

QUANTIZATION OF LOCALLY SYMMETRIC KÄHLER MANIFOLDS

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Abstract. We introduce noncommutative deformations of locally symmetric Kähler manifolds. A Kähler manifold M is said to be a locally symmetric Kähler manifold if the covariant derivative of the curvature tensor is vanishing. An algebraic derivation process to construct a locally symmetric Kähler manifold is given. As examples, star products for noncommutative Riemann surfaces and noncommutative $\mathbb{C}\mathbb{P}^N$ are constructed.

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1. Review of the Deformation Quantization with Separation of Variables

In this section, we review the deformation quantization with separation of variables to construct noncommutative Kähler manifolds.

An N -dimensional Kähler manifold M is described by using a Kähler potential. Let Φ be a Kähler potential and ω be a Kähler two-form

$$\omega := ig_{k\bar{l}}dz^k \wedge d\bar{z}^l, \quad g_{k\bar{l}} := \frac{\partial^2 \Phi}{\partial z^k \partial \bar{z}^l} \quad (1)$$

where z^i, \bar{z}^i ($i = 1, 2, \dots, N$) are complex local coordinates.

In this article, we use the Einstein summation convention over repeated indices. The $g^{\bar{k}l}$ is the inverse of the Kähler metric tensor $g_{k\bar{l}}$. That means $g^{\bar{k}l}g_{l\bar{m}} = \delta_{\bar{k}\bar{m}}$.

In the following, we use

$$\partial_k = \frac{\partial}{\partial z^k}, \quad \partial_{\bar{k}} = \frac{\partial}{\partial \bar{z}^k}. \tag{2}$$

Deformation quantization is defined as follows.

Definition 1 (Deformation quantization). *Deformation quantization of Poisson manifolds is defined as follows. \mathcal{F} is defined as a set of formal power series: $\mathcal{F} := \left\{ f \mid f = \sum_k f_k \hbar^k; f_k \in C^\infty(M) \right\}$. A star product is defined as*

$$f * g = \sum_k C_k(f, g) \hbar^k \tag{3}$$

such that the product satisfies the following conditions

1. $(\mathcal{F}, +, *)$ is a (noncommutative) algebra.
2. $C_k(\cdot, \cdot)$ is a bidifferential operator.
3. C_0 and C_1 are defined as $C_0(f, g) = fg$, $C_1(f, g) - C_1(g, f) = \{f, g\}$ where $\{f, g\}$ is the Poisson bracket.
4. $f * 1 = 1 * f = f$.

Karabegov introduced a method to obtain a deformation quantization of a Kähler manifold in [6]. His deformation quantization is called deformation quantizations with separation of variables

Definition 2 (A star product with separation of variables). *The operation $*$ is called a star product with separation of variables on a Kähler manifold when $a * f = af$ for an arbitrary holomorphic function a and $f * b = fb$ for an arbitrary anti-holomorphic function b .*

We use

$$D^{\bar{l}} = g^{\bar{l}k} \partial_k$$

and introduce $\mathcal{S} := \left\{ A; A = \sum_\alpha a_\alpha D^\alpha, a_\alpha \in C^\infty(M) \right\}$, where α is a multi-index $\alpha = (\alpha_1, \alpha_2, \dots, \alpha_n)$.

In this article, we also use the Einstein summation convention over repeated multi-indices and $a_\alpha D^\alpha := \sum_\alpha a_\alpha D^\alpha$.

There are some useful formulae. $D^{\bar{l}}$ satisfies the following equations.

$$[D^{\bar{l}}, D^{\bar{m}}] = 0, \quad [D^{\bar{l}}, \partial_{\bar{m}} \Phi] = \delta^{\bar{l}}_{\bar{m}}, \quad \text{for all } l, m \tag{4}$$

where $[A, B] = AB - BA$. Using them, one can construct a star product as a differential operator L_f such that $f * g = L_f g$.

Theorem 1. [Karabegov [6]]. For an arbitrary Kähler form ω , there exist a star product with separation of variables $*$ and it is constructed as follows. Let f be an element of \mathcal{F} and $A_n \in \mathcal{S}$ be a differential operator whose coefficients depend on f , i.e.,

$$A_n = a_{n,\alpha}(f)D^\alpha, \quad D^\alpha = \prod_{i=1}^n (D^{\bar{i}})^{\alpha_i}, \quad (D^{\bar{i}}) = g^{\bar{i}l} \partial_l \quad (5)$$

where α is an multi-index $\alpha = (\alpha_1, \alpha_2, \dots, \alpha_n)$. Then

$$L_f = \sum_{n=0}^{\infty} \hbar^n A_n \quad (6)$$

is uniquely determined such that it satisfies the following conditions.

1. For $R_{\partial_{\bar{l}}\Phi} = \partial_{\bar{l}}\Phi + \hbar \partial_{\bar{l}}$

$$[L_f, R_{\partial_{\bar{l}}\Phi}] = 0. \quad (7)$$

- 2.

$$L_f 1 = f * 1 = f. \quad (8)$$

Then the star products are given by

$$L_f g := f * g \quad (9)$$

and the star products satisfy the associativity

$$L_h(L_g f) = h * (g * f) = (h * g) * f = L_{L_h g} f. \quad (10)$$

Recall that each two of $D^{\bar{i}}$ commute each other, so if a multi index α is fixed then the A_n is uniquely determined. The equations (8)-(10) imply that $L_f g = f * g$ gives deformation quantization.

Definition 3. A map from differential operators to formal polynomials is defined as

$$\sigma(A; \xi) := \sum_{\alpha} a_{\alpha} \xi^{\alpha}$$

where

$$A = \sum_{\alpha} a_{\alpha} D^{\alpha}.$$

This map is called “twisted symbol”. It becomes easier to calculate commutators by using the following theorem.

Proposition 2 (Karabegov [6]). Let $a(\xi)$ be a twisted symbol of an operator A . Then the twisted symbol of the operator $[A, \partial_i \Phi]$ is equal to $\partial a / \partial \xi^{\bar{i}}$

$$\sigma([A, \partial_i \Phi]) = \frac{\partial}{\partial \xi^{\bar{i}}} \sigma(A).$$

This proposition follows from (4), i.e.,

$$\sigma([D^{\bar{l}}, \partial_i \Phi]) = \delta_{\bar{l}}^i.$$

2. Deformation Quantization with Separation of Variables for a Locally Symmetric Kähler Manifold

In this section, explicit formulas to obtain star products on local symmetric Kähler manifolds are constructed. A method of Karabegov in Section 1 is used for the constructing.

Operators $D^{\vec{\alpha}_n}$ and $D^{\vec{\beta}_n^*}$ are defined by using $D^k = g^{k\bar{m}} \partial_{\bar{m}}$ and $D^{\bar{j}} = g^{\bar{j}l} \partial_l$ as

$$D^{\vec{\alpha}_n} := D^{\alpha_1^n} D^{\alpha_2^n} \dots D^{\alpha_N^n}, \quad D^{\vec{\beta}_n} := D^{\beta_1} D^{\beta_2} \dots D^{\beta_N}$$

where

$$D^{\alpha_k} := (D^k)^{\alpha_k}, \quad D^{\beta_j} := (D^{\bar{j}})^{\beta_j}$$

and $\vec{\alpha}_n$ and $\vec{\beta}_n^*$ are N -dimensional vectors whose summation of their all elements are set to be n

$$\vec{\alpha}_n \in \left\{ (\gamma_1^n, \gamma_2^n, \dots, \gamma_N^n) \in \mathbb{Z}^N ; \sum_{k=1}^N \gamma_k^n = n \right\}$$

$$\vec{\beta}_n^* \in \left\{ (\gamma_1^n, \gamma_2^n, \dots, \gamma_N^n)^* \in \mathbb{Z}^N ; \sum_{k=1}^N \gamma_k^n = n \right\}$$

i.e.,

$$\vec{\alpha}_n := (\alpha_1^n, \alpha_2^n, \dots, \alpha_N^n), \quad |\vec{\alpha}_n| := \sum_{k=1}^N \alpha_k^n = n$$

$$\vec{\beta}_n^* := (\beta_1^n, \beta_2^n, \dots, \beta_N^n)^*, \quad |\vec{\beta}_n^*| := \sum_{k=1}^N \beta_k^n = n.$$

For $\vec{\alpha}_n \notin \mathbb{Z}_{\geq 0}^N$ we define $D^{\vec{\alpha}_n} := 0$.

For example, $D^{(1,2,3)} = D^1 (D^2)^2 (D^3)^3$, $D^{(2,4,0)^*} = (D^{\bar{1}})^2 (D^{\bar{2}})^4$ and $D^{(5,-2,3)} = 0$ for a three-dimensional manifolds case with $n = 6$.

\vec{e}_i is used as a N -dimensional vector

$$\vec{e}_i = (\delta_{1i}, \delta_{2i}, \dots, \delta_{Ni}). \quad (11)$$

A Riemannian (Kähler) manifold (M, g) is called a locally symmetric Riemannian (Kähler) manifold when $\nabla_m R_{ijk}{}^l = 0$ for all i, j, k, l, m . Only locally symmetric Kähler manifolds are discussed.

We assume that a star product with separation of variables for smooth functions f and g on a locally symmetric Kähler manifold M has a form

$$L_f g = f * g = \sum_{n=0}^{\infty} \sum_{\vec{\alpha}_n, \vec{\beta}_n^*} T_{\vec{\alpha}_n, \vec{\beta}_n^*}^n \left(D^{\vec{\alpha}_n} f \right) \left(D^{\vec{\beta}_n^*} g \right) \quad (12)$$

where $T_{\vec{\alpha}_n, \vec{\beta}_n^*}^n$ are covariantly constants. If $\vec{\alpha}_n \notin \mathbb{Z}_{\geq 0}^N$ or $\vec{\beta}_n \notin \mathbb{Z}_{\geq 0}^N$ then we define $T_{\vec{\alpha}_n, \vec{\beta}_n^*}^n := 0$. $\sum_{\vec{\alpha}_n, \vec{\beta}_n^*}$ is defined by the summation over all $\vec{\alpha}_n^*$ and $\vec{\beta}_n^*$ satisfying

$|\vec{\alpha}_n^*| = |\vec{\beta}_n^*| = n$. In brief

$$n = |\vec{\alpha}_n^*| := \sum_{i=1}^N \alpha_i^n, \quad n = |\vec{\beta}_n^*| := \sum_{i=1}^N \beta_i^n, \quad \sum_{\vec{\alpha}_n, \vec{\beta}_n^*} := \sum_{|\vec{\alpha}_n^*|=|\vec{\beta}_n^*|=n}.$$

Theorem 3. When the star product with separation of variables for smooth functions f and g on a local symmetric Kähler manifold is given as

$$f * g = \sum_{n=0}^{\infty} \sum_{\vec{\alpha}_n, \vec{\beta}_n^*} T_{\vec{\alpha}_n, \vec{\beta}_n^*}^n \left(D^{\vec{\alpha}_n} f \right) \left(D^{\vec{\beta}_n^*} g \right)$$

these smooth functions $T_{\vec{\alpha}_n, \vec{\beta}_n^*}^n$, which are covariantly constants, are determined by the following recurrence relations for all i

$$\begin{aligned} & \sum_{d=1}^N \hbar g_{id} T_{\vec{\alpha}_n - \vec{e}_d, \vec{\beta}_n^* - \vec{e}_i}^{n-1} \\ &= \beta_i T_{\vec{\alpha}_n, \vec{\beta}_n^*}^n + \sum_{k=1}^N \sum_{p=1}^N \frac{\hbar (\beta_k^n - \delta_{kp} - \delta_{ik} + 1) (\beta_k^n - \delta_{kp} - \delta_{ik} + 2)}{2} \\ & \quad \times R_{\vec{p}}^{\vec{k}\vec{k}} \bar{i} T_{\vec{\alpha}_n, \vec{\beta}_n^* - \vec{e}_{\vec{p}} + 2\vec{e}_{\vec{k}} - \vec{e}_i}^n + \sum_{k=1}^{N-1} \sum_{l=1}^{N-k} \sum_{p=1}^N \hbar (\beta_k^n - \delta_{kp} - \delta_{ik} + 1) \\ & \quad \times (\beta_{k+l}^n - \delta_{(k+l), p} - \delta_{i, (k+l)} + 1) R_{\vec{p}}^{\vec{k}+\vec{l}\vec{k}} \bar{i} T_{\vec{\alpha}_n, \vec{\beta}_n^* - \vec{e}_{\vec{p}} + \vec{e}_{\vec{k}} + \vec{e}_{\vec{k}+l} - \vec{e}_i}^n. \end{aligned}$$

Outline of Proof. Let f and g be smooth functions on a Kähler manifold M and L_f be a left star product by f given as (12). Then

$$\begin{aligned}\sigma([L_f, \partial_{\bar{i}}\Phi]) &= \frac{\partial\sigma(L_f)}{\partial\xi^{\bar{i}}} \\ &= \sum_{n=0}^{\infty} \sum_{\vec{\alpha}_n, \vec{\beta}_n^*} \beta_i^n T_{\vec{\alpha}_n, \vec{\beta}_n^*}^n \left(D^{\vec{\alpha}_n} f \right) \left(\xi^{\bar{1}\beta_1^n} \dots \xi^{\bar{i}\beta_i^n - 1} \dots \xi^{\bar{N}\beta_N^n} \right)\end{aligned}$$

or equivalently,

$$[L_f, \partial_{\bar{i}}\Phi]g = \sum_{n=0}^{\infty} \sum_{\vec{\alpha}_n, \vec{\beta}_n^*} \beta_i^n T_{\vec{\alpha}_n, \vec{\beta}_n^*}^n \left(D^{\vec{\alpha}_n} f \right) \left(D^{\vec{\beta}_n^* - \vec{e}_i} g \right). \quad (13)$$

The following formulas are given in [10]. For smooth functions f and g on a locally symmetric Kähler manifold, the following formulas are given.

$$\begin{aligned}\nabla_{\bar{j}_1} \dots \nabla_{\bar{j}_n} f &= g_{l_1 \bar{j}_1} \dots g_{l_n \bar{j}_n} D^{l_1} \dots D^{l_n} f \\ \nabla_{k_1} \dots \nabla_{k_n} g &= g_{\bar{m}_1 k_1} \dots g_{\bar{m}_n k_n} D^{\bar{m}_1} \dots D^{\bar{m}_n} g \\ D^{l_1} \dots D^{l_n} f &= g^{l_1 \bar{j}_1} \dots g^{l_n \bar{j}_n} \nabla_{\bar{j}_1} \dots \nabla_{\bar{j}_n} f \\ D^{\bar{m}_1} \dots D^{\bar{m}_n} g &= g^{\bar{m}_1 k_1} \dots g^{\bar{m}_n k_n} \nabla_{k_1} \dots \nabla_{k_n} g.\end{aligned}$$

If M is a locally symmetric Kähler manifold, these formulas derive

$$\begin{aligned}& [L_f, \hbar\partial_{\bar{i}}]g \\ &= \hbar \sum_{n=0}^{\infty} \sum_{\vec{\alpha}_n, \vec{\beta}_n^*} \sum_{k=1}^N \sum_{\vec{\alpha}_n, \vec{\beta}_n^*} \frac{\beta_k^n (\beta_k^n - 1)}{2} R_{\bar{\rho}}^{\bar{k}\bar{k}} \bar{\rho}^{\bar{i}} T_{\vec{\alpha}_n, \vec{\beta}_n^*}^n \left(D^{\vec{\alpha}_n} f \right) \left(D^{\vec{\beta}_n^* + \vec{e}_{\bar{\rho}} - \vec{e}_{\bar{k}}} g \right) \\ &+ \hbar \sum_{n=0}^{\infty} \sum_{k=1}^{N-1} \sum_{l=1}^{N-k} \sum_{\vec{\alpha}_n, \vec{\beta}_n^*} \beta_k^n \beta_{k+l}^n R_{\bar{\rho}}^{\bar{k}+\bar{l}\bar{k}} \bar{\rho}^{\bar{i}} T_{\vec{\alpha}_n, \vec{\beta}_n^*}^n \left(D^{\vec{\alpha}_n} f \right) \left(D^{\vec{\beta}_n^* + \vec{e}_{\bar{\rho}} - \vec{e}_{\bar{k}}} g \right) \\ &- \hbar \sum_{n=1}^{\infty} \sum_{\vec{\alpha}_{n-1}, \vec{\beta}_{n-1}^*} \sum_{d=1}^N g_{\bar{i}d} T_{\vec{\alpha}_{n-1}, \vec{\beta}_{n-1}^*}^{n-1} \left(D^{\vec{\alpha}_{n-1} + \vec{e}_{\bar{d}}} f \right) \left(D^{\vec{\beta}_{n-1}^*} g \right).\end{aligned}$$

■

Details of this proof are given in [5].

3. *–Products for Riemann Surfaces

*–products for Riemann surfaces are studied in this section for arbitrary Riemann surfaces regarded as locally symmetric Kähler manifold. Applying Theorem 3

for complex 1 dimensional case, $*$ -products for Riemann surfaces are obtained concretely. A formal discussions are given in [11], and star products are studied in [9].

The Scalar curvature R is defined as

$$R = g^{i\bar{j}} R_{i\bar{j}} = R_{i\bar{i}}^{\bar{j}j}.$$

Theorem 4. Let M be a one-dimensional locally symmetric Kähler manifold ($N = 1$) and f and g be smooth functions on M . The star product with separation of variables for f and g can be described as ¹

$$f * g = \sum_{n=0}^{\infty} \left[\left(g^{1\bar{1}} \right)^n \left\{ \prod_{k=1}^n \frac{2\hbar}{2k + \hbar k(k-1)R} \right\} \left\{ \left(g^{1\bar{1}} \frac{\partial}{\partial z} \right)^n f \right\} \left\{ \left(g^{1\bar{1}} \frac{\partial}{\partial \bar{z}} \right)^n g \right\} \right].$$

Example 1. Let (\mathbb{C}, g) be a complex plane as a one-dimensional locally symmetric Kähler manifold. The star product with separation of variables for f and g can be described as

$$f * g = \sum_{n=0}^{\infty} \left[\frac{\hbar^n}{n!} \left\{ \left(\frac{\partial}{\partial z} \right)^n f \right\} \left\{ \left(\frac{\partial}{\partial \bar{z}} \right)^n g \right\} \right].$$

Example 2. Wellknown flat torus embedding $X : S^1 \times S^1 \rightarrow \mathbb{R}^4$

$$X(u, v) = (\cos u, \sin u, \cos v, \sin v), \quad u = \operatorname{Re}(z), \quad v = \operatorname{Im}(z)$$

$$\implies R = \frac{-1}{\sqrt{EG}} \left\{ \frac{\partial}{\partial u} \left(\frac{1}{\sqrt{E}} \frac{\partial \sqrt{G}}{\partial u} \right) + \frac{\partial}{\partial v} \left(\frac{1}{\sqrt{G}} \frac{\partial \sqrt{E}}{\partial v} \right) \right\} = 0$$

where first fundamental forms are

$$E = \frac{\partial X}{\partial u} \cdot \frac{\partial X}{\partial u} = 1, \quad F = \frac{\partial X}{\partial u} \cdot \frac{\partial X}{\partial v} = 0, \quad G = \frac{\partial X}{\partial v} \cdot \frac{\partial X}{\partial v} = 1$$

hence u, v are isothermal coordinates on a torus and the pullback metric is defined as

$$\tilde{g}_{1\bar{1}} = E = G = 1.$$

If $(M, g) = (S^1 \times S^1, \tilde{g})$ then $R = R_{1\bar{1}}^{\bar{1}1} = 0$. Hence the star product with separation of variables for f and g can be described as also

$$f * g = \sum_{n=0}^{\infty} \left[\frac{\hbar^n}{n!} \left\{ \left(\frac{\partial}{\partial z} \right)^n f \right\} \left\{ \left(\frac{\partial}{\partial \bar{z}} \right)^n g \right\} \right].$$

¹Here we correct the typos in page 562 in [5]. $\prod_{k=1}^{n-1}$ in [5] should be $\prod_{k=1}^n$.

4. Projective Space Cases

In this section, we calculate star products of $\mathbb{C}\mathbb{P}^N$. These star products are also equal to the ones given in [1, 4, 10]. A projective space is a special Grassmann manifold and a Grassmann manifold is a special flag manifold. Deformation quantization of flag manifolds and Grassmann manifolds were studied in [2, 3, 7, 8]. At first, a function similar to the determinant is defined on the matrix space.

Definition 4 (permanent). *Let $C = (C_{k,l})_{1 \leq k \leq n, 1 \leq l \leq n}$ be a $n \times n$ matrix. We define $|\cdot|^+$ as a \mathbb{C} -valued function on $M(n, n; \mathbb{C})$ such that*

$$|C|^+ := \sum_{\sigma_n \in S_n} \prod_{k=1}^n C_{k, \sigma_n(k)}.$$

This is called “permanent”.

Definition 5. *A matrix $G^{\vec{\alpha}_n, \vec{\beta}_n^*}$ is defined by using the Hermitian metrics on M . Its elements are metrics on M and are located as follows. $\vec{\alpha}_n$ and $\vec{\beta}_n^*$ are elements of \mathbb{Z}^N*

$$G^{\vec{\alpha}_n, \vec{\beta}_n^*} = \begin{pmatrix} \tilde{G}_{11} & \cdots & \tilde{G}_{1n} \\ \vdots & \ddots & \vdots \\ \tilde{G}_{n1} & \cdots & \tilde{G}_{nn} \end{pmatrix}$$

where

$$\tilde{G}_{pq} =: g_{p\bar{q}} \begin{pmatrix} 1 & \cdots & 1 \\ \vdots & \ddots & \vdots \\ 1 & \cdots & 1 \end{pmatrix} \in M(\alpha_p^n, \beta_q^n; \mathbb{C}).$$

Theorem 5. *Let f and g be smooth functions on a projective space $\mathbb{C}\mathbb{P}^N$. A star product with separation of variables on a projective space $\mathbb{C}\mathbb{P}^N$ is given as*

$$f * g = f \cdot g + \sum_{n=1}^{\infty} \sum_{\vec{\alpha}_n, \vec{\beta}_n^*} |G^{\vec{\alpha}_n, \vec{\beta}_n^*}|^+ \left(\prod_{l=1}^N \frac{1}{\alpha_l^n! \beta_l^n!} \right) \prod_{k=1}^n \frac{\hbar}{(1 + \hbar - \hbar k)} (D^{\vec{\alpha}_n} f) (D^{\vec{\beta}_n^*} g). \tag{14}$$

Here, we correct the typos in (5.4) in [5].

Proof. *We show that*

$$T_{\vec{\alpha}_n, \vec{\beta}_n^*}^n = |G^{\vec{\alpha}_n, \vec{\beta}_n^*}|^+ \left(\prod_{l=1}^N \frac{1}{\alpha_l^n! \beta_l^n!} \right) \prod_{k=1}^n \frac{\hbar}{(1 + \hbar - \hbar k)}$$

satisfies (3)

$$\begin{aligned} & \sum_{d=1}^N \frac{\hbar g_{id}}{(1 + \hbar - \hbar n)} \beta_i^n T_{\vec{\alpha}_n - \vec{e}_d, \vec{\beta}_n^* - \vec{e}_i}^{n-1} \\ &= \sum_{d=1}^N g_{id} \alpha_d^n \left| G^{\vec{\alpha}_n - \vec{e}_d, \vec{\beta}_n^* - \vec{e}_i} \right|^+ \frac{\hbar}{(1 + \hbar - \hbar n)} \left(\prod_{l=1}^N \frac{1}{\alpha_l^n! \beta_l^n!} \right) \prod_{k=1}^n \frac{\hbar}{(1 + \hbar - \hbar k)}. \end{aligned}$$

Using cofactor expansion of permanent, the R.H.S. of the above is rewritten as

$$\left| G^{\vec{\alpha}_n, \vec{\beta}_n^*} \right|^+ \left(\prod_{l=1}^N \frac{1}{\alpha_l^n! \beta_l^n!} \right) \prod_{k=1}^n \frac{\hbar}{(1 + \hbar - \hbar k)}.$$

This shows the given $T_{\vec{\alpha}_n, \vec{\beta}_n^*}^n$ satisfies the recurrence relation (3). ■

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