

Summing over $n = 1, 2, \cdots$ results in the renewal equation

(3)
$$u(t) = u_1(t) + \int_0^t u(t-x) du_1(x), \qquad u(t) = \sum_{n=1}^\infty u_n(t).$$

Evidently, u(t) is continuously differentiable for $t \neq 1$; thus $P(\tau \leq t)$ is absolutely continuous on $(0, \infty)$. In fact, an easy computation with (3) leads to

$$Ee^{-\lambda \tau} = \frac{e^{-1}(1+\lambda)}{\lambda + e^{-1-\lambda}}, \qquad \lambda \geq 0.$$

(For intensity θ we get $e^{-\theta}(\theta + \lambda)/(\lambda + \theta e^{-\theta - \lambda})$.)

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