Commun. Math. Phys. 141, 41-61 (1991)



Condensation of a One-Dimensional Lattice Gas

Kurt Johansson*

Department of Mathematics, Royal Institute of Technology, Stockholm, Sweden**

Received February 6, 1990; in revised form April 16, 1991

Abstract. We consider a one-dimensional lattice gas in the canonical ensemble with interaction energy $1/r^{\alpha}$, $1 < \alpha \le 2$. Using an energy-entropy argument we show that the gas condenses at sufficiently low temperatures meaning that the gas has a non-uniform density in the thermodynamic limit.

0. Introduction

Consider a one-dimensional lattice gas in the canonical ensemble, i.e. the number of positions in $[1, L] \cap Z$ occupied by particles is a fixed number N. Let the interaction energy between two particles at distance r be $-1/r^{\alpha}$, $1 < \alpha \le 2$. We will show that, with appropriate boundary conditions, this lattice gas condenses at sufficiently low temperatures, meaning that the gas has a non-uniform density in the thermodynamic limit $N, L \to \infty, N/L \to \varrho, 0 < \varrho < 1$.

The proof is an energy-entropy argument which makes rigorous the following heuristic argument, see Landau-Lifshitz [6] abd Thouless [9]. Consider configurations that are not condensed already, say those that are approximately uniform. Suppose that we can partition these configurations into blocks $A_1, B_1, A_2, \ldots, A_k, B_k$, where almost every position in A_1, \ldots, A_k is occupied and almost every position in B_1, \ldots, B_k is empty. Also assume that all these blocks have length $\geq d$. Rearrange the blocks so that the condensed configuration $A_1A_2 \ldots A_kB_1 \ldots B_k$ is obtained. The change in energy is $\Delta E \sim -CkE_\alpha(d)$, where $E_\alpha(d) = d^{2-\alpha}$ if $1 < \alpha < 2$, $= \log d$ if $\alpha = 2$ and $\leq \text{const.}$ if $\alpha > 2$. The change in entropy due to many configurations being mapped to the same condensed configuration is roughly $\Delta S \sim -\log {\binom{L}{k}} \sim -k \log(L/k)$. The free energy thus changes by $\Delta F = \Delta E - \beta^{-1}\Delta S \sim k(-CE_\alpha(d) + \beta^{-1}\log(L/k))$. If $\alpha > 2$ and k = o(N),

^{*} Research supported by the Swedish research councils NFR and STUF

^{**} Mail address: Department of Mathematics, University of Uppsala, S-75238 Uppsala, Sweden

then we cannot have $\Delta F < 0$ with β finite and $N \to \infty$. If $\alpha = 2$ then $\Delta F < 0$ if $d \ge (L/k)^{1/\beta C}$, and thus in this case the system could condense. Similarly for $\alpha < 2$ with a smaller d. This explains why we can expect a condensation when $\alpha \le 2$ but not when a > 2. We see that in order for the above argument to work d must be sufficiently large. If we do rearrangements only locally a smaller d is possible. In the proof we will partition every configuration into blocks which are alternatingly dense and dilute. By making local rearrangements we obtain configurations with sufficiently long blocks to make the global rearrangement sketched above work.

The existence of a phase transition in a one-dimensional system with longrange interaction was first established by Dyson [2], who showed that there is spontaneous magnetization in the 1*d* Ising model with interaction $J_{ij} \sim 1/|i-j|^{\alpha}$, $1 < \alpha < 2$, at sufficiently low temperatures. That this also holds for $\alpha = 2$ was shown by Fröhlich and Spencer [4]. Recently the understanding of the case $\alpha = 2$ has improved considerably by use of the relation between the Ising model and certain percolation models. In particular the existence of the Thouless effect was proven. See Aizenman, Chayes and Newman [1]. Thus the existence of a phase transition in long-range 1*d* lattice models is not a new result. The novelty in this paper lies in the character of the proof, where we work in the canonical ensemble and establish the phase transition as a condensation. In an accompanying paper, Johansson [5], we will show how the same ideas can be extended to prove the existence of a liquid-gas transition in a one-dimensional continuous model.

The organization of the paper is as follows. In Sect. 1 we define the model and state our results. In Sect. 2 we define the partition of a configuration into blocks and the local rearrangement procedure. We then, in Sect. 3, state lemmas giving the change in energy and entropy due to the local rearrangement. These two lemmas are combined into an energy-entropy argument for the local rearrangement and we can then perform the global rearrangement and prove the main theorem. The proof of the entropy and energy lemmas occupy Sects. 4 and 5. The last section contains an argument which proves that, if the free energy is a strictly convex function of the density, then no condensation occurs. As shown in the first section this implies that at high temperatures the gas does not condense.

1. Definitions and Results

1.1. Definition of the Model

Let $\Lambda = [1, L] \cap Z$, $\Lambda^c = Z \setminus \Lambda$, where L is a positive integer. With each site $i \in \Lambda$ we associate a random variable n_i taking the values 0 and 1 and satisfying the constraint

$$\sum_{i\in\Lambda} n_i = N , \qquad (1.1)$$

where N is a positive integer. If $n_i = 1$ we say that the site *i* is occupied by a particle, otherwise it is empty. Equation (1.1) thus says that the total number of particles is fixed. As boundary conditions we set $n_i = 1$, if $i \in Z \cap (-\infty, 0]$ and $n_i = 0$ if $i \in Z \cap [L + 1, \infty)$. We let $\underline{n} = (n_i)_{i \in Z}$ and call this a *configuration* in Λ with the specified boundary conditions. $Q = Q_{N,L}$ denotes the set of all

configurations in Λ . The Hamiltonian is

$$H(\underline{n}) = H_{N,L}(\underline{n}) = -\sum_{\substack{i \in \mathcal{A}, j \in \mathbb{Z} \\ i \neq j}} \frac{n_i n_j}{|i - j|^{\alpha}}, \qquad (1.2)$$

where $\alpha > 1$. We will work with the canonical ensemble probability measure which is defined by

$$P(A) = P_{\beta,N,L}(A) = Z^{-1} \sum_{\underline{n} \in A} \exp(-\beta H(\underline{n})), \qquad (1.3)$$

where $A \subseteq Q$, β is the inverse temperature and

$$Z = Z(N, L) = \sum_{\underline{n} \in Q} \exp(-\beta H(\underline{n}))$$
(1.4)

is the partition function.

For $x, y \in Z$, $x \le y$ and a given configuration <u>n</u> we define

$$N(x, y)(\underline{n}) = \sum_{i: x \le i \le y} n_i.$$
(1.5)

 $N(x, y)(\underline{n})$ is the number of particles in [x, y].

Let $\varrho, 0 < \varrho < 1$, denote the asymptotic average density, and let $\Lambda \to Z^+$ denote the thermodynamic limit $N, L \to \infty$, $N/L \to \varrho$. For $0 < \tau_1 < \tau_2 < 1$ and $\underline{n} \in Q$ we define

$$d(\tau_1,\tau_2)(\underline{n}) = \frac{1}{L(\tau_2-\tau_1)} N([\tau_1 L], [\tau_2 L])(\underline{n}),$$

where $[\cdot]$ denotes the integer part.

Definition 1.1. We say that the gas has uniform density in the thermodynamic limit if for each $\varepsilon > 0$ and any fixed $0 < \tau_1 < \tau_2 < 1$,

$$\lim_{A\to Z^+} P\{\underline{n}\in Q; \varrho-\varepsilon \leq d(\tau_1,\tau_2) (\underline{n}) \leq \varrho+\varepsilon\} = 1.$$

Otherwise we will say that the gas has non-uniform density in the thermodynamic limit.

1.2. Results

For a given small $\delta > 0$ and density $\rho, 0 < \rho < 1$, we put

$$d_1 = (1 - \delta)^{-1} (\varrho - \delta), \quad d_2 = (1 - \delta)^{-1} \varrho.$$
 (1.6)

The main result of this paper is the following theorem.

Theorem 1.2. Let $1 < \alpha \le 2$. There exist positive constants K and κ , depending only on α , and $\beta_0 = \beta_0(\varrho, \alpha)$ such that if $\beta \ge \beta_0$ and $\delta = K \exp(-\kappa\beta)$, then for each $\varepsilon > 0$,

$$\lim_{A \to Z^+} P(\underline{n} \in Q; d(\tau_1, \tau_2) (\underline{n}) \ge 1 - 2\delta) = 1$$

for any fixed $0 \le \tau_1 < \tau_2 \le d_1 - \varepsilon$, and

$$\lim_{A \to Z^+} P(\underline{n} \in Q; d(\tau_1, \tau_2) (\underline{n}) \le 2\delta) = 1$$

for any fixed $d_2 + \varepsilon \leq \tau_1 < \tau_2 \leq 1$.

Thus for sufficiently high β the gas has non-uniform density in the thermodynamic limit. We also want to show that for sufficiently small β we have uniform density in the thermodynamic limit. To do this we will use the free energy of the lattice gas. The free energy, $f(\varrho, \beta)$, is defined by

$$\beta f(\varrho, \beta) = -\lim_{\Lambda \to Z^+} \frac{1}{L} \log Z(N, L), \qquad (1.7)$$

where Z(N,L) is the partition function defined above. It can be shown that this limit exists if $\alpha > 1$, is a convex function of ρ and does not depend on the choice of boundary conditions, see Ruelle [8]. In Sect. 6 we will prove the following theorem.

Theorem 1.3. If $f(\varrho, \beta)$ is a strictly convex function of ϱ , $0 < \varrho < 1$ for a fixed β , then the gas has uniform density in the thermodynamic limit for this β , and any ϱ , $0 < \varrho < 1$.

Using this we can now show

Theorem 1.4. If $\alpha > 1$ there is a β_1 such that if $\beta \leq \beta_1$ the gas has uniform density in the thermodynamic limit.

Proof. If we put $f(1,\beta) = \lim_{\varrho \to 1^-} f(\varrho,\beta)$ and $f(\varrho,\beta) = +\infty$ for $\varrho \notin [0,1]$, then by the lattice analogue of Theorem 3.4.4 in Ruelle [8], $f(\varrho,\beta)$ is a lower semicontinuous convex function of $\varrho \in R$. By the lattice analogue of Theorem 3.4.6 in Ruelle [8] the pressure $p(\mu,\beta)$ as a function of the chemical potential μ is given by the Legendre-Fenchel transform

$$\beta p(\mu, \beta) = \sup_{\varrho \in R} \left(\varrho \mu - \beta f(\varrho, \beta) \right).$$

Theorem 5.2.1 in Ruelle [8] says that there is a β_1 such that $p(\mu, \beta)$ is an analytic function of μ when $\beta \leq \beta_1$. By a theorem for convex functions, see Rockafellar [7] p. 253, $f(\varrho, \beta)$ is essentially strictly convex when $p(\mu, \beta)$ is differentiable as a function of μ . Thus $f(\varrho, \beta)$ is strictly convex in [0, 1) where it is finite. \Box

Corollary 1.5. If $\beta \ge \beta_0$, then $\partial p(\mu, \beta)/\partial \mu$ does not exist for all $\mu \in R$.

The lattice gas can be related to the Ising model by the substitution $n_i = \frac{1}{2}(s_i + 1)$, $s_i = \pm 1$. This gives the following relation between the pressure $p(\mu, \beta)$ for the lattice gas and the free energy $\psi(\beta, h)$ for the Ising model with interaction $J_{ij} = |i - j|^{-\alpha}$ and magnetic field h,

$$\psi(\beta, h) = 4\beta p(2h - 4\beta c, 4\beta) - h + \beta c,$$

where $c = 2\zeta(\alpha)$ and $\zeta(\alpha)$ is Riemanns ζ -function. Thus, by Corollary 1.5, when $\beta \ge \beta_0/4$, $(\partial \psi/\partial h)(\beta, h)$ does not exist for all $h \in R$. It is known that this can only happen when h = 0 and that we then have spontaneous magnetization, see Ellis [3], p. 105. The theorems above thus imply that the $1/r^{\alpha}$ -Ising model in one dimension has a phase transition when $\alpha \le 2$.

2. Blocks and Rearrangement

2.1. Definition of Blocks

We now begin the proof of Theorem 1. Basic for our proof is the partition of configurations into blocks which alternate between having a high and a low density of particles. Let $1 \le a < a' \le L + 1$ be two integers and $\underline{n} \in Q$ a

configuration. Then

$$A = \langle a, a' \rangle = (n_a, n_{a+1}, \ldots, n_{a'-1})$$

is called a block in <u>n</u>. A is an o-block if $n_a = n_{a'-1} = 1$ and an e-block if $n_a = n_{a'-1} = 0$. The infinite block $\langle -\infty, a \rangle$, $n_a = 1$, is always an o-block and $\langle a, \infty \rangle$, $n_a = 0$ is always an e-block, due to the boundary conditions. The length of the block A is |A| = a' - a, the number of positions in the block. Two o-(e-)blocks $A = \langle a, a' \rangle$ and $B = \langle a', a'' \rangle$ can be joined to a new o-(e-)block $AB = \langle a, a'' \rangle$. A set of integers $\gamma = \{a_1, \ldots, a_r\}$, $0 \le a_1 < a_2 < \cdots < a_r \le L+1$ defines a partition of a configuration <u>n</u> into blocks $\langle a_k, a_{k+1} \rangle$, $k = 0, \ldots, r$, where $a_0 = -\infty$, $a_{r+1} = \infty$. We will say that $\langle a_k, a_{k+1} \rangle$ is a block in (\underline{n}, γ) . The number of elements in γ will be denoted $|\gamma|$.

Let δ be as in Theorem 1.2. The constants K and κ will be specified later.

Definition 2.1. Let $\gamma(\underline{n})$ define a partition of the configuration \underline{n} into blocks. We will say that $(\underline{n}, \gamma(\underline{n}))$ has the *density property* if the blocks in $(\underline{n}, \gamma(\underline{n}))$ alternate between *o*- and *e*-blocks and for each *o*-(*e*-)block $A = \langle a, a' \rangle$ in $(\underline{n}, \gamma(\underline{n}))$, (i) $N(a, x) (\underline{n}) \ge 1 + (1 - \delta) (x - a)$ ($\le \delta (x - a)$),

(i) $N(a, x) (\underline{n}) \ge 1 + (1 - \delta) (x - a)$ $(\le \delta(x - a)),$ (ii) $N(x, a' - 1) (\underline{n}) \ge 1 + (1 - \delta) (a' - 1 - x)$ $(\le \delta(a' - 1 - x))$ if $a \le x \le a' - 1$.

Thus if a partition has the density property the *o*- and *e*-blocks that it defines have a high, respectively low, density of particles.

We will now define a partition $\gamma_1(\underline{n})$ for every $\underline{n} \in Q_1$. Put

$$u_k = 2^{k-1}, \quad v_k = \frac{2}{\delta} 4^{k-1}.$$

The blocks in $(\underline{n}, \gamma_1(\underline{n}))$ will be the basic building blocks in our argument. $\gamma_1(\underline{n})$ is defined inductively. Let

$$\gamma^{(0)}(\underline{n}) = \{i \in \{1, \ldots, L+1\}; |n_i - n_{i-1}| = 1\}.$$

The blocks defined by $\gamma^{(0)}(\underline{n})$ alternate between *o*- and *e*-blocks and in every block either all $n_i = 1$ or all $n_i = 0$. Suppose that $\gamma^{(k-1)}(\underline{n})$ has been defined and that the blocks defined by it alternate between *o*- and *e*-blocks. Consider all *o*-blocks in $(\underline{n}, \gamma^{(k-1)}(\underline{n}))$ of length $\geq v_k$, A_{j_1}, \ldots, A_{j_p} , $p \geq 1$, $(\langle -\infty, a_1 \rangle$ is always an *o*-block of length $\geq v_k$). If the *e*-blocks in $(\underline{n}, \gamma^{(k-1)}(\underline{n}))$ between $A_{j_r} = \langle a_{j_r}, a_{j_r+1} \rangle$ and $A_{j_{r+1}} = \langle a_{j_{r+1}}, a_{j_{r+1}+1} \rangle$ have total length $\leq u_k$ we put $\eta_r = [a_{j_r+1}, a_{j_{r+1}}] \cap \gamma^{(k-1)}(\underline{n})$. If the *e*-blocks have length $> u_k$ we put $\eta_r = \emptyset$. Do this for $r = 1, \ldots, p - 1$; if p = 1 there are no η_r 's. In the same way we define $\lambda_1, \ldots, \lambda_{q-1}$ with *e*-blocks and *o*-blocks exchanged. Put

$$\gamma^{(k)}(\underline{n}) = \gamma^{(k-1)}(\underline{n}) \setminus (\eta_1 \cup \cdots \cup \eta_{p-1} \cup \lambda_1 \cup \cdots \cup \lambda_{q-1}).$$

It is clear from the construction that the blocks defined by $\gamma^{(k)}(\underline{n})$ alternate between *o*- and *e*-blocks. We now let

$$\gamma_1(\underline{n}) = \gamma^{(2k_N)}(\underline{n}),$$

where

$$k_N = \left[\frac{1}{4}\log N\right].$$

Lemma 2.2. Consider two o-(e-)blocks of length $\geq v_k$ in $(\underline{n}, \gamma_1(\underline{n}))$, $\underline{n} \in Q_1$. Then the total length of the e-(o-)blocks between them is $> u_k$, $1 \leq k \leq 2k_N$.

Proof. Assume that the lemma is false. Then there is a $k, 1 \le k \le 2k_N$, and two o-(e-)blocks $\langle a, a' \rangle$ and $\langle b, b' \rangle$ in $(\underline{n}, \gamma_1(\underline{n}))$ of length $\ge v_k$, such that the length of the e-(o-)blocks between them is $\le u_k$ and all o-(e-)blocks between them have length $< v_k$. a' must have been a right endpoint of an o-(e-)block $\langle a'', a' \rangle$ and b a left endpoint of an o-(e-)block $\langle b, b'' \rangle$ in $(\underline{n}, \gamma^{(k-1)}(\underline{n}))$. Since b and a' were not removed at some later step the e-(o-)blocks between $\langle a'', a' \rangle$ and $\langle b, b'' \rangle$ in $(\underline{n}, \gamma^{(k-1)}(\underline{n}))$ have length $\le u_k$ and one of the blocks $\langle a'', a' \rangle$ and $\langle b, b'' \rangle$ in $(\underline{n}, \gamma^{(k-1)}(\underline{n}))$ have length $\le u_k$ and one of the blocks $\langle a'', a' \rangle$ and $\langle b, b'' \rangle$, say $\langle a'', a' \rangle$, must have length $< v_k$. But $a'' \notin \gamma_1(\underline{n})$, and hence a'' must have been removed at some later step, whereas a' remained. This is only possible if $\langle a'', a' \rangle$ has length $\ge v_k$ and we get a contradiction. \Box

Lemma 2.3. $(\underline{n}, \gamma_1(\underline{n}))$ has the density property for every $\underline{n} \in Q_1$.

The proof of this lemma will be given in Sect. 4.3.

2.2. The Elementary Rearrangement

Starting from a given configuration \underline{n} and a partition $\gamma = \{a_1, \ldots, a_{2r-1}\}$ of it, we will define a new configuration by letting two blocks in (\underline{n}, γ) change place. (\underline{n}, γ) consists of 2r blocks and the elementary rearrangement operation, S, consists of letting the shortest block in (\underline{n}, γ) change place with the shortest of its neighbouring blocks. To be more precise, let $S_{k,k+1}$ denote the operation where block number k changes place with block number k + 1. $S_{k,k+1}$ maps a triple $(\underline{n}, \gamma, \delta)$ to a triple $(\underline{n}', \gamma', \delta')$, where $\underline{n}, \underline{n}'$ are configurations, γ, γ' are partitions and δ, δ' are sets which give the positions of the "old partition points." We define

$$\underline{n}' = \langle a_0, a_1 \rangle \dots \langle a_{k-2}, a_{k-1} \rangle \langle a_k, a_{k+1} \rangle \langle a_{k-1}, a_k \rangle \langle a_{k+1}, a_{k+2} \rangle \dots \langle a_{2r-1}, a_{2r} \rangle , \gamma' = (\gamma \setminus \{a_{k-1}, a_k, a_{k+1}\}) \cup \{a_{k-1} + a_{k+1} - a_k\} , \delta' = (\delta \setminus (\delta \cap [a_{k-1}, a_k))) \cup (\delta \cap [a_{k-1}, a_k) + a_{k+1} - a_k) \cup (\delta \setminus (\delta \cap [a_k, a_{k+1})) \cup (\delta \cap [a_k, a_{k+1}) - (a_k - a_{k-1})) \cup \{a_{k-1}, a_{k+1}\} .$$

We can now define the elementary rearrangement operation S as follows. If $\langle a_{k-1}, a_k \rangle$ is the shortest block in (\underline{n}, γ) , the leftmost if there are several, then $S = S_{k-1,k}$ if $|a_{k-1} - a_{k-2}| \le |a_{k+1} - a_k|$ and $S = S_{k,k+1}$ if $|a_{k-1} - a_{k-2}| > |a_{k+1} - a_k|$. To motivate the definition of γ' and δ' , write $A_k = \langle a_{2(k-1)}, a_{2k-1} \rangle$, $B_k = \langle a_{2k-1}, a_{2k} \rangle$, $k = 1, \ldots, r$, and assume that A_p is the shortest block and $|B_{p-1}| \le |B_p|$. Then the partition points between A_{p-1}, B_{p-1} , between B_{p-1}, A_p and between A_p , B_p are removed from γ and the new partition point between $A_{p-1}A_p$ and $B_{p-1}B_p$ is added to γ , giving γ' . The "old partition points" tell us how $A_{p-1}A_p$ and $B_{p-1}B_p$ should be split up to give A_{p-1}, A_p, B_{p-1} and B_p ; they are added to δ . We also have to move previous "old partition points" lying inside B_{p-1} and A_p . This explains the definition of δ' . Compare Fig. 1.

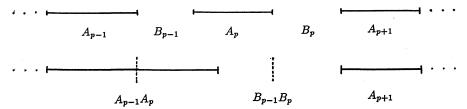


Fig. 1. The dotted lines indicate the locations of the old partition points which are added to δ

2.3. The Local Rearrangement

Starting with Q_1 we will now by successive elementary rearrangements define sets of configurations Q_j , $j = 1, ..., k_N$, and for each $\underline{n} \in Q_j$ a partition $\gamma_j(\underline{n})$. Assume that Q_{j-1} has been defined and that each $\underline{m} \in Q_{j-1}$ has been assigned a partition $\gamma_{j-1}(\underline{m})$. For each $\underline{m} \in Q_{j-1}$ we do successive elementary rearrangements until the shortest block is $\geq u_j$. This will surely happen after a finite number, say l, of elementary rearrangements, where l of course depends on \underline{m} . Let

$$(R(\underline{m}), R\gamma(\underline{m}), R\delta(\underline{m})) = S^{l}(\underline{m}, \gamma_{j-1}(\underline{m}), \emptyset).$$

We define Q_j and $\gamma_j(\underline{n})$ as follows

$$Q_{j} = \{R(\underline{m}); \underline{m} \in Q_{j-1}\},\$$

$$\gamma_{j}(\underline{n}) = \bigcup_{\underline{m} \in Q_{j-1}; R(\underline{m}) = \underline{n}} R\gamma(\underline{m}).$$
(2.1)

Lemma 2.4. $\underline{m} \in Q_{j-1}$, $j \ge 2$, is uniquely determined by $(R(\underline{m}), R\gamma(\underline{m}), R\delta(\underline{m}))$.

Proof. It is clear that the blocks in $R(\underline{m})$ defined by $R\gamma(\underline{m})$ alternate between o- and e-blocks. Let $R\gamma(\underline{m}) = \{a_1, \ldots, a_{2r-1}\}$ and $R\delta(\underline{m}) = \{b_1, \ldots, b_s\}$. Then $A_k = \langle a_{2(k-1)}, a_{2k-1} \rangle$, $k = 1, \ldots, r$, are o-blocks and $B_k = \langle a_{2k-1}, a_{2k} \rangle$, $k = 1, \ldots, r$, are e-blocks. $R\delta(\underline{m})$ tells us how to split these into smaller blocks, e.g. if $a_{2(k-1)} < b_i < b_{i+1} < \cdots < b_{i+p} < a_{2k-1}$, then A_k should be split up into o-blocks $\langle a_{2(k-1)}, b_i \rangle, \ldots, \langle b_{i+p}, a_{2k-1} \rangle$. Similarly the other blocks may be split up into smaller blocks of the same type. The o- and e-blocks thus obtained are the o- and e-blocks in $(\underline{m}, \gamma_{j-1}(\underline{m}))$. These blocks are correctly ordered since in the rearrangement procedure we do not change the order of the o-blocks among themselves or the e-blocks among themselves. We thus know the o- and e-blocks in $(\underline{m}, \gamma_{j-1}(\underline{m}))$ in the right order and this enables us to reconstruct \underline{m} uniquely. \Box

The proof of the following lemma is postponed to Sect. 4.3.

Lemma 2.5. $(\underline{n}, \gamma_j(\underline{n}))$ has the density property for every $\underline{n} \in Q_j$ and all blocks in $(\underline{n}, \gamma_j(\underline{n})), \underline{n} \in Q_j$, have length $\geq u_j$.

3. The Energy-Entropy Argument

3.1. The Entropy Lemma

As a consequence of the definition of the original partition and the rearrangement procedure there is a bound on how much a block can grow due to the rearrangements. This gives an estimate of the maximal distance from an element in $R\delta(\underline{m})$ to the nearest element in $R\gamma(\underline{m})$. More precisely we have the following lemma.

Lemma 3.1. Let $w_{j-1} = 9^{j-2}v_j$. Then for all $\underline{m} \in Q_{j-1}$ the distance from an element in $R\delta(\underline{m})$ to the closest element in $R\gamma(\underline{m})$ is $\leq 8w_{j-1}$, $2 \leq j \leq k_N$.

The proof of Lemma 3.1 is postponed to Sect. 4. We can now estimate the entropy loss due to the rearrangement.

Lemma 3.2. (The entropy lemma). Let $\underline{n} \in Q_j$, $2 \le j \le k_N$ and let $1 \le K \le |\gamma_j(\underline{n})|$, $M \ge 0$ be integers. If $L_{K,M}(\underline{n})$ denotes the number of configurations $\underline{m} \in Q_{j-1}$ such that $R(\underline{m}) = \underline{n}$, $|R\gamma(\underline{m})| = K$ and $|R\delta(\underline{m})| = M$, then

$$L_{K,M} \le e^{C_1 j M} \binom{|\gamma_j(\underline{n})|}{K} \sum_{r=0}^{\min(K,M)} \binom{K}{r}, \qquad (3.1)$$

where $C_1 = \log(72/\delta)$.

Proof. By Lemma 2.2 \underline{m} is uniquely determined if we know $\underline{n} = R(\underline{m}), R\gamma(\underline{m})$ and $R\delta(\underline{m})$. It follows from (2.1) that the number of possible $R\gamma(\underline{m})$ with $|R\gamma(\underline{m})| = K$ is not more than

$$\binom{|\gamma_j(\underline{n})|}{K}.$$

Imagine that each element in $R\delta(\underline{m})$ is assigned to the element in $R\gamma(\underline{m})$ which is closest to it. Suppose that r elements in $R\gamma(\underline{m})$ have been assigned elements from $R\delta(\underline{m}), r \leq K, r \leq M$. There are $\binom{K}{r}$ possibilities for these r elements. $R\delta(\underline{m}) = \{b_1, \ldots, b_M\}, b_1 < b_2 < \cdots < b_M$, is split up into r subsets $\{b_1, \ldots, b_{m_1}\}, \{b_{m_1+1}, \ldots, b_{m_1+m_2}\}, \ldots, \{b_{m_1+\dots+m_{r-1}+1}, \ldots, b_{m_1+\dots+m_r}\}, m_1 + \dots + m_r = M$. This can be done in $\binom{M}{r-1} \leq 2^M$ ways. By $(a)_j$ the distance from an element in $R\delta(\underline{m})$ to the element it has been assigned to is $\leq 8w_{j-1}$. Thus the number of possibilities for the positions of the elements in $R\delta(\underline{m})$, given m_1, \ldots, m_r and the elements in $R\gamma(\underline{m})$ they have been assigned to, is at most

$$\prod_{i=1}^{\prime} \binom{16w_{j-1}}{m_i} \le e^{M \log(16w_{j-1})} \le e^{jM \log(72/\delta)},$$

where we have used that $\binom{n}{k} \le \exp(k \log n)$. Combining all the combinatorial estimates we obtain (3.1).

3.2. The Energy Lemma

We also need an estimate of how much the energy changes when we perform an elementary rearrangement. This is provided by the next lemma which will be proved in Sect. 5.

Lemma 3.3. (*The energy lemma*). Let (\underline{n}, γ) satisfy the density property and assume that all blocks in (\underline{n}, γ) have length $\geq 2^{j-1}$. If \underline{n}' is obtained from \underline{n} by an elementary rearrangement of (\underline{n}, γ) , then

$$H(\underline{n}) - H(\underline{n}') \ge 2\kappa j, \qquad (3.2)$$

where κ is a positive constant that only depends on α .

Remark. When $\alpha < 2$ the right-hand side of (3.2) can be replaced by $\kappa' 2^{(2-\alpha)(j-1)}$.

From the energy lemma it follows that for all $\underline{m} \in Q_{j-1}$, $2 \le j \le k_N$ we have

$$H(\underline{m}) - H(R\underline{m})) \ge \kappa |R\delta(\underline{m})|j.$$
(3.3)

To see this note that the partitions we obtain in the intermediate steps as we go from \underline{m} to $T(\underline{m})$ by successive elementary rearrangements all satisfy the density

property. Thus we can use Lemma 3.2 repeatedly and since the total number of elementary rearrangements we do is $|R\delta(\underline{m})|/2$ and all blocks have length $\geq 2^{j-1}$ by Lemma 2.5 we get (3.3).

We can now specify the constants appearing in Theorem 1.2. Let

$$\delta_0 = \min\{\varrho, 1 - \varrho, 1/4, 2(\alpha - 1)\kappa/5\log 2)\}$$

and

$$\beta_0 = (3 + \log(72/\delta_0))/\kappa.$$

For a given $\beta \ge \beta_0$ we let δ be given by

$$\beta = (3 + \log(72/\delta))/\kappa, \qquad (3.4)$$

i.e. $\delta = K \exp(-\kappa\beta)$, where K is a numerical constant. Note that $\beta \ge \beta_0$ implies $\delta \le \delta_0$.

3.3. Energy-Entropy Argument for the Local Rearrangement

Above we have defined maps $R: Q_{j-1} \to Q_j, j = 2, ..., k_N$ and together they define a map $\mathscr{R}: Q_1 \to Q_{k_N}$. Define

$$H_1 = \{\underline{n} \in Q_1; |\gamma_{k_N}(\mathscr{R}(\underline{n}))| \ge 3\}$$

and $H_j = R(H_{j-1})$. Note that $H_j \subset Q_j$ so we have a partition $\gamma_j(\underline{n})$ for each $\underline{n} \in H_j$. We shall prove that

$$\lim_{A \to Z^+} P(H_1) = 0.$$
(3.5)

According to the next lemma, which will be proved in Sect. 4.2, this proves Theorem 1.2.

Lemma 3.4. Let d_1 , d_2 be defined by (1.6). If $0 \le \tau_1 < \tau_2 \le d_1 - \varepsilon$, $\varepsilon > 0$, then

$$\{\underline{n} \in Q; d(\tau_1, \tau_2) (\underline{n}) < 1 - 2\delta\} \subset H_1$$

and if $d_2 + \varepsilon \leq \tau_1 < \tau_2 \leq 1$, then

$$\{\underline{n} \in Q; d(\tau_1, \tau_2) (\underline{n}) > 2\delta\} \subset H_1$$

for all sufficiently large N, L.

We now turn to the proof of (3.5). The following inequality controls the energy-entropy balance in the local rearrangements. For $1 \le k \le k_N$ we have

$$\sum_{\underline{m}\in H_1} e^{-\beta H(\underline{n})} \le \eta_k \sum_{\underline{m}\in H_k} e^{-\beta H(\underline{m}) + \log(2k)|\gamma_k(\underline{m})|},\tag{3.6}$$

where $\eta_k \leq \eta$ and η is a numerical constant. We will prove (3.6) by induction. Assume that (3.6) is true for k = j - 1. Let $D_{j-1} = \log 2(j - 1)$. Then using Lemma 3.2 and (3.3) we find

$$\sum_{\underline{m}\in H_{j-1}} e^{-\beta H(\underline{m}) + D_{j-1}|\gamma_{j-1}(\underline{m})|} = \sum_{\underline{n}\in H_j} e^{-\beta H(\underline{n})} \sum_{\underline{m}\in H_{j-1}; R(\underline{m}) = \underline{n}} e^{-\beta (H(\underline{m}) - H(R(\underline{m})) + D_{j-1}|\gamma_{j-1}(\underline{m})|}$$
$$\leq \sum_{\underline{n}\in H_j} e^{-\beta H(\underline{n})} \sum_{K=1}^{|\gamma_j(\underline{n})|} \sum_{M=0}^{\infty} e^{-\beta KMj + D_{j-1}(M+K)} L_{K,M}(\underline{n}), \quad (3.7)$$

where we have also used that $|\gamma_{j-1}(\underline{m})| = |R\gamma(\underline{m})| + |R\delta(\underline{m})|$, which follows from the fact that in every elementary rearrangement, $S(\underline{n}, \gamma, \delta) = (\underline{n}', \gamma', \delta')$, $|\gamma'| = |\gamma| - 2$, $|\delta'| = |\delta| + 2$, and hence $|\gamma| + |\delta| = |\gamma'| + |\delta'|$. Using (3.1) we see that (3.7) can be estimated by

$$\sum_{\underline{n}\in H_j} e^{-\beta H(\underline{n})} \sum_{K=1}^{|\gamma_j(\underline{n})|} e^{D_{j-1}K} {|\gamma_j(\underline{n})| \choose K} G(K, \underline{n}), \qquad (3.8)$$

where

$$G(K,\underline{n}) = \sum_{M=0}^{\infty} \sum_{r=0}^{\min(K,M)} \binom{K}{r} e^{-(\beta \kappa - C_1)jM}.$$

According to (3.4) $\beta \kappa - C_1 = 3 \ge 1$. Thus

$$\begin{aligned} G(K,\underline{n}) &= \left(\sum_{M=0}^{K}\sum_{r=0}^{M} + \sum_{M=K+1}^{\infty}\sum_{r=0}^{K}\right) \binom{K}{r} e^{-jM} \\ &\geq \sum_{r=0}^{K}\binom{K}{r}\sum_{M=0}^{K} e^{-jM} + \sum_{M=K}^{\infty} 2^{K} e^{-jM} \\ &\leq \frac{1}{1-e^{-j}} \left[(1+e^{-j})^{K} + 2e^{-jK} \right] \\ &\leq \frac{1+2e^{-j}}{1-e^{-j}} \left(1+e^{-j} \right)^{K}. \end{aligned}$$

Inserting this in (3.8) and calculating the sum over K using the binomial theorem the expression in (3.8) can be estimated by

$$(1+2e^{-j})^2 \sum_{\underline{n}\in H_j} e^{-\beta H(\underline{n})} (1+e^{D_{j-1}}(1+e^{-j}))^{|\gamma_j(\underline{n})|}.$$

Since $\log(1 + e^{D_{j-1}}(1 + e^{-j})) \le \log(2j)$ we have completed the induction step if we put $\eta_j = \eta_{j-1}(1 + 2e^{-j})^2$. We see that

$$\eta_k = \prod_{j=2}^k 1 + 2e^{-j})^2 \le \prod_{j=1}^\infty (1 + 2e^{-j})^2 = \eta,$$

and (3.6) is proved.

3.4. Energy-Entropy Argument for the Global Rearrangement

We are now in position to rigorously carry out the argument sketched in the introduction. In H_{k_N} all blocks have length $\geq 2^{k_N-1} \sim N^{(\log 2)/4}$ and this is sufficient for that argument to work. For each $\underline{m} \in H_{k_N}$ we will perform successive elementary rearrangements starting from $(\underline{m}, \gamma_{k_N}(\underline{m}), \emptyset)$. Every elementary rearrangement reduces the number of elements in the partition with 2. We apply the S-operations repeatedly until we get a configuration $(G(\underline{m}), G\gamma(\underline{m}), G\delta(\underline{m}))$ with only one partition point, i.e. $|G\gamma(\underline{m})| = 1$. Just as in Lemma 2.2 we see that \underline{m} is uniquely determined by $(G(\underline{m}), G\gamma(\underline{m}), G\delta(\underline{m}))$. It is clear by the rearrangement procedure that $|G\gamma(\underline{m})| + |G\delta(\underline{m})| = |\gamma_{k_N}(\underline{m})| \cdot |\gamma_{k_N}(\underline{m})| = 2M - 1$ is odd and, by the

definition of H_1 , $M \ge 2$. Given $G(\underline{m})$ the number of possibilities for $G\gamma(\underline{m})$ and $G\delta(\underline{m})$ is less than

$$L\binom{L}{2M-2} \le e^{(2M-1)\log L}.$$
(3.9)

When we go from <u>m</u> to $G(\underline{m})$ we do $|G\delta(\underline{m})|/2 = M-1$ elementary rearrangements. Repeated application of Lemma 3.2 gives

$$H(\underline{m}) - H(G(\underline{m})) \ge 2\kappa(M-1)k_N \tag{3.10}$$

since all blocks in H_{k_N} have length $\geq 2^{k_N-1}$.

Thus by (3.6) with $k = k_N$, (3.9) and (3.10) we obtain

$$P(H_{1}) \leq \frac{\eta}{Z} \sum_{\underline{n} \in G(H_{k_{N}})} e^{-\beta H(\underline{n})} \sum_{M=2}^{\infty} e^{\log(2k_{N}L)(2M-1)-2\beta\kappa(M-1)k_{N}}$$
$$\leq \eta e^{\log(2k_{N}L)} \sum_{M=1}^{\infty} e^{(2\log(2k_{N}L)-\frac{1}{2}\beta\kappa\log N)M} \leq C \frac{(\log N)^{3}}{N}$$
(3.11)

for all sufficiently large N, L. Here we have used that $\beta \kappa \ge \beta_0 \kappa \ge 8$. It follows from (3.11) that (3.5) holds and we have proved Theorem 1.2.

4. Proof of the Entropy Lemma

4.1. Proof of Lemma 3.1

Call the assertions in Lemma 3.1 $(a)_j$ for a particular value of j. Let

$$w_{j,k} = 9^{j-1} v_k, \quad w_{j-1} = w_{j-1,j}.$$

Together with $(a)_j$ we will prove the following assertion which we call $(b)_j$, $2 \le j \le k_N$.

Consider two o-(e-)blocks in $(\underline{n}, \gamma_j(\underline{n}))$, $\underline{n} \in Q_j$, of length $\geq w_{j,k}$. Then the total length of the e-(o-)blocks between them is $\geq u_k$, $j \leq k \leq 2k_N - j$.

From Lemma 2.2 we know that $(b)_1$ is true. We will prove $(a)_j$ inductively by showing that $(b)_{j-1}$ implies $(a)_j$ and that $(b)_{j-1}$ and $(a)_j$ together imply $(b)_j$.

Assume that $(b)_{j-1}$ is true. $(R(\underline{m}), R\gamma(\underline{m}))$ has been obtained from $(\underline{m}, \gamma_{j-1}(\underline{m}))$ by successive elementary rearrangements. Consider a block A in $(R(\underline{m}), R\gamma(\underline{m}))$. Since there is a symmetry between o- and e-blocks we can assume that A is an o-block. A has been built up from o-blocks A_1, \ldots, A_r in $(\underline{m}, \gamma_{j-1}(\underline{m}))$. If we can show that

$$\min_{1 \le p \le r} \max\{|A_1 \dots A_{p-1}|, |A_{p+1} \dots A_r|\} \le 8w_{j-1}$$
(4.1)

then $(a)_i$ follows. In $(\underline{m}, \gamma_{i-1}(\underline{m}))$ we had *e*-blocks B_1, \ldots, B_{r-1} between A_1, \ldots, A_r ,

$$\underline{m} = \dots B_0 A_1 B_1 A_2 B_2 \dots A_{r-1} B_{r-1} A_r B_r \dots$$

 $R(\underline{m}) = \dots BAB' \dots$, where $B = \dots B_0 \dots B_{s-1}$ and $B' = B_s \dots B_r \dots$ for some s, $1 \le s \le r$. It follows that A_s must have been fixed throughout the rearrangement procedure, since if A_s has been involved in a rearrangement B_{s-1} and B_s would have been joined.

Claim 1. If the elementary rearrangement $\ldots A^1B^1A^2B^2 \ldots$ to $\ldots A^1A^2B^1B^2 \ldots$ is performed at some step, then $|A^2| \leq |A^1|$ and $|B^1| \leq |B^2|$, and one of B^1 and A^2 has length $< u_i$. Here A^1 can be an o- or an e-block.

This follows immediately from the definition of an elementary rearrangement.

Claim 2. Consider an e-block, B^p , to the left (right) of the o-block containing A_s at some step in the rearrangement procedure. Let A^1, \ldots, A^p be the o-blocks built up from A_1, \ldots, A_r to the left (right) of B^p . If $|B^p| \ge u_j$, then $|A^j| < u_j$, $j = 1, \ldots, p$ and $|A^1 \ldots A^p| \le 3w_{j-1}$.

This is seen as follows. We have the configuration

$$\dots A^1 B^1 \dots A^p B^p A^{p+1} \dots, \tag{4.2}$$

where B^1, \ldots, B^p are built up from some (or all) of B_1, \ldots, B_{s-1} . Assume that $|A^q| \ge u_j$, some $q, 1 \le q \le p$. Because B^p eventually ends up in B, at some step an *e*-block containing B^p must change place with an *o*-block containing A^q , but this is impossible according to Claim 1 since they both have length $\ge u_j$. Thus $|A^q| < u_j, 1 \le q \le p$, and (4.2) actually looks like

$$\ldots A_1 B_1 A_2 \ldots A_p B^p A^{p+1} \ldots$$

i.e. $A^i = A_i, i = 1, ..., p$.

If some o-block A_i is joined with A_{i-1} before it is joined with A_{i+1} (or A^{p+1} ...) we get an o-block of length $\geq u_j$ to the left of B^p and a similar contradiction. Thus A_p, \ldots, A_1 are joined successively to the o-block containing A^{p+1} . It follows from Claim 1 that when A_3 is joined with the o-block containing A_4 we must have

$$|B_2| \ge |B_2 \dots B_p|. \tag{4.3}$$

Similarly we must have

$$|B_1| \geq |B_2 \dots B_p|.$$

If $|B_2| \ge w_{j-1}$, then $|B_1| \ge w_{j-1}$ and by $(b)_{j-1}$, $|A_2| \ge u_j$ a contradiction. Hence $|B_2| < w_{j-1}$ and by (4.3)

 $w_{j-1} \geq |B_3 \dots B_p| \geq (p-2)u_{j-1}$,

and consequently

$$|A_1 \dots A_p| \le pu_j \le (w_{j-1}/u_{j-1} + 2)u_j \le 3w_{j-1}$$
.

Claim 3. Let $t \leq s$. If A_{t-1} is joined with A_t before A_{t-2} is joined with A_{t-1} , then $|A_1 \dots A_{t-2}| \leq 3w_{j-1}$. Similarly for $t \geq s$, if A_{t+1} is joined with A_t before A_{t+2} is joined with A_{t+1} , then $|A_{t+2} \dots A_r| \leq 3w_{j-1}$.

This follows immediately from Claim 2, since when A_{t-1} is joined with A_t we get an *e*-block of length $\ge u_j$ either to the left or to the right of the *o*-block containing A_{t-1} . Now A_{t-2} must lie to the left of this *e*-block and hence $|A_1 \dots A_{t-2}| \le 3w_{j-1}$.

We can now prove (4.1) by proving that $\max\{|A_1 \dots A_{s-1}|, |A_{s+1} \dots A_r|\} \le 8w_{j-1}$. Assume first that $|A_s| \ge w_{j-1}$. If $|B_s| \ge u_j$, then $|A_1 \dots A_{s-1}| \le 3w_{j-1}$ by Claim 2. Suppose that $|B_{s-1}| < u_j$. Then $(b)_{j-1}$ implies that $|A_{s-1}| < w_{j-1}$. If A_{s-1} and A_s are joined before A_{s-2} and A_{s-1} , then by Claim 3, $|A_1 \dots A_{s-2}| \le 3w_{j-1}$, and hence $|A_1 \dots A_{s-1}| \le w_{j-1}$. Otherwise, at some step an o-block $A_{s-j_1} \dots A_{s-2}$,

 $j_1 \ge 2$, is joined with A_{s-1} which is not joined with A_s . If it is A_{s-1} that changes place we get an *e*-block of length $\ge u_j$ to the right of A_{s-1} and hence $|A_1 \dots A_{s-1}| \le 3w_{j-1}$. In the other case we get $|A_{s-j_1} \dots A_{s-2}| \le |A_{s-1}| \le w_{j-1}$ and according to Claim 3, $|A_1 \dots A_{s-j_1}| \le w_{j-1}$. Thus $|A_1 \dots A_{s-1}| \le 3w_{j-1}$. By a symmetric argument we get $|A_{s+1} \dots A_r| \le 4w_{j-1}$.

Now assume that $|A_s| < w_{j-1}$ and that A_{s-1} is joined with A_s before A_{s+1} is joined with A_s , the other case is analogous. At some step an *o*-block $A_{j_1} \dots A_{s-1}$ is joined with A_s and by Claim 1, $|A_{j_1} \dots A_{s-1}| \le |A_s| < w_{j-1}$. By Claim 3, $|A_1 \dots A_{j_{1}-1}| \le 3w_{j-1}$ and hence $|A_1 \dots A_{s-1}| \le 4w_{j-1}$. At some later step an *o*-block $A_{s+1} \dots A_{k_1}$ is joined with an *o*-block $A_{j_2} \dots A_s$. It follows from Claim 1 that $|A_{s+1} \dots A_{k-1}| \le |A_{j_2} \dots A_s| \le 5w_{j-1}$. By Claim 3, $|A_{k_1+1} \dots A_{r_l}| \le 3w_{j-1}$ and thus $|A_{s+1} \dots A_r| \le 8w_{j-1}$. This concludes the proof of $(a)_j$.

We now turn to the proof that $(b)_{j-1}$ and $(a)_j$ imply $(b)_j$. Assume that $(b)_j$ is false. Then there is an $\underline{n} \in Q_j$ and two *e*-blocks, say, *B* and *B'* in $(\underline{n}, \gamma_j(\underline{n}))$ of length $\geq w_{j,k}$ such that the length of the *o*-blocks between *B* and *B'* is $\langle u_k, k \geq j$. We must have $k \geq j + 1$ since all blocks in $(\underline{n}, \gamma_j(\underline{n}))$ have length $\geq u_j$. Let A_1, \ldots, A_p be the *o*-blocks between $B = \langle b_0, a_1 \rangle$ and $B' = \langle b_p, a_{p+1} \rangle$. Write $A_i = \langle a_i, b_i \rangle$ and $B_i = \langle b_i, a_{i+1} \rangle$. By (2.1) there is an $\underline{m} \in Q_{j-1}$, $R(\underline{m}) = \underline{n}$ such that $a_1 \in R\gamma(\underline{m})$. It follows from the density property that at least one of the b'_i s must belong to $R\gamma(\underline{m})$, and we let b_q be the largest such b_i . In $(R(\underline{m}), R\gamma(\underline{m}))$ we have two *e*-blocks $C = \langle c, a_1 \rangle$, where $c \leq b_0$ and $C' = \langle b_q, c' \rangle$, where $c' \geq a_{p+1}$. Hence *C* and *C'* have length $\geq w_{j,k}$. Consider an *o*-block *A* in $(R(\underline{m}), R\gamma(\underline{m}))$ between *C* and *C'*. Then $A = \langle a_r, b_s \rangle$ for some *r*, *s*, $1 \leq r \leq s \leq q$. If r = s, then $|A| = |A_r|$. Assume that r < s. Then $A = A_r B_r \dots B_{s-1} A_s$. From the density property for $(R(\underline{m}), R\gamma(\underline{m}))$ we get that the number of occupied positions in *A* is

$$\geq (1 - \delta)|A| \geq (1 - \delta) \left(|A_r| + |B_r| + \dots + |B_{s-1}| + |A_s| \right).$$

On the other hand we get from the density property for $(\underline{n}, \gamma_j(\underline{n}))$ that the number of occupied positions in A is

$$\leq |A_r| + \cdots + |A_s| + \delta \left(|B_r| + \cdots + |B_{s-1}| \right).$$

It follows that

$$|A| \leq \left(1 + \frac{\delta}{1 - 2\delta}\right) (|A_r| + \dots + |A_s|) \,.$$

We have thus shown that the total length of the *o*-blocks in $(R(\underline{m}), R\gamma(\underline{m}))$ between *C* and *C'* is $\leq (1 + \delta(1 - 2\delta)^{-1}) \leq 2$ times the length of the *o*-blocks in $(\underline{n}, \gamma_j(\underline{n}))$ between *B* and *B'*, i.e. $< 2u_k = u_{k+1}$.

We will need the following two inequalities which are easily proved. If $k \ge j+1$, then

$$w_{i,k} > 32w_{i-1}$$
 and $w_{i,k} - 16w_{i-1} > w_{i-1,k+1}$.

It follows from $(a)_j$ that there must have been *e*-blocks C_1 and C'_1 in $(\underline{m}, \gamma_{j-1}(\underline{m}))$ of length $\geq w_{j,k} - 16w_{j-1} > w_{j-1,k+1}$ which are part of *C* respectively *C'*. The *o*blocks in $(R(\underline{m}), R\gamma(\underline{m}))$ between *C* and *C'* must have been built up from *o*-blocks in $(\underline{m}, \gamma_{j-1}(\underline{m}))$ between C_1 and C'_1 because if an *o*-block is moved across C_1 or C'_1 then they must at least double their length. Since $w_{j,k} - 16w_{j-1} > 16w_{j-1}$ this contradicts $(a)_j$. We thus have two *e*-blocks C_1 and C'_1 of length $\geq w_{j-1,k+1}$ in $(\underline{m}, \gamma_{j-1}(\underline{m}))$, such that the total length of the *o*-blocks between them is $< u_{k+1}$. This contradicts $(b)_{j-1}$ if $k+1 \leq 2k_N - (j-1)$, i.e. $k \leq 2k_N - j$. \Box

4.2. Proof of Lemma 3.4

We will need the following lemma, which we will prove in Sect. 4.3.

Lemma 4.1. Let $\gamma_1(\underline{n}) = \{a_1, \ldots, a_r\}, \underline{n} \in Q_1$. If $0 < x < x + z < a_1$ and $N(x, x + z)(\underline{n}) \leq (1 - 2\delta)z$, then $x + L \geq C_{\delta}z^2$, where C_{δ} is a constant that only depends on δ . Similarly, if $x + z > x > a_r$ and $N(x, x + z)(\underline{n}) \geq 2\delta z$, then $x - a_r \geq C_{\delta}z^2$.

Let $\underline{n} \in Q_1 \setminus H_1$ and let $\underline{m} = \mathscr{R}(\underline{n}) \in Q_{k_N}$. Then $\gamma_{k_N}(\underline{m})$ only contains one element which we denote by λ . Define

$$Q_j(\underline{m}) = \{\underline{n} \in Q_j; R^{k_N - j}(\underline{n}) = \underline{m}\}.$$

Then $\underline{n} \in Q_1(\underline{m}), Q_{k_N}(\underline{m}) = \{\underline{m}\}$ and $R(Q_{j-1}(\underline{m})) = Q_j(\underline{m})$. In $(\underline{m}, \gamma_{k_N}(\underline{m})), \langle -\infty, \lambda \rangle$ is an o-block and $\langle \lambda, \infty \rangle$ is an e-block. It follows from the density property that

$$N = N(1, \lambda - 1) (\underline{m}) + N(\lambda, L) (\underline{m}) \le \lambda + \delta(L - \lambda)$$

and

$$N = N(1, \lambda - 1) (\underline{m}) + N(\lambda, L) (\underline{m}) \ge (1 - \delta)\lambda$$

Thus

$$(1-\delta)^{-1}(N/L-\delta) \le \lambda/L \le (1-\delta)^{-1}N/L.$$

Given $\varepsilon > 0$ we see that for sufficiently large N, L,

$$d_1 - \varepsilon/2 \leq \lambda/L \leq d_2 + \varepsilon/2,$$

where d_1 and d_2 are given by (1.6). Define

$$\begin{aligned} \nu_j &= \min_{\underline{n} \in \mathcal{Q}_j(\underline{m})} (\min \gamma_j(\underline{n})), \quad j = 1, \dots, k_N, \\ \mu_j &= \max_{\underline{n} \in \mathcal{Q}_j(\underline{m})} (\max \gamma_j(\underline{n})), \quad j = 1, \dots, k_N. \end{aligned}$$

Let $\underline{n} \in Q_j(\underline{m})$. Then $\min \gamma_j(\underline{n}) = \min R\gamma(\underline{n}')$ for some $\underline{n}' \in Q_{j-1}(\underline{m})$ such that $R(\underline{n}') = \underline{n}$. It is clear from the definition of an elementary rearrangement that $\min R\gamma(\underline{n}') \ge \min \gamma_{j-1}(\underline{n}') \ge v_{j-1}$. Thus $v_j \le v_{j-1}$. There is an $\underline{n}'' \in Q_{j-1}(\underline{m})$ such that $v_{j-1} = \min \gamma_{j-1}(\underline{n}'')$. It follows from Lemma 3.1 that

 $\min R\gamma(\underline{n}'') - \min \gamma_{j-1}(\underline{n}'') \le 8w_{j-1}.$

Furthermore $\min R\gamma(\underline{n}'') \ge \min \gamma_j(R(\underline{n}'')) \ge v_j$ since $R(\underline{n}'') \in Q_j(\underline{m})$. Thus $v_j - v_{j-1} \le 8w_{j-1}$ and we obtain

$$\lambda - v_1 = v_{k_N} - v_1 \le \sum_{j=2}^{k_N} w_{j-1} \le CN^{\gamma},$$
(4.4)

where C is a constant and $\gamma = \frac{1}{4} \log 36 < 1$. Similarly we can show that $\mu_1 - \lambda \leq CN^{\gamma}$. Hence for N, L sufficiently large,

$$d_1 - \varepsilon < \frac{v_1}{L} \le \frac{\mu_1}{L} \le d_2 + \varepsilon.$$

If $0 \le \tau_1 \le \tau_2 \le d_2 - \varepsilon$ and $\underline{n} \in Q_1 \setminus H_1$ then $\langle [\tau_1 L], [\tau_2 L] \rangle$ lies inside an oblock in $(\underline{n}, \gamma_1(\underline{n}))$. It follows from Lemma 4.1 that if $d(\tau_1, \tau_2)(\underline{n}) \le 1 - 2\delta$, then $[\tau_1 L] + L \ge C_{\delta}([\tau_2 L] - [\tau_1 L])^2$, which is impossible if L is sufficiently large. Hence $d(\tau_1, \tau_2)(\underline{n}) > 1 - 2\delta$. Similarly, if $d_2 + \varepsilon \le \tau_1 < \tau_2 \le 1$, then $\langle [\tau_1 L], [\tau_1 L] \rangle$ is part of an *e*-block in $(\underline{n}, \gamma_1(\underline{n}))$ and Lemma 4.1 implies that we must have $d(\tau_1, \tau_2)(\underline{n}) \le 2\delta$. This completes the proof of Lemma 3.4. \Box

4.3. Proof of Lemmas 2.3, 2.5, and 4.1

Proof of Lemma 2.2. The proof of (i) and (ii) in Definition 2.1 for o- and e-blocks are similar and we will only give the proof of (i) for o-blocks. Let $\gamma^{(0)}, \ldots, \gamma^{(2k_N)} = \gamma_1$ be the partitions used in the inductive definition of γ_1 . Suppose that every o-block $A = \langle a, a' \rangle$ in $\gamma^{(k-1)}(\underline{n})$ satisfies $N(a, a+t) \ge 1 + (1 - \delta_{k-1})t$ for $0 \le t \le a' - a$ and some $\delta_{k-1} \ge 0$ independent of A. This is obviously true for k = 1 with $\delta_0 = 0$. Consider an o-block $B = \langle b, b' \rangle$ defined by $\gamma^{(k)}(\underline{n})$. Assume that the length of whole or parts of e- and o-blocks in $(\underline{n}, \gamma^{(k-1)}(\underline{n}))$ between b and b+s < b' is l and s-l respectively. The o-blocks between b and b+s contain at least $1 + (1 - \delta_{k-1})(s-l)$ particles according to our assumption. The construction of $\gamma^{(k)}(\underline{n})$ from $\gamma^{(k-1)}(\underline{n})$ shows that the length of the o-blocks between b and b+s must be at least $v_k l/u_k$. Hence $s - l \ge v_k l/u_k$ which gives $l \le u_k s/v_k$ and $(1 - \delta_{k-1})(s-l) \ge (1 - (\delta_k + u_k/v_k))s$. Thus we can take $\delta_k = \delta_{k-1} + u_k/v_k$ and we have proved (i) since

$$\delta = \sum_{k=1}^{\infty} \frac{u_k}{v_k}.$$

Proof of Lemma 2.5. We first prove that $(\underline{n}, \gamma_j(\underline{n}))$ has the density property for all $\underline{n} \in Q_j$ by induction on j. This is true for j = 1 by Lemma 2.3. Assume that $(\underline{m}, \gamma_{j-1}(\underline{m}))$ has the density property for all $\underline{m} \in Q_{j-1}$. Then clearly $(R(\underline{m}), R\gamma(\underline{m}))$ has the density property for all $\underline{m} \in Q_{j-1}$.

Let $\langle a, a' \rangle$ and $\langle a', a'' \rangle$ be two consecutive blocks in $(\underline{n}, \gamma_j(\underline{n}))$. Assume that $n_a = 1$. By (2.1), $a \in R\gamma(\underline{m}_1)$ and $a' \in R\gamma(\underline{m}_2)$ for some $\underline{m}_1, \underline{m}_2 \in Q_{j-1}, R(\underline{m}_1) = R(\underline{m}_2) = \underline{n}$. Let b be the element closest to the right of a in $R\gamma(\underline{m}_1)$, and b' the element closest to the left of a' in $R\gamma(\underline{m}_2)$. Then $b \ge a'$ and $b' \le a$. Since $n_a = 1$, $\langle a, b \rangle$ is an o-block in $(\underline{n}, R\gamma(\underline{m}_1))$ and thus N(a, a'-1) $(\underline{n}) \ge 1 + (1-\delta) (a'-1-a)$. We see that $\langle b', a' \rangle$ must be an o-block and hence $n_{a'-1} = 1$ and $n_{a'} = 0$. Consequently $\langle a, a' \rangle$ is an o-block and it satisfies (i) and (ii) in Definition 2.1 since $\langle a, b \rangle$ and $\langle b', a' \rangle$ is an o-block and that it satisfies (i) and (ii). Hence o- and e-blocks alternate in $(\underline{n}, \gamma_j(\underline{n}))$ and they satisfy (i) and (ii) in Definition 2.1.

Let $A = \langle a, a' \rangle$ be the shortest block in $(\underline{n}, \gamma_j(\underline{n})), a \in R\gamma(\underline{m}_1), R(\underline{m}_1) = \underline{n}$, where $\underline{m}_1 \in Q_{j-1}$. We take A to be an o-block, the other case is similar. If $|A| < u_j$ then $a' \notin R\gamma(\underline{m}_1)$ since all blocks in $(R(\underline{m}_1), R\gamma(\underline{m}_1))$ have length $\geq u_j$ by definition. To the right of A in $(\underline{n}, \gamma_j(\underline{n}))$ there is an e-block $B = \langle a', a'' \rangle$ of length $\geq |A|$. By (2.1) the next point to the right of a in $R\gamma(\underline{m}_1)$ is $\geq a''$. Hence

$$N(a, a'') \le a' - a + \delta(a'' - a) \le (1/2 + \delta)(a'' - a)$$

But since $\delta \leq 1/4$ when $\beta \geq \beta_0$, $1 - \delta > 1/2 + \delta$ and thus this contradicts the density property of $(\underline{n}, R\gamma(\underline{m}_1))$ which gives $N(a, a'') \geq (1 - \delta)(a'' - a)$. \Box

Proof of Lemma 4.1. We prove the first part of lemma, the second part is analogous. The assertion will be proved inductively for $(\underline{n}, \gamma^{(k)}(\underline{n}))$, $k = 0, \ldots, 2k_N$. The case k = 0 is trivial since N(x, x + s) $(\underline{n}) \leq (1 - 2\delta)s$ is impossible. Assume that the assertion is true in $(\underline{n}, \gamma^{(k-1)}(\underline{n}))$ and let $A = \langle -\infty, a \rangle$ be the first o-block in $(\underline{n}, \gamma^{(k)}(\underline{n}))$. If $\langle -\infty, a \rangle$ is an o-block also in $(\underline{n}, \gamma^{(k-1)}(\underline{n}))$ we are done, otherwise $A = A_1B_1A_2 \ldots B_{-1}A_p$, where A_1, \ldots, A_p are o-blocks and B_1, \ldots, B_{p-1} are eblocks in $(\underline{n}, \gamma^{(k-1)}(\underline{n}))$. It is clear from the construction that A_1 and A_p both have length $\geq v_k$. Let the length of the e-blocks which wholly or partly lie in $\langle x, x+s \rangle$ be y and suppose that $ju_k < y \leq (j+1)u_k$. It follows that $\langle x, x+s \rangle$ contains at least *j* o-blocks of length $\geq v_k$. Then $s - y \geq jv_k$ and N(x, x + s) (<u>n</u>) $\geq (1 - \delta_{k-1}) (s - y)$, where δ_{k-1} is the same as in the proof of Lemma 2.3. We get

$$N(x, x + s) (\underline{n}) \ge (1 - \delta_{k-1}) (1 - y/s)s$$

$$\ge (1 - \delta_{k-1}) (1 - 2u_k/v_k)s > (1 - 2\delta)s,$$

which contradicts our assumption. Hence $y \le u_k$, and from the inequalities $(1-\delta)(s-y) \le N(x, x+s)(\underline{n}) \le (1-2\delta)s$ we find

$$s \leq \delta^{-1}(1-\delta)y \leq \delta^{-1}(1-\delta)u_k$$

If $\langle x, x + s \rangle$ lies inside A_1 the result follows from the induction hypothesis. Otherwise the closest x can be to -L is $L - (s - y) \ge v_k - (s - y)$. Thus

$$\begin{aligned} x+L &\ge v_k - (1-2\delta)\delta^{-1}u_k = 2\delta^{-1}u_k^2 - (1-2\delta)\delta^{-1}u_k \\ &\ge 2u_k^2 \ge 2\left(\frac{\delta}{1-\delta}\right)^2 s^2. \end{aligned}$$

5. Proof of the Energy Lemma

In this section we will prove Lemma 3.3. Suppose that $\ldots A^1 B^1 A^2 B^2 \cdots \rightarrow \ldots A^1 A^2 B^1 B^2 \ldots$ is the elementary rearrangement that has been performed and that A^1 is an o-block. The case when A^1 is an e-block and thus B^2 an o-block is completely analogous by symmetry. We also assume to begin with that $\alpha < 2$. Write $A^1 = \langle a_1, b_1 \rangle$, $B^1 = \langle b_1, a_2 \rangle$, $A^2 = \langle a_2, b_2 \rangle$ and $B^2 = \langle b_2, a_3 \rangle$ and let $x_1 = |A^1|$, $y_1 = |B^1|$, $x_2 = |A^2|$, $y_2 = |B^2|$. The elementary rearrangement above is only performed if $x_1 \ge x_2$ and $y_1 \le y_2$.

Note that we gain energy when two blocks are moved closer to each other and loose energy when they are moved away from each other. If we take all energy losses into consideration but only include the energy gain due to A^1 and A^2 having been moved closer to each other, we get the following estimate

$$H(\underline{n}) - H(\underline{n}') \geq \sum_{\substack{a_1 \leq i < b_1 \\ a_2 \leq j < b_2}} n_i n_j \left(\frac{1}{(j - i - y_1)^{\alpha}} - \frac{1}{(j - i)^{\alpha}} \right) \\ - \sum_{\substack{a_2 \leq i < b_2 \\ j \geq b_2}} n_i n_j \left(\frac{1}{(j - i)^{\alpha}} - \frac{1}{(j - i + y_1)^{\alpha}} \right) - \sum_{\substack{b_1 \leq i < a_2 \\ a_2 \leq j < b_2}} \frac{n_i n_j}{(j - i)^{\alpha}} \\ - \sum_{\substack{a_2 \leq i < b_2 \\ b_2 \leq j < a_3}} n_i n_j \left(\frac{1}{(j - i)^{\alpha}} - \frac{1}{(j - i + y_1)^{\alpha}} \right) \\ - \sum_{\substack{b_1 \leq j < a_2}} n_i n_j \left(\frac{1}{(j - i)^{\alpha}} - \frac{1}{(j - i + x_2)^{\alpha}} \right) \\ = E_1 - E_2 - E_3 - E_4 - E_5.$$
(5.1)

Here E_1 is the energy gain due to A^1 and A^2 moving closer, E_2 the energy loss due to A^2 moving away from everything to the right of B^2 , E_3 the energy loss due to separating B^1 and A^2 , E_4 the energy loss due to A^2 moving away from B^2 and finally E_5 the energy loss due to B^1 moving away from A^1 and everything

to the left of A^1 . Our problem is to estimate E_1 to E_5 . Since n_i is 0 or 1 we have $n_i n_j \ge 1 - (1 - n_i) - (1 - n_j)$. If we use this in E_1 we get an estimate of E_1 from below by three sums $E_1^{(1)} - E_1^{(2)} - E_1^{(3)}$ corresponding to 1, $1 - n_i$, and $1 - n_j$ respectively. If we replace all n_i and n_j in E_2 by 1 we get an upper bound on E_2 . By rewriting the sums we get the following estimate:

$$E_{1}^{(1)} - E_{2} \ge \sum_{i=1}^{x_{1}} \sum_{j=0}^{x_{2}-1} \left(\frac{1}{(j+i)^{\alpha}} - \frac{1}{(j+i+y_{1})^{\alpha}} \right) - \sum_{i=1}^{x_{2}} \sum_{j=0}^{\infty} \left(\frac{1}{(j+i+y_{2})^{\alpha}} - \frac{1}{(j+i+y_{2}+y_{1})^{\alpha}} \right) \ge \sum_{i=1}^{x_{2}} \sum_{j=0}^{x_{2}-1} \left(\frac{1}{(j+i)^{\alpha}} - \frac{1}{(j+i+y_{1})^{\alpha}} \right) - \sum_{i=1}^{x_{2}} \sum_{j=0}^{\infty} \left(\frac{1}{(j+i+y_{1})^{\alpha}} - \frac{1}{(j+i+2y_{1})^{\alpha}} \right),$$
(5.2)

where in the second inequality we have used $x_1 \ge x_2$ and $y_1 \le y_2$. Assume first that $y_1 \ge x_2$. The terms in the first double sum in (5.2) are increasing functions of y_1 , whereas the terms in the second double sum are decreasing functions of y_1 . Hence (5.2) is smallest, for fixed x_2 and $y_1 \ge x_2$, when $y_1 = x_2$. Now assume that $x_1 > y_1$. If we use cancellation between terms in the double sums in (5.2) they can be rewritten as follows:

$$\begin{split} \sum_{i=1}^{x_2} \left(\sum_{j=0}^{y_1-1} \frac{1}{(i+j)^{\alpha}} - \sum_{j=x_2}^{x_2+y_1-1} \frac{1}{(i+j)^{\alpha}} \right) - \sum_{i=1}^{x_2} \sum_{j=0}^{y_1-1} \frac{1}{(i+j+y_1)^{\alpha}} \\ &= \sum_{i=1}^{x_2} \sum_{j=0}^{y_1-1} \left(\frac{1}{(i+j)^{\alpha}} - \frac{1}{(i+j+x_2)^{\alpha}} - \frac{1}{(i+j+x_2)^{\alpha}} \right) \\ &= \sum_{j=1}^{y_1} \left(\sum_{i=0}^{x_2-1} \frac{1}{(i+j)^{\alpha}} - \sum_{i=y_1}^{y_1+x_2-1} \frac{1}{(i+j)^{\alpha}} \right) - \sum_{j=1}^{y_1} \sum_{i=0}^{x_2-1} \frac{1}{(i+j+x_2)^{\alpha}} \\ &= \sum_{j=1}^{y_1} \left(\sum_{i=0}^{y_1-1} \frac{1}{(i+j)^{\alpha}} - \sum_{i=x_2}^{y_1+x_2-1} \frac{1}{(i+j)^{\alpha}} \right) \\ &- \sum_{j=1}^{y_1} \sum_{i=0}^{\infty} \left(\frac{1}{(i+j+x_2)^{\alpha}} - \frac{1}{(i+j+x_2)^{\alpha}} \right) \\ &= \sum_{j=1}^{y_1} \sum_{i=0}^{y_1-1} \left(\frac{1}{(i+j)^{\alpha}} - \frac{1}{(i+j+x_2)^{\alpha}} \right) \\ &- \sum_{j=1}^{y_1} \sum_{i=0}^{\infty} \left(\frac{1}{(i+j)^{\alpha}} - \frac{1}{(i+j+x_2)^{\alpha}} \right) \\ &- \sum_{j=1}^{y_1} \sum_{i=0}^{\infty} \left(\frac{1}{(i+j+x_2)^{\alpha}} - \frac{1}{(i+j+x_2)^{\alpha}} \right) . \end{split}$$
(5.3)

The same argument as above shows that if $x_2 \ge y_1$, y_1 fixed, the last expression in (5.3) is smallest when $x_2 = y_1$. If we put $x = \min(x_2, y_1)$ we obtain the estimate

$$E_1^{(1)} - E_2 \ge \sum_{i=1}^{x} \sum_{j=0}^{x-1} \left(\frac{1}{(i+j)^{\alpha}} - \frac{2}{(i+j+x)^{\alpha}} \right).$$
(5.4)

Split this sum into a sum over *i*, *j* such that $i + j \le x$ and a sum where i + j > x. If k = i + j we find

$$E_{1}^{(1)} - E_{2} \ge \sum_{k=1}^{x} k \left(\frac{1}{k^{\alpha}} - \frac{2}{(k+x)^{\alpha}} \right) + \sum_{k=1}^{x-1} k \left(\frac{1}{(2x-k)^{\alpha}} - \frac{2}{(3x-k)^{\alpha}} \right)$$
$$= \frac{1}{x^{\alpha}} \left(1 - 2^{1-\alpha} \right) + \sum_{k=1}^{x-1} \frac{1}{k^{\alpha-1}} \left(1 - f\left(\frac{x}{k};\alpha\right) \right), \tag{5.5}$$

where $f(t;\alpha) = 2(1+t)^{-\alpha} + 2(3t-1)^{-\alpha} - (2t-1)^{-\alpha}$. If x = 1 the sum in (5.5) is zero. At the end of this section we will show that

$$\sup_{t \ge 1} f(t; \alpha) = 1 - d_{\alpha} < 1$$
(5.6)

Define $g(x) = 1 + (2 - \alpha)^{-1}(x^{2-\alpha} - 1)$. If we use (5.6) in (5.5) we get

$$E_1^{(1)} - E_2 \ge \min\{1 - 2^{1-\alpha}, d_\alpha\} \sum_{k=1}^{x-1} \frac{1}{k^{\alpha-1}} \ge d'_{\alpha}g(x).$$
 (5.7)

We will now show how E_5 and $E_1^{(2)}$ are estimated. E_3 , E_4 and $E_1^{(3)}$ are estimated similarly. By the density property

$$\sum_{j=b_1}^{b_1+k} n_j \le \delta k, \quad 0 \le k \le y_1 - 1.$$

If we use $n_i \leq 1$, summation by parts and estimation of sums by integrals we get

$$E_{5} \leq \sum_{i < b_{1}} \sum_{j=b_{1}}^{a_{2}-1} n_{j} \left(\frac{1}{(j+i)^{\alpha}} - \frac{1}{(j+i+x_{2})^{\alpha}} \right)$$

$$\leq \delta \sum_{i=1}^{\infty} \sum_{j=1}^{y_{1}-1} \left(\frac{1}{(i+j)^{\alpha}} - \frac{1}{(i+j+x_{2})^{\alpha}} \right)$$

$$\leq \delta \sum_{i=1}^{x_{2}} \sum_{j=1}^{y_{1}} \frac{1}{(i+j)^{\alpha}} \leq \frac{\delta}{\alpha-1} g(x),$$

where $x = \min(x_2, y_1)$. The density property gives

$$\sum_{i=1}^{k} 1 - n_{b_1-1} \le \delta(k-1), \quad k = 1, \dots, x_1.$$

Using this estimate in a summation by parts we find

$$\begin{split} E_1^{(2)} &= \sum_{i=a_1}^{b_1-1} \sum_{j=a_2}^{b_2-1} (1-n_i) \left(\frac{1}{(j-i+y_1)^{\alpha}} - \frac{1}{(j-i)^{\alpha}} \right) \\ &= \sum_{i=1}^{x_1} \sum_{j=0}^{x_2-1} (1-n_{b_1-i}) \left(\frac{1}{(i+j)^{\alpha}} - \frac{1}{(i+j+y_1)^{\alpha}} \right) \\ &= \delta \sum_{i=2}^{\infty} \sum_{j=0}^{x_2-1} \left(\frac{1}{(i+j)^{\alpha}} - \frac{1}{(i+j+y_1)^{\alpha}} \right) \\ &= \delta \sum_{i=1}^{y_1} \sum_{j=1}^{x_2} \frac{1}{(i+j)^{\alpha}} = \frac{\delta}{\alpha-1} g(x) \,. \end{split}$$

A similar argument gives the same estimate for $E_1^{(3)}$, E_3 and E_4 , and together with (5.7) this gives

$$H(\underline{n}) - H(\underline{n}') \ge \left(d'_{\alpha} - \frac{5\delta}{\alpha - 1}\right).$$
(5.8)

Note that $g(x) \ge 1 + \log x$ and that $g(x) \to 1 + \log x$ as $\alpha \to 2-$. Put $d'_{\alpha} = 4\kappa/\log 2$ and assume that $\delta \le \frac{1}{5}(\alpha - 1)2\kappa/\log 2$. Since $x \ge 2^{j-1}$, (5.8) gives that $H(\underline{n}) - H(\underline{n}') \ge 2\kappa j$ which is what we wanted to prove.

We have proved Lemma 3.3 except for (5.6), the proof of which we outline below. We have

$$\frac{\partial f}{\partial \alpha} = -2 \frac{\log(1+t)}{(t+1)^{\alpha}} - 2 \frac{\log(3t-1)}{(3t-1)^{\alpha}} + \frac{\log(2t-1)}{(2t-1)^{\alpha}}$$

With some work the right inequality in

$$\frac{\log(2t-1)}{(2t-1)^{\alpha}} \le 2\frac{\log(1+t)}{(1+t)^{\alpha}} \Leftrightarrow \frac{\log(2t-1)}{\log(1+t)} \le 2\left(\frac{2t-1}{t+1}\right)^{\alpha}$$

can be poved for $t \ge 1$ and consequently $\partial f/\partial \alpha < 0$ if $t \ge 1$. When t > 1 it is not difficult to see that $f(t; \alpha) < 1$. When $1 < \alpha \le 2$, $f(1; \alpha) \ge 0$ and $(\partial f/\partial \alpha)(1; \alpha) > 0$. Hence $f(x; \alpha)$, for fixed α , attains its maximum at some point $t_{\alpha} > 1$; t_{α} only depends on α . Thus

$$f(t;\alpha) \le f(t_{\alpha};\alpha) < f(t_{\alpha};1) < 1$$

for all $t \ge 1$. If we put $1 - d_{\alpha} = f(t_{\alpha}; 1)$ we are done. \Box

6. Proof of Theorem 1.3

Fix a β such that the free energy $f(\rho, \beta)$ defined by (1.7) is a strictly convex function of ρ and fix N, L. Let $\underline{n} \in Q$ be a configuration and let $x_{\mu}(\underline{n})$ denote the μ :th occupied position in Λ , i.e. the position of the μ :th particle, $\mu = 1, ..., N$. For a given small $\varepsilon > 0$ we define

$$I_{\mu} = \{ x \in \Lambda; |x/L - \mu/N| > \varepsilon \}$$

$$M_{\varepsilon,\mu} = \{\underline{n} \in Q; x_{\mu}(\underline{n}) \in I_{\mu}\}.$$

K. Johansson

Let v = v(N, L) be such that

$$\sum_{\underline{n}\in M_{\varepsilon,\mu}}e^{-\beta H(\underline{n})}$$

is maximal when $\mu = v$ and let $\xi = \xi(N, L)$ be such that

$$\sum_{\underline{n}\in M_{\varepsilon,\mu};x_{\nu}(\underline{n})=x}e^{-\beta H(\underline{n})}$$

is maximal when $x = \xi$. Note that $\xi \in I_{\nu}$. If we define

$$H_1(\underline{n}) = -\sum_{\substack{i,j \in [1,\xi) \cap A \\ i \neq j}} \frac{n_i n_j}{|i-j|^{\alpha}}, \qquad H_2(\underline{n}) = -\sum_{\substack{i,j \in (\xi,L) \cap A \\ i \neq j}} \frac{n_i n_j}{|i-j|^{\alpha}},$$

then $H(\underline{n}) - H_1(\underline{n}) - H_2(\underline{n}) \ge o(N)$, where o(N) is positive and does not depend on \underline{n} . Let $M_{\varepsilon} = \bigcup_{\mu=1}^{N} M_{\varepsilon,\mu}$. Then

$$\sum_{\underline{n}\in M_{\varepsilon}} e^{-\beta H(\underline{n})} \leq \sum_{\mu=1}^{N} \sum_{\underline{n}\in M_{\varepsilon,\mu}} e^{-\beta H(\underline{n})}$$

$$\leq N \sum_{x\in I_{\nu}} \left(\sum_{\underline{n}\in M_{\varepsilon,\mu}; x_{\nu}(\underline{n})=x} e^{-\beta H(\underline{n})}\right)$$

$$\leq NL \sum_{\underline{n}\in M_{\varepsilon,\mu}; x_{\nu}(\underline{n})=\xi} e^{-\beta H(\underline{n})} \leq NL \sum_{\underline{n}; x_{\nu}(\underline{n})=\xi} e^{\beta H(\underline{n})}$$

$$\leq e^{o(N)} Z(\nu-1, \xi-1) Z(N-\nu, L-\xi). \tag{6.1}$$

We want to show that

$$\limsup_{\Lambda \to Z^+} P(M_{\varepsilon})^{1/L} \le c < 1,$$
(6.2)

where c is a constant (that depends on ε). Choose a subsequence $N_p \to \infty$, $L_p \to \infty$ and $N_p/L_p \to \varrho$ as $p \to \infty$ which gives the upper limit in (6.2) and such that

$$t = \lim_{p \to \infty} \frac{v(N_p, L_p)}{N_p}$$
 and $s = \lim_{p \to \infty} \frac{\xi(N_p, L_p)}{L_p}$

exist. Now, using (5.1) and the definition of free energy we see that

$$\limsup_{\Lambda \to Z^+} P(M_{\varepsilon})^{1/L} \le \exp \beta \left[f(\varrho, \beta) - sf\left(\frac{t}{s}\varrho, \beta\right) - (1-s)f\left(\frac{1-t}{1-s}\varrho, \beta\right) \right].$$

Observe that $\xi \in I_{\nu}$ implies $|\xi/L - \nu/N| \ge \varepsilon$ and consequently $|s - t| \ge \varepsilon$. Since $\xi \ge \nu$ and $N - \nu \le L - \xi$ we must have 0 < s < 1. Thus $\varrho t/s \ne \varrho(1 - t)/(1 - s)$ and the strict convexity of f gives

$$f(\varrho,\beta) - sf\left(\frac{t}{s}\,\varrho,\beta\right) - (1-s)f\left(\frac{1-t}{1-s}\,\varrho,\beta\right) < 0,$$

and we have proved (6.2). It follows from (6.2) that

$$\limsup_{\Lambda \to Z^+} P(M_{\varepsilon}) = 0 \tag{6.3}$$

60

for every fixed $\varepsilon > 0$. We will prove that

$$Q \setminus M_{\varepsilon'} \subset \{ \underline{n} \in Q; \varrho - \varepsilon \le d(\tau_1, \tau_2) (\underline{n}) \le \varrho + \varepsilon \},$$
(6.4)

where $\varepsilon' = \varepsilon/3\varrho$. Equations (6.3) and (6.4) together prove Theorem 1.3.

If $\underline{n} \in Q \setminus M_{\varepsilon'}$, then $|x_{\mu}(\underline{n})/L - \mu/N| \leq \varepsilon'$ for $\mu = 1, ..., N$. Thus for each τ , $0 < \tau < 1$, we have $\mu > (\tau + \varepsilon')N$, which gives $x_{\mu}(\underline{n}) > \tau L$, and hence $N(0, \tau L)(\underline{n}) \leq \mu - 1$. Consequently $N(0, \tau L) \leq (\tau + \varepsilon')N$. Similarly $\mu \leq (\tau - \varepsilon')N$ implies $N(0, \tau L)(\underline{n}) \geq [(\tau - \varepsilon')N]$. These two estimates of $N(0, \tau L)(\underline{n})$ give

$$\varrho(1-3\varepsilon') \le d(\tau_1,\tau_2) (\underline{n}) \le \varrho(1+3\varepsilon'),$$

and we have proved (6.4). \Box

References

- 1. Aizenman, M., Chayes, J.T., Chayes, L., Newman, C.M.: J. Stat. Phys. 50, 1 (1988)
- 2. Dyson, F.J.: Commun. Math. Phys. 12, 212 (1969) and 21, 269 (1971)
- 3. Ellis, R.S.: Entropy, large deviations and statistical mechanics. Berlin, Heidelberg, New York: Springer 1985
- 4. Fröhlich, J., Spencer, T.: Commun. Math. Phys. 84, 87-101 (1982)
- Johansson, K.: Separation of phases at low temperatures in a one-dimensional continuous gas. Commun. Math. Phys. 141, 259-278 (1991)
- 6. Landau, L.D., Lifshitz, E.M.: Statistical physics. London: Pergamon Press 1958
- 7. Rockafellar, R.T.: Convex analysis. Princeton, NJ: Princeton University Press 1970
- 8. Ruelle, D.: Statistical mechanics. New York: Benjamin 1969
- 9. Thouless, D.J.: Phys. Rev. 187, 732-733 (1969)

Communicated by M. Aizenman