

Embeddings of $U(1)$ -Current Algebras in Non-Commutative Algebras of Classical Statistical Mechanics*

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Abstract. Motivated by multiplicative K -homology, and understanding critical phenomena in some classical statistical mechanical models, we construct actions of $GL(\infty)$ on the operator algebras of V. Jones and Ocneanu, and analyse these in terms of embeddings of $U(1)$ -current algebras.

1. Introduction

In this paper we shall describe a construction of actions of the group $GL(\infty)$ on the V. Jones and Ocneanu algebras, to which both authors of this paper arrived at independently and for different reasons.

The first author was motivated by the search for a multiplicative analogue of the known additive K -homology theory. In this the Jones' index of subfactors would be the multiplicative analogue of the codimension of a subspace of a Hilbert space. This analogy would be in the same way as the classification of outer conjugacy classes of automorphisms of the hyperfinite factor of type II_1 is the multiplicative analogue of the classification of unitaries modulo the compact operators. The basic notion of a Fredholm module (\mathfrak{h}, D) on an algebra \mathfrak{A} gets replaced by a pair (A, σ_t) where A is a C^* -algebra, σ_t a one parameter group of automorphisms, while the unitary group $\mathfrak{U}(\mathfrak{A})$ acts on A by automorphisms. It is easy to get such a multiplicative module out of a Fredholm module (\mathfrak{h}, D) by the CAR construction, and the possibility of replacing the CAR algebra by the algebras of V. Jones and Ocneanu served as one motivation for the search of an action of $GL(\infty)$.

The second author was motivated by statistical mechanics, in particular by the transfer matrix method, and the problem of understanding critical phenomena and the continuum limit in classical models such as those of Potts and of Andrews et al.

The transfer matrix method allows us to reduce a two-dimensional classical or commutative statistical mechanical model to a one-dimensional quantum or non-

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commutative model. Subsequently, taking a continuum limit one obtains an Euclidean field theory, equipped with representations of the Virasoro algebra in the case of conformal invariant field theory. This framework is now well understood in the case of the nearest neighbour two dimensional Ising model [P, AE, EL, A, CE, K]. This paper is the first part of an investigation into understanding critical phenomena in other classical models, such as those of Potts [B] and Andrews et al. [ABF], and representations of the Virasoro algebra at the level of some one dimensional quantum model (before the scaling limit takes place) set in the operator algebras of Temperley and Lieb [TL], Jones [J], and Ocneanu [O]. Such algebras can be regarded as deformations of the Clifford algebra. This investigation begins here with a study of embeddings of $U(1)$ -current algebras in the one-dimensional quantum systems.

Consider the following two dimensional classical models on the square lattice \mathbb{Z}^2 :

(a) The Potts model [B]. At each site we have q -possible states, so that a configuration of the system is an element of $\{1, 2, \dots, q\}^{\mathbb{Z}^2}$. The Hamiltonian for the (standard) Potts model is $\mathcal{H}(\sigma) = - \sum_{\alpha, \beta} J \delta(\sigma_\alpha, \sigma_\beta)$, where the summation is over nearest neighbours.

(b) The Andrews-Baxter-Forrester model [ABF]. Again we have ℓ possible states at each site, but we only consider configurations $\sigma \in \{1, 2, \dots, \ell\}^{\mathbb{Z}^2}$ such that $|\sigma_\alpha - \sigma_\beta| = 1$ for nearest neighbours α, β . The Hamiltonian is obtained by considering interactions around faces.

In the algebraic approach, we reduce the problem set classically in $C(P)$, the continuous functions on a configuration space P to a non-commutative C^* -algebra A [P, AE, CE, K]. The programme is, given a classical equilibrium state μ at inverse temperature β , to find a linear map ϕ_μ on the C^* -algebra A , and a map $F \rightarrow F_\beta$ from local observables in $C(P)$ to A , such that a classical correlation function can be computed using a knowledge of the quantum system alone:

$$\phi(F) = \phi_\mu(F_\beta). \quad (1.1)$$

One expects that the map $F \rightarrow F_\beta$ depends only on the inverse temperature β , and that positivity of ϕ_μ is related to reflection positivity of μ [K].

Take for example the one dimensional 2-state Potts model [B]. This can be described using the algebra M_q of complex $q \times q$ matrices. The Hamiltonian for the Potts model with periodic boundary conditions is

$$\mathcal{H}_L(\sigma) = - \sum_{i=-L}^L J \delta(\sigma_i, \sigma_{i+1}). \quad (1.2)$$

Then the partition function is

$$Z_L = \sum_{\sigma} \exp -\beta \mathcal{H}(\sigma) = \sum_{\sigma} \exp(\beta J \delta(\sigma_i, \sigma_{i+1})) = \text{tr } T^{2L+1}, \quad (1.3)$$

if $T_{\sigma, \sigma'} = \exp(\beta J \delta(\sigma, \sigma'))$, a $q \times q$ matrix when $\sigma, \sigma' \in \{1, 2, \dots, q\}$. Similarly

$$\langle F \rangle_{\beta}^L = \sum_{\sigma} F(\sigma) \exp -\beta \mathcal{H}(\sigma) / Z_L = \text{tr } F_{\beta} T^{2L+1} / \text{tr } T^{2L+1}. \quad (1.4)$$

Consequently, letting $L \rightarrow \infty$, we have

$$\langle F \rangle_{\beta} = \langle F_{\beta} \Omega, \Omega \rangle, \quad (1.5)$$

where $\Omega = q^{-1/2}(1, \dots, 1)$ is the largest eigenvector of the transfer matrix $T = (e^K - 1)e^{K*f}$, if $(e^K - 1)(e^{K*} - 1) = q$, and f is the rank one projection on Ω .

One proceeds similarly in 2-dimensions by first taking a finite box $[-M, M] \times [-L, L]$ and pushing each column $\{i\} \times [-L, L]$ onto one site $(i, 0)$ and picking up one copy of M_q for each site $(i, 0)$ in the row $[-M, M] \times \{0\}$. Thus

$$Z_{LM} = \text{tr } T^{2L+1}, \langle F \rangle_\beta^{LM} = \text{tr } F_\beta^M T^{2L+1} / \text{tr } T^{2L+1}, \quad (1.6)$$

where now $T_M, F_\beta^M \in \bigotimes_{-M}^M M_q$, and one can proceed to take the thermodynamic limit as before. The transfer matrix T can be expressed, up to a scalar, as $T = V^{1/2} W V^{1/2}$. Here V comes from vertical interactions so that

$$V = \exp K_1^* \sum_i e_{2i},$$

and

$$e_{2i} = 1 \otimes \dots \otimes f \otimes 1 \dots \otimes 1_{i^{\text{th}} \text{ position}}. \quad (1.7)$$

We have taken different interaction constants J_1, J_2 in the vertical and horizontal directions with $K_i = \beta J_i$. Then W arising from horizontal interactions is diagonal, and

$$W = \exp K_2 \sum_i e_{2i+1}, e_{2i+1} = 1 \otimes \dots \otimes_{i, i+1} g \otimes 1 \dots \otimes 1, \quad (1.8)$$

where $g = \sum_{j=1}^q g_j \otimes g_j \in M_q \otimes M_q$, and g_j is the rank one projection on $\Omega_j = (\delta_{ij})_{i=1}^q$.

The family $\{e_i\}_{i=-\infty}^\infty$ of projections arising from the Potts model transfer matrix satisfy the relations $[J, B]$:

$$e_i e_j = e_j e_i, |i - j| \geq 2, \quad (1.9)$$

$$e_i e_{i \pm 1} e_i = \tau e_i, \quad (1.10)$$

where $\tau^{-1} = q$. Similarly, one can describe the transfer matrix of the ABF model (at least at criticality) with the aid of projections satisfying similar relations with $\tau^{-1} = 4 \cos^2 \pi / (\ell + 1)$, [KAW, P]. We note from Jones' classification of index of subfactors [J], that an infinite family $\{e_i\}$ of projections exists satisfying (1.9) and (1.10) if and only if $\tau^{-1} \in [4, \infty) \cup \{4 \cos^2 \pi / (\ell + 1) : \ell = 3, 4, \dots\}$. Observe that $2 = 4 \cos^2(\pi/4)$, $3 = 4 \cos^2(\pi/6)$. One can indeed generalise the ABF models [P] by taking any finite connected graph Γ , with vertices $\Gamma^{(0)}$ and edges $\Gamma^{(1)}$. Then consider the space of configurations $\sigma \in (\Gamma^{(0)})^{\mathbb{Z}^2}$ subject to the constraint that $(\sigma_\alpha, \sigma_\beta) \in \Gamma^{(1)}$ for all pairs α, β of nearest neighbours in \mathbb{Z}^2 . Then in this model, one gets a representation of the Jones' relations where $\tau^{-1} = \alpha^2$, and α is the largest eigenvalue of a certain incidence matrix associated with the graph Γ .

From a lattice model, one gets a field theory by taking a continuum limit; letting the lattice spacing $\varepsilon \rightarrow 0$, whilst simultaneously letting the temperature approach the critical temperature, in such a way that $\varepsilon \cdot (\text{correlation length})$ is constant.

This is well understood in the Ising model [SMJ, PT1, PT2, SO'C]. Belavin et al. [BPZ] suggested using conformal invariance to understand scale invariance of special critical points. Conformal invariance is understood via representations of the Virasoro algebras. A knowledge of which representations can appear tells us something about the field theory and the nature of the critical points. The Virasoro algebra is the unique central extension of the Witt algebra of vector fields on S^1 , generated by L_m , $m \in \mathbb{Z}$ and a central element c with

$$[L_m, L_n] = (m - n)L_{m+n} + \frac{c}{12}m(m^2 - 1)\delta_{m, -n}. \quad (1.11)$$

Unitarity of the field theory, (coming from reflection positivity of the statistical mechanical model) means $L_n^* = L_{-n}$. From physical considerations, one is also interested in highest weight representations, where the energy of L_0 is bounded below. There is a vector $|h\rangle$, the highest weight vector, a scalar h such that $L_0|h\rangle = h|h\rangle$, $L_n|h\rangle = 0$, $n > 0$. In an irreducible representation, c is a scalar, and so an irreducible highest weight unitary representation is determined by a pair (c, h) of real numbers. It is known that such representations exist if and only if either [FQS, GKO]:

$$c \geq 1, \quad h \geq 0, \quad (1.12)$$

$$\begin{aligned} c &= 1 - 6/m(m+1), \quad h = h_{p,q}(c) = [((m+1)p - mq)^2 - 1]/4m(m+1), \\ m &= 2, 3, 4, \dots, p = 1, 2, \dots, m-1, \quad q = 1, 2, \dots, p. \end{aligned} \quad (1.13)$$

The question then arises of which representations correspond to certain statistical mechanical models. The possible values of h are related to the scaling dimensions or critical exponents. In this way, comparing with known exponents, one can match [FQS, H] (1.13) $m = 3$ with the Ising model, $m = 4$ with the tricritical Ising model, $m = 5$ with the 3-state Potts model, and $m \geq 5$ with the ABF models. On the other hand we have noted how the Potts model and ABF models are related to the Jones' classification via the transfer matrix method. It is thus natural to look for a direct link between the FQS classification and that of Jones. In this present paper we construct $U(1)$ -currents directly from the Jones algebra, from which, using the Sugawara formula one can construct representations of the Virasoro algebra with central charge $c = 1$.

The operator algebras of Jones and Ocneanu provide a natural setting for deforming the Fermi-Dirac quantization of the classical fields on S^1 . The Jones algebra $A(\tau)$ is the C^* -algebra generated by projections $\{e_i\}_{i \in \mathbb{Z}}$ satisfying relations (1.9) and (1.10). The Ocneanu algebra (see Sect. 3) $A(\Gamma)$ is a slightly larger algebra associated to a connected graph Γ such that the largest eigenvalue of its associated incidence matrix is $\tau^{-1/2}$. In general $A(\Gamma)$ is generated by $A(\tau)$ and a finite dimensional C^* -algebra. In Sect. 2 we show how to embed a copy of an observable algebra in an operator algebra, simply by constructing certain representations of the Lie algebra $gl(\infty)$ from operators satisfying infinitesimal versions of the Jones' relations. This allows us in Sects. 3 and 4 to embed the observable algebra in the algebras of Jones and Ocneanu, with the Ocneanu algebra permitting a rich choice of embeddings. The Hamiltonian $L_0 = \sum ne_{2n}$ generates time evolution, and has a unique KMS state ω_β for each temperature β (if $\tau^{-1} < 4$), and a canonical ground

state ω_β (Sect. 5). The restriction of ω_β to any observable algebra constructed in Sect. 4 is quasi-free. In the state ω_∞ we can construct (Sect. 6) $U(1)$ -currents T_n satisfying $[T_m, T_n] = n\delta_{m, -n}$, to which one can apply the Sugawara formula.

What we have achieved is to construct representations of the Virasoro algebra in the setting of the algebras of V. Jones and Ocneanu, with central charge $c = 1$. It remains to construct representations of the Virasoro algebra with central charge $c = 1 - 6/m(m+1)$ in the setting of the algebras associated with index $\tau^{-1} = 4 \cos^2 \pi / (m+1)$. However, we have strong evidence from studying the characters of the discrete series of representations of the Virasoro algebra that this can be done. It also remains to understand how the families of $U(1)$ -currents which we have constructed, associated to a single graph Γ , fit together.

2. Representations of $gl(\infty)$ and of the Observable Algebra

If I is a countable set, let $gl(I)$ denote the Lie algebra of $|I| \times |I|$ matrices over \mathbb{C} , with only finitely many non-zero entries, generated by the elementary matrices $\{E_{ij} : i, j \in I\}$. The first lemma shows the equivalence between those representations of $gl(I)$, where E_{ii} are projections, with representations of the observable algebra over the Hilbert space $\ell^2(I)$.

Lemma 2.1. *Let I be a countable set, \mathcal{B} a unital C^* -algebra, and for each $i, j \in I$, let e_{ij} be elements of \mathcal{B} . Then the following pair of conditions are equivalent:*

- (2.1) (a) $e_{ii} = e_{ii}^2$,
 (b) $e_{ij}^* = e_{ji}$,
 (c) $[e_{ij}, e_{jl}] = e_{ii} - e_{jj}$,
 (d) $[e_{ij}, e_{jk}] = e_{ik}$ if $j \neq k$,
 (e) $[e_{ij}, e_{kl}] = 0$ if $j \neq k, l \neq i$.
- (2.2) *There is an unique unital $*$ -homomorphism π from the observable algebra over $\ell^2(I)$ into \mathcal{B} such that $\pi(a_i^* a_j) = e_{ij}$.*

Proof. (2.2) \Rightarrow (2.1) is clear as $E_{ij} \rightarrow a_i^* a_j$ extends to a Lie algebra representation of $gl(I)$.

(2.1) \Rightarrow (2.2). Suppose I is ordered, say an interval in \mathbb{Z} .

Let $f_i = e_{ii}$, $u_i = e_{i, i+1}$. From (2.1) (d) we have $f_i u_i - u_i f_i = u_i$ (*). Multiplying (*) on the left by f_i , we see $f_i u_i f_i = 0$, and then multiplying (*) on the right by f_i we see $u_i f_i = 0$, and so $f_i u_i = u_i$. Similarly, $u_i^* f_{i+1} = 0$, $u_i^* f_i = u_i^*$. Hence from (2.1) (c), $[u_i, u_i^*] = f_i - f_{i+1}$, we see $u_i^* u_i = f_{i+1}(1 - f_i)$, $u_i u_i^* = f_i(1 - f_{i+1})$. Then Lemma 2.1 will follow from:

Lemma 2.2. *Let I be a countable set, say an interval in \mathbb{Z} . For each $i \in I$, let f_i, u_i be elements of a unital C^* -algebra \mathcal{B} . Then the following three sets of conditions are equivalent:*

- (2.3) *Each f_i is a projection and the map $E_{ii} \rightarrow f_i$, $E_{i, i+1} \rightarrow u_i$ extends to a Lie algebra homomorphism from $gl(I)$ into \mathcal{B} .*
- (2.4) *There is an unique unital $*$ -homomorphism π from the observable algebra over $\ell^2(I)$ into \mathcal{B} such that $\pi(a_i^* a_i) = f_i$, $\pi(a_i^* a_{i+1}) = u_i$.*

- (a) f_i is a projection, (2.5)
 (b) $[f_i, f_j] = 0$,
 (c) $[u_i, f_j] = 0$ if $j \neq i, i+1$,
 (d) $[u_i, u_j^*] = 0$ if $|i-j| \geq 2$,
 (e) $u_i^* u_i = f_{i+1}(1-f_i)$, $u_i u_i^* = f_i(1-f_{i+1})$.

Proof. (2.5) \Rightarrow (2.4). Let $K = \{-1, 1\}^I$ with the equivalence relation \sim on K given by $x \sim y$ if $x = (x(i))$, $y = (y(i)) \in K$, $x(i) = y(i)$ except for finitely many coordinates, and $\sum_{i \in I} (x(i) - y(i)) = 0$. We identify the gauge invariant CAR algebra on $\ell^2(I)$ with the C^* -algebra $C^*(G)$ of the groupoid of the equivalence relation G of \sim [R]. In particular $C^*(G)$ contains $C(K)$, and so g_i , the projection given by the characteristic function of $\{x \in K : x(i) = 1\}$. Let v_i be the partial isometry with initial support $g_{i+1}(1-g_i)$ and final support $g_i(1-g_{i+1})$, which replaces $x(i) = -1$, $x(i+1) = 1$ with $x(i) = 1$, $x(i+1) = -1$, which does not affect $\sum x(i)$ and so v_i belongs to $C^*(G)$. The v_i generate the equivalence relation, since given any sequence $(-1, 1, -1, \dots, 1)$ of ± 1 , one can always put the -1 's on the left by replacing any pair $1, -1$ by $-1, 1$. Thus to show that $\pi : g_i, v_i \rightarrow f_i, u_i$ extends to a homomorphism between the C^* -algebras $C^*(G)$ and \mathcal{B} , one just has to check that it gives a covariant representation of a groupoid. It is clear that $\pi : g_i \rightarrow f_i$ extends to a homomorphism of $C(K)$ onto $C^*(f_i)$. To show that $U : v_i \rightarrow u_i$ extends to a representation of the inverse semigroup \mathcal{G} of partial isometries on K , one notes that \mathcal{G} is the universal inverse semigroup of partial isometries on generators $\{v_i, v_i^*, f_i, 1-f_i\}$ subject to the relations 2(a)–(e), which imply the infinitesimal Jones or braid relations:

$$v_i v_{i \pm 1} v_i = 0 \quad v_i^2 = 0 \quad (2.6)$$

together with $v_i v_{i \pm 1}^* = 0 = v_{i \pm 1}^* v_i$. To see that \mathcal{G} is universal one merely observes that any non-zero word in $\{v_i, v_i^*, f_i, 1-f_i\}$ can be uniquely expressed as

$$f(\varepsilon) v^\#(J_1) \dots v^\#(J_p), \quad (2.7)$$

where

- (a) $f(\varepsilon) = \prod_i f(\varepsilon_i)$, $f(\varepsilon_i) = 1$, f_i or $1-f_i$, and $\text{supp } \varepsilon = \{i : f_i \neq 1\}$ is finite.
 (b) $J_i = \{k_i, k_i+1, \dots, j_i\}$ $j_i \geq k_i$, $j_{i+1} > j_i$, $k_{i+1} > k_i$, $v^\#(J_i) = v_{k_i}^\# \dots v_{j_i}^\#$. We say $v^*(J_i)$ is starred, $v(J_i)$ is unstarred.
 (c) If $J_i^0 = \{k_i, \dots, j_i+1\}$, and one of $\{v^\#(J_i), v^*(J_{i+1})\}$ is starred and one is unstarred, then $J_i^0 \cap J_{i+1} = \emptyset$.
 (d) $\text{supp}(\varepsilon) \cap \bigcup_{i=1}^p J_i^0 = \emptyset$.

We can then deduce that (π, U) gives a covariant representation of the groupoid \mathcal{G} on K and so gives the required embedding of the observable algebra on $\ell^2(I)$ in \mathcal{B} .

Remark 2.3. The relations (2.6) give rise to Quantum Yang-Baxter equation in that $R_i(s) = \exp(sv_i)$ satisfy:

$$\begin{aligned} R_i(s) R_{i+1}(s+s') R_i(s') &= R_{i+1}(s') R_i(s+s') R_{i+1}(s), \\ R_i(s) R_j(s') &= R_j(s') R_i(s) \quad |i-j| \neq 1. \end{aligned} \quad (2.8)$$

3. Local Structure

We start with a C^* -algebra A with the following type of local structure over \mathbb{Z} . To every subset $A \subset \mathbb{Z}$ is associated a C^* -subalgebra $A(A) \subset A$ in such a way that

$$\text{if } A \subset A', \text{ then } A(A) \subset A(A'), \quad (3.1)$$

$$\text{if } d(A, A') \geq 1, \text{ then } A(A) \text{ and } A(A') \text{ commute.} \quad (3.2)$$

Here $d(A, A') = \min \{|i - j| : i \in A, j \in A'\}$. We are also given an automorphism σ of A such that

$$\sigma(A(A)) = A(A + 2), \text{ for any } A \in \mathbb{Z}. \quad (3.3)$$

Example 3.1. The Jones Algebra. If $\tau^{-1} \in [4, \infty) \cup \{4 \cos^2 \pi/m : m = 3, 4, \dots\}$, let $A = A_\tau = A(\tau)$ be the C^* -algebra generated by projections $e_n, n \in \mathbb{Z}$ satisfying the V. Jones' relations

$$e_n e_m = e_m e_n, \quad |m - n| \geq \tau, \quad (3.4)$$

$$e_n e_{n+1} e_n = \tau e_n. \quad (3.5)$$

For each subset $A \subset \mathbb{Z}$, let $A(A)$ be the C^* -subalgebra generated by $e_n, n \in A$. Let σ be the shift of 2, $\sigma(e_n) = e_{n+2}$, so that conditions (3.1)–(3.3) are obviously satisfied.

There is a unique positive trace Tr on $A(\tau)$ called the Markov trace, such that

$$\text{Tr}(x e_m) = \tau \text{Tr}(x) \quad x \in A(-\infty, m-1], \quad (3.6)$$

$$\text{Tr}(e_m) = \tau, \quad \text{Tr}(1) = 1. \quad (3.7)$$

In particular, for $\tau^{-1} = \frac{1}{2} = 4 \cos^2 \pi/4$, $A(\frac{1}{2})$ can be identified with the even part of the Clifford algebra over $l_{\mathbb{R}}^2(\mathbb{Z})$. More precisely let $\{\gamma_n : n \in \mathbb{Z}\}$ satisfy the Clifford relations $\gamma_n = \gamma_n^*, \{\gamma_n, \gamma_m\} = 2\delta_{nm}$ so that $U_n = i\gamma_n \gamma_{n+1}$ are self adjoint unitaries in the Clifford algebra $A^F = C^*(\gamma_n : n \in \mathbb{Z})$ satisfying the relations $U_n U_m = U_m U_n, |m - n| \geq 2$, and $U_n U_{n+1} = -U_{n+1} U_n$. Hence $e_n = (1 + U_n)/2 = (1 + i\gamma_n \gamma_{n+1})/2$ satisfy the V. Jones relations (3.4) and (3.5), so that $A(\frac{1}{2})$ can be identified with the even part of the Clifford algebra, namely $A_+^F = C^*(\gamma_i \gamma_j : i, j \in \mathbb{Z})$. (The algebra A^F is graded so that γ_n is odd.) Going from this real or Clifford picture to a CAR picture, one takes a complex structure on $l_{\mathbb{R}}^2(\mathbb{Z})$ with real basis η_n such that $i\eta_{2n} = \eta_{2n+1}, i\eta_{2n+1} = -\eta_{2n}$. Then the creation operators $a_n = (\gamma_{2n} + i\gamma_{2n+1})/2$, are odd, and satisfy the canonical commutation relations $\{a_n, a_m\} = 0, \{a_n, a_m^*\} = \delta_{nm}$. In this language we identify:

$$e_{2n} = a_n^* a_n, \quad (3.8)$$

$$e_{2n+1} = [1 + (a_n - a_n^*)(a_{n+1} + a_{n+1}^*)]/2. \quad (3.9)$$

Then $a_n^* a_n = e_{2n}, a_n^* a_{n+1} = e_{2n}(1 - 2e_{2n+1})e_{2n+2}$ defines a Lie algebra homomorphism of $gl(\infty)$ into $A(\frac{1}{2})$.

Example 3.2. The Ocneanu Algebra. Let Γ be a graph, with base point $*$, $\Gamma^{(0)}$ the set of vertices, $\Gamma^{(1)}$ the set of edges, so that $\Gamma^{(1)}$ consists of certain two element subsets of $\Gamma^{(0)}$. We say that $\gamma \in \Gamma^{(0)}$ is even (respectively odd) if it can be joined to $*$

by an even (respectively odd) number of edges. We let $\Gamma_+^{(0)}$ (respectively $\Gamma_-^{(0)}$) denote the even (respectively odd) vertices. Let K be the Cantor set of all sequences $(x_v)_{v \in \mathbb{Z}}$ with

$$\forall v, (x_v, x_{v+1}) \in \Gamma^{(1)}, \quad (3.10)$$

$$x_0 \text{ is even.} \quad (3.11)$$

On K consider the equivalence relation \sim with countable orbits given by $(x_v) \sim (y_v)$ if and only if $x_v = y_v$, except in some finite set of v 's. Let $A = A^\Gamma = A(\Gamma)$ be the corresponding C^* -algebra. For each (finite) subset $A \subset \mathbb{Z}$, let $A(A) = A^\Gamma(A)$ be the C^* -subalgebra generated by the following partial isometries $f_{\gamma, \gamma'}$: Both γ and γ' are elements of

$$\begin{aligned} \mathcal{A} = \{ \gamma \in (\Gamma^{(0)})^{A'} : (\gamma_i, \gamma_{i+1}) \in \Gamma^{(1)}, \gamma_{2m} \in \Gamma_+^{(0)}, \gamma_{2n+1} \in \Gamma_-^{(0)}, \\ \text{if } i, i+1, 2m, 2n+1 \in A' \}, \end{aligned} \quad (3.12)$$

where $A' = \{i : d(i, A) \leq 1\}$, and $\gamma(j) = \gamma'(j)$ if $j \notin A$. The partial isometry $f_{\gamma, \gamma'}$ has as initial domain the cylinder set $Z_\gamma = \{(x_v) : x_i = y_i, i \in A'\}$, and it replaces any such (x_v) in Z_γ by (y_v) in the cylinder set $Z_{\gamma'}$ where $y_v = x_v, v \notin A, y_v = \gamma'_v, v \in A$. The shift by 2 on K induces an automorphism σ of $A(\Gamma)$ satisfying (3.1)–(3.3).

Associated to the graph Γ is a symmetric 0–1 matrix A , indexed by the vertices $\Gamma^{(0)}$, with entries $A_{\varrho\eta}$ equal to 0 or 1, according as $(\varrho, \eta) \in \Gamma^{(1)}$ or $(\varrho, \eta) \notin \Gamma^{(1)}$. We let $t(\eta) = \{\varrho \in \Gamma^{(0)} : (\varrho, \eta) \in \Gamma^{(1)}\}$. If the graph is connected, then there exists a unique eigenvector $(v_\varrho)_{\varrho \in \Gamma^{(0)}}$ with strictly positive entries, and $\alpha \in (0, 1]$ such that

$$v_\eta = \sum_{\varrho \in t(\eta)} \alpha v_\varrho, \quad \sum_{\eta \text{ even}} v_\eta^2 = 1. \quad (3.13)$$

Then $X(\eta) = ((\alpha v_\varrho / v_\eta)^{1/2})_{\varrho \in \Gamma^{(0)}}$ defines a unit vector in $\ell^2(t(\eta))$. If $k \leq \ell \in \mathbb{Z}$, $\delta, \beta \in \Gamma^{(0)}$, let $\mathcal{G}_{\delta, \beta}^{[k, \ell]} = \{\gamma \in \mathcal{G}^{[k, \ell]} : \gamma(k-1) = \delta, \gamma(\ell+1) = \beta\}$. Then for each $n \in \mathbb{Z}$, let

$$e_n = \sum_{\eta} X(\eta)(\varrho) X(\eta)(\varrho') f_{\gamma, \gamma'}, \quad (3.14)$$

where the summation is over all $\eta \in \Gamma^{(0)}$ if n is even or all $\eta \in \Gamma_+^{(0)}$ if n is odd, $\gamma, \gamma' \in \mathcal{G}_{\eta, \eta}^{[n]}$, and $\gamma(n) = \varrho, \gamma'(n) = \varrho'$. Then e_n is a projection, being identified with a sum (over odd or even vertices depending on whether n is even or odd) of the rank one projection on $X(\eta)$ in $\text{End}(\ell^2(t(\eta)))$. The family e_n satisfy the Jones' relations (3.3) and (3.4) for $\tau = \alpha^2$, so that $A(\alpha^2) \subset A(\Gamma)$.

We define a trace Tr , called the Markov trace, on $A(\Gamma)$ to be the unique state on $A(\Gamma)$ such that

$$\text{Tr } f_{\gamma, \gamma'} = 0 \quad \text{if } \gamma \neq \gamma', \quad (3.15)$$

$$\text{Tr } f_{\gamma, \gamma} = \alpha^{\ell-k+2} v_\delta v_\beta, \quad \text{if } \gamma \in \mathcal{G}_{\delta, \beta}^{[k, \ell]}. \quad (3.16)$$

Then

$$\text{Tr}(xe_m) = \alpha^2 \text{Tr}(x), \quad x \in A(-\infty, m-1], \quad (3.17)$$

$$\text{Tr}(e_m) = \alpha^2, \quad \text{Tr}(1) = 1. \quad (3.18)$$

Consequently $\text{Tr}|A(\alpha^2)$ is the Markov trace on $A(\alpha^2)$.

We emphasize that $A(\Gamma) \neq A(\alpha^2)$ in general, and we shall indeed see in Sect. 6 that

$$A^{\alpha^2}\{0, 1, 2\}_{e_0(1-e_2)} = \mathbb{C} \oplus A^\Gamma\{0, 1, 2\}_{e_0(1-e_2)}.$$

However we claim that $A(\Gamma)$ is generated by $A(\alpha^2)$ and a finite dimensional algebra. Let ϱ (respectively π) denote the left (respectively right) action of A_k on A_k so that $\varrho(A_k)' = \pi(A_k)$. We first note that there is an action ϱ of $A_{k+1} = A[0, k+1]$ on A_k , whose commutant is the (right) action of A_{k-1} on A_k . The right action π of A_{k-1} on A_k is given by

$$\pi(f_{\gamma_1, \gamma'_1}) : f_{\gamma, \gamma'} \rightarrow f_{\gamma, \gamma'_1 \varrho}$$

if $\gamma' = \gamma_1 \varrho$, $\gamma_1 \gamma' \in \mathcal{G}^{[0, k]}$, $\gamma_1, \gamma'_1 \in \mathcal{G}^{[0, k-1]}$, $\varrho \in \mathcal{G}^{[k-1, k]}$, and $\pi(f_{\gamma_1, \gamma'_1}) f_{\gamma, \gamma'} = 0$ otherwise. To describe the commutant of the action of $\{f_{\gamma_1, \gamma'_1} : \gamma'_1(k-1) = \delta\}$ on $f_{\gamma, \gamma'}$ we note that we have the freedom to change $\tilde{\gamma} \in [0, k+2]$, where $\tilde{\gamma}|[0, k] = \gamma$, $\tilde{\gamma}(k+1) = \delta$, which can be described as an action ϱ of A_{k+1} as follows:

$$\varrho(f_{\tilde{\gamma}, \tilde{\gamma}'}): f_{\gamma, \gamma'} \rightarrow v_{\eta_1}^{1/2} v_{\eta_2}^{-1/2} f_{\tilde{\gamma}_0, \tilde{\gamma}'_0},$$

if $\gamma, \gamma' \in \mathcal{G}^{[0, k]}$, $\tilde{\gamma}, \tilde{\gamma}' \in \mathcal{G}^{[0, k+1]}$, $\gamma'_0|[0, k] = \gamma$, $\tilde{\gamma}_0 = \tilde{\gamma}|[0, k]$, $\gamma'_0|[0, k-1] = \gamma'|[0, k-1]$, $\gamma'_0(k) = \tilde{\gamma}(k)$, $\eta_1 = \gamma(k) = \gamma'(k)$, $\eta_2 = \tilde{\gamma}(k)$, and $\varrho(f_{\tilde{\gamma}, \tilde{\gamma}'}): f_{\gamma, \gamma'} = 0$ otherwise. The factor $v_{\eta_1}^{1/2} v_{\eta_2}^{1/2}$ is inserted to make ϱ a *-representation with respect to the inner product $\langle x, y \rangle = \text{tr } y^* x$ on A_{k-1} . Then $\varrho(A_{k+1})' = \pi(A_k)$, and ϱ is faithful if every vertex which is a right-hand end point of a path of length $k+1$, is also a right-hand end point of a path of length $k-1$ (with the same left-hand and points). We assume that this is the case. This will certainly be the case if the two restrictions of A^2 to even and odd vertices are irreducible matrices. Then $\varrho(e_{k+1}) = E_{k+1}$ is the orthogonal projection of A_k on A_{k-1} .

Lemma 3.3. *If k is sufficiently large, A_{k+1} is the C^* -algebra generated by A_k and e_{k+1} .*

Proof. We take k sufficiently large that ϱ is faithful. Let $x \in (\varrho(A_k), \varrho(e_{k+1}))' \subset \varrho(A_k)' = \pi(A_k)$. Then $x E_{k+1} = E_{k+1} x$ shows that $x = x E_{k+1} (I) = (E_{k+1} x) I = E_{k+1} (x) \in A_{k-1}$. Hence $(\varrho(A_k), \varrho(e_{k+1}))' \subset \pi(A_{k-1})$, and so $\varrho(A_{k+1}) \supset (\varrho(A_k), \varrho(e_{k+1}))'' \supset \pi(A_{k-1})' \supset \varrho(A_{k+1})$.

Corollary 3.4. *The C^* -algebra $A(\Gamma)$ is generated by $A(\alpha^2)$ and a finite dimensional C^* -algebra.*

4. Embedding the Observable Algebra in Local Algebras

Let A be a C^* -algebra with local structure as in (3.1)–(3.3) of Sect. 3.

Lemma 4.1. *Let $f_0 \in A\{0\}$ be a projection and $u_0 \in A\{0, 1, 2\}$ a partial isometry; with initial support $\sigma(f_0)(1 - f_0)$ and final support $f_0(1 - \sigma(f_0))$. Then there exists an*

unique $*$ -homomorphism j from the observable algebra \mathcal{O} over $\ell^2(\mathbb{Z})$ to A such that

$$j\sigma_{\mathcal{O}}(x) = \sigma_A(x), x \in \mathcal{O}, \quad (4.1)$$

$$j(a_0^*a_0) = f_0, \quad (4.2)$$

$$j(a_0^*a_1) = u_0. \quad (4.3)$$

Remark 4.2. Here $\sigma_{\mathcal{O}}$ is the automorphism of the observable algebra \mathcal{O} which is the restriction of the automorphism σ on the Fermion algebra given by $\sigma(\gamma_n) = \gamma_{n+2}$ or $\sigma(a_i) = a_{i+1}$ (see Example 3.1).

Proof of Lemma 4.1. Let $f_i = \sigma^i(f_0) \in A\{i\}$, $u_i = \sigma^i(u_0) \in A\{2i, 2i+1, 2i+2\}$ for $i \in \mathbb{Z}$. Then $\{f_i, u_i\}$ satisfy conditions (2.5) of Lemma 2.2.

Let us now examine Lemma 4.1 in the context of the Jones and Ocneanu algebras.

Lemma 4.2. Let A_{τ} be the V . Jones algebra as in Sect. 3.

(4.4) In $A\{0, 1, 2\}$ the projections $e_0(1 - e_2)$ and $(1 - e_0)e_2$ are equivalent.

(4.5) The reduced algebra of $A\{0, 1, 2\}$ by $e_0(1 - e_2)$ is one-dimensional.

Proof. (4.4). If $u = \tau^{-1}e_0(1 - e_2)e_1(1 - e_0)e_2 = e_0(1 - \tau^{-1}e_1)e_2 \in A\{0, 1, 2\}$, [cf. (Example 3.1)] then

$$u^*u = (1 - e_0)e_2, uu^* = e_2(1 - e_0).$$

(4.5) is clear from (4.4), $A\{0, 1, 2\} = C^*(e_0, e_1, e_2)$ and the Jones' relations.

Theorem 4.3. (a) For $\tau^{-1} \in [4, \infty) \cup \{4 \cos^2 \pi/m\}$, let $A = A(\tau)$ the V . Jones algebra. Then for each $t \in \mathbb{T}$, there exists an unique $*$ -endomorphism j_t of the observable algebra \mathcal{O} into $A(\tau)$ such that

$$j_t\sigma_{\mathcal{O}}(x) = \sigma_A j_t(x), \quad x \in \mathcal{O}, \quad (4.6)$$

$$j_t(a_i^*a_i) = e_{2i}, \quad i \in \mathbb{Z}, \quad (4.7)$$

$$j_t(a_i^*a_{i+1}) = \sigma_A^i(t) e_{2i}(1 - \tau^{-1}e_{2i+1})e_{2i+2}, \quad i \in \mathbb{Z}. \quad (4.8)$$

(b) For Γ a connected graph, let $A = A(\Gamma)$ be the Ocneanu algebra, and G^{Γ} be the unitary group of the reduced algebra of $A^{\Gamma}\{0, 1, 2\}$ by $e_0(1 - e_2)$. Then for each $t \in G^{\Gamma}$, there exists an unique $*$ -endomorphism j_t of the observable algebra \mathcal{O} into $A(\Gamma)$ satisfying (4.6)–(4.8).

Proof. Using Lemmas 2.2 and 4.2, it only remains to show that the embeddings are faithful. This can be shown by computing the Markov trace on the projection $f(\varepsilon)$ (in the notation of the proof of Lemma 2.2) in the centre $C((\pm 1)^{\mathbb{Z}})$ of \mathcal{O} , and observing that this is non-zero.

It only remains to compute the group G^{Γ} .

Proposition 4.4. For Γ a connected graph, let $A = A(\Gamma)$ be the Ocneanu algebra.

(4.9) Any function of $x_v, v \in A' \setminus A = \partial A$ lies in the centre of $A(A)$.

(4.10) Let $q_\eta^{(n)}$ be the function $q_\eta^{(n)}((x_v)) = 0$ if $x_n \neq \eta$, and $q_\eta^{(n)}((x_v)) = 1$ if $x_n = \eta$. Then $q_\eta^{(-1)}e_0q_\eta^{(1)} = 0$ if $\eta \neq \eta'$.

(4.11) Let $P^{(k)}(\eta, \eta')$ denote the paths of length k from η to η' . Then if $\eta \neq \eta'$, $q_\eta^{(-1)}e_0q_\eta^{(1)}(1 - e_2)q_\eta^{(3)}$ is a projection of dimension $\#P^{(2)}(\eta, \eta')$ in $A\{0, 1, 2\}$. (Here $\#(\emptyset) = 0$.)

(4.12) $q_\eta^{(-1)}e_0q_\eta^{(1)}(1 - e_2)q_\eta^{(3)}$ is a projection of dimension $\#t(\eta) - 1$ in $A\{0, 1, 2\}$.

Proof. (4.9): Any function $x_v, v \in \partial A$ is a linear combination of $\{g_\eta^{(v)} : \eta\}$. Then $g_\eta^{(v)} = \sum f_{\gamma, \gamma'} A(A)$, where the summation is over all $\gamma \in \mathcal{G}^A$ such that $\gamma(v) = \eta$. Moreover $g_\eta^{(v)}$ is a unit for $\{f_{\gamma, \gamma'} : \gamma, \gamma' \in \mathcal{G}^A, \gamma| \partial A = \gamma'| \partial A, \gamma(v) = \eta\}$ and annihilates $\{f_{\gamma, \gamma'} : \gamma, \gamma' \in \mathcal{G}^A, \gamma| \partial A = \gamma'| \partial A, \gamma(v) \neq \eta\}$. Hence $g_\eta^{(v)}$ commutes with $A(A) = \lim \{f_{\gamma, \gamma'} : \gamma, \gamma' \in \mathcal{G}^A, \gamma| \partial A = \gamma'| \partial A\}$.

(4.10): We have

$$A\{0\} = \bigoplus_{\eta, \eta'} \text{End } \ell^2(P^2(\eta, \eta')), \quad (4.13)$$

where the summation is over odd $\eta, \eta' \in \Gamma^{(0)}$. We have as in (4.9) that $q_\eta^{(-1)}$ (respectively $q_\eta^{(1)}$) is a unit for $\text{End } \ell^2(P^2(\alpha, \beta))$ if $\eta = \alpha$ (respectively $\eta = \beta$), and annihilates $\text{End } \ell^2(P^2(\alpha, \beta))$ if $\eta \neq \alpha$ (respectively $\eta \neq \beta$). Hence $q_\eta^{(-1)} \text{End } \ell^2 P(\alpha, \alpha) q_\eta^{(1)} = 0$ if $\eta \neq \eta'$. Consequently $q_\eta^{(-1)}e_0q_\eta^{(1)} = 0$ if $\eta \neq \eta'$ as $e_0 \in \bigoplus_{\alpha \text{ odd}} \text{End } \ell^2(P(\eta, \eta))$, [where we identify $P^2(\alpha, \alpha)$ with $t(\alpha)$].

(4.11): Take η, η' odd vertices in $\Gamma^{(0)}$, $\eta \neq \eta'$. Then

$$q_\eta^{(1)}q_\eta^{(3)} = \sum f_{\gamma', \gamma'},$$

where the summation is over all $\gamma' \in \mathcal{G}_{\eta, \eta'}^{(2)}$. Take $\gamma \in \mathcal{G}_{\eta}^{(0)}$, such that $q_{e_0}^{(-1)}e_0q_\eta^{(1)}$ being a rank one projection in $\text{End } \ell^2(t(\eta))$, there exists a unitary u in $\text{End } \ell^2(t(\eta)) \in A\{0\}$ such that $q_\eta^{(-1)}e_0q_\eta^{(1)} = u f_{\gamma, \gamma} u^*$. Hence, since $[u, f_{\gamma', \gamma'}] = 0$ for each $\gamma' \in \mathcal{G}_{\eta, \eta'}^{(2)}$, we see that

$$q_\eta^{(-1)}e_0q_\eta^{(1)}(1 - e_2)q_\eta^{(3)} = q_\eta^{(-1)}e_0q_\eta^{(1)}q_\eta^{(3)} = u(\sum f_{\gamma\gamma', \gamma\gamma'})u^*,$$

where the summation is over all $\gamma' \in \mathcal{G}_{\eta, \eta'}^{(2)}$, and which is of dimension $\#P^{(2)}(\eta, \eta')$ in $\text{End}(\ell^2(P^{(2)}(\eta, \eta')))$. (Here if $\gamma \in \mathcal{G}^{A_1}, \gamma' \in \mathcal{G}^{A_2}$, with $\gamma = \gamma'$ on $A'_1 \cap A'_2$, we let $\gamma\gamma'$ denote the obvious element of $\mathcal{G}^{A_1 \cup A_2}$.)

(4.12): As in the proof of (4.11), we see that the dimension of $q_\eta^{(-1)}e_0q_\eta^{(1)}(1 - e_2)q_\eta^{(3)}$ in $\text{End } \ell^2(P^{(4)}(\eta, \eta))$ is the same as that of $q_\eta^{(1)}(1 - e_2)q_\eta^{(3)}$ in $\text{End } \ell^2(t(\eta))$. Now e_2 is a rank one projection in $\text{End } \ell^2(t(\eta))$, and consequently its complement is of dimension $\#t(\eta) - 1$.

Corollary 4.5. If Γ is a connected graph, the reduction of the algebra $A^\Gamma\{0, 1, 2\}$ by $e_0(1 - e_2)$ can be identified with

$$\bigoplus_{\eta, \eta' \text{ odd}} M(N(\eta, \eta')), \quad (4.14)$$

where $M(N)$ denotes the algebra of complex $N \times N$ matrices, and the summation is over odd vertices η, η' in $\Gamma^{(0)}$ with

$$N(\eta, \eta') = \begin{cases} \#P^{(2)}(\eta, \eta') & \text{when } \eta \neq \eta' \\ \#t(\eta) - 1 & \text{when } \eta = \eta' \end{cases} \quad (4.15)$$

ω_{A_β} is the unique KMS state on the Clifford algebra $\text{Cliff}_{\mathbb{C}}(\ell^2(S^1)_{\mathbb{R}})$ for the Dirac Hamiltonian, with chemical potential $\beta^{-1} \log(\tau/(1-\tau))$, namely $D_\tau = D - \beta^{-1} \log \tau/(1-\tau)$. If $j = j_t$, for $t \in \mathbb{T} \subset G_I$, then E takes $A(\tau)$ onto $j(\mathcal{O}) \subset A(\tau)$.

Proof. The Markov property of the unique trace Tr on $A(\Gamma)$ shows that $\omega(x) = \lim_{m \rightarrow \infty} \text{Tr}(e^{-\beta H_m x})/\text{Tr}(e^{-\beta H_m})$ defines a state on $A(\Gamma)$, and which is KMS at inverse temperature β for σ_t . Let ϕ be any KMS state for σ_t at inverse temperature β , and $\phi_m = \phi/A_m$, $A_m = A^\Gamma[-2m, 2m]$. Then $\psi_m = \phi_m(e^{\beta H_m \cdot})/\phi_m(e^{\beta H_m})$ is a normalised positive trace on A_m . Let Δ_m denote the compact convex set of normalised positive traces on A_m , which can be identified with a subset of $\Delta_m = \{x \in \mathbb{R}^p : x_i \geq 0, \sum x_i = 1\}$, where $p = (\#(I^{(0)}))^2$. Then by [E, Sects. 5–6], $\Delta_m \supset \Delta_{m+1}$, $\text{diameter } \Delta_m \rightarrow 0$, and $\bigcap_{m=1}^{\infty} \Delta_m$ can be identified with the unique normalised trace Tr on $A(\Gamma)$. Hence given $\varepsilon > 0$, there exists m_0 , such that for $m \geq m_0$, $(1-\varepsilon)\text{Tr} \leq \psi_m \leq (1+\varepsilon)\text{Tr}$. Hence for $m \geq m_0$:

$$\begin{aligned} \phi(x) - \text{Tr}(e^{-\beta H_m x})/\text{Tr}(e^{-\beta H_m}) &= \psi_m(e^{-\beta H_m x})/\psi_m(e^{-\beta H_m}) - \text{Tr}(e^{-\beta H_m x})/\text{Tr}(e^{-\beta H_m}) \\ &\leq ((1+\varepsilon)(1-\varepsilon)^{-1} - 1) \text{Tr}(e^{-\beta H_m x})/\text{Tr}(e^{-\beta H_m}). \end{aligned}$$

Thus $\phi(x) = \lim_{m \rightarrow \infty} \text{Tr}(e^{-\beta H_m x})/\text{Tr}(e^{-\beta H_m}) = \omega(x)$, which shows uniqueness.

The state $\omega_\beta \circ j^{-1}|_{\mathcal{O}}$ on the observable algebra \mathcal{O} is independent of the choice of embedding j of \mathcal{O} in $A(\Gamma)$. In fact it is the restriction to \mathcal{O} of the quasi-free state ω_{A_β} , where $A_\beta = \tau e^{-\beta D}/(\tau e^{-\beta D} + (1-\tau)) = e^{-\beta D_\tau}(e^{-\beta D_\tau} + 1)$, $D_\tau = D - \beta^{-1} \log \tau/(1-\tau)$. In particular, taking $\beta = 0$, the restriction of the trace on $A(\Gamma)$ to \mathcal{O} is the restriction of the Powers state ϕ_λ on $\text{Cliff}_{\mathbb{C}}(\ell^2(\mathbb{Z})_{\mathbb{R}})$, for $\lambda = \tau/(1-\tau)$.

Next we show that there exists a projection of norm one from $A(\Gamma)$ onto the observable algebra $j(\mathcal{O})$.

Lemma 5.2. *Let Γ be a connected graph. Consider the type II_1 representation associated with the trace Tr on $A(\Gamma)$, and E the orthogonal projection on $j(\mathcal{O})$. Then $E(x) \in j\mathcal{O}[-2m, 2m]$, for all $x \in A^\Gamma(-2m, 2m)$, if we let $\mathcal{O}(A) = A^{1/2}(A)^\mathbb{T}$.*

Proof. We have $xy = yx$ for all $x \in A(-2m, 2m)$, $y \in \mathcal{O}((-\infty, -2m-1] \cup [2m+1, \infty))$. Hence $E(x)y = E(xy) = E(yx) = yE(x)$ for all $x \in A(-2m, 2m)$, $y \in \mathcal{O}((-\infty, -2m-1] \cup [2m+1, \infty))$. Consequently, if $R_\lambda(A)$ denotes the weak closure of $A^{1/2}(A)$ in the Powers state ϕ_λ , and $R_\lambda^\phi(A)$ denotes the centraliser $(A^{1/2}(A)^\mathbb{T})''$, then

$$\begin{aligned} E(x) &\in R_\lambda(-\infty, 2m-1] \otimes R_\lambda[-2m, 2m] \otimes R_\lambda[2m+1, \infty) \cap R_\lambda^\phi(-\infty, 2m+1] \\ &\quad \otimes \mathbb{C} \otimes R_\lambda^\phi[2m+1, \infty) \subset R_\lambda[-2m, 2m]. \end{aligned}$$

To calculate ω_β , take $x \in \mathcal{A}_n$, $m \geq n$, then

$$\omega_\beta(x) = \text{Tr}(e^{-\beta H_m x})/\text{Tr}(e^{-\beta H_m}) = \text{Tr} e^{-\beta H_m} E(x)/\text{Tr}(e^{-\beta H_m}) = \omega_\beta(E(x))$$

using $H_m \in j(\mathcal{O})$. Factoriality is clear, by uniqueness of KMS states for σ_t . Now $\omega_\beta(e_{2n}) \rightarrow 0$ as $n \rightarrow \infty$, and $\omega_\beta(1 - e_{2n}) \rightarrow 0$ as $n \rightarrow -\infty$. Thus it is convenient to introduce the projection in $\mathcal{M} = A(\Gamma)''$ given by

$$f = \lim_{N \rightarrow \infty} e_{-2N} e_{-2N+2} \cdots e_{-4} e_{-2} (1 - e_2) (1 - e_4) \cdots (1 - e_{2N}).$$

The automorphism group σ_t extends to a weakly continuous one parameter group of *-automorphism of \mathcal{M} . For $x \in A_m$, $N \geq m$,

$$\begin{aligned}\sigma_t(fxf) &= \lim_{N \rightarrow \infty} \sigma_t(f_N x f_N) = \lim_{N \rightarrow \infty} \text{Ad}(\exp it H_N)(f_N x f_N) \\ &= \lim_{N \rightarrow \infty} f_N x f_N = fxf \quad \text{as} \quad [H_N, f_N x f_N] = 0.\end{aligned}$$

Consequently, $\omega_\beta|_{\mathcal{M}_f}$ is a trace. But

$$\sigma(f(1-e_0)) = f e_0, \quad \text{and} \quad \omega_\beta(f e_0)/\omega_\beta(f(1-e_0)) = \tau/(1-\tau). \quad (5.4)$$

This means that \mathcal{M} is not of type II_1 for if ϕ is then the trace of \mathcal{M} , $\phi \circ \sigma = \phi$ and $\phi|_{\mathcal{M}_f} = \omega_\beta$ would contradict (5.4). Note moreover that

$$\sigma^n(f(1-e_0)) = \dots e_{-2N} \dots e_0 e_2 \dots e_{2n} (1-e_{2n+2}) \dots,$$

$n=0, 1, 2, \dots$ are a family of orthogonal projections whose sum

$$\sum_{n=0}^{\infty} \sigma^n(f(1-e_0)) = \dots e_{-2N} \dots e_{-4} e_2 e_4 \dots = e$$

has finite trace. Thus $(A(\Gamma))_e''$ is a factor of type II_1 . On the other hand note that $A(\Gamma)_e$ can be identified with $A^I[0, \infty)$.

6. Implementation of the Action of the Current Algebra at Zero Temperature

Let Γ be a connected graph, with Jones' index τ^{-1} , and j an embedding of the observable algebra \mathcal{O} in A^Γ as in Theorem 4.3. We first seek a canonical ground state for L_0 . In analogy with what was done in the Fermion algebra (see the appendix), we consider the dynamics $L_0^\varepsilon = L_0 + \varepsilon \delta_0$ corresponding to the perturbation $D + \varepsilon$ of the Dirac Hamiltonian, where δ_k is the derivation

$$\delta_k(x) = \sum_n i[j(a_{n+k}^* a_n), x], \quad (6.1)$$

defined on the elements x of finite support in A^Γ . The derivations δ_k leave the Jones' algebra A^τ and the observable algebra $j(\mathcal{O})$ invariant. Then L_0^ε has an unique KMS state ω_β^ε at inverse temperature β with

$$\omega_\beta^\varepsilon(x) = \text{Tr} e^{-\beta H_m^\varepsilon} x / \text{Tr} e^{-\beta H_m^\varepsilon} \quad (6.2)$$

for $x \in A^\Gamma[-2m, 2m]$, if $H_m^\varepsilon = \sum_{|n| \leq m} (n + \varepsilon) e_{2n}$.

Define

$$\omega_\infty(x) = \lim_{\varepsilon \downarrow 0} \lim_{\beta \rightarrow \infty} \omega_\beta^\varepsilon(x) = \lim_{N \rightarrow \infty} \text{Tr}(f_N x) / \text{Tr} f_N \quad (6.3)$$

for $x \in A^\Gamma$, where $f_N = e_{-2N} \dots e_{-4} e_{-2} (1-e_0) (1-e_2) \dots (1-e_{2N})$. Then ω_∞ is a ground state for L_0 on A^τ , and $\omega_\infty|_{\mathcal{O}} = \omega_{A^\infty}|_{\mathcal{O}}$ (for any embedding j of the observable algebra in A^Γ), where ω_{A^∞} is the ground state (Appendix) of $L_0|_{\mathcal{O}}$.

For $f \in C(S^1, \mathbb{R})$ consider the potential function $\Phi = \Phi^f$ defined on finite subsets $X \subset \mathbb{Z}$ by

$$\Phi(X) = \hat{f}_n j(a_{m+n}^* a_m) + \hat{f}_{-n} j(a_m^* a_{n+m}) \quad (6.4)$$

if $X = \{m, m+1, \dots, m+n\}$, $m \in \mathbb{Z}$, $n \in \mathbb{N}$, and $\Phi(X) = 0$ otherwise. Let

$$\mathcal{C} = \left\{ f \in C(S^1, \mathbb{R}) : \alpha = \alpha_f^2 = 2 \sum_{n=0}^{\infty} e^{\lambda n} (n+1) \hat{f}_{n+1} < \infty \text{ for some } \lambda > 0 \right\}, \quad (6.5)$$

and for $f \in \mathcal{C}$, construct the derivation $J(f)$ on the local elements A_0 in A^Γ , using the potential Φ^f defined above.

Thus

$$J(f)(x) = \sum_{m,n} i[\hat{f}_n j(a_{m+n}^* a_m), x], \quad x \in A_0, \quad (6.6)$$

and by [BR, Theorem 6.2.4], A_0 consists of analytic elements for $J(f)$, $J(f)$ is closable and its closure, also denoted by $J(f)$ generates a strongly continuous one-parameter group of $*$ -automorphisms θ_t^f of A^Γ such that

$$\theta_t^f(x) = \sum_{n=0}^{\infty} (J(f))^n(x)/n!, \quad x \in A_0 \quad (6.7)$$

$$= \lim_{n \rightarrow \infty} \text{Ad}(\exp(iH_n^f t))(x), \quad x \in A, \quad (6.8)$$

where $H_n^f = \sum_{|\alpha|, |\beta| \leq n} \hat{f}_{\alpha-\beta} j(a_\alpha^* a_\beta)$.

Moreover, by [BR, Theorem 6.2.4],

$$\|J(f)^n(x)\| \leq (2\alpha_f^2/\lambda)^n \|x\| n! \exp \lambda(2m+1) \quad (6.9)$$

for $x \in A[-2m, 2m]$, and it is clear from (6.7) that $\theta_t^f \theta_t^g = \theta_t^{f+g}$, $f, g \in \mathcal{C}$, $t \in \mathbb{R}$.

Let (π_∞, K, Ω) denote the GNS triple of ω_∞ on A^Γ . For $f \in \mathcal{C}$, let

$$\mathcal{D}(:J(f):) = \left\{ \zeta \in K : \lim_{N \rightarrow \infty} \sum_{|\alpha|, |\beta| \leq n} \hat{f}_{\alpha-\beta} \pi_\infty j(a_\alpha^* a_\beta) \zeta \text{ exists} \right\}, \quad (6.10)$$

and let

$$:J(f):\zeta = \lim_{N \rightarrow \infty} \sum_{|\alpha|, |\beta| \leq n} \hat{f}_{\alpha-\beta} \pi_\infty j(a_\alpha^* a_\beta) \zeta \quad \text{for } \zeta \in \mathcal{D}(:J(f):). \quad (6.11)$$

Then $x\Omega \in \mathcal{D}(:J(f):)$ for all $x \in A_0$, as

$$\lim_{N \rightarrow \infty} \sum \hat{f}_{\alpha-\beta} \pi_\infty j(a_\alpha^* a_\beta) \Omega \text{ in the Hilbert space } [\pi_\infty j(\mathcal{C})\Omega]^- \quad (6.12)$$

exists [and is $:J(f):\Omega$] by the classical theory in the appendix, and

$$J(f)(x) = \lim_{n \rightarrow \infty} \sum_{|\alpha|, |\beta| \leq n} i[\hat{f}_{\alpha-\beta} j(a_\alpha^* a_\beta), x] \quad (6.13)$$

by (6.6). Consequently for $x \in A_0$:

$$i : J(f) : x\Omega = J(f)(x)\Omega + ix : J(f) : \Omega. \quad (6.14)$$

Inductively we see that $A_0\Omega \subset \mathcal{D}((:J(f):)^n)$ and

$$(i : J(f) :)^n x\Omega = \sum_{r=0}^n {}^nC_r J(f)^{n-r}(x) (i : J(f) :)^r \Omega. \quad (6.15)$$

Hence by (6.9) and (Appendix), for $x \in \mathcal{A}[-2m, 2m]$:

$$\begin{aligned} & \| :J(f): \|^n x \Omega \| \\ & \leq \sum_{r=0}^n {}^n C_n (2\alpha_f^2/\lambda)^{n-r} \|x\| (n-r)! \exp \lambda(2m+1) ((2r)!/r!)^{1/2} (|f|/2)^r \\ & = \sum_{r=0}^n (2\alpha_f^2/\lambda)^{n-r} ((2r)!)^{1/2} (r!)^{-3/2} (|f|/2)^r n! \|x\| \exp \lambda(2m+1). \end{aligned} \quad (6.16)$$

Thus

$$\begin{aligned} & \sum_{n=0}^{\infty} t^n \| :J(f): \|^n x \Omega \| / n! \\ & \leq \|x\| \exp \lambda(2m+1) \sum_{r=0}^{\infty} (t|f|/2)^r (2r)!^{1/2} (r!)^{-3/2} \sum_{n=r}^{\infty} (2t\alpha_f^2/\lambda)^{n-r} \\ & = \|x\| \exp \lambda(2m+1) (1 - 2t\alpha_f^2/\lambda)^{-1} \sum_{r=0}^{\infty} (t|f|/2)^r (2r)!^{-1/2} (r!)^{-3/2}, \end{aligned} \quad (6.17)$$

$< \infty$ for sufficiently small t , $(2t\alpha_f^2/\lambda < 1)$.

Hence the symmetric operator $:J(f):$ has a dense set of analytic elements and so is essentially self adjoint. We again denote its closure by $:J(f):$.

The derivations $J(f)$ and $J(g)$ commute (for $f, g \in \mathcal{C}$). Thus for $x \in A_0$, we have

$$[[i:J(f):, i:J(g):], x] \Omega = [[J(f), J(g)], x] \Omega = 0.$$

Hence

$$[:J(f):, :J(g):] x \Omega = x[:J(f):, J(g):] \Omega = x \frac{1}{2\pi i} \int f dg \Omega,$$

where in the last equality we have used the classical theory of the Appendix in the Hilbert space $[\pi_{\infty} j(\mathcal{O}) \Omega]^-$.

We summarize this as follows, using the notation

$$\tilde{U}(C^\infty(S^1)) = \{h \in U(C^\infty(S^1)) : h = U^m \exp ig, m \in \mathbb{Z}, g \in \mathcal{C}\}.$$

Theorem 6.1. *Let Γ be a connected graph with Jones index τ^{-1} .*

(6.18) $J(f) = \sum_{n=-\infty}^{\infty} \hat{f}_n \delta_n$, $f \in \mathcal{C}$ defines a distribution with values in the derivations of A^Γ (respectively A^τ) and a representation of $\tilde{U}(C^\infty(S^1))$ in the group of automorphisms of A^Γ (respectively A^τ).

(6.19) The time evolution σ_t satisfies $J(\dot{f}) = -J(f')$.

(6.20) The action of the current algebra leaves the state ω_∞ :

$$\omega_\infty(x) = \lim_{N \rightarrow \infty} \frac{\text{Tr}(e_{-N} \dots e_{-2}(1-e_0) \dots (1-e_N)x)}{\text{Tr}(e_{-N} \dots e_{-2}(1-e_0) \dots (1-e_N))}$$

on A^Γ (respectively A^τ) quasi-invariant, and is implemented by self adjoint operators $:J(f):$ ($f \in \mathcal{C}$) satisfying

$$[:J(f):, J(g):] = \frac{1}{2\pi i} \int f dg. \quad (6.20)$$

To get a representation of the Virasoro algebra, with $c=1$, from the currents $T_n = :J_n:$, one uses the Sugawara formula:

$$L_n = \sum \frac{1}{2} \circ T_{m+n} T_{-m} \circ, \quad (6.21)$$

where

$$\circ T_m T_n \circ = \begin{cases} T_m T_n & \text{if } m < 0. \\ T_n T_m & \text{if } m > 0. \end{cases}$$

Appendix

We recall the Fermi-Dirac quantisation of the classical fields on S^1 . If H is a real Hilbert space, the Clifford algebra $\text{Cliff}_{\mathbb{C}}(H)$ is the universal C^* -algebra generated by the range of a linear map γ on H with the relations $\gamma(h)^* = \gamma(h)$, $\gamma(h)^2 = \|h\|^2$, $h \in H$. $\text{Cliff}_{\mathbb{C}}(H)$ has a $\mathbb{Z}/2$ grading so that each $\gamma(h)$ is odd. If i is a complex structure on H then $a(f) = [\gamma(f) - i\gamma(if)]/2$ for $f \in H_{\mathbb{C}}$, the complexification of H , defines a representation of the canonical anti-commutation relations. If $\{e_i\}_{i \in I}$ is a complete orthonormal basis for $H_{\mathbb{C}}$, we let $a_i = a(e_i)$. The field algebra \mathcal{A} is $\text{Cliff}_{\mathbb{C}}(\ell^2(\mathbb{Z})_{\mathbb{R}})$, and $\psi(\theta) = \sum a_n e^{-in\theta}$ gives the Fermi quantum field.

The gauge group $U(1)$ of the first kind acts on \mathcal{A} by $\alpha_s(a_n) = e^{-is} a_n$. The corresponding derivation $\delta = (\alpha_s)'_{s=0}$ takes a_n to $-ia_n$, and is given by $\delta(x) = i \sum [a_n^* a_m, x]$. The algebra of observable \mathcal{O} is the α -fixed point algebra, and is generated by $\{a_n^* a_m, 1\}$.

The gauge group $U(C(S^1))$ of the second kind acts on \mathcal{A} by $f: \psi(h) \rightarrow \psi(fh)$. Moreover $C^\infty(S^1, U(1)) (\subset C(S^1, U(1)) = U(C(S^1)))$ is generated by the shift U and the Lie algebra of real valued smooth maps on S^1 , in that we can write any $f \in C^\infty(S^1, U(1))$ uniquely as $f = U^m f_0$, where m is the winding number of f and $f_0 = \exp ig$ for some $g \in C^\infty(S^1, \mathbb{R})$. Then $g \in C^\infty(S^1, \mathbb{R})$ gives a $*$ -derivation $J(g)$ on \mathcal{A}_0 the $*$ -algebra generated by $\{a_m\}$. Complexifying, and taking $g = e^{im\theta}$, there is a derivation δ_m such that $\delta_m(a_n) = -ia_{n+m}$, $\delta_m(a_n^*) = ia_{n-m}^*$. Then $\delta_m^* = \delta_{-m}$, and

$$\delta_m(x) = i \sum_n [a_{n+m}^* a_m, x].$$

For $g = \sum \hat{g}_m e^{im\theta}$, $J(g) = \sum \hat{g}_m \delta_m = \frac{1}{2\pi} \int f(\theta) J(\theta) d\theta$, where $J(\theta) = \sum e^{-im\theta} \delta_m$ is a distribution taking values in the derivations on \mathcal{A}_0 .

Time evolution is given by $\sigma_t = \text{Cliff}(e^{itD})$ on \mathcal{A} , where D is the Dirac operator $-\text{id}/d\theta$ on S^1 , or $\sigma_t(a_n) = e^{-int} a_n$. The $*$ -derivation $L_0 = (\sigma_t)'_{t=0}$ takes a_n to ina_n and

$$L_0(x) = \sum in [a_n^* a_n, x].$$

Conservation of the current is expressed by $[L_0, J(\theta)] = -J'(\theta)$.

If C is a positive contraction on a complex Hilbert space K , we let ω_C denote the quasi-free state on $\text{Cliff}_{\mathbb{C}}(K_{\mathbb{R}})$ such that $\omega(a^*(f) a(g)) = \langle Cf, g \rangle$. If D is a self-adjoint operator, then for all $0 \leq \beta < \infty$, there exists a unique KMS state ω_β at inverse temperature β for $\sigma_t = \text{Cliff}(e^{itD})$, which is quasi-free with covariance operator $A_\beta = e^{-\beta D} (e^{-\beta D} + 1)^{-1}$. The Dirac Hamiltonian has a kernel, there is no unique

ground state, so we take the canonical ground state ω_∞ with covariance operator

$$A_\infty = \lim_{\varepsilon \downarrow 0} \lim_{\beta \rightarrow \infty} e^{-\beta(D+\varepsilon)} [e^{-\beta(D+\varepsilon)} + 1]^{-1} = (1 - F)/2,$$

where $F = \lim_{\varepsilon \downarrow 0} (D + \varepsilon) |D + \varepsilon|^{-1}$ is the “phase” of D . In the state ω_∞ using the normal ordering $:a_i^* a_j: = a_i^* a_j - \omega_\infty(a_i^* a_j)$, L_0 is implemented by $\sum n :a_n^* a_n:$ and δ_k by $J_k = \sum :a_{n+k}^* a_n:$. Then $J_k^* = J_{-k}$, and $[J_m, J_n] = n\delta_{m, -n}$. Consequently $J(g)$ is implemented by $:J(g): = \sum \hat{g}_m J_m$, and

$$[:J(f):, :J(g):] = \frac{1}{2\pi i} \int_{S^1} f dg$$

for $f, g \in \mathcal{A} = \{h \in C(S^1) : |h|^2 \equiv \sum |k|^2 |\hat{h}_k|^2 < \infty\}$. Thus $f \rightarrow \exp i :J(f):$ is a representation of the canonical commutation relations on \mathcal{A} . Moreover, the vacuum vector Ω gives the Boson Fock state

$$\langle \exp i :J(f): \Omega, \Omega \rangle = \exp -|f|^2/4.$$

Thus $f \rightarrow :J(f):$ is a free Bose field and

$$\langle :J(f):^{2n} \Omega, \Omega \rangle = \frac{(2n)!}{n!} \left(\frac{|f|}{2} \right)^{2n}.$$

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