Asymptotic behavior of Lévy measure density corresponding to inverse local time

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Abstract: For a one dimensional diffusion process $\mathbf{D}_{s,m}^*$ and the harmonic transformed process \mathbf{D}_{s_h,m_h}^* , the asymptotic behavior of the Lévy measure density corresponding to the inverse local time at the regular end point is investigated. The asymptotic behavior of n^* , the Lévy measure density corresponding to $\mathbf{D}_{s,m}^*$, follows from asymptotic behavior of the speed measure m. However, that of n^{h*} , the Lévy measure density corresponding to \mathbf{D}_{s_h,m_h}^* , is given by a simple form, n^* multiplied by an exponential decay function, for any harmonic function h based on the original diffusion operator.

Key words: Lévy measure density; asymptotic behavior; inverse local time.

1. Inverse local time and Lévy measure density. Let s be a continuous increasing function on an open interval $I = (l_1, l_2)$, where $-\infty < l_1 < l_2 \le \infty$, and let m be a right continuous increasing function on I. We assume

(1)
$$|s(l_1)| + |m(l_1)| < \infty$$
,

where we set $u(l_i) = \lim_{x \to l_i, x \in I} u(x)$, i = 1, 2, if there exist the limits, for functions u on I. (1) implies that the end point l_1 is regular in the sense of Feller [2]. We pose the reflecting or absorbing boundary condition at l_i (i = 1, 2) if it is regular. Let $\mathcal{G}_{s,m}$ be a one dimensional diffusion operator on I with scale function s, speed measure m, and null killing measure. We denote by $\mathbf{D}_{s,m}^* = [X(t), P_x^*]$ [resp. $\mathbf{D}_{s,m}^o = [X(t), P_x^o]$] the one dimensional diffusion process on I with $\mathcal{G}_{s,m}$ as the generator and with l_1 being reflecting [resp. absorbing]. Let denote by $l^*(t,\xi)$ the local time of $\mathbf{D}_{s,m}^*$, that is,

$$\int_0^t f(X(u)) \, du = \int_I l^*(t,\xi) f(\xi) \, dm(\xi), \quad t > 0,$$

for bounded continuous functions f on I. Since $l^*(t,\xi)$ is continuous and nondecreasing in t P_x^* -a.s., there is the right continuous inverse function $l^{*-1}(t,\xi)$. Note that there exists the inverse local time $l^{*-1}(t,l_1)$ at the end point l_1 , which is denoted

by $\tau^*(t)$. Combining Lévy formulas due to R. M. Blumenthal and R. K. Getoor ([1], Chapter V, Theorem 3.21) and those due to K. Itô and H. P. McKean ([3], Section 6.2), we obtain the following result. We give the proof in another paper.

Proposition 1. The Laplace transform of the distribution of $[\tau^*(t), t \geq 0]$ is given by the following

(2)
$$E_{l_1}^*[e^{-\lambda \tau^*(t)}]$$

= $\exp\left\{-\gamma^* t - t \int_0^\infty (1 - e^{-\lambda \xi}) n^*(\xi) d\xi\right\},$

(3)
$$\gamma^* = \begin{cases} 0 & \text{if } s(l_2) = \infty, \text{ or} \\ l_2 & \text{is regular and reflecting,} \\ 1/\{s(l_2) - s(l_1)\} & \text{if } s(l_2) < \infty, \end{cases}$$

(4)
$$n^*(\xi) = \lim_{x \to l_1} q^*(\xi, x) / \{s(x) - s(l_1)\},$$

where $E_{l_1}^*$ stands for the expectation with respect to P_l^* .

(5)
$$\int_0^t q^*(\xi, x) d\xi = P_x^*(\sigma_{l_1} < t), \quad x \in I, \ t > 0,$$

and σ_{l_1} is the first hitting time for l_1 .

We note a representation of $q^*(\xi, x)$ in terms of the transition probability density $p^o(t, x, y)$ of $\mathbf{D}_{s,m}^o$, that is,

(6)
$$q^*(\xi, x) = \lim_{z \to l_1} p^o(\xi, z, x) / \{s(z) - s(l_1)\},$$
$$\xi > 0, \ x \in I.$$

Here $p^{o}(t, x, y)$ is the transition probability density with respect to dm for $\mathbf{D}_{s,m}^{o}$, that is,

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$$P_x^o(X(t) \in E) = \int_E p^o(t, x, y) \, dm(y),$$

for $x \in I$, $E \in \mathcal{B}(I)$, where $\mathcal{B}(I)$ stands for the set of all Borel sets of I. It is well known that $p^o(t, x, y)$ is represented as

$$(7) \quad p^{o}(t,x,y)$$

$$= \int_{(0,\infty)} e^{-\lambda t} \psi^{o}(x,-\lambda) \psi^{o}(y,-\lambda) d\sigma^{o}(\lambda),$$

$$t > 0, \ x,y \in I,$$

where $d\sigma^{o}(\lambda)$ is a Borel measure on $(0, \infty)$ satisfying

(8)
$$\int_{(0,\infty)} e^{-\lambda t} d\sigma^{o}(\lambda) < \infty, \quad t > 0,$$

and $\psi^o(x, \alpha)$, $x \in I$, $\alpha \in \mathbb{C}$, is the unique solution of the following integral equation (9).

(9)
$$\psi^{o}(x,\alpha) = s(x) - s(l_{1})$$

 $+ \alpha \int_{(l_{1},x]} \{s(x) - s(y)\} \psi^{o}(y,\alpha) dm(y).$

By means of (4), (6), (7) and (9), we find (10) $n^{*}(\xi) = \lim_{x,y\to l_{1}} D_{s(x)} D_{s(y)} p^{o}(\xi, x, y)$ $= \int_{(0,\infty)} e^{-\lambda \xi} d\sigma^{o}(\lambda),$

where $D_{s(x)}$ denotes the right derivative with respect to s(x). $n^*(\xi)$ is the Lévy measure density of the inverse local time $[\tau^*(t), t \geq 0]$.

Example 2. Let $l_1 = 0$, $l_2 = l < \infty$, s(x) = x and $m(x) = C(l-x)^{-(1+1/\rho)}$, where C is a positive number and $0 < \rho < 1$. (1) is satisfied. By virtue of Proposition 1, the Laplace transform of the distribution of $[\tau^*(t), t \ge 0]$ is given by

(11)
$$E_0^*[e^{-\lambda \tau^*(t)}] = \exp\left\{-t/l - t \int_0^\infty (1 - e^{-\lambda \xi}) n^*(\xi) \, d\xi\right\},$$
(12)
$$n^*(\xi) = \int_0^\infty e^{-\lambda t} \sigma_o^o(\lambda) \, d\lambda,$$

where

(13)
$$p^{o}(t, x, y)$$

= $\int_{(0,\infty)} e^{-\lambda t} \psi^{o}(x, -\lambda) \psi^{o}(y, -\lambda) \sigma_{o}^{o}(\lambda) d\lambda$,

(14)
$$\psi^{o}(x, -\lambda) = \rho \pi \sqrt{l(l-x)} \times \{-N_{\rho}(c_{\rho}l^{-1/2\rho}\sqrt{\lambda})J_{\rho}(c_{\rho}(l-x)^{-1/2\rho}\sqrt{\lambda}) + J_{o}(c_{\rho}l^{-1/2\rho}\sqrt{\lambda})N_{o}(c_{\rho}(l-x)^{-1/2\rho}\sqrt{\lambda})\},$$

(15)
$$\sigma_o^o(\lambda) = (l\rho\pi^2)^{-1} \times \{J_\rho(c_\rho l^{-1/2\rho}\sqrt{\lambda})^2 + N_\rho(c_\rho l^{-1/2\rho}\sqrt{\lambda})^2\}^{-1}.$$

Here $c_{\rho} = 2\{C\rho(1+\rho)\}^{1/2}$, and $J_{\rho}(z)$ are $N_{\rho}(z)$ are Bessel functions. We prove (13) with (14) and (15) in another paper. Noting the asymptotic behavior of Bessel functions, we have

$$\sigma_o^o(\lambda) \sim rac{C^
ho C_1(
ho)}{l^2 \Gamma(1+
ho)} \, \lambda^
ho \quad ext{as} \quad \lambda o 0,$$

where $f(t) \sim g(t)$ as $t \to 0$ [resp. $t \to \infty$] stands for $\lim_{t \to 0} [\text{resp.} t \to \infty] f(t)/g(t) = 1$ for positive functions f(t) and g(t), and $C_1(\rho)$ is a positive number given by

(16)
$$C_1(\rho) = {\rho(1+\rho)}^{\rho}/\Gamma(\rho).$$

Therefore we find

(17)
$$n^*(\xi) \sim l^{-2} C^{\rho} C_1(\rho) \xi^{-(1+\rho)}$$
 as $\xi \to \infty$.

Example 3. Let $l_1=0, l_2=\infty, s(x)=x$ and $m(x)=Cx^{-1+1/\rho}$, where C is a positive number and $0<\rho<1$. (1) is satisfied. By virtue of Proposition 1, the Laplace transform of the distribution of $[\tau^*(t),\ t\geq 0]$ is given by

(18)
$$E_0^*[e^{-\lambda \tau^*(t)}]$$

= $\exp\left\{-t \int_0^\infty (1 - e^{-\lambda \xi}) n^*(\xi) d\xi\right\},$

(19)
$$n^*(\xi) = \int_0^\infty e^{-\lambda t} \sigma_o^o(\lambda) d\lambda = C^\rho C_2(\rho) \xi^{-(1+\rho)},$$

where $C_2(\rho)$ is a positive number given by

(20)
$$C_2(\rho) = {\rho(1-\rho)}^{\rho}/\Gamma(\rho).$$

(19) follows from the following representation of $p^o(t, x, y)$.

(21)
$$p^{o}(t, x, y)$$

= $\int_{(0,\infty)} e^{-\lambda t} \psi^{o}(x, -\lambda) \psi^{o}(y, -\lambda) \sigma_{o}^{o}(\lambda) d\lambda$,

(22)
$$\psi^{o}(x, -\lambda) = \frac{\Gamma(1+\rho)}{\{C\rho(1-\rho)\lambda\}^{\rho/2}} \sqrt{x} \times J_{\rho}(2\sqrt{C\rho(1-\rho)\lambda}x^{1/2\rho}),$$

(23)
$$\sigma_o^o(\lambda) = \frac{C^\rho C_2(\rho)}{\Gamma(1+\rho)} \lambda^\rho.$$

2. Asymptotic behavior of Lévy measure densities. In this section we consider asymptotic behavior of Lévy measure densities. We assume one

of the following (A1), (A2) and (A3), where 0 < $\rho < 1$ and L(x) is a slowly varying function.

(A1) $l_1 = 0, l_2 = l < \infty, s(x) = x \text{ and } m(x) \text{ sat-}$ isfies $|m(0)| < \infty$ and

(24)
$$m(l-1/x) \sim x^{1+1/\rho} L(x)$$
 as $x \to \infty$.

(A2) $l_1 = 0$, $l_2 = \infty$, s(x) = x and m(x) satisfies $|m(0)| < \infty$ and

(25)
$$m(x) \sim x^{-1+1/\rho} L(x)$$
 as $x \to \infty$.

(A3) $l_1 = 0$, $l_2 = \infty$, s(x) = x and m(x) satisfies $\lim_{x\to\infty} m(x) = \infty$ and

(26)
$$m(x) \sim x^{-1+1/\rho} L(x)$$
 as $x \to 0$.

Since $l_1 = 0$ is regular, we can define the inverse local time $\tau^*(t)$ at 0 by putting the reflecting boundary condition. We obtain the following asymptotic behavior of Lévy measure densities. Let K(x) be another slowly varying function such that

$$\begin{split} (27) & \lim_{x\to\infty} K(x)^{1/\rho} L(x^\rho K(x)) \\ &= \lim_{x\to\infty} L(x)^\rho K(x^{1/\rho} L(x)) = 1, \\ \text{where } x\to\infty \text{ should be read as } x\to0 \text{ when } (\mathbf{A3}) \text{ is} \end{split}$$

satisfied.

Theorem 4. Assume (A1). Then the Laplace transform of the distribution of $[\tau^*(t), t \geq 0]$ is given by the same formula as (11) and the Lévy $measure\ density\ satisfies$

(28)
$$n^*(\xi) \sim l^{-2} C_1(\rho) \xi^{-(1+\rho)} K(\xi)^{-1}$$

 $as \quad \xi \to \infty.$

Theorem 5. Assume (A2) [resp. (A3)]. Then the Laplace transform of the distribution of $[\tau^*(t), t \geq 0]$ is given by the same formula as (18) and the Lévy measure density satisfies

(29)
$$n^*(\xi) \sim C_2(\rho) \xi^{-(1+\rho)} K(\xi)^{-1}$$
$$as \quad \xi \to \infty \quad [resp. \ \xi \to 0].$$

Proof of Theorem 4. The assumption (A1) implies that (A.1) with $\theta = 0$ of [10] is satisfied, where we should replace the role of l_1 by that l_2 in (A.1) of [10]. Since $l_1 = 0$ is regular, we can put $l_1 = 0$ in (3.1) of [10]. Thus, by means of (5.11) of [10], we have

$$\int_{(0,\infty)} e^{-\lambda t} d\sigma^{o}(\lambda) \sim l^{-2} C_1(\rho) t^{-(1+\rho)} K(t)^{-1}$$

Combining this with (10), we obtain (28).

Theorem 5 follows from some results on Krein's correspondence. The arguments of Krein's correspondence are due to [4] and [6]. Let denote by \mathcal{M} the totality of nonnegative right continuous nondecreasing functions $\mu(x)$ on $[0,\infty]$ such that $\mu(x) \not\equiv \infty \text{ and } \mu(\infty) = \infty. \text{ For } \mu \in \mathcal{M} \text{ set } \mu(0-) = 0$ and let $\varphi(x,\lambda)$ be the solution of the integral

$$\varphi(x,\lambda) = 1 + \lambda \int_{[0,x]} (x-y) \varphi(y,\lambda) \, d\mu(y), \quad x \in [0,l),$$

where $\lambda \in \mathbf{C}$ and $l = \sup\{x : \mu(x) < \infty\}$. We set

$$\kappa(\alpha) = \int_0^l \varphi(x, \alpha)^{-2} dx, \quad \alpha > 0.$$

 κ is called the *characteristic function* of μ and the correspondence $\mu \in \mathcal{M} \to \kappa$ is called Krein's correspondence. Let \mathcal{K} be the set of functions κ on $(0, \infty)$ such that

$$\kappa(\alpha) = c + \int_{[0,\infty)} (\alpha + \lambda)^{-1} d\sigma(\lambda), \quad \alpha > 0,$$

for some $c \ge 0$ and some nonnegative Borel measure σ on $[0,\infty)$ satisfying $\int_{[0,\infty)} (1+\lambda)^{-1} d\sigma(\lambda) < \infty$. It is well known that Krein's correspondence is a one to one map from \mathcal{M} onto \mathcal{K} (see [4], e.g.). From now on we denote by $\mu \in \mathcal{M} \leftrightarrow \kappa \in \mathcal{K}$ Krein's correspondence. In [5] Kasahara proved the following asymptotic theorem on Krein's corespondence, where $0 < \rho < 1$, L(x) and K(x) are slowly varying functions satisfying (27), and $C_3(\rho) = \rho/\{\Gamma(1 - \rho)\}$ ρ) $C_2(\rho)$ }.

Theorem 6 ([5]). $\mu \in \mathcal{M} \leftrightarrow \kappa \in \mathcal{K}$ and l = ∞ . Then the following (30), (31) and (32) are equivalent each other.

(30)
$$\mu(x) \sim x^{-1+1/\rho} L(x)$$
 as $x \to \infty / x \to 0$ /.

(31)
$$\kappa(\alpha) \sim C_3(\rho)\alpha^{-\rho}K(1/\alpha)$$
as $\alpha \to 0 /\alpha \to \infty$.

(32)
$$\sigma(\lambda) \sim \{C_3(\rho)/\Gamma(\rho)\Gamma(2-\rho)\}\lambda^{1-\rho}K(1/\lambda)$$

 $as \quad \lambda \to 0 \ [\lambda \to \infty].$

Now we show Theorem 5.

Proof of Theorem 5. Assume (A2) [resp. (A3)]. Since $m \in \mathcal{M}$, there is the characteristic function $\kappa \in \mathcal{K}$ such that $m \leftrightarrow \kappa$. By means of Theorem 6, $\kappa(\alpha)$ satisfies (31). As we saw in Lemma 3 of [7], $1/\alpha\kappa(\alpha) \in \mathcal{K}$ and the corresponding spectral measure $d\sigma_*(\lambda)$ coincides with $d\sigma^o(\lambda)/\lambda$ for $\lambda > 0$ and $\sigma_*(\{0\}) = 0$. Noting $1/\alpha \kappa(\alpha) \sim$ $C_3(\rho)^{-1}\alpha^{\rho-1}/K(1/\alpha)$ as $\alpha \to 0$ [resp. $\alpha \to \infty$], and the relation between (31) and (32), we get

$$\sigma_*(\lambda) \sim \{\Gamma(1-\rho)\Gamma(1+\rho)C_3(\rho)\}^{-1}\lambda^{\rho}/K(1/\lambda)$$

as $\lambda \to 0$ [resp. $\lambda \to \infty$],

and hence

$$\sigma^{o}(\lambda) = \int_{(0,\lambda]} \xi \, d\sigma_{*}(\xi) \sim \frac{\rho}{1+\rho} \, \lambda \sigma_{*}(\lambda)$$
as $\lambda \to 0$ [resp. $\lambda \to \infty$].

Thus we obtain

$$\int_{(0,\infty)} e^{-\lambda t} d\sigma^{o}(\lambda) \sim C_{2}(\rho) t^{-(1+\rho)} K(t)^{-1}$$
as $t \to \infty$ [resp. $t \to 0$].

Combining this with (10), we obtain (29).

3. Inverse local time of harmonic transformed diffusion processes. In this section we consider inverse local times of harmonic transformed diffusion processes and the corresponding Lévy measure densities. Let $\mathbf{D}_{s,m}^*$ and $\mathbf{D}_{s,m}^o$ be diffusion processes on I as in Section 1. For both diffusion processes we pose the absorbing boundary condition at l_2 whenever it is regular, that is, $|s(l_2)| + |m(l_2)| < \infty$.

For $\beta \geq 0$, let h be a positive continuous function on I satisfying $\mathcal{G}_{s,m}h = \beta h$. We set

$$s_h(x) = \int_{(c_0, x]} h(y)^{-2} ds(y),$$

 $m_h(x) = \int_{(c_0, x]} h(y)^2 dm(y),$

where $c_0 \in I$ is fixed arbitrarily. Let us consider a harmonic transformed diffusion process on I whose generator is given by \mathcal{G}_{s_h,m_h} . It is known that h(x) is represented as a linear combination of $g_i(x,\beta)$ (i=1,2) such that $g_i(x,\beta)$ is positive and continuous in x, $g_1(x,\beta)$ is nondecreasing in x, $g_2(x,\beta)$ is non-increasing in x, $g_i(l_i,\beta) = 0$ if $|s(l_i)| < \infty$, and $\mathcal{G}_{s,m}g_i = \beta g_i$. Note that there exist such functions $g_i(\beta)$, i=1,2 ([3]). In the following we set

(33)
$$h(x) = B_1 g_1(x, \beta) + B_2 g_2(x, \beta),$$

where $B_1 \geq 0$, $B_2 > 0$. Since $g_1(l_1, \beta) = 0$, (33) implies $h(l_1) \in (0, \infty)$, and by virtue of Theorem 1.1 of [8], $|s_h(l_1)| + |m_h(l_1)| < \infty$, that is, l_1 is regular for harmonic transformed diffusion processes. Let $\mathbf{D}_{s_h,m_h}^* = [X(t), P_x^{h*}]$ [resp. $\mathbf{D}_{s_h,m_h}^o = [X(t), P_x^{ho}]$] the one dimensional diffusion process on I with \mathcal{G}_{s_h,m_h} as the generator and with l_1 being reflecting [resp.

absorbing]. For both diffusion processes we pose the absorbing boundary condition at l_2 whenever it is regular, that is, $|s_h(l_2)| + |m_h(l_2)| < \infty$. We denote by $[\tau^{(h*)}(t), t \geq 0]$ the inverse local time of \mathbf{D}_{s_h,m_h}^* at the end point l_1 .

We derive the following result from Proposition 1, Theorem 1.1 of [8] and Theorem 3.2 of [9].

Theorem 7. The Laplace transform of the distribution of $[\tau^{h*}(t), t \geq 0]$ is given by the following

(34)
$$E_{l_1}^{h*} \left[e^{-\lambda \tau^{h*}(t)} \right]$$

= $\exp \left\{ -\gamma^{h*} t - t \int_0^\infty (1 - e^{-\lambda \xi}) n^{h*}(\xi) d\xi \right\},$

(35)
$$\gamma^{h*} = \begin{cases} 0 & \text{if } B_1 = 0, \\ 1/\{s_h(l_2) - s_h(l_1)\} & \text{if } B_1 > 0, \end{cases}$$

(36)
$$n^{h*}(\xi) = (B_2 g_2(l_1, \beta))^2 e^{-\beta \xi} n^*(\xi),$$

where $E_{l_1}^{h*}$ stands for the expectation with respect to $P_{l_1}^{h*}$ and $n^*(\xi)$ is given by (4).

We should note that n^{h^*} is independent of B_1 . Finally we study asymptotic behavior of Lévy measure density $n^{h^*}(\xi)$. Assume that $\mathbf{D}_{s,m}^*$ satisfies one of $(\mathbf{A1})$, $(\mathbf{A2})$ and $(\mathbf{A3})$. We might suppose that the asymptotic behavior of $n^{h^*}(\xi)$ depends on those of $s_h(x)$ and $m_h(x)$ as $x \to l_2$, and hence that of h(x) as $x \to l_2$. However the asymptotic behavior of $n^{h^*}(\xi)$ is given by a quite simple form $n^*(\xi)$ multiplied by $e^{-\beta\xi}$.

Theorem 8. Assume one of (A1), (A2) and (A3). Let h be given by (33). Then (34), (35) and (36) hold. In particular, the asymptotic behavior of Lévy measure density n^{h*} is given by (36) with $n^*(\xi)$ satisfying (28) [resp. (29)] if (A1) [resp. (A2) or (A3)] is satisfied.

References

- [1] R. M. Blumenthal and R. K. Getoor, Markov processes and potential theory, Pure and Applied Mathematics, Vol. 29, Academic Press, New York, 1968.
- [2] W. Feller, The parabolic differential equations and the associated semi-groups of transformations, Ann. of Math. (2) 55 (1952), 468–519.
- [3] K. Itô and H. P. McKean, Jr., Diffusion processes and their sample paths, Springer-Verlag, New York, 1974.
- [4] I. S. Kac and M. G. Krein, On the spectral functions of the string, American Mathematical Society Translations, Series 2. Vol. 103, Amer. Math. Soc., Providence, RI, 1974.
- [5] Y. Kasahara, Spectral theory of generalized sec-

- ond order differential operators and its applications to Markov processes, Japan. J. Math. (N.S.) ${\bf 1}$ (1975/76), no. 1, 67–84.
- [6] S. Kotani and S. Watanabe, Krein's spectral theory of strings and generalized diffusion processes, in Functional analysis in Markov processes (Katata/Kyoto, 1981), 235–259, Lecture Notes in Math., 923, Springer, Berlin, 1982.
- [7] N. Minami, Y. Ogura and M. Tomisaki, Asymptotic behavior of elementary solutions of one-dimensional generalized diffusion equations, Ann. Probab. 13 (1985), no. 3, 698–715.
- [8] T. Takemura, State of boundaries for harmonic transforms of one-dimensional generalized diffusion processes, Annual Reports of Graduate School of Humanities and Sciences, Nara Women's University, 25 (2010), 285–294.
 [9] T. Takemura and M. Tomisaki, Lévy measure
- [9] T. Takemura and M. Tomisaki, Lévy measure density corresponding to inverse local time, Publ. Res. Inst. Math. Sci. 49 (2013), no. 3, 563-599.
- [10] M. Tomisaki, Asymptotic behavior of elementary solutions of transient generalized diffusion equations, J. Math. Soc. Japan **40** (1988), no. 4, 561–581.