

Research Article

The Yang-Laplace Transform for Solving the IVPs with Local Fractional Derivative

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The IVPs with local fractional derivative are considered in this paper. Analytical solutions for the homogeneous and nonhomogeneous local fractional differential equations are discussed by using the Yang-Laplace transform.

1. Introduction

In recent years, the ordinary and partial differential equations have found applications in many problems in mathematical physics [1, 2]. Initial value problems (IVPs) for ordinary and partial differential equations have been developed by some authors in [3–6]. There are analytical methods and numerical methods for solving the differential equations, such as the finite element method [6], the harmonic wavelet method [7–9], the Adomian decomposition method [10–12], the homotopy analysis method [13, 14], the homotopy decomposition method [15, 16], the heat balance integral method [17, 18], the homotopy perturbation method [19], the variational iteration method [20], and other methods [21].

Recently, owing to limit of classical and fractional differential equations, the local fractional differential equations have been applied to describe nondifferentiable problems for the heat and wave in fractal media [22, 23], the structure relation in fractal elasticity [24], and Fokker-Planck equation in fractal media [25]. Some methods were utilized to solve the local fractional differential equations. For example, the local fractional variation iteration method was used to solve the heat conduction in fractal media [26, 27]. The local fractional decomposition method for solving the local fractional diffusion and heat-conduction equations was considered in [28, 29]. The local fractional series expansion method for solving the Schrödinger equation with the local

fractional derivative was presented [30]. The Yang-Laplace transform structured in 2011 [22] was suggested to deal with local fractional differential equations [31, 32]. The coupling method for variational iteration method within Yang-Laplace transform for solving the heat conduction in fractal media was proposed in [33].

In this paper, our aim is to use the Yang-Laplace transform to solve IVPs with local fractional derivative. The structure of the paper is as follows. In Section 2, some definitions and properties for the Yang-Laplace transform are given. Section 3 is devoted to the solutions for the homogeneous and nonhomogeneous IVPs with local fractional derivative. Finally, conclusions are presented in Section 4.

2. Yang-Laplace Transform

In this section we show some definitions and properties for the Yang-Laplace transform.

The local fractional integral operator is defined as [22, 23, 26–33]

$$\begin{aligned} {}_a I_b^{(\alpha)} f(x) &= \frac{1}{\Gamma(1+\alpha)} \int_a^b f(t) (dt)^\alpha \\ &= \frac{1}{\Gamma(1+\alpha)} \lim_{\Delta t \rightarrow 0} \sum_{j=0}^{N-1} f(t_j) (\Delta t_j)^\alpha, \end{aligned} \quad (1)$$

where $\Delta t_j = t_{j+1} - t_j$, $\Delta t = \max\{\Delta t_0, \Delta t_1, \dots, \Delta t_j, \dots\}$, $[t_j, t_{j+1}]$, $j = 0, \dots, N - 1$, $t_0 = a, t_N = b$, is a partition of the interval $[a, b]$.

As the inverse operator of (1), the local fractional derivative operator is given by [22, 23, 26–33]

$$f^{(\alpha)}(x_0) = \frac{d^\alpha f(x)}{dx^\alpha} \Big|_{x=x_0} = \lim_{x \rightarrow x_0} \frac{\Delta^\alpha (f(x) - f(x_0))}{(x - x_0)^\alpha}, \quad (2)$$

with $\Delta^\alpha (f(x) - f(x_0)) \cong \Gamma(1 + \alpha)\Delta(f(x) - f(x_0))$.

The Yang-Laplace transform is expressed by [22, 31–33]

$$\tilde{L}_\alpha \{f(x)\} = f_s^{\tilde{L}, \alpha}(s) = \frac{1}{\Gamma(1 + \alpha)} \int_0^\infty E_\alpha(-s^\alpha x^\alpha) f(x) (dx)^\alpha, \quad (3)$$

$0 < \alpha \leq 1,$

where $f(x)$ is a local fractional continuous function.

The inverse Yang-Laplace transform reads as [22, 31–33]

$$f(x) = \tilde{L}_\alpha^{-1} \{f_s^{\tilde{L}, \alpha}(s)\} = \frac{1}{(2\pi)^\alpha} \times \int_{\beta-i\infty}^{\beta+i\infty} E_\alpha(s^\alpha x^\alpha) f_s^{\tilde{L}, \alpha}(s) (ds)^\alpha, \quad (4)$$

where $s^\alpha = \beta^\alpha + i^\alpha \infty^\alpha$ and $\text{Re}(s^\alpha) = \beta^\alpha$.

Some properties for Yang-Laplace transform are presented as follows [21, 22, 22–33]:

$$\tilde{L}_\alpha \{af(x) + bg(x)\} = a\tilde{L}_\alpha \{f(x)\} + b\tilde{L}_\alpha \{g(x)\}, \quad (5)$$

$$\tilde{L}_\alpha \{f^{(n\alpha)}(x)\} = s^{n\alpha} \tilde{L}_\alpha \{f(x)\} - \sum_{k=1}^n s^{(k-1)\alpha} f^{(n-k)\alpha}(0), \quad (6)$$

$$\lim_{x \rightarrow 0} f(x) = \lim_{s \rightarrow \infty} s^\alpha F(s), \quad (7)$$

$$\lim_{x \rightarrow \infty} f(x) = \lim_{s \rightarrow 0} s^\alpha F(s), \quad (8)$$

$$\tilde{L}_\alpha \{f(ax)\} = \frac{1}{a^\alpha} f_s^{\tilde{L}, \alpha} \left(\frac{s}{a} \right), \quad a > 0, \quad (9)$$

$$\tilde{L}_\alpha \{x^{k\alpha} f(x)\} = (-1)^k \frac{d^{k\alpha} f_s^{\tilde{L}, \alpha}(s)}{ds^{k\alpha}}, \quad (10)$$

$$\tilde{L}_\alpha \{f(x - c)\} = f_s^{\tilde{L}, \alpha}(s) E_\alpha(-c^\alpha s^\alpha), \quad (11)$$

$$\tilde{L}_\alpha \{f(x) E_\alpha(c^\alpha x^\alpha)\} = f_s^{\tilde{L}, \alpha}(s - c), \quad (12)$$

$$\tilde{L}_\alpha \{x^{k\alpha} E_\alpha(c^\alpha x^\alpha)\} = \frac{\Gamma(1 + k\alpha)}{(s - c)^{(k+1)\alpha}}, \quad (13)$$

$$\tilde{L}_\alpha \{\sin_\alpha(c^\alpha x^\alpha)\} = \frac{c^\alpha}{s^{2\alpha} + c^{2\alpha}}, \quad (14)$$

$$\tilde{L}_\alpha \{\cos_\alpha(c^\alpha x^\alpha)\} = \frac{s^\alpha}{s^{2\alpha} + c^{2\alpha}}, \quad (15)$$

$$\tilde{L}_\alpha \{x^{k\alpha}\} = \frac{\Gamma(1 + k\alpha)}{s^{(k+1)\alpha}}. \quad (16)$$

3. IVPs with Local Fractional Derivatives

In this section we handle the homogeneous and non-homogeneous IVPs with local fractional derivative.

3.1. Homogeneous IVPs with Local Fractional Derivative

Example 1. The homogeneous IVPs with local fractional derivative are expressed by

$$\frac{d^{2\alpha} y}{d^{2\alpha} x} - \frac{d^\alpha y}{d^\alpha x} + 2y = 0. \quad (17)$$

The initial boundary conditions are presented as

$$y(0) = 1, \quad y^{(\alpha)}(0) = 0. \quad (18)$$

From (6) we have

$$\tilde{L}_\alpha \{y^{(\alpha)}(x)\} = s^\alpha \tilde{L}_\alpha \{y(x)\} - y(0), \quad (19)$$

$$\tilde{L}_\alpha \{y^{(2\alpha)}(x)\} = s^{2\alpha} \tilde{L}_\alpha \{y(x)\} - s^\alpha y(0) - f^{(\alpha)}(0). \quad (20)$$

Hence, making use of (19) and (20), (19) can be written as

$$s^{2\alpha} \tilde{L}_\alpha \{y(x)\} - s^\alpha y(0) - f^{(\alpha)}(0) - \{s^\alpha \tilde{L}_\alpha \{y(x)\} - y(0)\} + 2\tilde{L}_\alpha \{y(x)\} = 0. \quad (21)$$

Hence, we obtain

$$\tilde{L}_\alpha \{y(x)\} = \frac{1}{s^\alpha + 2} y(0) = \frac{1}{s^\alpha + 2}. \quad (22)$$

So, making use of (13), we get the solution of (17):

$$y(x) = E_\alpha(-2x^\alpha). \quad (23)$$

The solution of (17) for $\alpha = \ln 2 / \ln 3$ is shown in Figure 1.

Example 2. Let us consider the homogeneous IVPs with local fractional derivative in the form

$$\frac{d^{4\alpha} y}{d^{4\alpha} x} - y = 0 \quad (24)$$

subject to initial boundary conditions

$$y(0) = 0, \quad y^{(\alpha)}(0) = 0, \quad (25)$$

$$y^{(2\alpha)}(0) = 0, \quad y^{(3\alpha)}(0) = 1.$$

From (6) we have

$$\tilde{L}_\alpha \{y^{(4\alpha)}(x)\} = s^{4\alpha} \tilde{L}_\alpha \{y(x)\} - s^{3\alpha} y(0) - s^{2\alpha} y^{(\alpha)}(0) - s^\alpha y^{(2\alpha)}(0) - f^{(3\alpha)}(0),$$

so that

$$s^{4\alpha} \tilde{L}_\alpha \{y(x)\} - s^{3\alpha} y(0) - s^{2\alpha} y^{(\alpha)}(0) - s^\alpha y^{(2\alpha)}(0) - f^{(3\alpha)}(0) - \tilde{L}_\alpha \{y(x)\} = 0. \quad (27)$$

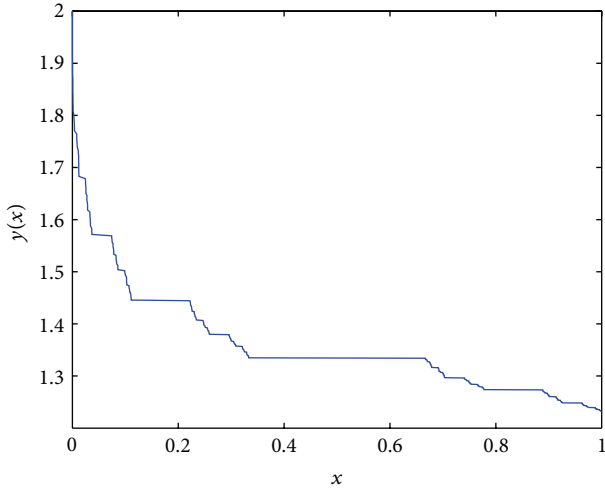


FIGURE 1: Graph of $y(x)$ for $\alpha = \ln 2/\ln 3$.

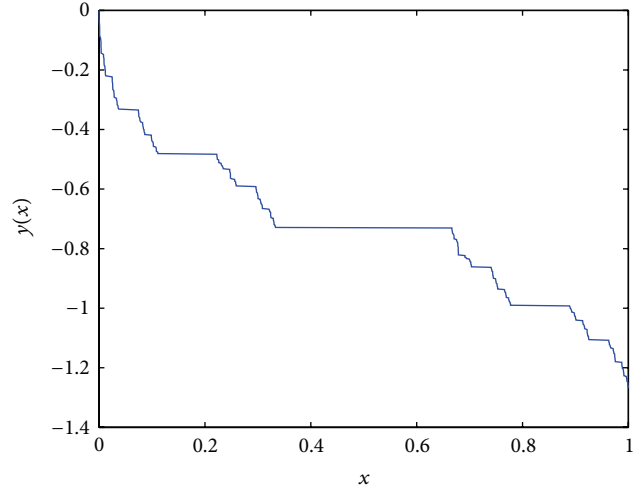


FIGURE 2: Graph of $y(x)$ for $\alpha = \ln 2/\ln 3$.

Hence, (27) can be written as

$$s^{4\alpha} \tilde{L}_\alpha \{y(x)\} - 1 - \tilde{L}_\alpha \{y(x)\} = 0, \tag{28}$$

which leads to

$$\tilde{L}_\alpha \{y(x)\} = \frac{1}{s^{4\alpha} - 1}. \tag{29}$$

Therefore, we get

$$\begin{aligned} y(x) &= \tilde{L}_\alpha^{-1} \left\{ \frac{1}{s^{4\alpha} - 1} \right\} \\ &= \tilde{L}_\alpha^{-1} \left\{ \frac{1}{2} \left(\frac{1}{2s^\alpha - 1} - \frac{1}{2s^\alpha + 1} - \frac{1}{s^{2\alpha} + 1} \right) \right\} \tag{30} \\ &= \frac{1}{4} E_\alpha(-x^\alpha) - \frac{1}{4} E_\alpha(x^\alpha) - \frac{1}{2} \sin_\alpha(x^\alpha). \end{aligned}$$

The exact solution of (24) for $\alpha = \ln 2/\ln 3$ is shown in Figure 2.

3.2. Nonhomogeneous IVPs with Local Fractional Derivative

Example 3. We now consider the non-homogeneous IVPs with local fractional derivative

$$\frac{d^{2\alpha} y}{d^{2\alpha} x} - y = \sin_\alpha(x^\alpha) \tag{31}$$

subject to initial boundary conditions

$$y(0) = 0, \quad y^{(\alpha)}(0) = 1. \tag{32}$$

By using (6), we have

$$\begin{aligned} \tilde{L}_\alpha \{y^{(2\alpha)}(x)\} &= s^{2\alpha} \tilde{L}_\alpha \{y(x)\} - s^\alpha y(0) - f^{(\alpha)}(0), \\ \tilde{L}_\alpha \{\sin_\alpha(x^\alpha)\} &= \frac{1}{s^{2\alpha} + 1} \end{aligned} \tag{33}$$

so that

$$\tilde{L}_\alpha \{y(x)\} = \frac{3}{4} \left(\frac{1}{s^\alpha - 1} - \frac{1}{s^\alpha + 1} \right) - \frac{1}{2} \frac{1}{s^{2\alpha} + 1}. \tag{34}$$

So,

$$y(x) = \frac{3}{4} E_\alpha(-x^\alpha) - \frac{3}{4} E_\alpha(x^\alpha) - \frac{1}{2} \sin_\alpha(x^\alpha). \tag{35}$$

The exact solution of (31) for $\alpha = \ln 2/\ln 3$ is shown in Figure 3.

Example 4. The non-homogeneous IVPs with local fractional derivative are

$$\frac{d^{2\alpha} y}{d^{2\alpha} x} + y = E_\alpha(x^\alpha). \tag{36}$$

The initial boundary conditions are

$$y(0) = 1, \quad y^{(\alpha)}(0) = 0. \tag{37}$$

In view of (6), we give

$$\tilde{L}_\alpha \{y(x)\} = \frac{1}{(s^\alpha + 1)(s^{2\alpha} + 1)} + \frac{s^\alpha}{s^{2\alpha} + 1}. \tag{38}$$

So, we obtain

$$\begin{aligned} y(x) &= \cos_\alpha(x^\alpha) + \frac{1}{\Gamma(1 + \alpha)} \int_0^x E_\alpha(x-t)^\alpha \sin_\alpha(t^\alpha) (dt)^\alpha \\ &= \cos_\alpha(x^\alpha) + \frac{1}{\Gamma(1 + \alpha)} \\ &\quad \times \int_0^x E_\alpha(t^\alpha) (\sin_\alpha(x^\alpha) \cos_\alpha(t^\alpha) \\ &\quad \quad - \cos_\alpha(x^\alpha) \sin_\alpha(t^\alpha)) (dt)^\alpha \end{aligned}$$

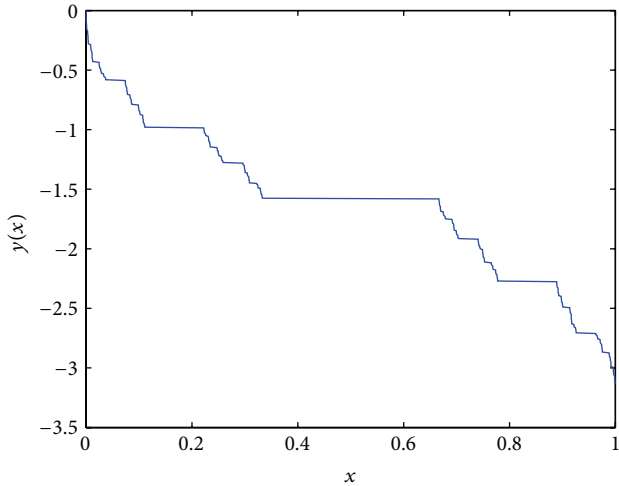


FIGURE 3: Graph of $y(x)$ for $\alpha = \ln 2 / \ln 3$.

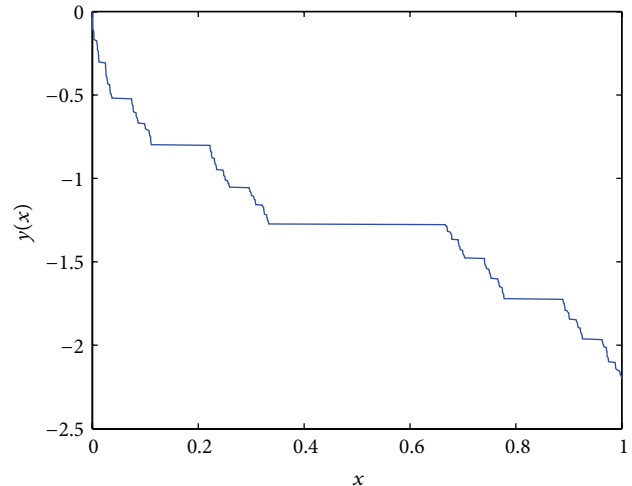


FIGURE 4: Graph of $y(x)$ for $\alpha = \ln 2 / \ln 3$.

$$\begin{aligned}
 &= \cos_\alpha(x^\alpha) \\
 &+ \sin_\alpha(x^\alpha) \left\{ \frac{1}{\Gamma(1+\alpha)} \int_0^x E_\alpha(t^\alpha) \cos_\alpha(t^\alpha) (dt)^\alpha \right\} \\
 &- \cos_\alpha(x^\alpha) \left\{ \frac{1}{\Gamma(1+\alpha)} \int_0^x E_\alpha(t^\alpha) \sin_\alpha(t^\alpha) (dt)^\alpha \right\} \\
 &= \cos_\alpha(x^\alpha) \\
 &+ \frac{\sin_\alpha(x^\alpha) \{E_\alpha(x^\alpha) [\cos_\alpha(x^\alpha) + \sin_\alpha(x^\alpha)] - 1\}}{2} \\
 &- \frac{\cos_\alpha(x^\alpha) \{E_\alpha(x^\alpha) [\sin_\alpha(x^\alpha) - \cos_\alpha(x^\alpha)] + 1\}}{2} \\
 &= \frac{1}{2} [\cos_\alpha(x^\alpha) - \sin_\alpha(x^\alpha) + E_\alpha(x^\alpha)].
 \end{aligned}
 \tag{39}$$

The exact solution of (36) for $\alpha = \ln 2 / \ln 3$ is shown in Figure 4.

4. Conclusions

In this work we have used the Yang-Laplace transform to handle the homogeneous and non-homogeneous IVPs with looselocal fractional derivative. Some illustrative examples of approximate solutions for local fractional IVPs are discussed. The nondifferentiable solutions for fractal dimension $\alpha = \ln 2 / \ln 3$ are shown graphically. The obtained results illustrate that the Yang-Laplace transform is an efficient mathematical tool to solve the homogeneous and non-homogeneous IVPs with local fractional derivative.

Conflict of Interests

The authors declare that there is no conflicts of interests regarding publication of this paper.

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