

Vol. 15 (2010), Paper no. 44, pages 1394–1414.

Journal URL

<http://www.math.washington.edu/~ejpecp/>

Universal Behavior of Connectivity Properties in Fractal Percolation Models

Erik I. Broman * Federico Camia †

Abstract

Partially motivated by the desire to better understand the connectivity phase transition in fractal percolation, we introduce and study a class of continuum fractal percolation models in dimension $d \geq 2$. These include a scale invariant version of the classical (Poisson) Boolean model of stochastic geometry and (for $d = 2$) the Brownian loop soup introduced by Lawler and Werner. The models lead to random fractal sets whose connectivity properties depend on a parameter λ . In this paper we mainly study the transition between a phase where the random fractal sets are totally disconnected and a phase where they contain connected components larger than one point. In particular, we show that there are connected components larger than one point at the unique value of λ that separates the two phases (called the critical point). We prove that such a behavior occurs also in Mandelbrot's fractal percolation in all dimensions $d \geq 2$. Our results show that it is a generic feature, independent of the dimension or the precise definition of the model, and is essentially a consequence of scale invariance alone. Furthermore, for $d = 2$ we prove that the presence of connected components larger than one point implies the presence of a unique, unbounded, connected component .

Key words: random fractals, fractal percolation, continuum percolation, Mandelbrot percolation, phase transition, crossing probability, discontinuity, Brownian loop soup, Poisson Boolean Model.

AMS 2000 Subject Classification: Primary 60D05, 28A80, 60K35.

Submitted to EJP on March 10, 2010, final version accepted August 31, 2010.

*Department of Mathematics, Chalmers University of Technology, S-412 96 Göteborg, Sweden. E-mail: broman@chalmers.se

†Department of Mathematics, Vrije Universiteit Amsterdam, De Boelelaan 1081a, 1081 HV Amsterdam, The Netherlands. E-mail: fede@few.vu.nl

1 Introduction

Many deterministic constructions generating fractal sets have random analogues that produce *random fractals* which do not have the self-similarity of their non-random counterpart, but are *statistically self-similar* in the sense that enlargements of small parts have the same statistical distribution as the whole set.

Random fractals can have complex topological structure, for example they can be highly multiply connected, and can exhibit *connectivity phase transitions*, corresponding to sudden changes of topological structure as a continuously varying parameter goes through a critical value.

In this paper, we introduce and study a natural class of random fractals that exhibit, in dimension $d \geq 2$, such a connectivity phase transition: when a parameter increases continuously through a critical value, the connectivity suddenly breaks down and the random fractals become totally disconnected with probability one. (We remind the reader that a set is called totally disconnected if it contains no connected component larger than one point.) The fractals we study are defined as the complement of the union of sets generated by a Poisson point process of intensity λ times a *scale invariant* measure on a space of subsets of \mathbb{R}^d (see Section 2).

Examples of such random fractals include a scale invariant version of the classical (Poisson) Boolean model of stochastic geometry (see [29; 32] and [24] for a multiscale version of the model), the Brownian loop soup [18] (both will be discussed in some more detail in the next section), and the models studied in [27]. The scale invariant (Poisson) Boolean model is a natural model for a porous medium with cavities on many different scales (but it has also been used as a simplified model in cosmology — see [14]). It is obtained via a Poisson point process in $(d+1)$ -dimensional space, where the first d coordinates of the points give the locations of the centers of d -dimensional balls whose radii are given by the last coordinate. The distribution of the radii r has density $(1/r)^{d+1}$, which ensures scale invariance. There is no reason, except simplicity, for using balls, and the model can be naturally generalized by associating random shapes to the points of the Poisson process. Another natural way to generalize the model is obtained by considering a Poisson point process directly in the space of “shapes,” i.e., subsets of \mathbb{R}^d . In dimension $d = 2$, this is how the Brownian loop soup is defined, with the distribution of the random shapes given by the distribution of Brownian loops. In this paper we consider this type of models with general scale invariant distributions on shapes (see Definition 2.1). The reason is that we want to study what features in the behavior of fractal percolation models are a consequence of scale invariance alone.

Our main result consists in showing that, when the intensity λ of the Poisson process is at its critical value, the random fractals are in the connected phase in the sense that they contain connected components larger than one point with probability one (see Theorem 2.4). This is reminiscent of the nature of the phase transition in Mandelbrot’s fractal percolation [21; 22], which is discussed in more detail in Section 4.

Our proof of Theorem 2.4 is interesting in that it shows that the nature of the connectivity phase transition described in the theorem is essentially a consequence of scale invariance alone, and in particular does not depend on the dimension d . The same proof applies to other models as well, including Mandelbrot’s fractal percolation. (We note that the proofs of Theorems 5.3 and 5.4 of [10] also show the importance of scale invariance, but are very two-dimensional). Along the way, we prove a discontinuity result for the probability that a random fractal contains a connected component crossing a “shell-like domain,” which is interesting in its own right (see Corollary 2.6).

The main result is stated in Section 2 while Sections 3 and 4 contain additional two-dimensional results and results concerning Mandelbrot's fractal percolation model respectively.

1.1 Two Motivating Examples

Two prototypical examples of the type of models that we consider in this paper are a fully scale invariant version of the multiscale Poisson Boolean model studied in Chapter 8 of [24] and in [25; 26] (see [2] for recent results on that model, whose precise definition is given below) and, in two dimensions, the Brownian loop soup of Lawler and Werner [18].

The Brownian loop soup with density $\lambda > 0$ is a realization of a Poisson point process with intensity λ times the *Brownian loop measure*, where the latter is essentially the only measure on loops that is conformally invariant (see [18] and [17] for precise definitions). A sample of the Brownian loop soup is a countable family of unrooted Brownian loops in \mathbb{R}^2 (there is no non-intersection condition or other interaction between the loops). The Brownian loop measure can also be considered as a measure on hulls (i.e., compact connected sets $K \subset \mathbb{R}^2$ such that $\mathbb{R}^2 \setminus K$ is connected) by filling in the bounded loops.

The scale invariant Boolean model in d dimensions is a Poisson point process on $\mathbb{R}^d \times (0, \infty)$ with intensity $\lambda r^{-(d+1)} dr dx$, where $\lambda \in (0, \infty)$, dr is Lebesgue measure on \mathbb{R}^+ and dx is the d -dimensional Lebesgue measure. Each realization \mathcal{P} of the point process gives rise to a collection of balls in \mathbb{R}^d in the following way. For each point $\xi \in \mathcal{P}$ there is a corresponding ball $b(\xi)$. The projection on \mathbb{R}^d of ξ gives the position of the center of the ball and the radius of the ball is given by the value of the last coordinate of ξ .

Since we want to show the analogy between the two models, and later generalize them, we give an alternative description of the scale invariant Boolean model. One can obtain the random collection of balls described above as a realization of a Poisson point process with intensity $\lambda \mu^{Bool}$, where μ^{Bool} is the measure defined by $\mu^{Bool}(\tilde{E}) = \int_D \int_a^b r^{-(d+1)} dr dx$, for all sets \tilde{E} that are collections of balls of radius $r \in (a, b)$ with center in an open subset D of \mathbb{R}^d . (Denoting by $\tilde{\mathcal{E}}$ the collection of sets \tilde{E} used in the definition of μ^{Bool} , it is easy to see that $\tilde{\mathcal{E}}$ is closed under pairwise intersections. Therefore, the probabilities of events in $\tilde{\mathcal{E}}$ determine μ^{Bool} uniquely as a measure on the σ -algebra $\sigma(\tilde{\mathcal{E}})$. This choice of σ -algebra is only an example, later we will work with different measurable sets.) Here, μ^{Bool} plays the same role as the Brownian loop measure in the definition of the Brownian loop soup. Note that μ^{Bool} is scale invariant in the following sense. Let \tilde{E}' denote the collection of balls with center in sD and radius $r \in (sa, sb)$, for some scale factor s . Then $\mu^{Bool}(\tilde{E}') = \int_{sD} \int_{sa}^{sb} r^{-(d+1)} dr dx = \int_D \int_a^b (sr)^{-(d+1)} s dr s^d dx = \mu(\tilde{E})$.

We are interested in fractal sets obtained by considering the complement of the union of random sets like those produced by the scale invariant Boolean model or the Brownian loop soup, possibly with a cutoff on the maximal diameter of the random sets. Fractals have frequently been used to model physical systems, such as porous media, and in that context the presence of a cutoff is a very natural assumption. Furthermore, it will be easy to see from the definitions that without a cutoff or some other restriction, the complement of the union of the random sets is a.s. empty. An alternative possibility is to consider the restriction of the scale invariant Boolean model or the Brownian loop soup to a bounded domain D . By this we mean that one keeps only those balls or Brownian loops that are contained in D , which automatically provides a cutoff on the size of the retained sets. This approach is particularly natural in the Brownian loop soup context, since then, when λ is below a

critical value, the boundaries of clusters of filled Brownian loops form a realization of a Conformal Loop Ensemble (see [33; 34; 30] and [31] for the definition and properties of Conformal Loop Ensembles).

All proofs are contained in Section 5.

2 Definitions and Main Results

We first remind the reader of the definition of Poisson point process. Let (M, \mathcal{M}, μ) be a measure space, with M a topological space, \mathcal{M} the Borel σ -algebra, and μ a σ -finite measure. A *Poisson point process with intensity measure μ* is a collection of random variables $\{N(E) : E \in \mathcal{M}, \mu(E) < \infty\}$ satisfying the following properties:

- With probability 1, $E \mapsto N(E)$ is a counting measure (i.e., it takes only nonnegative integer values).
- For fixed E , $N(E)$ is a Poisson random variable with mean $\mu(E)$.
- If E_1, E_2, \dots, E_n are mutually disjoint, then $N(E_1), N(E_2), \dots, N(E_n)$ are independent.

The random set of points $\mathcal{P} = \{\xi \in M : N(\{\xi\}) = 1\}$ is called a *Poisson realization* of the measure μ . In the rest of the paper, $d \geq 2$, and D will always denote a bounded, open subset of \mathbb{R}^d , which will be called a *domain*, and \bar{D} will denote the closure of D . If K is a subset of \mathbb{R}^d , we let $sK = \{x \in \mathbb{R}^d : x/s \in K\}$. Here M will be the set of connected, compact subsets of \mathbb{R}^d with nonempty interior. For our purposes, we need not specify the topology, but we require that the Borel σ -algebra contains all sets of the form $E(B; a, b) = \{K \in M : a < \text{diam}(K) \leq b, K \subset B\}$ for all $0 \leq a < b$ and all Borel sets $B \subset \mathbb{R}^d$. If we denote the collection of sets $E(B; a, b)$ by \mathcal{E} , the latter is a π -system (i.e., closed under finite intersections), and one may set $\mathcal{M} = \sigma(\mathcal{E})$.

We now give a precise definition of scale invariance, followed by the main definitions of the paper.

Definition 2.1. We say that an infinite measure μ on (M, \mathcal{M}) is scale invariant if, for any $E \in \mathcal{M}$ with $\mu(E) < \infty$ and any $0 < s < \infty$, $\mu(E') = \mu(E)$, where $E' = \{K : K/s \in E\}$.

Definition 2.2. A scale invariant (Poissonian) soup in D with intensity $\lambda\mu$ is the collection of sets from a Poisson realization of $\lambda\mu$ that are contained in D , where μ is a translation and scale invariant measure.

Note that the soup inherits the scale invariance of the measure μ , so that soup realizations in domains related by uniform scaling are statistically self-similar. For instance, if $0 < s < 1$ and D are such that $sD = \{x \in \mathbb{R}^d : x/s \in D\} \subset D$, and \mathcal{K}_D denotes a realization of a scale invariant soup in D , then the collection of sets from \mathcal{K}_D contained in sD is distributed like a scaled version $s\mathcal{K}_D$ of \mathcal{K}_D (where the elements of $s\mathcal{K}_D$ are the sets $sK = \{x \in \mathbb{R}^d : x/s \in K\}$ with $K \in \mathcal{K}_D$).

Definition 2.3. A full space (Poissonian) soup with intensity $\lambda\mu$ and cutoff $\delta > 0$ is a Poisson realization from a measure $\lambda\mu_\delta$, where μ_δ is the measure induced by μ on sets of diameter at most δ , and μ is a translation and scale invariant measure.

The scale invariance of the soup can now be expressed in the following way. Let $0 < s < 1$, and let D and D' be two disjoint domains such that $s^{-1}\text{diam}(D) = \text{diam}(D') \leq \delta$ (where $\text{diam}(\cdot)$ denotes Euclidean diameter) and with D' obtained by translating sD . Then, as before, the sets that are contained in D' are distributed like a copy scaled by s of the sets contained in D . In other words, the soup is statistically self-similar at all scales smaller than the cutoff. Note that the full space soup is also stationary due to the translation invariance of μ .

Clearly, the value δ of the cutoff is not important, since one can always scale space to make it become 1. In the rest of the paper, when we talk about the full space soup without specifying the cutoff δ , we implicitly assume that $\delta = 1$.

We will consider translation and scale invariant measures μ that satisfy the following condition.

- (\star) Given a domain D and two positive real numbers $d_1 < d_2$, let $F = F(D; d_1, d_2)$ be the collection of compact connected sets with nonempty interior that intersect D and have diameters $> d_1$ and $\leq d_2$; then $\mu(F) < \infty$.

Remarks Condition (\star) is very natural and is clearly satisfied by μ^{Bool} and by the Brownian loop measure (which are also translation and scale invariant). Its purpose is to ensure that $\lambda_c > 0$ in Theorem 2.4 below and to ensure the left-continuity of certain crossing probabilities (see the beginning of Section 5). Note that the set F can be written as $F(D; d_1, d_2) = E(D'; d_1, d_2) \cap E(D' \setminus D; d_1, d_2)^c$, where D' is the (Euclidean) d_2 -neighborhood of D and the superscript c denotes the complement. Therefore, F is measurable by our assumptions on \mathcal{M} .

We are now ready to state the main results of the paper.

Theorem 2.4. *For every translation and scale invariant measure μ satisfying (\star), there exists $\lambda_c = \lambda_c(\mu)$, with $0 < \lambda_c < \infty$, such that, with probability one, the complement of the scale invariant soup with density $\lambda\mu$ contains connected components larger than one point if $\lambda \leq \lambda_c$, and is totally disconnected if $\lambda > \lambda_c$. The result holds for the full space soup and for the soup in any domain D with the same λ_c .*

We say that a random fractal *percolates* if it contains connected components larger than one point. (This is not the definition typically used for Mandelbrot’s fractal percolation, which involves a certain crossing event — see Section 4 below — but we think it is more natural and “canonical,” at least in the present context, precisely because it does not involve an arbitrary crossing event.) Theorem 2.4 therefore says that for the class of models included in the statement, with probability one the system *percolates at criticality*. We will show that this is equivalent to having positive probability for certain crossing events involving “shell-like” (deterministic) domains.

Remark As pointed out to us by an anonymous referee, a standard example of percolation at criticality is the appearance of a k -ary subtree inside a Galton-Watson tree (e.g., in Bernoulli percolation on a b -ary tree), with $k \geq 2$. This example has in fact played a role in fractal percolation (e.g., it is used in the proof of Theorem 1 of [6]). It would be interesting to determine whether there is a connection between that example and the class of models treated in this paper.

Corollary 2.5. *Consider a full space soup in \mathbb{R}^d with density $\lambda\mu$, where μ is a translation and scale invariant measure satisfying (\star). If $\lambda \leq \lambda_c(\mu)$, with probability one, the complement of the soup*

contains arbitrarily large connected components. Moreover, if μ is invariant under rotations, any two open subsets of \mathbb{R}^d are intersected by the same connected component of the complement of the soup with positive probability.

Corollary 2.5 leaves open the question of existence, and possibly uniqueness, of an unbounded connected component. We are able to address this question only for $d = 2$ (see Theorem 3.2).

Remark The measure μ does not need to be completely scale invariant for our results above to hold. As it will be clear from the proofs, it suffices that there is an infinite sequence of scale factors $s_j \downarrow 0$ such that μ is invariant under scaling by s_j , in the sense described above. This is the case, for instance, for the multiscale Boolean model studied in Chapter 8 of [24] and in [25; 26].

Indeed, our definition of self-similar soups is aimed at identifying a natural class of models that is easy to define and contains interesting examples; we did not try to define the most general class of models to which our methods apply. In Section 4 we will use Mandelbrot’s fractal percolation to illustrate how our techniques can be easily applied to an even larger class of models.

The main technical tool in proving Theorem 2.4 and Corollary 2.5 is Lemma 5.3, presented in Section 5. The lemma implies that the probability that the complement in a “shell-like” domain A of a full space soup contains a connected component that touches both the “inner” and the “outer” boundary of the domain has a discontinuity at some $0 < \lambda_c^A < \infty$, jumping from a positive value at λ_c^A to zero for $\lambda > \lambda_c^A$. It is then easy to see that the complement of the soup must be totally disconnected for $\lambda > \lambda_c^A$ (Lemma 5.2), which implies that λ_c^A is the same for all “shell-like” domains and coincides with the λ_c of Theorem 2.4.

For future reference we define what we mean by a *shell* and a *simple shell*. We call a set A a *shell* if it can be written as $A = D \setminus \overline{D'}$, where D and D' are two non-empty, bounded, d -dimensional open sets with $\overline{D'} \subset D$. A shell A is *simple* if D and D' are open, concentric (d -dimensional) cubes. We will denote by Φ_A the probability that the complement of a full space soup contains a connected component that touches both the “inner” and the “outer” boundary of A .

We note that the proof of Lemma 5.3 makes essential use of the shell geometry and would not work in the case, for instance, of crossings of cubes. Throughout the proof of Theorem 2.4, we choose to use simple shells because they are easier to work with. However, all our results can be readily generalized to any shell. In particular, we have the following discontinuity result, which is interesting in its own right.

Corollary 2.6. *For all $d \geq 2$, all shells A , and all translation and scale invariant measures μ satisfying (\star) , the following holds:*

- $\Phi_A(\lambda) > 0$ if $\lambda \leq \lambda_c(\mu)$,
- $\Phi_A(\lambda) = 0$ if $\lambda > \lambda_c(\mu)$.

3 Two-Dimensional Soups

In two dimensions one can obtain additional information and show that, like in Mandelbrot’s fractal percolation, a unique infinite connected component appears as soon as there is positive probability of having connected components larger than one point (that is, at and below the critical point λ_c).

To prepare for our main result of this section, Theorem 3.2 below, consider a self-similar soup in the unit square $(0, 1)^2$, and let $g(\lambda)$ be the probability that the complement of the soup contains a connected component that crosses the square in the first coordinate direction, connecting the two opposite sides of the square. We then have the following result.

Theorem 3.1. *For every translation and scale invariant μ in two dimensions which satisfies condition (\star) and is invariant under reflections through the coordinate axes and rotations by 90 degrees, $g(\lambda_c(\mu)) > 0$.*

The invariance under reflections through the coordinate axes is required because $g(\lambda)$ is defined using crossings of the unit square $(0, 1)^2$.

In the case of the Brownian loop soup, Werner states a version of Theorem 3.1 in [33; 34]. The choice of the unit square in Theorem 3.1 is made only for convenience, and similar results can be proved in the same way for more general domains. The reflection invariance is a technical condition needed in the proof in order to apply a technique from [10] (see the proof of Lemma 5.1 there). The same technique, combined with Theorem 3.1, can be used to prove the next theorem, which is our main result of this section.

Theorem 3.2. *For every translation and scale invariant μ in two dimensions which satisfies condition (\star) and is invariant under reflections through the coordinate axes and rotations by 90 degrees, if $\lambda \leq \lambda_c(\mu)$, the complement of the full plane soup with density $\lambda\mu$ has a unique unbounded component with probability one.*

The informed reader might believe that the uniqueness result would follow from a version of the classical Burton-Keane argument (see [4]). However, in the Burton-Keane argument it is crucial, for instance, that a path from the inside to the outside of a cube of side length n uses at least (roughly) a portion $1/n^{d-1}$ of the “surface” of the cube (e.g., the number of sites of \mathbb{Z}^d on the boundary, for a lattice model defined on \mathbb{Z}^d), so that there is enough space for at most $O(n^{d-1})$ disjoint paths. There is clearly no analogue of this for the continuous models in this paper (nor for Mandelbrot percolation), since the relevant paths have no “thickness.”

4 Applications to Mandelbrot’s Fractal Percolation

The method of proof of Theorem 2.4 works in greater generality than the class of scale invariant soup models introduced in this paper. In order for the method to work, it suffices to have some form of scale invariance. To illustrate this fact, we will consider a well-known model, called *fractal percolation*, that was introduced by Mandelbrot [21; 22] and is defined by the following iterative procedure.

For any integers $d \geq 2$ and $N \geq 2$, and real number $0 < p < 1$, one starts by partitioning the unit cube $[0, 1]^d \subset \mathbb{R}^d$ into N^d subcubes of equal size. Each subcube is independently retained with probability p and discarded otherwise. This produces a random set $\mathcal{C}_N^1 = \mathcal{C}_N^1(d, p) \subset [0, 1]^d$. The same procedure is then repeated inside each retained subcube, generating the random set $\mathcal{C}_N^2 \subset \mathcal{C}_N^1$. Iterating the procedure ad infinitum yields an infinite sequence of random sets $[0, 1]^d \supset \mathcal{C}_N^1 \supset \mathcal{C}_N^2 \supset \mathcal{C}_N^3 \supset \dots$. It is easy to see that the *limiting retained set* $\mathcal{C}_N := \bigcap_{k=1}^{\infty} \mathcal{C}_N^k$ is well defined.

Several authors studied various aspects of Mandelbrot’s fractal percolation, including the Hausdorff dimension of \mathcal{C}_N , as detailed in [9], and the possible existence of paths [6; 10; 23; 7; 35; 5; 12;

13; 3] and $(d - 1)$ -dimensional “sheets” [7; 28; 3] traversing the unit cube between opposite faces. Dekking and Meester [10] proposed a “morphology” of more general “random Cantor sets,” obtained by generalizing the successive “deletion of middle thirds” construction using random substitutions, and showed that there can be several critical points at which the connectivity properties of a set change. Accounts of fractal percolation can be found in [8] and Chapter 15 of [11].

In this section we define three potentially different critical points. Theorem 4.1 shows that two of them are in fact the same. Furthermore, we prove that the third one is equal to the other two for N large enough, and conjecture that they are in fact the same for all N .

The first critical point is

$$\tilde{p}_c = \tilde{p}_c(N, d) := \sup\{p : \mathcal{C}_N \text{ is totally disconnected with probability one}\}.$$

To define the second critical point we focus on a specific shell. This choice is convenient but arbitrary and unnecessarily restrictive. Indeed, the proof of the next theorem shows that we could have chosen any other shell, so that \hat{p}_c defined below is independent of the choice of shell. Let $A \subset [0, 1]^d$ be the domain obtained by removing from the open unit cube $(0, 1)^d$ the cube $(1/2, \dots, 1/2) + [0, 1/3]^d$ of side length $1/3$, centered at $(1/2, \dots, 1/2)$. Denote by $\phi_A(p)$ the probability that the limiting retained set contains a connected component that intersects both the “inner” and the “outer” boundary of A . The second critical point is $\hat{p}_c = \hat{p}_c(N, d) := \inf\{p : \phi_A(p) > 0\}$. Our first result of this section concerns $\hat{p}_c(N, d)$ and $\tilde{p}_c(N, d)$.

Theorem 4.1. *For all $d \geq 2$ and $N \geq 2$, $\hat{p}_c = \hat{p}_c(N, d)$ satisfies $0 < \hat{p}_c < 1$. Moreover $\phi_A(\hat{p}_c) > 0$, while \mathcal{C}_N is a.s. totally disconnected when $p < \hat{p}_c$. Hence, $\hat{p}_c(N, d) = \tilde{p}_c(N, d)$ for every N and d .*

Mandelbrot’s fractal percolation can be extended to a full space model by tiling \mathbb{R}^d with independent copies of the system in the natural way. We call this model *full space fractal percolation*. As a consequence of the previous theorem we have the following result.

Corollary 4.2. *Consider full space fractal percolation with $d \geq 2$ and $N \geq 2$. With probability one, the limiting retained set contains arbitrarily large connected components for $p \geq \hat{p}_c$, and is totally disconnected for $p < \hat{p}_c$.*

We say that there is a (left to right) crossing of the unit cube if \mathcal{C}_N contains a connected component that intersects both $\{0\} \times [0, 1]^{d-1}$ and $\{1\} \times [0, 1]^{d-1}$. Let $p_c(N, d)$ be the infimum over all p such that there is a crossing of the unit cube with positive probability. Sometimes the system is said to percolate when such a crossing occurs. For $d = 2$ and all $N \geq 2$, Chayes, Chayes and Durrett [6] discovered that, at the critical point $p_c(N, 2)$, the probability of a crossing is strictly positive (see [10] for a simple proof). A slightly weaker result in three dimensions was obtained in [7]. Broman and Camia [3] were able to extend the result of Chayes, Chayes and Durrett to all $d \geq 2$, but only for sufficiently large N . However, the same is conjectured to hold for all N .

It is interesting to notice that in two dimensions one can prove that, for $p = p_c(N, 2)$, \mathcal{C}_N contains an infinite connected component with probability one [6]. This is in sharp contrast with lattice percolation, where it has been proved that, with probability one, the system does not have an infinite cluster at the critical point in dimensions 2 and ≥ 19 . (The same is believed to hold in all dimensions — see [15] for a general account of percolation theory.)

By Theorem 4.1, \mathcal{C}_N is totally disconnected with probability one when $p < \hat{p}_c$, so that $\hat{p}_c(N, d) \leq p_c(N, d)$ for all N and d . Furthermore we have the following result.

Theorem 4.3. For $d = 2$ we have that $\hat{p}_c(N, 2) = p_c(N, 2)$ for all $N \geq 2$. Furthermore, for every $d \geq 3$, there exists $N_0 = N_0(d)$ such that, for all $N \geq N_0$, $\hat{p}_c(N, d) = p_c(N, d)$.

Remark We conjecture that $\hat{p}_c(N, d) = p_c(N, d)$ for all $N \geq 2$ and all $d \geq 2$.

5 Proofs

5.1 Proofs of the Main Results

Before we can give the actual proofs, we need some definitions. Let $A = D \setminus \overline{D'}$ be a shell. Given a set K , if $A \setminus K$ contains a connected component that connects the boundary of D with that of D' (in other words, if K does not disconnect $\partial D'$ from ∂D), we say that the complement of K crosses A , or that there is a crossing of A in $A \setminus K$. We let $\Phi_A(\lambda)$ denote the probability that the complement of a full space soup crosses A . If μ satisfies condition (\star) , the function $\Phi_A(\lambda)$ is left-continuous in λ . To see this, consider $\Phi_A^\varepsilon(\lambda)$, the analogous crossing probability obtained by disregarding sets in the soup of diameter smaller than ε . A standard coupling of Poisson processes (between different values of λ) shows that, if condition (\star) is satisfied, $\Phi_A^\varepsilon(\lambda)$ is continuous in λ for any $\varepsilon > 0$. Furthermore, $\Phi_A(\lambda)$ corresponds to the $\varepsilon \rightarrow 0$ limit of $\Phi_A^\varepsilon(\lambda)$, and is therefore left-continuous in λ , since $\Phi_A^\varepsilon(\lambda)$ is nonincreasing in λ .

We now start with the proof of our main theorem, leaving out some lemmas that will be proved later.

Proof of Theorem 2.4

Full Space Soup We will first prove the result for the full space soup. Define $\lambda_c = \lambda_c(\mu)$ to be the infimum of all λ such that with probability one the complement of the full space soup contains at most isolated points. Clearly, $\Phi_A(\lambda) = 0$ for $\lambda > \lambda_c$. The following lemma, whose proof is standard and deferred till later on, holds.

Lemma 5.1. For any translation and scale invariant measure μ satisfying condition (\star) , we have $0 < \lambda_c(\mu) < \infty$.

In order to conclude the proof for the full space soup, it suffices to show that $\Phi_A(\lambda_c) > 0$ for some simple shell A , since that would imply that the complement of the soup cannot be totally disconnected with probability one at $\lambda = \lambda_c$. In order to achieve that, we combine the left-continuity of $\Phi_A(\lambda)$ with the two following lemmas.

Lemma 5.2. If λ is such that $\Phi_A(\lambda) = 0$ for some simple shell A , then the complement of the full space soup with density $\lambda\mu$ is totally disconnected with probability one.

Lemma 5.3. For any simple shell A , there exists an $\varepsilon > 0$ such that, if $\Phi_A(\lambda) \leq \varepsilon$, then $\Phi_A(\lambda) = 0$.

Lemma 5.2 implies that for all simple shells A , $\Phi_A(\lambda) > 0$ for $\lambda < \lambda_c$. Together with the left-continuity of $\Phi_A(\lambda)$ and Lemma 5.3, this implies that $\Phi_A(\lambda_c) > 0$, which concludes this part of the proof.

Soup in a Bounded Domain We now prove the result for the soup in a domain D . Let λ_c^D be the infimum of the set of λ 's such that the complement of the soup with intensity $\lambda\mu$ in D is totally disconnected with probability one. Coupling the soup in D with a full space soup with cutoff larger than the maximum radius allowed in D by using the same Poisson realization for both before applying the cutoff or the condition that sets be contained in D , one can easily see that $\lambda_c^D \geq \lambda_c$. Indeed, more sets are “discarded” in the case of the soup in D , meaning that the intersection with D of the complement of the full space soup is contained in the complement of the soup in D . For any $\lambda < \lambda_c$, the complement of the full space soup intersected with D contains a connected component larger than one point with positive probability. This is because we can cover \mathbb{R}^d with a countable number of copies of D and use translation invariance. It follows that the complement of the soup in D , must also contain a connected component larger than one point with positive probability showing that $\lambda \leq \lambda_c^D$ so that $\lambda_c^D \geq \lambda_c$.

On the other hand, for any closed set $G \subset D$, when $\lambda > \lambda_c$, the intersection with G of the complement of the soup in D is easily seen to be totally disconnected by comparing it with the intersection of G with the complement of a full space soup with cutoff smaller than $\frac{1}{2}\text{dist}(G, \partial D)$, coupled to the soup in D in the same way as before. Therefore $\lambda_c^D \leq \lambda_c$, and we conclude that $\lambda_c^D = \lambda_c$ for all domains D .

It remains to check that at the critical point λ_c , the complement of the soup in D contains connected components larger than one point. This can be done by coupling the soup in D , as before, to a full space soup with cutoff larger than the maximum radius allowed in D . The intersection with D of the complement of such a soup is contained in the complement of the soup in D . For $\lambda = \lambda_c$, the complement of the full space soup intersected with D contains a connected component larger than one point with positive probability, and so does the complement of the soup in D . The proof of Theorem 2.4 is therefore complete. \square

We now turn to the proofs of Lemmas 5.1, 5.2 and 5.3. Lemma 5.1 can be proved in various, rather standard, ways. A detailed proof in the context of the multiscale Boolean model can be found in Chapter 8 of [24], while a different proof in the context of the Brownian loop soup is sketched in [34]. Both proofs are given in two dimensions, but the dimensionality of the space is irrelevant and the same arguments work in all dimensions. Since the same ideas work for all scale invariant soups, we only sketch the proof of the lemma, and refer the interested reader to [24] or [34] for more details.

Proof of Lemma 5.1 Since the whole space can be partitioned in a countable number of cubes, in order to prove that $\lambda_c < \infty$, it suffices to show that, for some λ sufficiently large, the unit cube $[0, 1]^d$ is completely covered with probability one.

Let n be a positive integer. Given μ , choose α small enough so that, with strictly positive probability, the cube $(-\alpha, \alpha)^d$ is covered by a set from the soup of intensity μ contained inside the cube $(-1, 1)^d$. Note that, if the probability of the event described above is strictly positive for a soup of intensity μ , and if we denote by $a(\lambda)$ the probability of the same event for a soup of intensity $\lambda\mu$, we have $a(\lambda) \rightarrow 1$ as $\lambda \rightarrow \infty$.

We can cover the unit cube with order 2^n cubes of side length $\alpha 2^{-n+1}$. Each such cube is contained inside order n nested simple shells with diameter $\leq 2\sqrt{2}$ and constant ratio α between the side length of the outer cube and that of the inner cube. Therefore, by scale invariance the probability

that the unit cube is not covered by sets from a soup with intensity $\lambda\mu$ is bounded above by some constant times

$$2^n(1 - a(\lambda))^n = 2^{(1 - |\log_2(1 - a(\lambda))|)n}.$$

For λ so large that $|\log_2(1 - a(\lambda))| > 1$, the exponent in the bound is negative and so the bound tends to zero as $n \rightarrow \infty$, which concludes the proof of the first part of the lemma.

Now let \mathcal{K} denote the collection of sets (from a soup) that intersect the unit cube $[0, 1]^d$. For each set $K \in \mathcal{K}$, define $q(K) \in \mathbb{N}$ in such a way that $\text{diam}(K) \in (2^{-q(K)-1}, 2^{-q(K)}]$. Partition the unit cube in cubes of side length 2^{-n} , and for each cube C of the partition Π_n define

$$\tilde{X}(C) = 1_{\{\exists K: q(K)=n \text{ and } K \cap C \neq \emptyset\}}.$$

The random variables $\{\tilde{X}(C)\}_{C \in \Pi_n}$ are not independent, since two adjacent elements of Π_n can intersect the same set K . It is in fact easy to see that they are positively correlated. However, it is possible to couple the collection of random variables $\{\tilde{X}(C)\}_{C \in \Pi_n}$ with a collection of independent Bernoulli random variables $\{X(C)\}_{C \in \Pi_n}$ such that, for each C , $\tilde{X}(C) \geq X(C)$ almost surely (see, e.g., [20]). Moreover, for any $\delta > 0$ we can take $P(X(C) = 1) \geq 1 - \delta$, if $P(\tilde{X}(C) = 1)$ is close enough to 1.

Let $F_n(C)$ denote the collection of compact sets with nonempty interior that intersect C and have diameters in $(2^{-n-1}, 2^{-n}]$. Then, for all $C \in \Pi_n$,

$$P(\tilde{X}(C) = 1) = e^{-b\lambda},$$

where $b = \mu(F_n(C)) < \infty$ by scale (and translation) invariance and condition (\star) . Therefore, by taking λ small enough, we can make $P(\tilde{X}(C) = 1)$, and thus also $P(X(C) = 1)$, arbitrarily close to 1. Since the collection of random variables $\{X(C) : C \in \Pi_n, n \in \mathbb{N}\}$ defines a Mandelbrot percolation process in $[0, 1]^d$ with retention probability equal to $P(X(C) = 1)$, and whose retained set is contained in $[0, 1]^d \setminus \mathcal{K}$, for sufficiently small λ , $[0, 1]^d \setminus \mathcal{K}$ contains connected components larger than one point with positive probability. \square

Proof of Lemma 5.2 Let $\lambda > 0$ and the simple shell $A = B^{out} \setminus \overline{B^{in}}$ be such that $\Phi_A(\lambda) = 0$. Because of scale and translation invariance, $\Phi_{A'}(\lambda) = 0$ for any A' obtained by translating a scaled shell sA with $s \leq 1$.

Given $\varepsilon > 0$, take $s = s(\varepsilon)$ such that $0 < s < 1$ and $\text{diam}(sB^{out}) = s \text{diam}(B^{out}) < \varepsilon$, and consider the simple shell sA . Consider a tiling of \mathbb{R}^d with non-overlapping (except along the boundaries) translates of $s\overline{B^{in}}$ such that the centers of the cubes form a regular lattice isomorphic to \mathbb{Z}^d . If the full space soup contains a connected component of diameter larger than ε , such a component must intersect some of the cubes from the tiling.

For any cube from the tiling, the probability that it is intersected by a connected component of the complement of the soup of diameter larger than ε is bounded above by the probability that a translate of sA is crossed by the complement of the soup. This follows from the fact that the diameter of the connected component is strictly larger than the diameter of sA . Since the probability of crossing sA is zero, and the number of cubes in the tiling is countable, we conclude that the full space soup cannot contain a connected component of diameter larger than ε , for any $\varepsilon > 0$. \square

Proof of Lemma 5.3 Let B^{in} and B^{out} be the two d -dimensional, concentric, open cubes such that $A = B^{out} \setminus \overline{B^{in}}$. For $0 < s < 1$, consider a tiling of \mathbb{R}^d with non-overlapping (except along the boundaries) translates of sB^{in} such that the centers of the cubes form a regular lattice isomorphic to \mathbb{Z}^d . We can use this isomorphism to put the translates of sB^{in} in a one-to-one correspondence with the vertices of \mathbb{Z}^d , and thus index them via the vertices of \mathbb{Z}^d .

For each translate $B_{x,s}^{in}$, $x \in \mathbb{Z}^d$, of sB^{in} , consider the translate $B_{x,s}^{out}$ of sB^{out} concentric to $B_{x,s}^{in}$. The two define the simple shell $A_{x,s} = B_{x,s}^{out} \setminus \overline{B_{x,s}^{in}}$. We then have a collection $\{A_{x,s}\}_{x \in \mathbb{Z}^d}$ of (overlapping) simple shells indexed by \mathbb{Z}^d .

Let \mathcal{K}_s denote the collection of sets from the full space soup with diameter at most s . Obviously, \mathcal{K}_s is distributed like a full space soup with cutoff $\delta = s$. Let $K_s = \bigcup_{K \in \mathcal{K}_s} K$, and denote by $\Psi_x(\lambda, s)$ the probability that there is a crossing of $A_{x,s}$ in the complement of K_s . It immediately follows from scale and translation invariance and the way $A_{x,s}$ has been defined that $\Psi_x(\lambda, s) = \Phi_{s^{-1}A_{x,s}}(\lambda) = \Phi_A(\lambda)$.

We now introduce the graph \mathbb{M}^d whose set of vertices is \mathbb{Z}^d and whose set of edges is given by the adjacency relation: $x \sim y$ if and only if $\|x - y\| = 1$, where $\|x\| = \|(x_1, \dots, x_d)\| := \max_{1 \leq i \leq d} |x_i|$ and $|\cdot|$ denotes absolute value. Next, for each $0 < s < 1$, we define the random variables $\{X_s(x)\}_{x \in \mathbb{Z}^d}$ by letting $X_s(x) = 1$ if there is a crossing of $A_{x,s}$ in the complement of K_s , and $X_s(x) = 0$ otherwise. By construction, the probability that $X_s(x) = 1$ equals $\Psi_x(\lambda, s) = \Phi_{s^{-1}A_{x,s}}(\lambda) = \Phi_A(\lambda) < 1$.

Note that, if $\|x - y\| > [\text{diam}(B^{out}) + 2]/l$, where l is the Euclidean side length of B^{in} , then $X_s(x)$ and $X_s(y)$ are independent of each other. This implies that we can apply Theorem B26 of [19] (see p. 14 there; the result first appeared in [20]) to conclude that there exist i.i.d. random variables $\{Y_s(x)\}_{x \in \mathbb{Z}^d}$ such that $Y_s(x) = 1$ with probability $p < 1$ and $Y_s(x) = 0$ otherwise, and $Y_s(x) \geq X_s(x)$ for every $x \in \mathbb{Z}^d$. Moreover, one can let $p \rightarrow 0$ as $\Phi_A(\lambda) \rightarrow 0$.

For each $0 < s < 1$, using the random variables $\{Y_s(x)\}_{x \in \mathbb{Z}^d}$, we can define a Bernoulli site percolation model on \mathbb{M}^d by declaring $x \in \mathbb{Z}^d$ open if $Y_s(x) = 1$ and closed if $Y_s(x) = 0$. We denote by $p_c(d)$ the critical value for Bernoulli site percolation on \mathbb{M}^d . (See [15] for a general account on percolation theory.)

Let $G_s := \{x \in \mathbb{Z}^d : A_{x,s} \subset A\}$, i.e., the set of vertices of \mathbb{Z}^d corresponding to simple shells $A_{x,s}$ contained in A . Note that, if s is sufficiently small, $\mathbb{Z}^d \setminus G_s$ contains two components, of which one is unbounded. These components are connected in terms of the adjacency relation \sim used to define \mathbb{M}^d when considered as subsets of the vertex set of \mathbb{M}^d .

When this is the case, if there is a crossing of A in the complement of the full space soup, then the percolation process on \mathbb{M}^d defined above has an open cluster that connects the bounded component to the unbounded component of $\mathbb{Z}^d \setminus G_s$, “crossing” G_s . The reason is that if the crossing of A intersects a box $B_{x,s}^{in}$, then $A_{x,s}$ must be also be crossed and so $Y_s(x) = 1$. The diameter of such an open cluster is at least of order $\text{dist}(\partial B^{out}, \partial B^{in})/sl$, and the cluster is contained in G_s , whose diameter is of the order of $\text{diam}(B^{out})/sl$. (Here, for $K_1, K_2 \subset \mathbb{R}^d$, $\text{dist}(K_1, K_2) := \inf\{|x - y| : x \in K_1, y \in K_2\}$.)

We are now ready to conclude the proof. Assume that $\Phi_A(\lambda) < \varepsilon$ with $0 < \varepsilon < 1$ so small that one can choose $p = P(Y_s(x) = 1)$ so that $p < p_c(d)$. Take s so small that G_s contains two connected components, as explained above. Then, for every s sufficiently small, $\Phi_A(\lambda)$ is bounded above by the probability that the Bernoulli percolation process defined via the random variables $\{Y_s(x)\}_{x \in \mathbb{Z}^d}$ contains an open cluster of diameter at least L/s inside a region of linear size at most L'/s , for some $L, L' < \infty$. Since $p < p_c(d)$, it follows from standard percolation results (see, e.g., [15]) that the

probability of such an event goes to zero as $s \rightarrow 0$, proving the lemma. \square

Proof of Corollary 2.6 Let A be a shell and let A' be a simple shell such that $A \subset A'$. Using Theorem 2.4 we conclude that $\Phi_A(\lambda) = 0$ if $\lambda > \lambda_c(\mu)$. Furthermore, we have that

$$0 < \Phi_{A'}(\lambda_c(\mu)) \leq \Phi_A(\lambda_c(\mu)),$$

since a crossing of A' implies a crossing of A . \square

Remark It is possible to define the notion of $(d - 1)$ -dimensional crossings of general shells. For example, a crossing could be, informally, a connected subset of the complement of the soup which divides the shell into two disjoint parts, both touching the “inner” and “outer” boundary of the shell. Using this definition of crossing it is possible to show results analogous to Theorem 2.4 and Corollary 2.6. Note that, for $d = 2$ this is not the definition that we use, but it is easy to see that our results would still be true with this definition of crossing.

The proof of our second main result is now easy.

Proof of Corollary 2.5 The first claim can be proved using Theorem 2.4 and a simple scaling argument, but it is also an immediate consequence of Corollary 2.6 combined with translation invariance.

The proof of the second claim uses rotation invariance. Let us consider, without loss of generality, two disjoint open balls, B_1 and B_2 , of radii r_1 and r_2 and centered at x_1 and x_2 , respectively. (If the balls are not disjoint, the complement of the soup intersects $B_1 \cap B_2$ with a connected component larger than one point with positive probability.) Let $d_{12} = |x_1 - x_2|$, and consider the shell $A = \{x : |x - x_1| < d_{12}\} \setminus \{x : |x - x_1| \leq r_1/2\}$. From Corollary 2.6 we know that $\Phi_A(\lambda) > 0$ for $\lambda \leq \lambda_c$. It then follows from rotation invariance that there is positive probability that the complement of the soup contains a connected component that connects the sphere $\{x : |x - x_1| = r_1/2\}$ with the surface $\{x : |x - x_1| = d_{12}\} \cap B_2$. Such a connected component must intersect both B_1 and B_2 . \square

5.2 Proofs of the Additional Two-Dimensional Results

Proof of Theorem 3.1 Let $\tilde{f}(\lambda)$ denote the probability that the complement of the full plane soup with density $\lambda\mu$ and cutoff $\delta = 2$ contains a connected component that crosses the rectangle $[0, 1] \times [0, 2]$ horizontally, and $\tilde{g}(\lambda)$ the probability that it contains a connected component that crosses the square $[0, 1]^2$ horizontally. Consider the annulus $A = [-3/2, 3/2]^2 \setminus [-1/2, 1/2]^2$. By Corollary 2.6, $\Phi_A(\lambda) > 0$ if $\lambda \leq \lambda_c(\mu)$. It is easy to see (Figure 4) that any crossing of A must cross either a square of side length 1 or a rectangle of side lengths 1 and 2 in the “easy” direction. Using translation and rotation invariance, this implies $\tilde{f}(\lambda_c) > 0$.

Let us now couple the full plane soup with cutoff $\delta = 2$ with full plane fractal percolation with $N = 3$ in such a way that at level $k = 0, 1, \dots$ of the fractal percolation construction, a square of

side length 3^{-k} is discarded if and only if it is covered by sets of the soup of diameter between $2/3^k$ and $2/3^{k+1}$, which happens with positive probability q . It is immediate that the limiting retained set of the full plane fractal percolation process contains the complement of the full plane soup. Therefore, $\tilde{f}(\lambda_c) > 0$ implies that there is positive probability that the limiting retained set of the fractal percolation process contains a connected component that crosses the rectangle $[0, 1] \times [0, 2]$ horizontally.

Note that in the fractal percolation process defined above, squares are not discarded independently, due to the presence of sets that can intersect two or more squares (up to four). However, two level- k squares of side length 3^{-k} at distance larger than $2/3^k$ are retained or discarded independently. In particular, if one marks every third square in a line of level- k squares, all marked squares are discarded with probability $q > 0$, independently of each other. This observation implies that we can apply the proof of Lemma 5.1 of [10] to the fractal percolation process defined above (as the reader can easily check).

The proof of the lemma shows that horizontal crossings of the rectangle $[0, 1] \times [0, 2]$ cannot be “too straight,” they must possess a certain “wavyness” so that, using invariance under reflections through the y -axis and translations, and the fact that crossing events are positively correlated, the horizontal crossings in five partially overlapping 1×2 rectangles “hook up” with positive probability to form a horizontal crossing of the rectangle $[0, 3] \times [0, 2]$ (see Figure 5 of [10] and the discussion in the proof of Lemma 5.1 there).

Since the horizontal crossings of the rectangle $[0, 1] \times [0, 2]$ in the complement of the soup form a subset of the fractal percolation crossings, they must possess the same “wavyness” property. In our setting, positive correlation of crossing events follows, for instance, from [16] and the fact that crossing events for the complement of the soup are decreasing. (Let $\mathcal{H} \subset \mathcal{H}'$ denote two soup realizations; an event \mathcal{A} is decreasing if $\mathcal{H} \notin \mathcal{A}$ implies $\mathcal{H}' \notin \mathcal{A}$.) Therefore, using the same “hook up” technique as in the proof of Lemma 5.1 of [10], but with crossings in the complement of the soup, we can conclude that there is positive probability that the complement of the full plane soup contains a horizontal crossing of the rectangle $[0, 3] \times [0, 2]$, and thus also of the square $[0, 3] \times [0, 3]$. By scaling, this implies that $\tilde{g}(\lambda_c) > 0$.

To conclude the proof, we couple the soup in the unit square $(0, 1)^2$ with density $\lambda_c(\mu)\mu$ with the full space soup with the same density and cutoff $\delta = 2$ by using the same Poisson realization for both before applying the cutoff or the condition that discs be contained in $(0, 1)^2$. Clearly, the intersection with $(0, 1)^2$ of the complement of the full space soup is contained in the complement of the soup in $(0, 1)^2$. Therefore, $\tilde{g}(\lambda_c) > 0$ implies $g(\lambda_c) > 0$, as required. \square

Proof of Theorem 3.2 The proof of Theorem 3.1 shows that, for all $\lambda \leq \lambda_c(\mu)$, there is positive probability that the complement of the full plane soup with density $\lambda\mu$ and cutoff $\delta = 1$ contains a horizontal crossing of the rectangle $[0, 3] \times [0, 2]$. Crossing events like the one just mentioned are decreasing and are therefore positively correlated (see, e.g., Lemma 2.2 of [16]).

Let us call $a(\lambda)$ the probability that the complement of the full plane soup with density $\lambda\mu$ and cutoff $\delta = 1$ contains a horizontal crossing of the rectangle $[0, 3] \times [0, 2]$, and $b(\lambda)$ the probability that it contains a vertical crossing of the square $[0, 2] \times [0, 2]$. Positive correlation of crossing events, combined with a standard “pasting” argument (see Fig. 1), implies that the probability that the complement of the full plane soup with intensity $\lambda\mu$ and cutoff $\delta = 1$ contains a horizontal crossing of the rectangle $[0, 6] \times [0, 2]$ is bounded below by $a(\lambda)^4 b(\lambda)^3$. We denote by h_0 this probability,

and by h_n the probability of the event \mathcal{B}_n that the complement of the soup contains a horizontal crossing of $[0, 2 \cdot 3^{n+1}] \times [0, 2 \cdot 3^n]$.

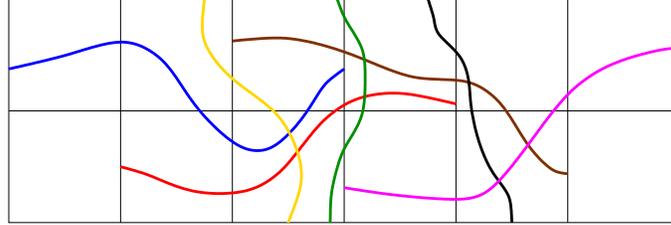


Figure 1: Pasting horizontal crossings of 3 by 2 rectangles and vertical crossings of 2 by 2 squares, a crossing of a 6 by 2 rectangle is obtained. (The different colors for the crossings serve only to enhance the visibility of the figure.)

Consider the full plane soup with cutoff $\delta_n = 3^{-n}$ obtained by a “thinning” of the soup with cutoff $\delta = 1$ that consists in removing from it all the sets with diameter larger than 3^{-n} . By scaling, the probability that the complement of the soup with cutoff δ_n contains a horizontal crossing of the rectangle $[0, 3] \times [0, 2]$ is equal to h_n . The complements of the soups with cutoffs $\{\delta_n\}_{n \in \mathbb{N}}$ form an increasing (in the sense of inclusion of sets) sequence of nested sets. Therefore, the limit of h_n as $n \rightarrow \infty$ is the probability of $\bigcup_{n \geq 0} \mathcal{B}_n$. By Kolmogorov’s zero-one law, the latter probability is either 0 or 1. However, since it cannot be smaller than $h_0 > 0$, it must necessarily be 1.

Having established that $\lim_{n \rightarrow \infty} h_n = 1$, both the existence and the uniqueness of an unbounded component with probability one follow from standard pasting arguments. To prove existence one can use the event depicted in Fig. 2 (and its rotation by 90 degrees) to couple the complement of the soup to a one-dependent bond percolation process on a square grid with parameter $p_n \rightarrow 1$ as $n \rightarrow \infty$ (see Fig. 2). Thus, choosing n sufficiently large implies percolation in the bond percolation process and, by the coupling, existence of an unbounded component in the complement of the soup. Now note that the event \mathcal{A}_n that the complement of the soup contains a circuit inside $[-3^{n+1}, 3^{n+1}]^2 \setminus [-3^n, 3^n]^2$ surrounding $[-3^n, 3^n]^2$ has probability bounded below by h_n^4 (see Fig. 3). Hence, \mathcal{A}_n occurs for infinitely many n , which implies the uniqueness of the unbounded component. \square

5.3 Proofs Concerning Mandelbrot’s Fractal Percolation Model

Sketch of the Proof of Theorem 4.1 The fact that $0 < \hat{p}_c(N, d) < 1$ follows from the same arguments that show that $0 < p_c(N, d) < 1$, where $p_c(N, d)$ is the critical probability defined in terms of crossings of a cube (see Section 4).

Showing that $\phi_A(\hat{p}_c) > 0$ follows the strategy of the proof of Lemma 5.3. In fact, the proof is even easier in this case, since the strategy is particularly well-suited for Mandelbrot’s fractal percolation model. The reason for this lies in the geometry of the fractal construction.

Using the fact that $\phi_A(p) = 0$ for $p < \hat{p}_c$, the proof that the limiting retained set \mathcal{C}_N is totally disconnected if $p < \hat{p}_c$ is essentially the same as the proof of Lemma 5.2, to which we refer the

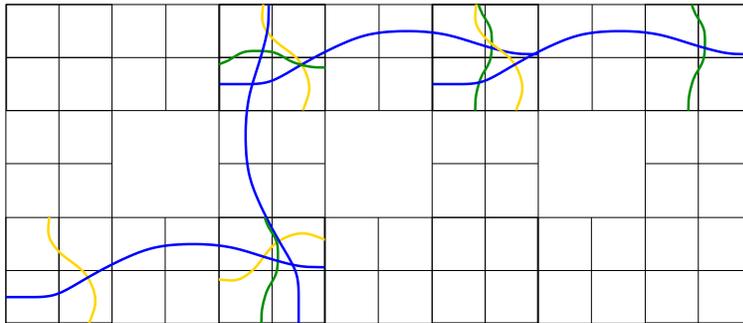
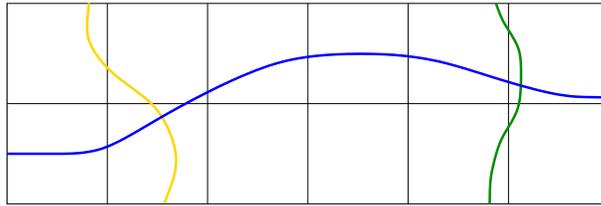


Figure 2: The event depicted above and its rotation by 90 degrees can be used to couple the complement of the soup to a one-dependent percolation process on a square grid whose open edges correspond to rectangles where the event occurs. We denote by p_n the probability of the event when the elementary squares in the figure have side length 3^n , i.e., the probability that an edge is open in the corresponding bond percolation process. (As in Fig. 1, the different colors for the crossings serve only to enhance the visibility of the figure.)

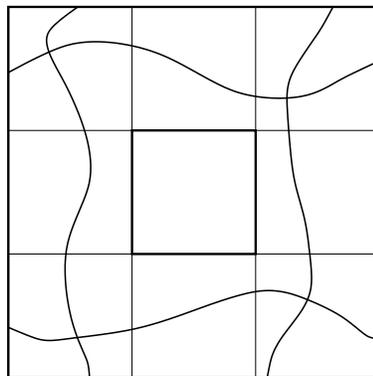


Figure 3: Four crossings of rectangles forming a circuit inside an annulus.

reader. This also shows that $\hat{p}_c(N, d) = \tilde{p}_c(N, d)$. \square

Proof of Corollary 4.2 Since full space fractal percolation is obtained by tiling \mathbb{R}^d with independent

copies of fractal percolation in $[0, 1]^d$, it immediately follows from Theorem 4.1 that the system is totally disconnected when $p < \hat{p}_c = \hat{p}_c(N, d)$.

Let us now show that, when $p \geq \hat{p}_c$, for any $M > 0$, the system contains a connected component with diameter larger than M with probability one. Consider the event that the unit cube $[0, 1]^d$ contains a connected component of diameter larger than ε , and let $\pi(p, \varepsilon)$ denote its probability. Fix an $0 < \varepsilon_0 < 1$ and observe that by Theorem 4.1 $\pi(p, \varepsilon_0) > 0$. Let k_0 be the smallest integer such that $\varepsilon_0 N^{k_0} > M$.

Consider the full space fractal percolation process obtained by conditioning on total retention in the unit cube $[0, 1]^d$ of the first k_0 iterations. Since the limiting retained set of this process is clearly stochastically larger than the limiting retained set of the original one, the probability that the conditioned process contains a connected component of diameter larger than ε_0 in the unit cube $[0, 1]^d$ is at least $\pi(p, \varepsilon_0) > 0$. However, by scaling, this probability is the same as the probability that the original process contains a connected component of diameter larger than $N^{k_0} \varepsilon_0 > M$ in the cube $N^{k_0} [0, 1]^d$. We can now use translation invariance to conclude that a connected component with diameter larger than M is present in the original full space fractal percolation system with probability one. \square

The next proof relies on ideas developed in [3], and is similar to the proof of Theorem 1.1 there. For this reason, we present here only a sketch of the proof, referring the interested reader to [3] for more details.

Sketch of the Proof of Theorem 4.3 To prove the first statement, it suffices to show that if the shell A in the definition of \hat{p}_c is crossed with positive probability, then the unit square $[0, 1]^2$ is also crossed with positive probability. It is easy to see (Fig. 4) that any crossing of A must cross either a square of side length $1/3$ or a rectangle of side lengths $1/3$ and $2/3$ in the “easy” direction. Using Theorem 4.1, this implies that when $p \geq \hat{p}_c$, there is positive probability of having a crossing of the rectangle $[0, 1/3] \times [0, 2/3]$ in the horizontal direction. According to results from [10] (see Lemma 5.1 there) this implies that there is a crossing of $[0, 1] \times [0, 2/3]$ with positive probability, which implies the first statement of the theorem.

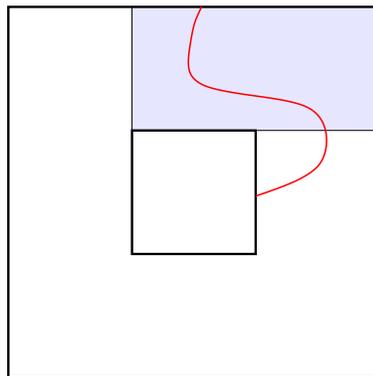


Figure 4: The shaded rectangle is crossed in the “easy” direction.

Since the limiting retained set \mathcal{C}_N is totally disconnected with probability one when $p < \hat{p}_c(N, d)$, it is immediate that $\hat{p}_c(N, d) \leq p_c(N, d)$ for all N . Assume that $p_c(N, d) > \hat{p}_c(N, d)$ for all N . Then,

for each N , we can choose $p_0 = p_0(N, d)$ such that $\hat{p}_c(N, d) < p_0(N, d) < p_c(N, d)$. We will show that this leads to a contradiction for N large enough.

Consider Mandelbrot's fractal percolation process in $[0, 1]^d$ with retention probability p_0 , and denote by \mathcal{A}_k the event that there is complete retention up to the k -th iteration, i.e., $\mathcal{C}_N^k = [0, 1]^d$. Let \mathbb{L}^d be the d -dimensional lattice with vertex set \mathbb{Z}^d and with edge set given by the adjacency relation: $(x_1, \dots, x_d) = x \sim y = (y_1, \dots, y_d)$ if and only if $x \neq y$, $|x_i - y_i| \leq 1$ for all i and $x_i = y_i$ for at least one value of i .

Conditioned on \mathcal{A}_{k-1} , we can couple level k of the fractal percolation process to a diminishment percolation process (see [1; 15]) on \mathbb{L}^d with the following diminishment rule: for a vertex $x \in \mathbb{Z}^d$, if all \mathbb{L}^d -neighbors of x are closed, except possibly two nearest neighbors in \mathbb{Z}^d , we make x closed, regardless of its state before the diminishment. Note that this has the effect of "diminishing" the percolation configuration by changing the state of some vertices from open to closed. The diminishment is *essential* in the language of [1] (see also [3]). The coupling is exactly the same as in the proof of Theorem 1.1 of [3] (although the diminishment rule is different), therefore the interested reader can check the details in [3].

Let $\psi_k(p_0)$ denote the probability that there is an open crossing between the inner and outer boundaries of $N^k A$ in the diminishment percolation process with initial density p_0 of open vertices. A feature of the coupling is that the diminishment percolation process dominates the fractal percolation process in the sense that $\phi_A(p_0) \leq \psi_k(p_0)$ (see [3]).

From [12; 13] we know that, for all $d \geq 2$, $p_c(N, d) \rightarrow p'_c(d)$ as $N \rightarrow \infty$, where $p'_c(d)$ is the critical value for Bernoulli site percolation on \mathbb{L}^d . Furthermore, for the diminished percolation model, the critical value $p''_c(d)$ satisfies $p''_c(d) > p'_c(d)$ (which follows from the diminishment being essential, see [1; 15] for more details on enhancement and diminishment percolation).

This implies that, for fixed d and N sufficiently large, the diminishment percolation process with initial density $p_0 = p_0(N, d) < p_c(N, d) < p''_c(d)$. Observe that it is not subcritical in the sense that $p_0 < p'_c(d)$, rather it is the *diminished* percolation process that is subcritical. Therefore, for N sufficiently large, $\lim_{k \rightarrow \infty} \psi_k(p_0) = 0$, which implies that $\phi_A(p_0) = 0$. However, since $p_0 > \hat{p}_c(N, d)$, $\phi_A(\hat{p}_c(N, d)) > 0$ by Theorem 4.1, and crossing probabilities are increasing in p , this leads to a contradiction. \square

Remark The result from [10] that is used for the above proof when $d = 2$ uses planar arguments and therefore cannot be readily generalized to $d \geq 3$.

Acknowledgements The second author thanks the Department of Mathematics of Chalmers University of Technology for the hospitality during a visit in 2008 when this work was started. He also thanks the Netherlands Organization for Scientific Research (NWO) for making the visit possible and for financial support through a VENI grant while part of this work was carried out. The first author thanks the Department of Mathematics of the Vrije Universiteit Amsterdam for the hospitality during a subsequent visit. Both authors thank Institut Mittag-Leffler for the hospitality during the Spring 2009 program on Discrete Probability when this paper was completed. They also thank Ronald Meester for suggesting to look at the multiscale Boolean model studied in [24] and for giving many useful comments on the manuscript. Finally they thank an anonymous referee who provided many good comments and useful suggestions.

References

- [1] M. Aizenmann and G. Grimmett, Strict Monotonicity for Critical Points in Percolation and Ferromagnetic Models, *J. Stat. Phys.* **63**, 817–835 (1991). MR1116036
- [2] H. Biermé and A. Estrade, Covering the whole space with Poisson random balls, in preparation (2010).
- [3] E.I. Broman and F. Camia, Large- N limit of crossing probabilities, discontinuity, and asymptotic behavior of threshold values in Mandelbrot’s fractal percolation process, *Electron. J. Probab.* **13**, 980–999 (2008). MR2413292
- [4] R.M. Burton and M. Keane, Density and Uniqueness in Percolation, *Comm. Math. Phys.* **121**, 501–505 (1989). MR0990777
- [5] J.T. Chayes and L. Chayes, The large- N limit of the threshold values in Mandelbrot’s fractal percolation process, *J. Phys. A: Math. Gen.* **22**, L501–L506 (1989).
- [6] J.T. Chayes, L. Chayes and R. Durrett, Connectivity Properties of Mandelbrot’s Percolation Process, *Probab. Theory Relat. Fields* **77**, 307–324 (1988). MR0931500
- [7] J.T. Chayes, L. Chayes, E. Grannan and G. Swindle, Phase transitions in Mandelbrot’s percolation process in three dimensions, *Probab. Theory Relat. Fields* **90**, 291–300 (1991). MR1133368
- [8] L. Chayes, Aspects of the fractal percolation process, *Progress in Probability* **37**, 113–143 (1995). MR1391973
- [9] E.M. Dekking and G.R. Grimmett, Superbranching processes and projections of random Cantor sets, *Probab. Theory Relat. Fields* **78**, 335–355 (1988). MR0949177
- [10] E.M. Dekking and R.W.J. Meester, On the structure of Mandelbrot’s percolation process and other random Cantor sets, *J. Stat. Phys.* **58**, 1109–1126 (1990). MR1049059
- [11] K. Falconer, *Fractal Geometry* (2nd ed.), Wiley, Chichester, 2003. MR2118797
- [12] K.J. Falconer and G.R. Grimmett, The critical point of fractal percolation in three and more dimensions, *J. Phys. A: Math. Gen.* **24**, L491–L494 (1991). MR1117858
- [13] K.J. Falconer and G.R. Grimmett, On the geometry of Random Cantor Sets and Fractal Percolation, *J. Theor. Probab.* **5**, 465–485 (1992). MR1176432
- [14] B. Freivogel and M. Kleban, A Conformal Field Theory for Eternal Inflation?, *JHEP12(2009)019*. MR2593019
- [15] G. Grimmett, *Percolation* (2nd ed.), Springer-Verlag, Berlin, 1999. MR1707339
- [16] S. Janson, Bounds on the distribution of extremal values of a scanning process, *Stochastic Process. Appl.* **18**, 313–328 (1984). MR0770197
- [17] G.F. Lawler, *Conformally Invariant Processes in the Plane*, Mathematical Surveys and Monographs **114**, American Mathematical Society, Providence, RI (2005). MR2129588

- [18] G.F. Lawler and W. Werner, The Brownian loop soup, *Probab. Theory Relat. Fields* **128**, 565–588 (2004). MR2045953
- [19] T.M. Liggett *Stochastic Interacting Systems: Contact, Voter and Exclusion Processes* Springer-Verlag, Berlin, 1999. MR1717346
- [20] T.M. Liggett, R.H. Schonmann and A.M. Stacey, Domination by product measures, *Ann. Probab.* **25**, 71–95 (1997). MR1428500
- [21] B.B. Mandelbrot, Intermittent turbulence in self-similar cascades: divergence of high moments and dimension of the carrier, *J. Fluid Mech.* **62**, 331–358 (1974).
- [22] B.B. Mandelbrot, *The Fractal Geometry of Nature*, W.H. Freeman, San Francisco, 1983. MR0665254
- [23] R.W.J. Meester, Connectivity in fractal percolation, *J. Theor. Prob.* **5**, 775–789 (1992). MR1182680
- [24] R. Meester and R. Roy, *Continuum Percolation*, Cambridge University Press, New York, 1996. MR1409145
- [25] M.V. Menshikov, S.Yu. Popov and M. Vachkovskaia, On the connectivity properties of the complementary set in fractal percolation models, *Probab. Theory Relat. Fields* **119**, 176–186 (2001). MR1818245
- [26] M.V. Menshikov, S.Yu. Popov and M. Vachkovskaia, On a multiscale continuous percolation model with unbounded defects, *Bull. Braz. Math. Society* **34**, 417–435 (2003). MR2045167
- [27] S. Nacu and W. Werner, Random soups, carpets and dimensions, *J. London Math. Soc.*, to appear (2010).
- [28] M.E. Orzechowski, On the Phase Transition to Sheet Percolation in Random Cantor Sets, *J. Stat. Phys.* **82**, 1081–1098 (1996). MR1372435
- [29] R. Schneider and W. Weil, *Stochastic and Integral Geometry*, Springer-Verlag, Berlin, 2008. MR2455326
- [30] S. Sheffield and W. Werner, Conformal Loop Ensembles: Construction via Loop-soups, arXiv:1006.2373v1 (2010).
- [31] S. Sheffield and W. Werner, Conformal Loop Ensembles: The Markovian Characterization, arXiv:1006.2374v1 (2010).
- [32] D. Stoyan, W.S. Kendall and J. Mecke, *Stochastic geometry and its applications* (2nd ed.), Wiley, Chichester, 1985. MR0895588
- [33] W. Werner, SLEs as boundaries of clusters of Brownian loops, *C. R. Math. Acad. Sci. Paris* **337**, 481–486 (2003). MR2023758
- [34] W. Werner, Some recent aspects of random conformally invariant systems, in *Les Houches School Proceedings: Session LXXXII, Mathematical Statistical Physics* (A. Bovier, F. Dunlop, A. van Enter, J. Dalibard editors), pp. 57–98, Elsevier, Amsterdam, 2006. MR2581883

- [35] D.G. White, On fractal percolation in \mathbb{R}^2 , *Statist. Probab. Lett.* **45**, 187–190 (1999).
MR1718427