23. The W^{k,p}-continuity of Wave Operators for Schrödinger Operators

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1. Introduction. Theorems. For Schrödinger operators $H = H_0 + V(x)$ and $H_0 = D_1^2 + \cdots + D_m^2$, $D_j = -i\partial/\partial x_j$, the wave operators W_{\pm} and Z_{\pm} are defined by the limits in $L^2 \equiv L^2(\mathbb{R}^m)$:

(1.1)
$$W_{\pm}u = \lim_{t \to \pm \infty} e^{itH} e^{-itH_0} u, \quad Z_{\pm}u = \lim_{t \to \pm \infty} e^{itH_0} e^{-itH} P_c(H) u,$$

where $P_c(H)$ is the orthogonal projection onto the continuous spectral subspace $L_c^2(H)$ for H. We assume that V(x) satisfies the following condition, where $m_* = (m-1)/(m-2)$, $\langle x \rangle = (1+|x|^2)^{1/2}$ and \mathcal{F} is the Fourier transform. We take and fix $\sigma > 2/m_*$, $\delta > \max(m+2, 3m/2-2)$ and an integer $l \ge 0$.

Assumption 1.1. V(x) is a real valued function on \mathbb{R}^m , $m \ge 3$, such that $\mathcal{F}(\langle x \rangle^{\sigma} D_x^{\alpha} V) \in L^{m_*}$ for any $|\alpha| \le l$ and satisfies either (1) $|| \mathcal{F}(\langle x \rangle^{\sigma} V) ||_{L_{m_*}} \equiv C(V)$ is sufficiently small or (2) m = 2m' - 1 is odd and $|D^{\alpha}V(x)| \le C_{\alpha} \langle x \rangle^{-\delta}$ for any $|\alpha| \le \max\{l, m' - 4 + l\}$.

Under the assumption, V is H_0 -bounded and is short-range in the sense of Agmon [1]. Hence H with domain $D(H) = D(H_0) = W^{2,2}$ is selfadjoint and both limits in (1.1) exist ([1], [8]); W_{\pm} are partial isometries from L^2 onto $L_c^2(H)$ and $Z_{\pm} = W_{\pm}^*$. Consequently, the continuous part H_c of H is unitarily equivalent to H_0 and, for any Borel function f, $f(H)P_c(H) = W_{\pm}f(H_0)W_{\pm}^*$, $f(H_0) = W_{\pm}^*f(H)P_c(H)W_{\pm}$. The main result of this paper is the following

Theorem 1.1. Let V satisfy Assumption 1.1 and let 0 be neither eigenvalue nor resonance of H. Then, for any $1 \le p \le \infty$ and integral $0 \le k \le l$, W_{\pm} and Z_{\pm} originally defined on $L^2 \cap W^{k,p}$ can be extended to bounded operators in $W^{k,p}$.

Remark 1.1. We say 0 is resonance of H if $-\Delta u(x) + V(x)u(x) = 0$ has a solution u such that $\langle x \rangle^{-\gamma}u(x) \in L^2$ for any $\gamma > 1/2$ but not for $\gamma = 0$. Under the assumption, 0 is not resonance if $m \ge 5$, and is neither eigenvalue nor resonance if C(V) is small enough.

Remark 1.2. If 0 is resonance, Theorem 1.1 never holds. If 0 is eigenvalue of H, then it does not hold in general. This can be seen by comparing the results of [3] or [9] with Theorem 1.3 below.

In the sequel, we always assume that the condition of Theorem 1.1 is satisfied. For Banach spaces X and Y, B(X, Y) is the space of bounded operators from X to Y, B(X) = B(X, X). Theorem 1.1 yields the following

Theorem 1.2. Let $1 \le p$, $q \le \infty$ and let $0 \le k$, $k' \le l$ be integers. Then: $C^{-1} \| f(H_0) \|_{B(W^{k,p}, W^{k',q})} \le \| f(H) P_c(H) \|_{B(W^{k,p}, W^{k',q})} \le C \| f(H_0) \|_{B(W^{k,p}, W^{k',q})}$

where the constant C > 0 is independent of Borel functions f.

An immediate corollary of Theorem 1.2 is the $L^{p} - L^{q}$ estimate for time dependent Schrödinger, Klein-Gordon and wave equations. Under slightly different conditions on V, such estimate has been proven recently for Schrödinger and wave equations ([5], [2]). See [4] for related results.

Theorem 1.3. Let $0 \le k \le l$ be integral, $2 \le p \le \infty$ and 1/p + 1/q = 1. Then:

 $\|e^{-itH}P_{c}(H)u\|_{W^{k,p}} \leq C_{kp} |t|^{m(1/p-1/2)} \|u\|_{W^{k,q}}, \quad u \in L^{2} \cap W^{k,q}.$

Theorem 1.4. Let $0 \le k \le l$ be integral, $2 \le p \le 2(m+1)/(m-1)$ and 1/p + 1/q = 1. Then, there exists a constant $C_{kp} > 0$ such that for any $\phi, \phi \in L^2_c(H) \cap W^{k,q}$ the solution u(t, x) of the Cauchy problem $\partial^2 u / \partial t^2 = \Delta u - \mu^2 u - Vu$, $u(0, x) = \phi(x)$, $u_t(0, x) = \psi(x)$ satisfies

$$\| u(t, \cdot) \|_{W^{k,p}} \le C_{kp} \| t \|^{1+m(1/p-1/q)} (\| \phi \|_{W^{k,q}} + \| \sqrt{H + \mu^2} \phi \|_{W^{k,q}}).$$
we if $k \le l-1$ $\| \sqrt{H + \mu^2} \phi \|_{W^{k,q}}$ may be replaced by $\| \phi \|_{W^{k,q}}$.

Moreover, if $k \leq l-1$, $\|\sqrt{H} + \mu^2 \phi\|_{W^{k,q}}$ may be replaced by $\|\phi\|_{W^{k+1,q}}$.

Another consequence of Theorem 1.1 is on a multiplier theorem for the generalized Fourier transforms. We assume (2) of Assumption 1.1. Then, for any $\xi \in \mathbb{R}^m \setminus \{0\}$, there exists a unique solution of $(-\Delta + V(x))\phi_{\pm}(x, \xi) = |\xi|^2 \phi_{\pm}(x, \xi)$ satisfying the radiation condition: $\phi_{\pm}(x, \xi) = e^{ix\cdot\xi} + e^{\pm i|x||\xi|} |x|^{-(m-1)/2} (g(\hat{x}, \xi) + O(|x|^{-1}))$. Define

$$\mathscr{F}_{\pm}u(\xi) = (2\pi)^{-m/2} \int_{\mathbb{R}^m} \overline{\phi_{\pm}(x,\xi)} u(x) dx.$$

 $\begin{aligned} \mathscr{F}_{\pm} & \text{ are unitary from } L^2_c(H) \text{ onto } L^2(R^m) \text{ ; } W_{\pm} = \mathscr{F}_{\pm}^*\mathscr{F} \text{ ; and } \mathscr{F}_{\pm} \text{ diagonalize} \\ H_c : \mathscr{F}_{\pm}H_c\mathscr{F}_{\pm}^*g(\xi) = |\xi|^2g(\xi). \text{ Identity } \mathscr{F}_{\pm}^*\mathscr{F}_{\pm} = P_c(H) \text{ is equivalent to the} \\ \text{generalized eigenfunction expansions: } u(x) = (2\pi)^{-m/2} \int_{R^m} \phi_{\pm}(x,\xi) \\ \mathscr{F}_{\pm}u(\xi)d\xi, \ u \in L^2_c(H). \end{aligned}$

For a function f on \mathbb{R}^m , M_f is the multiplication operator with $f(\xi)$ and $f(D) = \mathcal{F}^* M_f \mathcal{F}$.

Theorem 1.5. Let V satisfy Assumption 1.1, (2) and let $1 \le p, q \le \infty$. Then for any Borel function f, we have

 $C^{-1} \| f(D) \|_{B(W^{1,p},W^{1,q})} \leq \| \mathcal{F}_{\pm}^* M_f \mathcal{F}_{\pm} \|_{B(W^{1,p},W^{1,q})} \leq C \| f(D) \|_{B(W^{1,p},W^{1,q})},$ where the constant C is independent of f.

Combining Theorem 1.5 with known Fourier multiplier theorems, we can give explicit conditions on $f(\xi)$ for $\mathscr{F}_{\pm}^*M_{f}\mathscr{F}_{\pm}$ to be bounded in $W^{l,p}$.

2. Outline of the proof of Theorem 1.1. We prove the case l = 0 first. We treat W_+ only. $R(z) = (H-z)^{-1}$ and $R_0(z) = (H_0 - z)^{-1}$ are the resolvents. Assumption 1.1 implies $||V||_{L^{(m\pm\epsilon)/2}} \leq CC(V)$, $\varepsilon > 0$. We decompose V(x) = A(x)B(x) with $A, B \in L^{m+\epsilon} \cap L^{m-\epsilon}$. Then A and B are supersmooth ([7]) and Kato's theory of smooth operator ([6]) implies that W_+f can be written as $W_+f = \sum_{n=0}^{N} (-1)^n W_n f + L_N f$ for $N = 0, 1 \cdots$. Here W_0 is the identity operator, W_n^* , $n \geq 1$, which we estimate rather than W_n itself, and L_N are written as

(2.1)
$$W_n^* f = \frac{-1}{2\pi i} \int_{-\infty}^{\infty} R_0 (\lambda + i0) \left(V R_0 (\lambda - i0) \right)^n f d\lambda$$

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$$= (-i)^{n} \int_{[0,\infty)^{n}} e^{itH_{0}} V e^{-it_{1}H_{0}} V \cdots V e^{-it_{n}H_{0}} f dt_{1} \cdots dt_{n}, \quad t = t_{1} + \cdots + t_{n},$$
(2.2) $L_{N}f = \frac{1}{2\pi i} \int_{0}^{\infty} (R_{0}(\lambda - i0)V)^{N} R(\lambda - i0) \times V\{R_{0}(\lambda + i0) - R_{0}(\lambda - i0)\} f d\lambda.$

Moreover, $||W_n||_{B(L^2)} \leq (CC(V))^n$ and $||L_n||_{B(L^2)} \leq (CC(V))^{n+1}$ for $n = 0, 1, \cdots$ and, when C(V) is small, $W_+ = \sum_{n=0}^{\infty} (-1)^n W_n$ converges in the operator norm in $B(L^2)$.

Suppose for the moment $\hat{V} \in C_0^{\infty}$. Set $K_n(k_1, \ldots, k_n) = i^n (2\pi)^{-nm/2} 2^{-n}$ $\prod_{j=1}^n \hat{V}(k_j - k_{j-1})$, where $k_0 = 0$. Define for $(t_1, \ldots, t_n) \in \mathbb{R}^n$ and $(\omega_1, \ldots, \omega_n) \in \Sigma^n$, Σ is the unit sphere of \mathbb{R}^m : $\hat{K}_n(t_1, \ldots, t_n, \omega_1, \ldots, \omega_n) =$

$$\int_{[0,\infty)^n} e^{-i\Sigma_{j=1}^n t_j s_j/2} (s_1 \cdots s_n)^{m-2} K(s_1 \omega_1, \dots, s_n \omega_n) ds_1 \cdots ds_n.$$

Lemma 2.1. For $n = 1, 2, ..., W_n^{\dagger} f(x)$ can be written in the form

(2.3)
$$W_n^* f(x) = \int_{(0,\infty)^{n-1} \times I \times \Sigma^n} \hat{K}_n(t_1,\ldots,t_{n-1},\tau,\omega_1,\ldots,\omega_n) \times f(\bar{x}+\rho) dt_1 \cdots dt_{n-1} d\tau d\omega_1 \cdots d\omega_n$$

where, $\bar{y} = x - 2(\omega_n \cdot y)\omega_n$ is the reflection of y along the ω_n axis, $\rho = t_1\bar{\omega}_1 + \cdots + t_{n-1}\bar{\omega}_{n-1} - \tau\omega_n$, and where $I = (-\infty, 2\omega_n x - \sigma)$, $\sigma = 2\omega_n(x + t_1\omega_1 + \cdots + t_{n-1}\omega_{n-1})$, is the range of the integration by the variable τ . Note that ρ does not depend on x. Extending the range of integration by τ to

the whole line after taking the absolute values of both sides of (2.3), we have

$$|W_n^*f(x)| \leq \int_{[0,\infty)^{n-1} \times R \times \Sigma^n} |\hat{K}_n(t_1,\ldots,t_{n-1},\tau,\omega_1,\ldots,\omega_n) \times f(\bar{x}+\rho)| dt_1 \cdots dt_{n-1} d\tau d\omega_1 \cdots d\omega_n.$$

Noting that $x \rightarrow \bar{x}$ is an isometry, we obtain by applying Minkowski's inequality:

Lemma 2.2. We have $\| W_n^* f \|_{L^p} \leq 2 \| \hat{K}_n \|_{L^{1}([0,\infty)^n \times \Sigma^n)} \| f \|_{L^p}, 1 \leq p \leq \infty$, for $n = 1, 2, \cdots$. Let $X = L^{m-1}([0, \infty)^n, L^1(\Sigma^n))$. We have $\| \hat{K}_n \|_X \leq C^n \| K_n \|_{L^{m*}(\mathbb{R}^{mn})} \leq C^n \| \hat{V} \|_{L^{m*}}^n$ by Hölder and Hausdorff-Young inequalities and likewise $\| (\Pi_{j=1}^n \langle t_j \rangle) \hat{K}_n \|_X \leq C^n \| \mathcal{F}(\langle x \rangle^2 V) \|_{L^{m*}}^n$. Hence, by the multilinear interpolation inequality and the assumption $\sigma > 2/m_*$, we obtain

$$\|\hat{K}_{n}\|_{L^{1}([0,\infty)^{n},L^{1}(\Sigma^{n}))} \leq C^{n} \|(\prod_{j=1}^{n} \langle t_{j} \rangle^{\sigma/2}) \hat{K}_{n}\|_{X} \leq (C_{1}C(V))^{n}.$$

Combining this with Lemma 2.2, we thus have

Lemma 2.3. We have $|| W_n^* f ||_{L^p} \le C_1 (C_2 C(V))^n || f ||_{L^p}, 1 \le p \le \infty$.

It is easy to see by the density argument that Lemma 2.3 extends to any V with $\mathscr{F}(\langle x \rangle^{\sigma} V) \in L^{m_*}$ and, if C(V) is small, the series $W_+ = \sum_{n=0}^{\infty} (-1)^n W_n$ converges in the operator norm in $B(L^{\flat})$. This proves Theorem 1.1 in the case C(V) is small.

Suppose Assumption 1.1, (2) now. m = 2m' - 1. Since W_j , $j = 1, \ldots, m$, are bounded in L^p as shown above, we have only to show that L_m is also

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bounded in L^{\flat} . We estimate its integral kernel L(x, y). Setting $N(k) = \{R_0(k^2 - i0)V\}^{m'-1}R(k^2 - i0)\{VR_0(k^2 - i0)\}^{m'-1}$, we rewrite: (2.4) $L_m = \frac{1}{\pi i} \int_0^{\infty} R_0(k^2 - i0)VN(k)V\{R_0(k^2 + i0) - R_0(k^2 - i0)\}kdk$. Let $G_{\pm}(x, k) = \pm (i/4(2\pi)^{\nu})|x|^{2-m}(k|x|)^{\nu}H_{\nu}^{(j)}(k|x|)$ be the outgoing (incoming) fundamental solutions of $-\Delta - k^2$, where $H_{\nu}^{(j)}(r), \nu = (m-2)/2$, is the Hankel function of *j*-th kind and j = 1 for + case and j = 2 for -. The integral kernel of $R_0(k^2 \pm i0)$ is given by $G_{\pm,x,k}(y) = G_{\pm}(x - y, k)$, and, consequently, that of the integrand of (2.4) is given by $(N(k)V(G_{+,y,k} - G_{-,y,k}), VG_{+,x,k})$. Set $G_{\pm,x,k}(y) = e^{\pm ik|x|}\tilde{G}_{\pm,x,k}(y)$ and define $T_{\pm}(x, y, k) = (N(k)V\tilde{G}_{\pm,y,k}, V\tilde{G}_{+,x,k})$ and

(2.5)
$$L_{\pm}(x, y) = \frac{1}{i\pi} \int_{0}^{\infty} e^{-ik(|x|\mp|y|)} T_{\pm}(x, y, k) k dk,$$

so that $L(x, y) = L_+(x, y) - L_-(x, y)$. The mapping properties of $R(\lambda \pm i0)$ and $R_0(\lambda \pm i0)$ given in [9] and [3] imply

Lemma 2.4. Let $j = 0, \ldots, m' + 1, \gamma, \gamma' > j + 1/2$ and $0 \le s, s' \le m' - 2$. Then $\langle x \rangle^{-\gamma} N(k) \langle x \rangle^{-\gamma'}$ is a $B(W^{-s',2}, W^{s,2})$ -valued C^{j} -function of k and

(2.6) $\| (d/dk)^{j} \langle x \rangle^{-\tau} N(k) \langle x \rangle^{-\tau'} \|_{B(W^{-s',2},W^{s,2})} \leq C \langle k \rangle^{-(m-s-s')}.$

We estimates $L_{\pm}(x, y)$ by performing integration by parts in (2.5), using Lemma 2.4, the decay property as $|x| \to \infty$ of $||\langle y \rangle^{-r} (d/dx)^{j} \tilde{G}_{+,x,k} ||_{L^{q}(\mathbb{R}^{m})}$ for suitable γ , j and q and the fact that $e^{\pm ir} r^{\nu} H_{\nu}^{(j)}(r)$ are polynomials of degree m'-1 for odd m = 2m'-1. Cancellations of the boundary terms of the integral (2.5) for $L_{+}(x, y)$ and $L_{-}(x, y)$ occur and we obtain

Lemma 2.5. We have $\sup_{y \in R^m} \int_{R^m} |L(x, y)| dx < \infty$ and $\sup_{x \in R^m} \int_{R^m} |L(x, y)| dy < \infty$.

This is a well known criterion for L_m to be bounded in L^{\flat} and concludes the proof of Theorem 1.1 for the case l = 0.

When l = 1, we compute $D_j W_n f$ and $D_j L_m f$, j = 1, ..., n. For example, using the first equation of (2.1), $D_j W_n f - W_n D_j f$ can be computed as

$$\sum_{j=0}^{n} \frac{1}{2\pi i} \int_{-\infty}^{\infty} (R_0(\lambda - i0) V)^j R_0(\lambda - i0) (D_j V) (R_0(\lambda - i0) V)^{n-j-1} R_0(\lambda + i0) f d\lambda.$$

This is a sum of terms which have exactly the same form as the adjoint of the second of (2.1) except that one of V is replaced by $D_j V$. Thus, the argument which leads to Lemma 2.3 above implies

 $||D_{j}W_{n}f||_{L^{p}} \leq (n+1)(C_{1}C(V))^{n-1}(C(V) + C(D_{j}V))||f||_{W^{1,p}}, 1 \leq p \leq \infty$. This shows that the series $\sum_{n=1}^{\infty} W_{n}$ converges in the operator norm in $B(W^{1,p})$ for the same value of C(V) for which it converges in $B(L^{p})$. Thus W_{+} is bounded in $W^{1,p}$. The argument for $D_{j}L_{m}f$ is similar and Theorem 1.1 holds for l = 1. For general $l \geq 2$, we repeat this argument. The detail will appear elsewhere.

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