

116. A Stationary Free Boundary Problem for a Circular Flow with or without Surface Tension^{*)}

By Hisashi OKAMOTO

Department of Mathematics, University of Tokyo

(Communicated by Kôzaku YOSIDA, M. J. A., Dec. 13, 1982)

§ 1. In this note we are concerned with a free boundary problem which is a model for a flow around a planet. The problem is stated as follows.

Problem. Given a unit circle Γ in \mathbb{R}^2 , find a closed Jordan curve γ outside Γ and a function V such that

$$(1.1) \quad \Delta V = 0 \quad \text{in } \Omega_\gamma,$$

$$(1.2) \quad V|_\Gamma = 0, \quad V|_\gamma = a,$$

$$(1.3) \quad \frac{1}{2} |\nabla V|^2 + Q + \sigma K_\gamma = \text{unknown constant on } \gamma,$$

$$(1.4) \quad |\Omega_\gamma| = \omega_0.$$

Here Ω_γ is a doubly connected domain between Γ and γ (see Fig. 1). Constants $a > 0$, $\omega_0 > 0$ and $\sigma \geq 0$ are given. σ is the surface tension coefficient. Q is a given function defined outside Γ . K_γ is the curvature of γ ($K_\gamma > 0$ if γ is convex). $|\Omega_\gamma|$ denotes the area of Ω_γ .

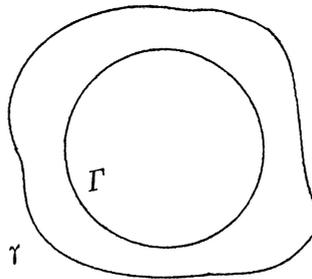


Fig. 1

Remark. We have assumed that the fluid is perfect, irrotational and that V is a stream function for the flow. Ω_γ is the flow region.

The more precise physical meaning of this problem will be explained in a forthcoming paper where we will give proofs of theorems in § 2.

Trivial solution. If Q is radially symmetric, i.e., $Q = Q_0(r)$ ($r = (x^2 + y^2)^{1/2}$), then there exists the following trivial solution. Take a number $r_0 > 1$ satisfying $\pi r_0^2 - \pi = \omega_0$. Then a circle γ_0 of radius r_0 with the origin as its center is a solution for any $\sigma \geq 0$. In fact the

^{*)} Partially supported by the Fûjukai.

corresponding stream function V is given by

$$(1.5) \quad V = V(r) = \frac{a}{\log r_0} \log r \quad (1 < r < r_0).$$

The unknown constant in (1.3) is given by

$$\frac{1}{2} \left(\frac{a}{r_0 \log r_0} \right)^2 + Q_0(r_0) + \frac{\sigma}{r_0}.$$

Our aim is to investigate perturbation and bifurcation of this trivial solution.

§ 2. **Mathematical formulation and results.** We prepare some symbols.

$$\begin{aligned} \Omega &= \{(x, y) \in \mathbf{R}^2; 1 < x^2 + y^2 < \infty\}, \\ S^1 &= \{(x, y) \in \mathbf{R}^2; 1 = x^2 + y^2\}, \\ C^{m+\alpha}(\bar{\Omega}), C^{m+\alpha}(S^1) & \quad (m = 0, 1, 2, \dots, 0 < \alpha < 1); \end{aligned}$$

the Hölder spaces with the usual norm $\| \cdot \|_{m+\alpha, \Omega}$ or $\| \cdot \|_{m+\alpha, S^1}$, resp. We fix a number $\alpha \in]0, 1[$ and a $Q_0 \in C^{2+\alpha}([1, \infty))$. The case of $Q_0(r) = -g/r$ with a positive constant g or $Q_0 \equiv 0$ is physically typical.

When a small $u \in C^{3+\alpha}(S^1)$ is given, we denote by γ_u a closed Jordan curve which is parametrized in the polar coordinates as $(r_0 + u(\theta), \theta)$ ($0 \leq \theta < 2\pi$). Hereafter we identify a function on S^1 with a 2π -periodic function on \mathbf{R} . We denote a domain between Γ and γ_u by Ω_u . The curvature of γ_u is denoted by K_u . V_u is the unique solution of the Dirichlet problem

$$(2.1) \quad \Delta V_u = 0 \quad \text{in } \Omega_u,$$

$$(2.2) \quad V_u|_{\Gamma} = 0, \quad V_u|_{\gamma_u} = a.$$

For $u \in C^{3+\alpha}(S^1)$, $Q \in C^{2+\alpha}(\bar{\Omega})$ and $\xi \in \mathbf{R}$, we put

$$(2.3) \quad F_1(a, Q; u, \xi) = \left(\frac{1}{2} |\nabla V|^2 + Q \right) \Big|_{\gamma_u} + \sigma K_u - \xi_0 - \xi,$$

where

$$(2.4) \quad \begin{aligned} \xi_0 &= \frac{1}{2} \left(\frac{a}{r_0 \log r_0} \right)^2 + Q_0(r_0) + \sigma/r_0, \\ F_2(a, Q; u, \xi) &= \frac{1}{2} \int_0^{2\pi} (r_0 + u(\theta))^2 d\theta - \pi - \omega_0, \end{aligned}$$

$$(2.5) \quad F(a, Q; u, \xi) = (F_1(a, Q; u, \xi), F_2(a, Q; u, \xi)).$$

Then it is easy to see that $F(a, Q_0; 0, 0) = (0, 0)$ and that γ_u is a solution for Q if and only if $F(a, Q; u, \xi) = (0, 0)$ for some $\xi \in \mathbf{R}$. Note that $F(a, \cdot; \cdot, \cdot)$ is a well-defined continuous mapping from a neighborhood of $(Q_0; 0, 0)$ in $C^{2+\alpha}(\bar{\Omega}) \times C^{3+\alpha}(S^1) \times \mathbf{R}$ into $C^{1+\alpha}(S^1) \times \mathbf{R}$.

Now perturbation of the trivial solution is possible in the following sense.

Theorem 1. *Assume that $\sigma > 0$. Define a_n by*

$$a_n = \left(\frac{\sigma(n^2 - 1)/r_0^2 + (\partial Q_0/\partial r)(r_0)}{r_0^{-1} + nr_0^{-1}(r_0^n + r_0^{-n})/(r_0^n - r_0^{-n})} \right)^{1/2} r_0 \log r_0$$

for

$$n \in \tilde{N} = \left\{ n \in N; \sigma(n^2 - 1)/r_0^2 + \frac{\partial Q_0}{\partial r}(r_0) \geq 0 \right\}.$$

Then, for any $a \notin \{a_n\}_{n \in \tilde{N}}$, there exists a positive constant δ such that for any $Q \in C^{2+\alpha}(\bar{\Omega})$ satisfying $\|Q - Q_0\|_{2+\alpha, \Omega} < \delta$ we have a solution $\{u, \xi\}$ of the equation $F(a, Q; u, \xi) = (0, 0)$. The solution is unique in some neighborhood of the origin.

Theorem 2. Assume that $\sigma = 0$. Let $Q_0 \in C^{1+\alpha}(\bar{\Omega})$ satisfy

$$(2.6) \quad \frac{\partial Q_0}{\partial r}(r_0) < a^2/r_0^3 (\log r_0)^2.$$

Then there exists a positive constant ε such that for any $Q \in C^{10+(1/2)+\alpha}(\bar{\Omega})$ satisfying $\|Q - Q_0\|_{10+(1/2)+\alpha} < \varepsilon$ we have a solution $\{u, \xi\} \in C^{6+\alpha}(S^1) \times \mathbf{R}$ of the equation $F(a, Q; u, \xi) = (0, 0)$.

The following two theorems are concerned with bifurcation phenomena.

Theorem 3. Fix a natural number n . Assume that $\sigma > 0$. Assume also that

$$(2.7) \quad a_m \neq a_n \quad \text{for all } m \neq n.$$

Then there exists a branch of non-trivial solution of $F(a, -g/r; u, \xi) = (0, 0)$ through $(a_n, 0, 0)$. If n is sufficiently large, then the bifurcation occurs subcritically.

Theorem 4. Assume that $\sigma = 0$. Assume also (2.6). Then, in some $C^{1+\alpha}$ -neighborhood of r_0 , there exists no solution other than r_0 .

Remark. The condition (2.7) is satisfied if n is sufficiently large or if $2\sigma \geq g$.

Theorem 1 is proved by a classical implicit function theorem, while a Nash-Moser implicit function theorem is used to prove Theorem 2. Theorem 3 is a consequence of a bifurcation theory due to Sattinger, Golubitsky and Schaeffer. Theorem 4 is proved by the maximum principle. Details will be published in the future.

Acknowledgment. The writer would like to express his deep gratitude to Profs. H. Fujita and H. Kawarada for their valuable advice and helpful discussions.

References

- [1] Golubitsky, M., and Schaeffer, D.: Imperfect bifurcation in the presence of symmetry. *Comm. Math. Phys.*, **67**, 205-232 (1979).
- [2] Landau, L. D., and Lifschitz, E. M.: *Fluid Mechanics*. Pergamon, London (1959).
- [3] Sattinger, D. H.: Group theoretic methods in bifurcation theory. *Lect. Notes in Math.*, vol. 762, Springer-Verlag, Berlin, Heidelberg, New York (1979).