

Spatial asymptotics for the parabolic Anderson model driven by a Gaussian rough noise

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Abstract

The aim of this paper is to establish the almost sure asymptotic behavior as the space variable becomes large, for the solution to the one spatial dimensional stochastic heat equation driven by a Gaussian noise which is white in time and which has the covariance structure of a fractional Brownian motion with Hurst parameter $H \in (\frac{1}{4}, \frac{1}{2})$ in the space variable.

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1 Introduction

This article is concerned with a linear stochastic heat equation on $\mathbb{R}_+ \times \mathbb{R}$, formally written as

$$\frac{\partial u}{\partial t} = \frac{1}{2} \frac{\partial^2 u}{\partial x^2} + u \dot{W}, \quad t \geq 0, \quad x \in \mathbb{R}, \quad (1.1)$$

where \dot{W} is a Gaussian noise which is white in time and colored in space, and we are interested in regimes where the spatial behavior of \dot{W} is rougher than white noise. More

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specifically, our noise can be seen as the formal space-time derivative of a centered Gaussian process whose covariance is given by

$$\mathbb{E}[W(s, x)W(t, y)] = \frac{1}{2} (|x|^{2H} + |y|^{2H} - |x - y|^{2H}) (s \wedge t), \tag{1.2}$$

where $\frac{1}{4} < H < \frac{1}{2}$. That is, W is a standard Brownian motion in time and a fractional Brownian motion with Hurst parameter H in the space variable. Notice that the spatial covariance of \dot{W} , which is formally given by $\gamma(x - y) = H(2H - 1)|x - y|^{2H-2}$, is not locally integrable when $H < \frac{1}{2}$. It is in fact a nonpositive distribution, and therefore the stochastic integration with respect to W cannot be handled by classical theories (see e.g. [11, 10, 16]). However, one has recently been able (cf. [14]) to give a proper definition of equation (1.1) and to solve it in a space of Hölder continuous processes (see also the recent work [2], covering the linear case (1.1)). We shall take those results for granted.

Let us now highlight the fact that space-time asymptotics for stochastic heat equations like (1.1) have attracted a lot of attention in the recent past. This line of research stems from different motivations, and among them let us quote the following. For a fixed $t > 0$, the large scale behavior of the function $x \mapsto u(t, x)$ is dramatically influenced by the presence of the noise \dot{W} in (1.1) (as opposed to a deterministic equation with no noise). One way to quantify this assertion is to analyze the asymptotic behavior of $x \mapsto u(t, x)$ as $|x| \rightarrow \infty$. Results in this sense include intermittency results, upper and lower bounds for $M_R \equiv \sup_{|x| \leq R} u(t, x)$ contained in [8], and culminate in the sharp results obtained in [6]. Roughly speaking, in case of a white noise in time like in (1.2), those articles establish that $\log(M_R)$ behaves like $[\log(R)]^\psi$, for an exponent ψ which depends on the spatial covariance structure of \dot{W} . In particular, if the spatial covariance of \dot{W} is described by the Riesz kernel $|x|^{-\alpha}$ for $\alpha \in (0, 1)$, one gets $\psi = \frac{2}{4-\alpha}$. This interpolates between a regular situation in space ($\alpha = 0$ and $\psi = 1/2$) and the white noise setting ($\alpha = 1$ and $\psi = 2/3$). In any case those results are in sharp contrast with the deterministic case, for which $x \mapsto u(t, x)$ stays bounded.

With these preliminaries in mind, the current contribution completes the space-time asymptotic picture for the stochastic heat equation, covering very rough situations like the ones described by (1.2). Namely, we shall get the following spatial asymptotics.

Theorem 1.1. *Let W be the Gaussian field given by the covariance (1.2), and assume $\frac{1}{4} < H < \frac{1}{2}$. Let u be the unique solution to equation (1.1) driven by \dot{W} with initial condition $u_0 = 1$ (see Proposition 2.3 for a precise statement), and consider $t > 0$. Then*

$$\lim_{R \rightarrow \infty} (\log R)^{-\frac{1}{1+H}} \log \left(\max_{|x| \leq R} u(t, x) \right) = c_0(H) (t\mathcal{E})^{\frac{H}{1+H}} \quad \text{a.s.}, \tag{1.3}$$

where \mathcal{E} is the variational constant to be defined below in Proposition 3.1 and $c_0(H)$ is a constant depending only on H , given by

$$c_0(H) = (1 + H) \left(\frac{c_H}{2} \right)^{\frac{1}{1+H}} \left(\frac{1}{H} \right)^{\frac{H}{1+H}}, \tag{1.4}$$

with c_H defined in (2.2).

Notice that the value of the limit in (1.3) coincides with the corresponding limit in the case $H > \frac{1}{2}$, obtained in [6]. In fact, with the change of variables $H = 1 - \frac{\alpha}{2}$ and assuming that $c_H = 1$ (that is, we take $|\xi|^{1-2H} d\xi$ as the spectral measure), this limit equals to

$$\frac{4 - \alpha}{4} \left(\frac{4\mathcal{E}}{2 - \alpha} \right)^{\frac{2-\alpha}{4-\alpha}} t^{\frac{2-\alpha}{4-\alpha}},$$

which is equal to the right-hand side of (1.20) in [6], when $\alpha_0 = 1$, $d = 1$ and $\theta = 1$.

Let us say a few words about our strategy in order to prove Theorem 1.1. It can be roughly divided in two main steps:

(i) *Tail estimate for $u(t, x)$.* Let us fix $t \in \mathbb{R}_+$ and $x \in \mathbb{R}$. We will see (cf Corollary 4.5) that for large a , we have

$$\mathbb{P}(\log(u(t, x)) \geq a) \simeq \exp\left(-\frac{\hat{c}_{H,t} a^{1+H}}{t^H}\right), \tag{1.5}$$

where $\hat{c}_{H,t}$ is determined by a variational problem. This stems, via some large deviation arguments, from a sharp analysis of the high moments of $u(t, x)$. Namely, our main effort in order to get the tail behavior will be to prove that for large $m \in \mathbb{N}$, we have (see Theorem 4.4)

$$\mathbb{E}[(u(t, x))^m] \simeq \exp\left(c_H t m^{1+\frac{1}{H}}\right), \tag{1.6}$$

with a variational expression for c_H . Towards this aim, we resort to a Feynman-Kac representation for the moments of $u(t, x)$, which involves a kind of intersection local time for an m -dimensional Brownian motion weighted by a singular potential. We are thus able to relate the quantity $\mathbb{E}[(u(t, x))^m]$ to a semi-group on $L^2(\mathbb{R}^m)$, and this semi-group admits a generator A_m which can be expressed as the Laplace operator on \mathbb{R}^m perturbed by a singular distributional potential. Then we shall get our asymptotic result (1.6) thanks to a careful spectral analysis of A_m . The technicalities related to this step are detailed in Sections 3 and 4.

(ii) *Spatial behavior.* Once the tail of $\log(u(t, x))$ has been sharply estimated, we can complete the study of the asymptotic behavior in the following way: on the interval $[-M, M]$ for large M , we are able to produce some random variables X_1, \dots, X_N such that:

- N is of order $2M$.
- X_1, \dots, X_N are i.i.d, and satisfy approximately (1.5). More precisely, for any $\delta > 0$, we shall prove that one can choose λ conveniently so that the inequality $\mathbb{P}(\log(|X_i|) > \lambda(\log R)^{\frac{1}{1+H}}) > R^{-(1-\frac{\delta}{2})}$ holds true for large R .
- X_1, \dots, X_N are approximations of $u(t, x_1), \dots, u(t, x_N)$ with $x_1, \dots, x_N \in [-M, M]$.
- Fluctuations of u in small boxes around x_1, \dots, x_N are small.

With this information in hand, the behavior $(\log R)^{\frac{1}{1+H}}$ in Theorem 1.1 can be heuristically understood as follows: for an additional parameter λ , we have

$$\mathbb{P}\left(\max_{j \leq N} \log(|X_j|) \leq \lambda[\log R]^{\frac{1}{1+H}}\right) = \left[1 - \mathbb{P}\left(\log(|X_j|) > \lambda[\log(R)]^{\frac{1}{1+H}}\right)\right]^N,$$

and thanks to the tail estimate (1.5), we obtain

$$\mathbb{P}\left(\max_{j \leq N} \log(|X_j|) \leq \lambda[\log R]^{\frac{1}{1+H}}\right) \simeq \left[1 - \exp\left(-\hat{c}_{H,t} \lambda^{1+H} \log R\right)\right]^N.$$

With some elementary calculus considerations, and playing with the extra parameter λ , one can now easily check that for R large enough

$$\mathbb{P}\left(\max_{j \leq N} \log(|X_j|) \leq \lambda[\log R]^{\frac{1}{1+H}}\right) \leq \exp(-R^\nu),$$

with a positive ν . Otherwise stated, we obtain an exponentially small probability of having $\log(|X_j|)$ of order less than $[\log R]^{\frac{1}{1+H}}$. Using a Borel-Cantelli type argument and the fact that fluctuations of u in small boxes around x_1, \dots, x_N are small, we thus prove Theorem 1.1.

As already mentioned, the spatial covariance γ of the noise \dot{W} driving equation (1.1) is a nonpositive distribution. With respect to smoother cases such as the ones treated in [6], this induces some serious additional difficulties which can be summarized as follows. First, the variational asymptotic results involving the generator A_m cannot be reduced to a one-dimensional situation due to the singularities of γ . We thus have to handle a family of optimization problems in $L^2(\mathbb{R}^m)$ for arbitrarily large m . Then, still in the part concerning the asymptotic behavior of $m \mapsto \mathbb{E}[(u(t, x))^m]$, the upper bound obtained in [6] relied heavily on a compactification by folding argument for which the positivity of γ was essential. This approach is no longer applicable here, and we have to replace it by a coarse graining procedure. Finally, the localization procedure and the study of fluctuations in the spatial behavior step of our proof, though similar in spirit to the one in Conus et al. [9], is more involved in its implementation. More specifically, in our case the moment estimates cannot be obtained by using sharp Burkholder inequalities, because of the roughness of the noise. For this reason we use Wiener chaos expansions and hypercontractivity, which are more suitable methods in our context. The fluctuation estimates alluded to above are also obtained through chaos expansions.

The paper is organized as follows. Section 2 contains some preliminaries on stochastic integration with respect to the rough noise \dot{W} and the mild formulation of equation (1.1). We introduce the variational quantities and their asymptotic behavior when time is large in Section 3. Section 3.3 deals with Feynman-Kac semigroups and in Section 4 we derive the precise moments asymptotics which are required to show Theorem 1.1. The proof of Theorem 1.1 is given in Section 5. A technical lemma is proved in the appendix.

2 Multiplicative stochastic heat equation

This section is devoted to recall the basic existence and uniqueness results for the stochastic equation with rough space-time noise.

2.1 Structure of the noise

Recall that we are considering a Gaussian field W whose covariance structure is given by (1.2). As mentioned above, the stochastic integration with respect to W has only been introduced recently in [2, 14], and we proceed now to a brief review of the results therein.

Let us start by introducing our basic notation on Fourier transforms of functions. The space of Schwartz functions is denoted by \mathcal{S} . Its dual, the space of tempered distributions, is \mathcal{S}' . The Fourier transform of a function $g \in \mathcal{S}$ is defined by

$$\mathcal{F}g(\xi) = \int_{\mathbb{R}} e^{-i\xi x} g(x) dx,$$

so that the inverse Fourier transform is given by $\mathcal{F}^{-1}g(\xi) = (2\pi)^{-1}\mathcal{F}g(-\xi)$.

Let $\mathcal{D}((0, \infty) \times \mathbb{R})$ denote the space of real-valued infinitely differentiable functions with compact support on $(0, \infty) \times \mathbb{R}$. Taking into account the spectral representation of the covariance function of the fractional Brownian motion in the case $H < \frac{1}{2}$ proved in [15, Theorem 3.1], we represent our noise W by a zero-mean Gaussian family $\{W(\varphi), \varphi \in \mathcal{D}((0, \infty) \times \mathbb{R})\}$ defined on a complete probability space $(\Omega, \mathcal{F}, \mathbb{P})$, whose covariance structure is given by

$$\mathbb{E}[W(\varphi)W(\psi)] = c_H \int_{\mathbb{R}_+ \times \mathbb{R}} \mathcal{F}\varphi(s, \xi) \overline{\mathcal{F}\psi(s, \xi)} \mu(d\xi) ds, \tag{2.1}$$

where the Fourier transforms $\mathcal{F}\varphi, \mathcal{F}\psi$ are understood as Fourier transforms in space

only and where we have set

$$c_H = \frac{1}{2\pi} \Gamma(2H + 1) \sin(\pi H), \quad \text{and} \quad \mu(d\xi) = |\xi|^{1-2H} d\xi. \tag{2.2}$$

We denote by γ the Fourier transform of the measure $\mu(d\xi)$. Formally $c_H \gamma(x) \delta_0(s)$ is the covariance function of the generalized noise \dot{W} . However, notice that γ is a generalized function and the integral $\gamma(x) = \int_{\mathbb{R}} e^{-i\xi x} \mu(d\xi)$ is not defined pointwise. Rather, it is defined as a linear functional given by

$$\int_{\mathbb{R}} \gamma(x) \varphi(x) dx = \int_{\mathbb{R}} \mathcal{F}\varphi(\xi) \mu(d\xi),$$

for any $\varphi \in \mathcal{S}(\mathbb{R})$. As a generalized function, γ is non-negative definite in the sense that

$$\int_{\mathbb{R} \times \mathbb{R}} \gamma(x - y) \varphi(x) \varphi(y) dx dy = \int_{\mathbb{R}} |\mathcal{F}\varphi(\xi)|^2 \mu(d\xi) \geq 0, \quad \varphi \in \mathcal{S}(\mathbb{R}).$$

On the other hand, $\gamma(x)$ (or more precisely, its truncated form) takes negative values somewhere. As mentioned in the introduction, this fact makes the problem of spatial asymptotics for equation (1.1) much harder than in [6].

Let \mathcal{H} be the closure of $\mathcal{D}((0, \infty) \times \mathbb{R})$ under the semi-norm induced by the right-hand side of (2.1). The Gaussian family W can be extended as an isonormal Gaussian process $W = \{W(\varphi), \varphi \in \mathcal{H}\}$ indexed by the Hilbert space \mathcal{H} . The space \mathcal{H} can be identified with the homogenous Sobolev space of order $\frac{1}{2} - H$ of functions with values in $L^2(\mathbb{R}_+)$, namely $\mathcal{H} = \dot{H}^{\frac{1}{2}-H}(L^2(\mathbb{R}_+))$ (see [1] for the definition of $\dot{H}^{\frac{1}{2}-H}$).

Let us close this subsection by the definition of Itô's type integral in our context, which will play a crucial role in the sequel. We will make use of the notation for any $t \geq 0$ and $\varphi \in \mathcal{S}(\mathbb{R})$

$$W(t, \varphi) = W(\mathbf{1}_{[0,t]} \varphi). \tag{2.3}$$

For each $t \geq 0$, we denote by \mathcal{F}_t the σ -field generated by the random variables $\{W(s, \varphi) : s \in [0, t], \varphi \in \mathcal{S}(\mathbb{R})\}$. The following proposition is borrowed from [14].

Proposition 2.1. *Let L_a^2 be the space of predictable processes g defined on $\mathbb{R}_+ \times \mathbb{R}$ such that almost surely $g \in \mathcal{H}$ and $\mathbb{E}[\|g\|_{\mathcal{H}}^2] < \infty$. Then, the stochastic integral $\int_{\mathbb{R}_+} \int_{\mathbb{R}} g(s, x) W(ds, dx)$ is well defined for $g \in L_a^2$. Furthermore, the following isometry property holds true*

$$\begin{aligned} \mathbb{E} \left[\left(\int_{\mathbb{R}_+} \int_{\mathbb{R}} g(s, x) W(ds, dx) \right)^2 \right] &= \mathbb{E} [\|g\|_{\mathcal{H}}^2] \\ &= c_H \int_{\mathbb{R}_+ \times \mathbb{R}} \mathbb{E} [|\mathcal{F}g(s, \xi)|^2] |\xi|^{1-2H} d\xi ds. \end{aligned} \tag{2.4}$$

2.2 Stochastic heat equation with rough multiplicative noise

Recall that we are considering equation (1.1) driven by the noise described in Section 2.1. For the sake of simplicity, we shall moreover choose $u(0, \cdot) = 1$ as the initial condition.

Definition 2.2. *Let $u = \{u(t, x), t \geq 0, x \in \mathbb{R}\}$ be a real-valued predictable stochastic process. Assume that for all $t \geq 0$ and $x \in \mathbb{R}$ the process $\{p_{t-s}(x - y)u(s, y)\mathbf{1}_{[0,t]}(s), 0 \leq s \leq t, y \in \mathbb{R}\}$ is an element of L_a^2 , where $p_t(x)$ is the heat kernel on the real line related to $\frac{1}{2}\Delta$ and L_a^2 is defined in Proposition 2.1. We say that u is a mild solution of (1.1) if for all $t \geq 0$ and $x \in \mathbb{R}$ we have*

$$u(t, x) = 1 + \int_0^t \int_{\mathbb{R}} p_{t-s}(x - y) u(s, y) W(ds, dy) \quad a.s., \tag{2.5}$$

where the stochastic integral is understood in the Itô sense of Proposition 2.1.

Let $\{B_j, j \geq 1\}$ be a collection of independent standard Brownian motions, all independent of W . For all $t \geq 0$ and $j < k$, we can define (see [14] again) the functional

$$V_{j,k}(t) := \int_0^t \gamma(B_j(s) - B_k(s)) ds, \tag{2.6}$$

which is interpreted as the following limit in $L^2(\Omega)$

$$V_{j,k}(t) = \lim_{\varepsilon \rightarrow 0} \int_0^t \int_{\mathbb{R}} e^{-\varepsilon|\xi|^2} e^{i\xi(B_j(s) - B_k(s))} \mu(d\xi) ds.$$

With these notations in mind, let us quote an existence and uniqueness result for our equation of interest.

Proposition 2.3. *There is a unique nonnegative mild solution u to equation (2.5), understood as in Definition 2.2. Moreover, recalling our notation (2.6) above, we have:*

(i) *The following Feynman-Kac formula for the moments of u holds true for $m \geq 2$*

$$\mathbb{E} [(u(t, x))^m] = \mathbb{E}_x [\exp(c_H Q_m(t))], \quad \text{with } Q_m(t) = \sum_{1 \leq j < k \leq m} V_{j,k}(t), \tag{2.7}$$

where $\{B_j, j = 1, \dots, m\}$ is a family of independent standard Brownian motions starting from $x \in \mathbb{R}$, and \mathbb{E}_x denotes the expected value with respect to the Wiener measure shifted by x .

(ii) *For any $m \geq 1$ and any $\alpha > 0$ there exist some constants c_1 and c_2 such that*

$$\mathbb{E} [\exp(\alpha Q_m(t))] \leq c_1 \exp\left(c_2 m^{1+\frac{1}{H}} t\right). \tag{2.8}$$

In particular, for $t \geq 0$ and $x \in \mathbb{R}$ we have the existence of two constants c_3 and c_4 satisfying

$$\mathbb{E} [(u(t, x))^m] \leq c_3 \exp\left(c_4 m^{1+\frac{1}{H}} t\right).$$

One of the main steps in our estimates will be to obtain a sharp refinement of inequality (2.8).

3 Preliminaries on Dirichlet forms and semigroups

As mentioned in the Introduction, the semigroup related to a certain operator A_m will play a prominent role in our analysis of the spatial behavior of u . The current section defines and analyzes these objects.

3.1 Variational quantities

We will see in Section 4 that our sharp asymptotic estimates rely on an optimization problem for some variational quantities related to equation (1.1). We now derive some analytic properties of those quantities.

3.1.1 A variational form on \mathbb{R}

Let us consider the following general problem: let \mathcal{K}_1 be the space defined by

$$\mathcal{K}_1 = \{g \in L^2(\mathbb{R}) : \|g\|_2 = 1 \text{ and } g' \in L^2(\mathbb{R})\}. \tag{3.1}$$

Next, for a given parameter $\theta > 0$ and $g \in \mathcal{K}_1$ set

$$H_\theta(g) = \theta \int_{\mathbb{R}} |\mathcal{F}g^2(\xi)|^2 |\xi|^{1-2H} d\xi - \frac{1}{2} \int_{\mathbb{R}} |g'(x)|^2 dx. \tag{3.2}$$

We are interested in optimizing this kind of variational quantity, and here is a first result in this direction.

Proposition 3.1. For $\theta > 0$ and $g \in \mathcal{K}_1$ consider the variational quantity $H_\theta(g)$ defined by (3.2) and set

$$\mathcal{E}_\theta := \sup \{H_\theta(g) : g \in \mathcal{K}_1\}. \tag{3.3}$$

Then the following holds true:

- (i) The quantity \mathcal{E}_θ is finite for any $\theta > 0$.
- (ii) Setting $\mathcal{E} = \mathcal{E}_1$, we have $\mathcal{E}_\theta = \theta^{\frac{1}{2H}} \mathcal{E}$.

Proof. Let us first focus on item (i). For any $g \in \mathcal{K}_1$ and $\xi \in \mathbb{R}$ we have

$$|\mathcal{F}g^2(\xi)| = \left| \int_{\mathbb{R}} e^{-i\xi x} g^2(x) dx \right| \leq \int_{\mathbb{R}} |g^2(x)| dx = 1.$$

In addition, an elementary integration by parts argument shows that

$$\int_{\mathbb{R}} e^{-i\xi x} g^2(x) dx = -i \int_{\mathbb{R}} \left(\frac{1}{\xi} \frac{dg^2}{dx} \right) e^{-i\xi x} dx.$$

Hence, for any $\xi \in \mathbb{R}$ and $g \in \mathcal{K}_1$ we get

$$|\mathcal{F}g^2(\xi)| \leq |\xi|^{-1} \int_{\mathbb{R}} \left| \frac{dg^2}{dx}(x) \right| dx = 2|\xi|^{-1} \int_{\mathbb{R}} |g(x)||g'(x)| dx \leq 2|\xi|^{-1} \|g'\|_2,$$

where the last relation is due to Cauchy-Schwarz' inequality plus the fact that $\|g\|_{L^2} = 1$ for $g \in \mathcal{K}_1$. Consider now an additional parameter $R > 0$. Gathering the two bounds we have obtained for $|\mathcal{F}g^2(\xi)|$, we end up with

$$\int_{\mathbb{R}} |\mathcal{F}g^2(\xi)|^2 |\xi|^{1-2H} d\xi \leq \int_{-R}^R |\xi|^{1-2H} d\xi + 4\|g'\|_2^2 \int_{\{|\xi| \geq R\}} |\xi|^{-(1+2H)} d\xi = \frac{R^{2-2H}}{1-H} + \frac{4\|g'\|_2^2}{H R^{2H}}.$$

We thus take $R = R_\theta$ large enough, such that $\frac{4\theta}{H R_\theta^{2H}} \leq \frac{1}{2}$. We get

$$\theta \int_{\mathbb{R}} |\mathcal{F}g^2(\xi)| |\xi|^{1-2H} d\xi - \frac{1}{2} \int_{\mathbb{R}} |g'(x)|^2 dx \leq \frac{\theta R_\theta^{2-2H}}{1-H}.$$

Since the quantity R_θ does not depend on g , the above inequality is valid for all $g \in \mathcal{K}_1$. The proof of item (i) is thus finished.

In order to check item (ii), consider a parameter $a > 0$, and for $g \in \mathcal{K}_1$ set $g_a(x) = a^{1/2}g(ax)$. It is readily checked that $g_a \in \mathcal{K}_1$ whenever $g \in \mathcal{K}_1$. In addition, we have

$$\mathcal{F}g_a^2(\xi) = \mathcal{F}g^2\left(\frac{\xi}{a}\right), \quad \text{and} \quad g'_a(x) = a^{3/2}g'(ax).$$

Plugging this information into the definition (3.2), we obtain

$$H_\theta(g_a) = \theta a^{2-2H} \int_{\mathbb{R}} |\mathcal{F}g^2(\xi)|^2 |\xi|^{1-2H} d\xi - \frac{a^2}{2} \int_{\mathbb{R}} |g'(x)|^2 dx.$$

We now choose $a = \theta^{\frac{1}{2H}}$, which yields $H_\theta(g_a) = \theta^{\frac{1}{H}} H_1(g)$. Our point (ii) is now easily derived. □

Remark 3.2. Notice that Proposition 3.1 holds for any $H \in (0, \frac{1}{2})$.

Remark 3.3. The variational quantity H_θ is defined in (3.2) appealing to Fourier coordinates. We could also have tried to introduce it in direct coordinates as

$$\hat{H}_\theta(g) = \int_{\mathbb{R} \times \mathbb{R}} \frac{g^2(x)g^2(y)}{|x-y|^{2-2H}} dx dy - \frac{1}{2} \int_{\mathbb{R}} |g'(x)|^2 dx. \tag{3.4}$$

However, this quantity blows up for a broad class of functions in \mathcal{K}_1 . Indeed, for $g \in \mathcal{K}_1$ such that $g \geq \alpha > 0$ on a neighborhood of 0, we have

$$\int_{\mathbb{R} \times \mathbb{R}} \frac{g^2(x)g^2(y)}{|x-y|^{2-2H}} dx dy \geq \alpha^4 \int_{\{|x| \leq \varepsilon\} \times \{|y| \leq \varepsilon\}} \frac{1}{|x-y|^{2-2H}} dx dy = \infty.$$

3.1.2 A variational form on \mathbb{R}^m

Fix $m \geq 2$. Our future computations also rely on the following variational quantity on \mathbb{R}^m :

$$K_{\theta,m}(g) = \frac{\theta}{m} \sum_{1 \leq j < k \leq m} \int_{\mathbb{R}^m} \gamma(x_j - x_k) g^2(x) dx - \frac{1}{2} \int_{\mathbb{R}^m} |\nabla g(x)|^2 dx. \tag{3.5}$$

Notice that $K_{\theta,m}$ can also be interpreted as a Dirichlet form related to a Schrödinger type generator $A_{\theta,m}$, that is,

$$K_{\theta,m}(g) = \langle A_{\theta,m}g, g \rangle_{L^2(\mathbb{R}^m)}, \quad \text{with} \quad A_{\theta,m} = \frac{1}{2}\Delta + \frac{\theta}{m} \sum_{1 \leq j < k \leq m} \gamma(x_j - x_k). \tag{3.6}$$

Observe, however, that $K_{\theta,m}$ and $A_{\theta,m}$ are only defined for smooth test functions, due to the fact that γ is a distribution.

Remark 3.4. The quantity $K_{\theta,m}(g)$ can also be expressed in Fourier modes. Indeed, the inverse Fourier transform of $x \in \mathbb{R}^m \mapsto \gamma(x_j - x_k)$ is given by

$$\varphi_{jk}(\xi) = |\xi_j|^{1-2H} \delta_0(\xi_j + \xi_k) \prod_{l \neq j,k} \delta_0(\xi_l).$$

Hence for $j < k$ we end up with

$$\int_{\mathbb{R}^m} \gamma(x_j - x_k) g^2(x) dx = \int_{\mathbb{R}^m} |\lambda|^{1-2H} \hat{g}_{jk}(\lambda) d\lambda,$$

where we define

$$\hat{g}_{jk}(\lambda) = \mathcal{F}g^2(0, \dots, 0, \overbrace{\lambda}^j, 0, \dots, 0, \overbrace{-\lambda}^k, 0, \dots, 0). \tag{3.7}$$

Summarizing, we have obtained

$$K_{\theta,m}(g) = -\frac{1}{2(2\pi)^m} \int_{\mathbb{R}^m} |\xi|^2 |\mathcal{F}g(\xi)|^2 d\xi + \frac{\theta}{m} \sum_{1 \leq j < k \leq m} \int_{\mathbb{R}} |\lambda|^{1-2H} \hat{g}_{jk}(\lambda) d\lambda. \tag{3.8}$$

3.2 Asymptotic results for principal eigenvalues

With those preliminary notions in hand, we now relate the principal eigenvalue of $A_{\theta,m}$ with the quantity \mathcal{E}_θ following the methodology introduced in [7].

Proposition 3.5. Consider $\theta > 0$ and the quantity $K_{\theta,m}(g)$ given by (3.5). Define the set \mathcal{K}_m (which is the equivalent of \mathcal{K}_1 for functions defined on \mathbb{R}^m) as follows

$$\mathcal{K}_m = \{g \in L^2(\mathbb{R}^m) : \|g\|_2 = 1 \text{ and } \nabla g \in L^2(\mathbb{R}^m)\}. \tag{3.9}$$

We define the principal eigenvalue of the operator $A_{\theta,m}$ by

$$\lambda_{\theta,m} = \sup \{K_{\theta,m}(g); g \in \mathcal{K}_m\}. \tag{3.10}$$

Then the following asymptotic behavior holds true

$$\lim_{m \rightarrow \infty} \frac{\lambda_{\theta,m}}{m} = \mathcal{E}_{\theta/2} = \left(\frac{\theta}{2}\right)^{\frac{1}{H}} \mathcal{E},$$

where we recall that \mathcal{E}_θ is defined by (3.3).

Proof. We mostly focus on the upper bound, the lower bound being easier to obtain. To this aim we divide the proof in several steps.

Step 1: Cutoff procedure. Observe that the results in [7] only hold for a pointwise defined function $\gamma(x)$. For this reason we introduce the decomposition

$$\gamma(x) = \int_{-M}^M e^{-i\xi x} \mu(d\xi) + \int_{\mathbb{R} \setminus [-M, M]} e^{-i\xi x} \mu(d\xi) = \gamma_M^1(x) + \gamma_M^2(x), \tag{3.11}$$

where the above identity (namely the second term $\gamma_M^2(x)$) is understood in the distribution sense. Also notice that for $j = 1, 2$, the function γ_M^j can be seen as the Fourier transform of the measure μ_M^j , where μ_M^1 and μ_M^2 are defined as follows:

$$\mu_M^1(d\xi) = |\xi|^{1-2H} \mathbf{1}_{[-M, M]}(\xi) d\xi, \quad \text{and} \quad \mu_M^2(d\xi) = |\xi|^{1-2H} \mathbf{1}_{[-M, M]^c}(\xi) d\xi. \tag{3.12}$$

Then, for any $\delta \in (0, 1)$, we can write

$$\sup_{g \in \mathcal{K}_m} K_{\theta, m}(g) \leq B_{m, M}^1 + B_{m, M}^2, \tag{3.13}$$

where

$$B_{m, M}^1 = \sup_{g \in \mathcal{K}_m} \left\{ \frac{\theta}{m} \sum_{1 \leq j < k \leq m} \int_{\mathbb{R}^m} \gamma_M^1(x_j - x_k) g^2(x) dx - \frac{1-\delta}{2} \int_{\mathbb{R}^m} |\nabla g(x)|^2 dx \right\}$$

and

$$B_{m, M}^2 = \sup_{g \in \mathcal{K}_m} \left\{ \frac{\theta}{m} \sum_{1 \leq j < k \leq m} \int_{\mathbb{R}^m} \gamma_M^2(x_j - x_k) g^2(x) dx - \frac{\delta}{2} \int_{\mathbb{R}^m} |\nabla g(x)|^2 dx \right\}.$$

The term $B_{m, M}^1$ is handled by [7] and we obtain

$$\lim_{m \rightarrow \infty} \frac{1}{m} B_{m, M}^1 = \mathfrak{E}_{M, \delta, \theta}, \quad \text{with} \quad \lim_{M \rightarrow \infty} \mathfrak{E}_{M, \delta, \theta} = \left(\frac{\theta}{2}\right)^{\frac{1}{H}} (1-\delta)^{1-\frac{1}{H}} \mathcal{E}, \tag{3.14}$$

where the limiting behavior for $\mathfrak{E}_{M, \delta, \theta}$ is a direct consequence of Proposition 3.1.

Step 2: Upper bound for $B_{m, M}^2$. We claim that the following inequality holds

$$B_{m, M}^2 \leq \frac{m-1}{2} \sup_{g \in \mathcal{K}_1} \left\{ \theta \int_{\mathbb{R}} \gamma_M^2(x) g^2(x) dx - \frac{\delta}{2} \int_{\mathbb{R}} |g'(x)|^2 dx \right\}. \tag{3.15}$$

To show this inequality, we write

$$\begin{aligned} & \sum_{1 \leq j < k \leq m} \int_{\mathbb{R}^m} \gamma_M^2(x_j - x_k) g^2(x_1, \dots, x_m) dx_1 \cdots dx_m \\ &= \frac{1}{2} \sum_{j=1}^m \sum_{k: k \neq j} \int_{\mathbb{R}^m} \gamma_M^2(x_j - x_k) g^2(x_1, \dots, x_m) dx_1 \cdots dx_m, \end{aligned}$$

and

$$\int_{\mathbb{R}^m} |\nabla g(x)|^2 dx = \sum_{j=1}^m \|\nabla_j g\|_2^2 = \frac{1}{m-1} \sum_{j=1}^m \sum_{k: k \neq j} \|\nabla_k g\|_2^2.$$

Replacing the supremum of the sum by a sum of supremums in the expression defining $B_{m, M}^2$, we obtain

$$\begin{aligned} B_{m, M}^2 &\leq \frac{1}{2m} \sum_{j=1}^m \sup_{g \in \mathcal{K}_m} \left\{ \theta \sum_{k: k \neq j} \int_{\mathbb{R}^m} \gamma_M^2(x_j - x_k) g^2(x) dx - \frac{m\delta}{m-1} \sum_{k: k \neq j} \|\nabla_k g\|_2^2 \right\} \\ &= \frac{1}{2} \sup_{g \in \mathcal{K}_m} D_{m, M}(g), \end{aligned} \tag{3.16}$$

where we have set

$$D_{m,M}(g) = \theta \sum_{k=2}^m \int_{\mathbb{R}^m} \gamma_M^2(x_k - x_1) g^2(x) dx - \frac{m\delta}{m-1} \sum_{k=2}^m \|\nabla_k g\|_2^2. \tag{3.17}$$

We now proceed to estimate the terms $\int_{\mathbb{R}^m} \gamma_M^2(x_k - x_1) g^2(x) dx$ and $\sum_{k=2}^m \|\nabla_k g\|_2^2$ above. In order to estimate the terms $\int_{\mathbb{R}^m} \gamma_M^2(x_k - x_1) g^2(x) dx$, notice that

$$\begin{aligned} & \int_{\mathbb{R}^{m-1}} \gamma_M^2(x_k - x_1) g^2(x_1, \dots, x_m) dx_2 \cdots dx_m \\ &= \int_{\mathbb{R}^{m-1}} \gamma_M^2(x_k) g^2(x_1, x_2 + x_1, \dots, x_m + x_1) dx_2 \cdots dx_m \\ &= \int_{\mathbb{R}^{m-1}} \gamma_M^2(x_k) \tilde{g}^2(x_1, x_2, \dots, x_m) dx_2 \cdots dx_m, \end{aligned}$$

where the function $\tilde{g} : \mathbb{R}^m \rightarrow \mathbb{R}$ is defined by

$$\tilde{g}(x_1, x_2, \dots, x_m) = g(x_1, x_2 + x_1, \dots, x_m + x_1).$$

Now observe that \tilde{g} belongs to the space \mathcal{K}_m defined by (3.5). We thus obtain

$$\sum_{k=2}^m \int_{\mathbb{R}^m} \gamma_M^2(x_k - x_1) g^2(x_1, \dots, x_m) dx = \sum_{k=2}^m \int_{\mathbb{R}^m} \gamma_M^2(x_k) \tilde{g}^2(x_1, \dots, x_m) dx.$$

As far as the terms $\|\nabla_k g\|_2^2$ in the definition of $D_{m,M}(g)$ are concerned, we just notice that $\|\nabla_k \tilde{g}\|_2^2 = \|\nabla_k g\|_2^2$ for every $2 \leq k \leq m$. Hence, recalling the definition (3.17) of $D_{m,M}(g)$, for any $g \in \mathcal{K}_m$ we end up with

$$\begin{aligned} D_{m,M}(g) &= \theta \sum_{k=2}^m \int_{\mathbb{R}^m} \gamma_M^2(x_k) \tilde{g}^2(x) dx - \frac{m\delta}{m-1} \sum_{k=2}^m \|\nabla_k \tilde{g}\|_2^2 \\ &\leq \sup_{g \in \mathcal{K}_m} \left\{ \theta \sum_{k=2}^m \int_{\mathbb{R}^m} \gamma_M^2(x_k) g^2(x) dx - \frac{m\delta}{m-1} \sum_{k=2}^m \|\nabla_k g\|_2^2 \right\}, \end{aligned}$$

and thus

$$D_{m,M}(g) \leq \sup_{g \in \mathcal{K}_m} \left\{ \theta \sum_{k=2}^m \int_{\mathbb{R}} \gamma_M^2(x_k) g_k^2(x_k) dx_k - \frac{m\delta}{m-1} \sum_{k=2}^m \|\nabla_k g\|_2^2 \right\}, \tag{3.18}$$

where for $k = 2, \dots, m$ we have set

$$g_k(x_k) = \left(\int_{\mathbb{R}^{m-1}} g^2(x_1, \dots, x_m) \prod_{1 \leq j \leq m, j \neq k} dx_j \right)^{1/2}.$$

Let us further analyze the term $\|\nabla_k g\|_2^2 = \|\nabla_k g\|_{L^2(\mathbb{R}^m)}^2$ above, and relate it with $\|g'_k\|_2^2 = \|g'_k\|_{L^2(\mathbb{R})}^2$. To this aim, it is readily checked that $g_k \in \mathcal{K}_1$ whenever $g \in \mathcal{K}_m$. Furthermore, we have

$$g'_k(x_k) = \left(\int_{\mathbb{R}^{m-1}} g^2(x) \prod_{j \neq k} dx_j \right)^{-\frac{1}{2}} \int_{\mathbb{R}^{m-1}} g(x) \nabla_k g(x) \prod_{j \neq k} dx_j.$$

By Cauchy-Schwarz' inequality, we thus get

$$|g'_k(x_k)|^2 \leq \int_{\mathbb{R}^{m-1}} |\nabla_k g(x)|^2 \prod_{j \neq k} dx_j, \tag{3.19}$$

which yields

$$\|g'_k\|_2^2 = \int_{\mathbb{R}} |g'_k(x_k)|^2 dx_k \leq \|\nabla_k g\|_2^2, \quad k = 2, \dots, m.$$

Plugging this inequality into (3.18), we have obtained

$$D_{m,M}(g) \leq \sup_{g \in \mathcal{K}_m} \left\{ \theta \sum_{k=2}^m \int_{\mathbb{R}} \gamma_M^2(x_k) g_k^2(x_k) dx_k - \frac{m\delta}{m-1} \sum_{k=2}^m \|g'_k\|_2^2 \right\},$$

and recalling relation (3.16), this yields

$$\begin{aligned} B_{m,M}^2 &\leq \frac{1}{2} \sup_{g \in \mathcal{K}_m} \left\{ \theta \sum_{k=2}^m \int_{\mathbb{R}} \gamma_M^2(x_k) g_k^2(x_k) dx_k - \frac{m\delta}{m-1} \sum_{k=2}^m \|g'_k\|_2^2 \right\} \\ &\leq \frac{m-1}{2} \sup_{g \in \mathcal{K}_1} \left\{ \theta \int_{\mathbb{R}} \gamma_M^2(x) g^2(x) dx - \delta \int_{\mathbb{R}} |g'(x)|^2 dx \right\}, \end{aligned}$$

which shows (3.15).

Step 3: Limiting behavior for $B_{m,M}^2$. The estimate (3.15) implies

$$\limsup_{m \rightarrow \infty} \frac{B_{m,M}^2}{m} \leq \frac{1}{2} \sup_{g \in \mathcal{K}_1} \left\{ \theta \int_{\mathbb{R}} \gamma_M^2(x) g^2(x) dx - \delta \int_{\mathbb{R}} |g'(x)|^2 dx \right\}. \quad (3.20)$$

We now show that this variation goes to zero as M tends to ∞ . This can be seen by a simple integration by parts argument. For any $g \in \mathcal{K}_1$ we have

$$\mathcal{F}g^2(\xi) = \int_{\mathbb{R}} e^{-i\xi x} g^2(x) dx = -\frac{i}{\xi} \int_{\mathbb{R}} e^{-i\xi x} (g^2)'(x) dx = -\frac{i}{\xi} [\mathcal{F}(g^2)'](\xi).$$

Hence Plancherel's identity yields

$$\begin{aligned} \int_{\mathbb{R}} \gamma_M^2(x) g^2(x) dx &= \int_{\{|\xi| \geq M\}} |\mathcal{F}g^2(\xi)| \mu(d\xi) = \int_{\{|\xi| \geq M\}} |[\mathcal{F}(g^2)'](\xi)| |\xi|^{-2H} d\xi \\ &\leq \left(\int_{\{|\xi| \geq M\}} |\xi|^{-4H} d\xi \right)^{1/2} \left(\int_{\mathbb{R}} |[\mathcal{F}(g^2)'](\xi)|^2 d\xi \right)^{1/2} = \frac{\left(\int_{\mathbb{R}} |[\mathcal{F}(g^2)'](\xi)|^2 d\xi \right)^{1/2}}{(4H-1)^{1/2} M^{2H-1/2}}, \end{aligned} \quad (3.21)$$

and it should be stressed that we are using our assumption $H > 1/4$ in order to get nondivergent integrals above. In addition, another use of Plancherel's identity enables us to write

$$\int_{\mathbb{R}} |[\mathcal{F}(g^2)'](\xi)|^2 d\xi = 2\pi \int_{\mathbb{R}} |(g^2)'(x)|^2 dx = 4\pi \int_{\mathbb{R}} g^2(x) |g'(x)|^2 dx. \quad (3.22)$$

We now get a uniform bound on g^2 . Since $g \in \mathcal{K}_1$ we have $\|g\|_2 = 1$ and thus

$$g^2(x) = 2 \int_{-\infty}^x g(y) g'(y) dy \leq 2 \|g\|_2 \|g'\|_2 = 2 \|g'\|_2.$$

Plugging this information into (3.22) we obtain

$$\int_{\mathbb{R}} |[\mathcal{F}(g^2)'](\xi)|^2 d\xi \leq 4\pi \|g'\|_2^3.$$

Finally, recalling (3.21) we can write

$$\theta \int_{\mathbb{R}} \gamma_M^2(x) g^2(x) dx - \delta \int_{\mathbb{R}} |g'(x)|^2 dx \leq \varphi(\|g'\|_2^{1/2}),$$

where $\varphi(x) = ax^3 - \delta x^4$ and $a = \frac{2\pi^{1/2}\theta}{(4H-1)^{1/2}M^{2H-1/2}}$. Optimizing over the values of $x \geq 0$ we get a bound which is uniform over $g \in \mathcal{K}_1$

$$\theta \int_{\mathbb{R}} \gamma_M^2(x)g^2(x)dx - \delta \int_{\mathbb{R}} |g'(x)|^2 dx \leq \frac{c\theta^4}{\delta^3 M^{2(4H-1)}},$$

with a (nonrelevant) constant $c = \frac{27\pi^2}{16}$. Plugging this uniform bound into (3.20), we thus easily end up with

$$\lim_{M \rightarrow \infty} \limsup_{m \rightarrow \infty} \frac{B_{m,M}^2}{m} = 0. \tag{3.23}$$

Step 4: Conclusion for the upper bound. Let us report our estimates (3.14) and (3.23) into the upper bound (3.13). This trivially yields

$$\limsup_{m \rightarrow \infty} \frac{1}{m} \sup_{g \in \mathcal{K}_m} K_{\theta,m}(g) \leq \left(\frac{\theta}{2}\right)^{\frac{1}{H}} (1-\delta)^{1-\frac{1}{H}} \mathcal{E}.$$

Step 5: Lower bound. In order to get the lower bound, we proceed to a direct verification, replacing the class \mathcal{K}_m by the smaller class

$$\mathcal{K}_0 \equiv \left\{ g(x) = \prod_{j=1}^m g_0(x_j), \quad g_0 \in \mathcal{K}_1 \right\}.$$

Towards this aim, we resort to the expression (3.8) of $K_{\theta,m}(g)$ in Fourier modes. Then, for $g \in \mathcal{K}_0$, it is readily checked that $\mathcal{F}g(\xi) = \prod_{k=1}^m \mathcal{F}g_0(\xi_k)$ for all $\xi = (\xi_1, \dots, \xi_k)$. Invoking this relation, plus the fact that $g_0 \in \mathcal{K}_1$, we get

$$\int_{\mathbb{R}^m} |\xi|^2 |\mathcal{F}g(\xi)|^2 d\xi = (2\pi)^{m-1} m \int_{\mathbb{R}} |\lambda|^2 |\mathcal{F}g_0(\lambda)|^2 d\lambda = (2\pi)^m m \int_{\mathbb{R}} |g'(x)|^2 dx,$$

and recalling that the functions \hat{g}_{jk} are introduced in (3.7),

$$\int_{\mathbb{R}} |\lambda|^{1-2H} \hat{g}_{jk}(\lambda) d\lambda = \int_{\mathbb{R}} |\lambda|^{1-2H} |\mathcal{F}g_0^2(\lambda)|^2 d\lambda.$$

Plugging those identities into the expression (3.8), we get

$$K_{\theta,m}(g) = -\frac{m}{2} \int_{\mathbb{R}} (g'_0(x))^2 dx + \frac{\theta(m-1)}{2} \int_{\mathbb{R}} |\lambda|^{1-2H} |\mathcal{F}g_0^2(\lambda)|^2 d\lambda.$$

This quantity is easily related to the variational expression $H_{\theta}(g_0)$ defined by (3.2), from which the identity

$$\lim_{m \rightarrow \infty} \sup_{g \in \mathcal{K}_0} \frac{K_{\theta,m}(g)}{m} = \left(\frac{\theta}{2}\right)^{\frac{1}{H}} \mathcal{E}$$

is readily checked. This shows our lower bound and finishes the proof. □

3.3 Feynman-Kac semi-groups

For $m \geq 1$ and $t \geq 0$, consider the quantity $Q_m(t)$ defined by (2.7). Our moment study for the solution to (1.1) will rely on the spectral behavior of the following semi-group acting on $L^2(\mathbb{R}^m)$

$$T_{m,t}g(x) = \mathbb{E}_x \left[g(B(t)) \exp\left(\frac{c_H Q_m(t)}{m}\right) \right], \tag{3.24}$$

where c_H is the constant defined in (2.2) and $x \in \mathbb{R}^m$ represents the initial condition for the m -dimensional Brownian motion B appearing in (2.7). We now establish some basic properties of those operators.

Proposition 3.6. Let $\{T_{m,t}, t \geq 0\}$ be the family of operators introduced in (3.24). Then the following properties hold true:

(i) The family $\{T_{m,t}, t \geq 0\}$ defines a semi-group of bounded self-adjoint operators on the space $L^2(\mathbb{R}^m)$.

(ii) The generator of $T_{m,t}$ is given on test functions by $A_m = A_{c_H,m}$ defined by (3.6), that is

$$A_m g = \frac{1}{2} \Delta g + \frac{c_H}{m} \sum_{1 \leq j < k \leq m} \gamma(x_j - x_k).$$

There exists an integer $m_0 \geq 1$ such that for $m \geq m_0$, A_m admits a self-adjoint extension (still denoted by A_m), defined on a domain $\text{Dom}(A_m) \subset L^2(\mathbb{R}^m)$.

Proof. Let us first prove the boundedness of $T_{m,t}$. To this aim, notice that for $g \in L^2(\mathbb{R}^m)$ we have

$$|T_{m,t} g(x)|^2 = \mathbb{E}_x^2 \left[g(B(t)) \exp \left(\frac{c_H Q_m(t)}{m} \right) \right].$$

We now apply Cauchy-Schwarz' inequality, together with relation (2.8), in order to get

$$|T_{m,t} g(x)|^2 \leq \mathbb{E}_x [g^2(B(t))] \mathbb{E}_x \left[\exp \left(\frac{2 c_H Q_m(t)}{m} \right) \right] \leq c_{m,H} \mathbb{E}_x [g^2(B(t))].$$

Taking into account the fact that $p_t(x - \cdot)$ is the density of B_t , we obtain

$$\begin{aligned} \|T_{m,t} g\|_{L^2(\mathbb{R}^m)}^2 &\leq c_{m,H} \int_{\mathbb{R}^m} \mathbb{E}_x [g^2(B_t)] dx = c_{m,H} \int_{\mathbb{R}^m} \left(\int_{\mathbb{R}^m} p_t(x - y) g^2(y) dy \right) dx \\ &= c_{m,H} \int_{\mathbb{R}^m} \left(\int_{\mathbb{R}^m} p_t(x - y) dx \right) g^2(y) dy = c_{m,H} \|g\|_{L^2(\mathbb{R}^m)}^2, \end{aligned}$$

which proves boundedness in $L^2(\mathbb{R}^m)$. The self-adjointness of $T_{m,t}$ is then easily derived.

It is also readily checked that the infinitesimal generator of $T_{m,t}$, acting on test functions in $\mathcal{S}(\mathbb{R}^m)$, is given by A_m . In addition, this operator is obviously symmetric. In order to show that it admits a self-adjoint extension, it is sufficient (thanks to the classical Freidrichs extension theorem) to show that

$$\langle A_m g, g \rangle \leq c \|g\|_{L^2(\mathbb{R}^m)}^2, \tag{3.25}$$

for all $g \in \text{Dom}(A_m)$ and for a constant $c > 0$. We now prove this inequality for m large enough: indeed, for all test functions we have $\langle A_m g, g \rangle = K_m(g)$, where we have set $K_m(g) = K_{c_H,m}(g)$ and $K_{\theta,m}(g)$ is defined by (3.5). The fact that there exists an $m_0 \geq 1$ such that relation (3.25) holds for $m \geq m_0$ is then an immediate consequence of Proposition 3.5. This finishes our proof. \square

As in the proof of Proposition 3.5, our future considerations will also rely on a truncated version of the operators $T_{m,t}$. Let us label their definition for further use.

Definition 3.7. Consider a parameter $M > 0$, and the function γ_M^1 defined by (3.11). For $m \geq 1$ we introduce the following quantities, defined similarly to (2.7) but replacing γ by the smoothed function γ_M^1 given by (3.11):

$$Q_{m,M}^1(t) = \sum_{1 \leq j < k \leq m} \int_0^t \gamma_M^1(B_j(s) - B_k(s)) ds, \tag{3.26}$$

$$\hat{Q}_{m,M}^1(t) = \sum_{1 \leq j, k \leq m} \int_0^t \gamma_M^1(B_j(s) - B_k(s)) ds. \tag{3.27}$$

Related to these Feynman-Kac functionals, we consider a family of operators $\{\hat{T}_{m,M,t}, t \geq 0\}$ acting on $L^2(\mathbb{R}^m)$, indexed by a parameter $\theta > 0$

$$\hat{T}_{m,M,t}g(x) = \mathbb{E}_x \left[g(B(t)) \exp \left(\frac{\theta \hat{Q}_{m,M}^1(t)}{m} \right) \right]. \tag{3.28}$$

The family of operators we have just introduced enjoys the following property.

Proposition 3.8. *Let $\{\hat{T}_{m,M,t}, t \geq 0\}$ be the families of operators introduced in (3.28). The conclusions of Proposition 3.6 remain true for this semi-group, with a generator $\hat{A}_{m,M}$ given by*

$$\hat{A}_{m,M}g = \frac{1}{2}\Delta g + \frac{\theta}{m} \sum_{1 \leq j,k \leq m} \gamma_M^1(x_j - x_k).$$

Furthermore, the following limiting behavior holds true for all $x \in \mathbb{R}^m$

$$\lim_{t \rightarrow \infty} \frac{\log \left(\hat{T}_{m,M,t} \mathbf{1}(x) \right)}{t} = \lambda_{m,M}, \tag{3.29}$$

where $\lambda_{m,M}$ is the principal eigenvalue of $\hat{A}_{m,M}$, defined similarly to (3.10) by

$$\lambda_{m,M} = \sup \left\{ \theta \sum_{j,k=1}^m \int_{\mathbb{R}^m} \gamma_M^1(x_j - x_k) g^2(x) dx - \frac{1}{2} \int_{\mathbb{R}^m} |\nabla g(x)|^2 dx; g \in \mathcal{K}_m \right\}, \tag{3.30}$$

and where we recall that \mathcal{K}_m is introduced in (3.9).

Proof. The self-adjointness of $T_{m,M,t}$ is completely classical and left to the reader. In order to get the self-adjointness of $A_{m,M}$, we just realize that $\gamma_M^1 : \mathbb{R} \rightarrow \mathbb{R}$ is the inverse Fourier transform of a function which is in $L^2(\mathbb{R})$ with compact support. Therefore γ_M^1 admits bounded derivatives of all order. The desired self-adjointness property is thus a consequence of classical results, which are summarized e.g in [3]. Finally relation (3.29) is a classical Feynman-Kac limit, for which we refer to [3, Theorem 4.1.6]. \square

4 Asymptotic properties of the moments

As in [4, 6, 8], the spatial asymptotics for equation (2.5) will be established thanks to a sharp estimate of the tail of $u(t, 0)$ for $t \geq 0$ fixed. As we will see later, this can be related to some estimates on the moments $u(t, 0)$, and our next step will be to obtain the exact asymptotic behavior (as $m \rightarrow \infty$) of these moments. Before we begin with this task, we will reduce our problem thanks to a series of lemmas.

First let us observe that the generator we have considered for a proper normalization procedure in Section 3.1 involves a sum of the form $\frac{1}{m} \sum_{1 \leq j < k \leq m} \gamma(x_j - x_k)$, where we emphasize the normalization by m . However, the quantity we manipulate in our Feynman-Kac representation (2.7) is $Q_m(t)$, which does not exhibit this normalizing term. We will introduce the missing normalization by a simple scaling argument.

Lemma 4.1. *Let u be the solution to (1.1). Then the moments of u admit the following representation*

$$\mathbb{E} [(u(t, 0))^m] = \mathbb{E}_0 \left[\exp \left(\frac{c_H Q_m(t_m)}{m} \right) \right], \tag{4.1}$$

where $t_m = m^{1/H}t$.

Proof. Start from expression (2.7) for $\mathbb{E}[(u(t, 0))^m]$, and recall that

$$Q_m(t) = \sum_{1 \leq j < k \leq m} \int_0^t \int_{\mathbb{R}} e^{i\xi(B_j(s) - B_k(s))} |\xi|^{1-2H} d\xi ds.$$

Setting $v = m^{1/H} s$ in the time integral, and invoking the fact that $B_j(m^{-1/H} \cdot)$ is equal in law to $m^{-1/2H} B_j(\cdot)$, we get

$$Q_m(t) \stackrel{(d)}{=} \frac{1}{m^{1/H}} \sum_{1 \leq j < k \leq m} \int_0^{t_m} \int_{\mathbb{R}} e^{i \frac{\xi(B_j(v) - B_k(v))}{m^{1/2H}}} |\xi|^{1-2H} d\xi dv.$$

We now set $m^{-1/2H} \xi = \lambda$ in the space integral above. This easily yields an equality in law between $Q_m(t)$ and $m^{-1} Q_m(t_m)$, and thus

$$\mathbb{E}[(u(t, 0))^m] = \mathbb{E}_0 \left[\exp \left(\frac{c_H Q_m(t_m)}{m} \right) \right], \tag{4.2}$$

which corresponds to our claim. □

We now establish a couple of simple monotonicity properties for the quantity Q_m which will feature in the sequel. The first one is related to our regularization procedure.

Lemma 4.2. *Let $Q_{m,M}^1$ be the quantity defined by (3.26) and let α be a positive constant. Then, the map $M \mapsto \mathbb{E}_0[\exp(\alpha Q_{m,M}^1(t))]$ is increasing.*

Proof. Consider $M > 0$. Then we have

$$\mathbb{E}_0 [\exp(\alpha Q_{m,M}^1(t))] = \sum_{n=0}^{\infty} \frac{\alpha^n \mathbb{E}_0 [(Q_{m,M}^1(t))^n]}{n!}. \tag{4.3}$$

We now show that, for any fixed $n \geq 1$, the map $M \mapsto \mathbb{E}_0[(Q_{m,M}^1(t))^n]$ is increasing. Indeed, recalling our notation (3.12) for μ_M^1 , it is readily checked that

$$\mathbb{E}_0 [(Q_{m,M}^1(t))^n] = \prod_{l=1}^n \sum_{1 \leq j_l < k_l \leq m} \int_{[0,t]^n} \int_{\mathbb{R}^n} \mathbb{E}_0 \left[e^{i \sum_{l=1}^n \xi_l (B_{j_l}(s_l) - B_{k_l}(s_l))} \right] [\mu_M^1]^{\otimes n}(d\xi) ds, \tag{4.4}$$

where we use the simple convention $d\xi = d\xi_1 \cdots d\xi_n$ and $ds = ds_1 \cdots ds_n$. Now notice that for all j, k, l we have

$$\mathbb{E}_0 \left[e^{i \sum_{l=1}^m \xi_l (B_{j_l}(s_l) - B_{k_l}(s_l))} \right] > 0.$$

This easily yields the fact that $M \mapsto \mathbb{E}_0[(\alpha Q_{m,M}^1(t))^n]$ is increasing. Going back to our decomposition (4.3), we have thus obtained that $M \mapsto \mathbb{E}_0[\exp(\alpha Q_{m,M}^1(t))]$ is increasing. □

The second monotonicity property we need concerns the dependence with respect to the initial condition for our underlying Brownian motion.

Lemma 4.3. *Recall that $B = \{B(s) = (B_1(s), \dots, B_m(s)), s \geq 0\}$ is an m -dimensional Brownian motion. For any $x \in \mathbb{R}^m$, recall that we use \mathbb{E}_x to denote the mathematical expectation with respect to B with $B(0) = x$. Then the following relation is verified for all $m \geq 1, t \geq 0, \theta \geq 0$ and $x \in \mathbb{R}^m$*

$$\mathbb{E}_0 \left[\exp \left(\frac{\theta Q_m(t)}{m} \right) \right] \geq \mathbb{E}_x \left[\exp \left(\frac{\theta Q_m(t)}{m} \right) \right].$$

For any fixed $M > 0$, the same property holds true for $Q_{m,M}^1(t)$ and $Q_{m,M}^2(t)$.

Proof. We focus on the property for Q_m , the equivalent for $Q_{m,M}^i$ ($i = 1, 2$) being shown exactly in the same way. Furthermore, resorting to expansion (4.3) as in the previous proof, our claim can be reduced to show that

$$\mathbb{E}_0 [(Q_m(t))^n] \geq \mathbb{E}_x [(Q_m(t))^n], \quad \text{for all } n \geq 1, \tag{4.5}$$

We now focus on relation (4.5).

In order to show (4.5), we first decompose the moments of $Q_m(t)$ similarly to (4.4)

$$\mathbb{E}_x [(Q_m(t))^n] = \int_{[0,t]^n} \int_{(\mathbb{R}^m)^n} \mathbb{E}_0 [D_{m,n}] \mu^{\otimes n}(d\xi) ds, \tag{4.6}$$

where we have set $D_{m,n} = \prod_{l=1}^n \sum_{1 \leq j_l < k_l \leq m} e^{i\xi_l(x_j - x_k)} e^{i\xi_l(B_{j_l}(s_l) - B_{k_l}(s_l))}$. Furthermore, one can reorder terms in the quantity $D_{m,n}$, and obtain

$$D_{m,n} = \sum_{\substack{j_1, \dots, j_n, k_1, \dots, k_n=1 \\ j_l < k_l, \forall l=1, \dots, n}}^m C(j_1, \dots, j_n, k_1, \dots, k_n) e^{i \sum_{l=1}^n \xi_l (B_{j_l}(s_l) - B_{k_l}(s_l))},$$

where the constants $C(j_1, \dots, j_n, k_1, \dots, k_n)$ satisfy $|C(j_1, \dots, j_n, k_1, \dots, k_n)| = 1$. It should also be noticed that

$$\sum_{\substack{j_1, \dots, j_n, k_1, \dots, k_n=1 \\ j_l < k_l, \forall l=1, \dots, n}}^m e^{i \sum_{l=1}^n \xi_l (B_{j_l}(s_l) - B_{k_l}(s_l))} = \prod_{l=1}^n \sum_{1 \leq j < k \leq m} e^{i \xi_l (B_j(s_l) - B_k(s_l))}. \tag{4.7}$$

Hence, taking the mathematical expectation yields

$$\begin{aligned} \mathbb{E}_0 [D_{m,n}] &= \left| \sum_{\substack{j_1, \dots, j_n, k_1, \dots, k_n=1 \\ j_l < k_l, \forall l=1, \dots, n}}^m C(j_1, \dots, j_n, k_1, \dots, k_n) \mathbb{E}_0 \left[e^{i \sum_{l=1}^n \xi_l (B_{j_l}(s_l) - B_{k_l}(s_l))} \right] \right| \\ &\leq \sum_{\substack{j_1, \dots, j_n, k_1, \dots, k_n=1 \\ j_l < k_l, \forall l=1, \dots, n}}^m \mathbb{E}_0 \left[e^{i \sum_{l=1}^n \xi_l (B_{j_l}(s_l) - B_{k_l}(s_l))} \right], \end{aligned}$$

where the inequality follows from the fact that $\mathbb{E}[e^{i \sum_{l=1}^n \xi_l (B_{j_l}(s_l) - B_{k_l}(s_l))}] > 0$. Plugging this information into (4.6) and taking (4.7) into account, we thus end up with.

$$\begin{aligned} \mathbb{E}_x [(Q_m(t))^n] &\leq \int_{[0,t]^n} \int_{(\mathbb{R}^m)^n} \mathbb{E}_0 \left[\prod_{l=1}^n \sum_{1 \leq j < k \leq m} e^{i \xi_l (B_j(s_l) - B_k(s_l))} \right] \mu^{\otimes n}(d\xi) ds \\ &= \mathbb{E}_0 [(Q_m(t))^n]. \end{aligned}$$

This proves (4.5), and thus our claim. □

We are now ready to prove our main asymptotic theorem for moments of u .

Theorem 4.4. *Let $t \in \mathbb{R}_+$ and $x \in \mathbb{R}$. Then the following asymptotic result holds true*

$$\lim_{m \rightarrow \infty} \frac{1}{m^{1+\frac{1}{H}}} \log (\mathbb{E} [(u(t, x))^m]) = \left(\frac{c_H}{2}\right)^{\frac{1}{H}} \mathcal{E} t, \tag{4.8}$$

where $\mathcal{E} = \mathcal{E}_1$ denotes the constant given by relation (3.3).

Proof. Recall that the law of $u(t, x)$ does not depend on x , and we thus start from expression (4.1) for $\mathbb{E}[(u(t, 0))^m]$. With a lower bound in mind, our first task will be to relate this expression to the semi-group $T_{m,t}$ introduced in (3.24).

Step 1: Spectral representation. Let us consider a function $g : \mathbb{R} \rightarrow \mathbb{R}$ with support in a given interval $[-a, a]$. For $m \geq 1$ we define $g_m = g^{\otimes m}$, that is,

$$g_m : \mathbb{R}^m \longrightarrow \mathbb{R}^m, \quad x = (x_1, \dots, x_m) \longmapsto \prod_{j=1}^m g(x_j).$$

Owing to Lemma 4.3 plus the trivial relation $\frac{g(x)}{\|g\|_\infty} \leq 1$, we have

$$\begin{aligned} \mathbb{E}[(u(t, 0))^m] &\geq (2a)^{-m} \int_{[-a, a]^m} \mathbb{E}_x \left[\exp \left(\frac{c_H Q_m(t_m)}{m} \right) \right] dx \\ &\geq (2\|g\|_\infty^2 a)^{-m} \int_{[-a, a]^m} g_m(x) \mathbb{E}_x \left[g_m(B(t_m)) \exp \left(\frac{c_H Q_m(t_m)}{m} \right) \right] dx. \end{aligned}$$

Invoking our definition (3.24), we have thus obtained

$$\mathbb{E}[(u(t, 0))^m] \geq \frac{\langle T_{m, t_m} g_m, g_m \rangle_{L^2(\mathbb{R}^m)}}{(2\|g\|_\infty^2 a)^m}. \tag{4.9}$$

We now consider a small parameter $\varepsilon > 0$. We also recall our definitions (3.5) and (3.6) for $K_{\theta, m}$ and $A_{\theta, m}$, and set $K_m := K_{c_H, m}$ (resp. $A_m := A_{c_H, m}$) to alleviate notations. Owing to our computations in the proof of Proposition 3.5 (Step 5), we can choose g satisfying $\|g\|_{L^2(\mathbb{R})} = 1$ and such that

$$\liminf_{m \rightarrow \infty} \frac{K_m(g_m)}{m} \geq \mathcal{E}_H - \varepsilon, \quad \text{where } \mathcal{E}_H = \left(\frac{c_H}{2}\right)^{\frac{1}{H}} \mathcal{E}. \tag{4.10}$$

This quantity can be related to the semi-group $T_{m, t}$ in the following way. Due to the fact that A_m is self-adjoint (see Proposition 3.6), $T_{m, t}$ admits a spectral representation related to its generator. Hence, since $\|g\|_{L^2(\mathbb{R})} = 1$, there exists a probability measure ν_g on \mathbb{R} such that:

$$\langle T_{m, t_m} g_m, g_m \rangle_{L^2(\mathbb{R}^m)} = \int_{\mathbb{R}} \exp(t_m \lambda) \nu_g(d\lambda) \geq \exp \left(t_m \int_{\mathbb{R}} \lambda \nu_g(d\lambda) \right),$$

which yields

$$\langle T_{m, t_m} g_m, g_m \rangle_{L^2(\mathbb{R}^m)} \geq \exp(t_m \langle A_m g, g \rangle_{L^2(\mathbb{R}^m)}) = \exp(t_m K_m(g_m)). \tag{4.11}$$

Plugging this information into (4.9) and recalling that g satisfies (4.10), we end up with the following inequality, valid for m large enough

$$\mathbb{E}[(u(t, 0))^m] \geq \frac{\exp(t_m K_m(g_m))}{(2\|g\|_\infty^2 a)^m} \geq \frac{\exp(mt_m (\mathcal{E}_H - \varepsilon))}{(2\|g\|_\infty^2 a)^m}. \tag{4.12}$$

Step 2: Lower bound. Starting from (4.12), the desired lower bound is now easily derived: taking into account the fact that $mt_m = m^{1+\frac{1}{H}} t$, we get (for m large enough)

$$\frac{1}{m^{1+\frac{1}{H}}} \log(\mathbb{E}[(u(t, 0))^m]) \geq (\mathcal{E}_H - \varepsilon) t - \frac{\log(2\|g\|_\infty^2 a)}{m^{\frac{1}{H}}}.$$

Taking limits in m and recalling that we have chosen an arbitrarily small ε , we have proved that

$$\liminf_{m \rightarrow \infty} \frac{1}{m^{1+\frac{1}{H}}} \log(\mathbb{E}[(u(t, 0))^m]) \geq \mathcal{E}_H t,$$

which corresponds to our claim.

Step 3: Cut-off in the Feynman-Kac representation. Let us go back to the Feynman-Kac representation of moments for $u(t, x)$ given by (2.7), and write the quantity $Q_m(t)$ therein as

$$Q_m(t) = \sum_{1 \leq j < k \leq m} \int_0^t \int_{\mathbb{R}} e^{i\xi(B_j(s) - B_k(s))} \mu(d\xi) ds.$$

In order to get the upper bound part of our theorem, we shall replace the quantity $Q_m(t)$, with its diverging high frequency modes, by the quantity $Q_{m,M}^1(t)$ defined in (3.26). To this aim, recall that $\gamma_M^1, \gamma_M^2, \mu_M^1, \mu_M^2$ are defined in (3.11)–(3.12), and notice that we have

$$Q_m(t) = Q_{m,M}^1(t) + Q_{m,M}^2(t), \quad \text{with} \quad Q_{m,M}^1(t) = \sum_{1 \leq j < k \leq m} \int_0^t \gamma_M^1(B_j(s) - B_k(s)) ds. \tag{4.13}$$

Our cut-off procedure is now expressed in the following form. For two conjugate exponents $p, q > 1$, Hölder’s inequality yields

$$\mathbb{E} [(u(t, 0))^m] \leq \mathbb{E}_0^{1/p} \left[\exp \left(\frac{p c_H Q_{m,M}^1(t_m)}{m} \right) \right] \mathbb{E}_0^{1/q} \left[\exp \left(\frac{q c_H Q_{m,M}^2(t_m)}{m} \right) \right]. \tag{4.14}$$

We will prove that our study can be reduced to analysis of the term $Q_{m,M}^1$ by showing that for an arbitrary $q \geq 1$

$$\lim_{M \rightarrow \infty} \limsup_{m \rightarrow \infty} \frac{1}{m^{1+\frac{1}{H}}} \log \left(\mathbb{E}_0^{1/q} \left[\exp \left(\frac{q c_H Q_{m,M}^2(t_m)}{m} \right) \right] \right) \leq 0. \tag{4.15}$$

Step 4: Proof of (4.15). Let us generalize somehow our problem, and show that for any $\theta > 0$ we have

$$\lim_{M \rightarrow \infty} \limsup_{m \rightarrow \infty} \frac{1}{m^{1+\frac{1}{H}}} \log (R_{\theta, M, m}) \leq 0, \quad \text{with} \quad R_{\theta, M, m} = \mathbb{E}_0 \left[\exp \left(\frac{\theta Q_{m,M}^2(t_m)}{m} \right) \right]. \tag{4.16}$$

Also notice that

$$Q_{m,M}^2(t_m) = \frac{1}{2} \sum_{j=1}^m \sum_{k \neq j} \int_0^{t_m} \int_{\mathbb{R}} e^{i\xi(B_j(s) - B_k(s))} \mu_M^2(d\xi) ds.$$

A rough bound on $R_{\theta, M, m}$ is thus obtained by applying Hölder’s inequality, similarly to [7, Section 3]

$$\begin{aligned} R_{\theta, M, m} &\leq \prod_{j=1}^m \mathbb{E}_0^{1/m} \left[\exp \left(\frac{\theta}{2} \sum_{k \neq j} \int_0^{t_m} \int_{\mathbb{R}} e^{i\xi(B_j(s) - B_k(s))} \mu_M^2(d\xi) ds \right) \right] \\ &= \mathbb{E}_0 \left[\exp \left(\frac{\theta}{2} \sum_{j=2}^m \int_0^{t_m} \int_{\mathbb{R}} e^{i\xi(B_j(s) - B_1(s))} \mu_M^2(d\xi) ds \right) \right]. \end{aligned} \tag{4.17}$$

We shall now prove that

$$R_{\theta, M, m} \leq \mathbb{E}_0^{m-1} \left[\exp \left(\frac{\theta}{2} \int_0^{t_m} \int_{\mathbb{R}} e^{i\xi B(s)} \mu_M^2(d\xi) ds \right) \right]. \tag{4.18}$$

Indeed, a simple series expansion for the exponential function in (4.17) reveals that $R_{\theta, M, m} \leq \sum_{n \geq 0} \frac{\theta^n}{2^n n!} L_n$, with

$$L_n = \sum_{j_1, \dots, j_n=2}^m \int_{[0, t_m]^n} \int_{\mathbb{R}^n} \mathbb{E}_0 \left[e^{i\xi \sum_{l=1}^n (B_{j_l}(s) - B_1(s))} \right] \prod_{l=1}^n \mu_M^2(d\xi_l) ds_l.$$

In addition, observe that

$$\mathbb{E}_0 \left[e^{i\xi \sum_{l=1}^n (B_{j_l}(s) - B_1(s))} \right] = \mathbb{E}_0 \left[e^{i\xi \sum_{l=1}^n B_{j_l}(s)} \right] \mathbb{E}_0 \left[e^{-i\xi \sum_{l=1}^n B_1(s)} \right].$$

Owing to the fact that $\mathbb{E}_0[e^{i\xi \sum_{l=1}^n B_{j_l}(s)}] \geq 0$ and $\mathbb{E}_0[e^{-i\xi \sum_{l=1}^n B_1(s)}] \in (0, 1)$, we end up with

$$\begin{aligned} L_n &\leq \sum_{j_1, \dots, j_n=2}^m \int_{[0, t_m]^n} \int_{\mathbb{R}^n} \mathbb{E}_0 \left[e^{i\xi \sum_{l=1}^n B_{j_l}(s)} \right] \prod_{l=1}^n \mu_M^2(d\xi_l) ds_l \\ &= \mathbb{E}_0 \left[\left| \sum_{j=2}^m \int_0^{t_m} \int_{\mathbb{R}} e^{i\xi B_j(s)} \mu_M^2(d\xi) ds \right|^n \right]. \end{aligned}$$

Invoking the series expansion of the exponential function again, this last bound easily yields our claim (4.18).

Starting from (4.18), we can prove (4.15). To this aim, a direct consequence of (4.18) is the following inequality

$$\frac{1}{m^{1+\frac{1}{H}}} \log(R_{\theta, M, m}) \leq \frac{1}{m^{\frac{1}{H}}} \log \left(\mathbb{E}_0 \left[\exp \left(\frac{\theta}{2} \int_0^{t_m} \int_{\mathbb{R}} e^{i\xi B(s)} \mu_M^2(d\xi) ds \right) \right] \right). \quad (4.19)$$

Furthermore, some positivity arguments similar to the one we resorted to in Lemma 4.2 show that

$$\mathbb{E}_0 \left[\exp \left(\frac{\theta}{2} \int_0^{t_m} \int_{\mathbb{R}} e^{i\xi B(s)} \mu_M^2(d\xi) ds \right) \right] \leq \mathbb{E}_0 \left[\exp \left(\frac{\theta}{2} \int_0^{\lceil t_m \rceil + 1} \int_{\mathbb{R}} e^{i\xi B(s)} \mu_M^2(d\xi) ds \right) \right],$$

where $\lceil x \rceil$ denotes the integer part of a real number x . Hence, applying successively Markov's property for B and Lemma 4.3, we get

$$\begin{aligned} \mathbb{E}_0 \left[\exp \left(\frac{\theta}{2} \int_0^{t_m} \int_{\mathbb{R}} e^{i\xi B(s)} \mu_M^2(d\xi) ds \right) \right] &\leq \sup_{x \in \mathbb{R}} \mathbb{E}_x^{\lceil t_m \rceil + 1} \left[\exp \left(\frac{\theta}{2} \int_0^1 \int_{\mathbb{R}} e^{i\xi B(s)} \mu_M^2(d\xi) ds \right) \right] \\ &\leq \mathbb{E}_0^{\lceil t_m \rceil + 1} \left[\exp \left(\frac{\theta}{2} \int_0^1 \int_{\mathbb{R}} e^{i\xi B(s)} \mu_M^2(d\xi) ds \right) \right]. \end{aligned}$$

Plugging this inequality into (4.19) yields

$$\begin{aligned} \frac{1}{m^{1+\frac{1}{H}}} \log(R_{\theta, M, m}) &\leq \frac{\lceil t_m \rceil + 1}{m^{1+\frac{1}{H}}} \log \left(\mathbb{E}_0 \left[\exp \left(\frac{\theta}{2} \int_0^1 \int_{\mathbb{R}} e^{i\xi B(s)} \mu_M^2(d\xi) ds \right) \right] \right) \\ &\leq 2t \log \left(\mathbb{E}_0 \left[\exp \left(\frac{\theta}{2} \int_0^1 \int_{\mathbb{R}} e^{i\xi B(s)} \mu_M^2(d\xi) ds \right) \right] \right). \end{aligned}$$

Summarizing, we have obtained that

$$\lim_{M \rightarrow \infty} \limsup_{m \rightarrow \infty} \frac{1}{m^{1+\frac{1}{H}}} \log(R_{\theta, M, m}) \leq \lim_{M \rightarrow \infty} \log \left(\mathbb{E}_0 \left[\exp \left(\frac{\theta}{2} \int_0^1 \int_{\mathbb{R}} e^{i\xi B(s)} \mu_M^2(d\xi) ds \right) \right] \right) = 0,$$

where the last relation can be seen by means of some elementary computations, very similar to the ones displayed in [14, p. 50]. This finishes the proof of (4.15).

Step 5: An expression with diagonal terms. Let us now focus on the behavior of the quantity $Q_{m, M}^1(t)$ defined by (4.13). Specifically, having (4.13) and (4.14) in mind, we wish to find the asymptotic behavior of

$$\mathcal{A}_{m, M}^1(\theta) \equiv \mathbb{E}_0 \left[\exp \left(\frac{\theta Q_{m, M}^1(t_m)}{m} \right) \right],$$

for a given parameter $\theta > 0$. We wish to reduce this asymptotic study to an evaluation involving Feynman-Kac semigroups. A first step in this direction is to realize that, since the cut-off measure μ_M^1 is now finite, the asymptotic behavior of $m \mapsto \mathbb{E}[(u(t, x))^m]$ is not perturbed by adding the diagonal terms corresponding to $j = k$ in the sum defining $Q_{m,M}^1(t)$. That is, one can replace $Q_{m,M}^1(t)$ by

$$\hat{Q}_{m,M}^1(t) = \sum_{j,k=1}^m \int_0^t \int_{\mathbb{R}} e^{i\xi(B_j(s)-B_k(s))} \mu_M^1(d\xi) ds = \int_0^t \int_{\mathbb{R}} |h_B(s, \xi)|^2 \mu_M^1(d\xi) ds, \quad (4.20)$$

where we have set $h_B(s, \xi) = \sum_{j=1}^m e^{i\xi B_j(s)}$.

Indeed, let us recast Definition (3.26) as

$$Q_{m,M}^1(t) = \sum_{1 \leq j < k \leq m} q_{m,M}^{jk}(t), \quad \text{where} \quad q_{m,M}^{jk}(t) = \int_0^t \int_{\mathbb{R}} e^{i\xi(B_j(s)-B_k(s))} \mu_M^1(d\xi) ds.$$

Now it is readily checked that the diagonal terms $q_{m,M}^{jj}$ can be expressed as

$$q_{m,M}^{jj}(t) = \int_0^t \int_{-M}^M |\xi|^{1-2H} d\xi ds = c_{1,H} t M^{2-2H}.$$

Thus, recalling the definition (4.20) of $\hat{Q}_{m,M}^1(t) = \sum_{j,k=1}^m q_{m,M}^{jk}(t)$, we have

$$\hat{\mathcal{A}}_{n,M}^1(\theta) := \mathbb{E}_0 \left[\exp \left(\frac{\theta \hat{Q}_{m,M}^1(t_m)}{m} \right) \right] = \exp \left(\theta t M^{2-2H} m^{\frac{1}{H}} \right) \mathcal{A}_{n,M}^1(2\theta), \quad (4.21)$$

where we have replaced θ by 2θ as a parameter of $\mathcal{A}_{n,M}^1$, due to repetitions of off diagonal terms in the definition of $\hat{Q}_{m,M}^1$. This easily entails that

$$\limsup_{m \rightarrow \infty} m^{-\frac{1+H}{H}} \log \left(\mathcal{A}_{n,M}^1(\theta) \right) = \limsup_{m \rightarrow \infty} m^{-\frac{1+H}{H}} \log \left(\hat{\mathcal{A}}_{n,M}^1 \right). \quad (4.22)$$

Recalling once again our formula (2.7), we now focus on the evaluation of $\hat{\mathcal{A}}_{n,M}^1$.

Step 6: A coarse graining procedure. Thanks to relation (4.22), our problem is reduced to a Feynman-Kac asymptotics for the semi-group related to $\hat{\mathcal{A}}_{n,M}^1$. However, the semi-group $\hat{T}_{m,M,t}$ is considered in Proposition 3.8 with m, M fixed and $t \rightarrow \infty$. In contrast, we consider here a situation where both m and t_m are going to ∞ . We solve this problem by a coarse graining type procedure which is described below.

To this aim, consider a fixed $\rho \in \mathbb{N}$ and let us decompose m as $m = n\rho + r$ with $n, r \in \mathbb{N}$ and $0 \leq r \leq n - 1$. We can write

$$\hat{Q}_{m,M}^1(t) = \int_0^t \int_{\mathbb{R}} \left| \sum_{j=0}^n \sum_{k=0}^{\rho_j-1} e^{i\xi B_{j\rho+k}(s)} \right|^2 \mu_M^1(d\xi) ds,$$

where $\rho_j = \rho$ for $0 \leq j \leq n - 1$ and $\rho_n = r + 1$. Then invoke the elementary relation $|\sum_{j=0}^n z_j|^2 \leq (n + 1) \sum_{j=0}^n |z_j|^2$, valid on \mathbb{C} , to get

$$\hat{Q}_{m,M}^1(t) \leq (n + 1) \sum_{j=0}^n \int_0^t \int_{\mathbb{R}} \left| \sum_{k=0}^{\rho_j-1} e^{i\xi B_{j\rho+k}(s)} \right|^2 \mu_M^1(d\xi) ds.$$

Thanks to the independence of the Brownian motions and the definition (4.21) of $\hat{\mathcal{A}}_{n,M}^1$, we thus obtain

$$\hat{\mathcal{A}}_{m,M}^1 = \mathbb{E}_0 \left[\exp \left(\frac{\theta \hat{Q}_{m,M}^1(t_m)}{m} \right) \right] \leq \left| \hat{I}_{\rho,M}(t_m) \right|^n \hat{R}_{m,M}(t_m), \quad (4.23)$$

where we set $\theta_\rho = \frac{\theta(n+1)}{m}$ and

$$\hat{I}_{\rho,M}(t_m) = \mathbb{E}_0 \left[\exp \left(\theta_\rho \int_0^{t_m} \int_{\mathbb{R}} \left| \sum_{k=1}^{\rho} e^{i\xi B_k(s)} \right|^2 \mu_M^1(d\xi) ds \right) \right].$$

In relation (4.23), the remainder term $\hat{R}_{m,M}(t_m)$ is defined by

$$\hat{R}_{m,M}(t_m) = \mathbb{E}_0 \left[\exp \left(\theta_\rho \int_0^{t_m} \int_{\mathbb{R}} \left| \sum_{k=1}^r e^{i\xi B_k(s)} \right|^2 \mu_M^1(d\xi) ds \right) \right],$$

and notice that we also have

$$\left| \theta_\rho - \frac{\theta}{\rho} \right| \leq \frac{\theta}{\rho n}. \tag{4.24}$$

Let us now evaluate the term $\hat{I}_{\rho,M}(t_m)$ above. To this aim, we linearize our expression again and write

$$\int_0^t \int_{\mathbb{R}} \left| \sum_{k=1}^{\rho} e^{i\xi B_k(s)} \right|^2 \mu_M^1(d\xi) ds = \sum_{j,k=1}^{\rho} \int_0^t \gamma_M^1(B_j(s) - B_k(s)) ds.$$

Related to this expression, consider the semi-group $\hat{T}_{\rho,M,t}$ defined by (3.28), for a parameter $\theta = \theta_\rho$. Then it is readily checked that $\hat{I}_{\rho,M}(t_m) = \hat{T}_{\rho,M,t} \mathbf{1}(0)$. We can now apply Proposition 3.8 and relation (4.24), which yields

$$\lim_{m \rightarrow \infty} \frac{\log \left(\hat{I}_{\rho,M}(t_m) \right)}{t_m} = \lambda_{\rho,M}, \tag{4.25}$$

where $\lambda_{\rho,M}$ is given by (3.30) with θ replaced by $\frac{\theta}{\rho}$. Along the same lines, one can also check that

$$\lim_{m \rightarrow \infty} \frac{\log \left(\hat{R}_{m,M}(t_m) \right)}{t_m} = 0. \tag{4.26}$$

Step 7: Conclusion for the upper bound. Let us go back to decomposition (4.23) and invoke relation (4.26) in order to write

$$\begin{aligned} \frac{1}{m^{1+\frac{1}{H}}} \log \left(\hat{\mathcal{A}}_{m,M}^1(\theta) \right) &\leq \frac{n}{m^{1+\frac{1}{H}}} \log \left(\hat{I}_{\rho,M}(t_m) \right) + \varepsilon_{m,M} \\ &= \left(\frac{t}{\rho} \right) \left(\frac{m-r}{m} \right) \frac{1}{t_m} \log \left(\hat{I}_{\rho,M}(t_m) \right) + \varepsilon_{m,M}, \end{aligned}$$

where $\lim_{m \rightarrow \infty} \varepsilon_{m,M} = 0$. Plugging the limiting behavior (4.25) into this inequality, we end up with

$$\limsup_{m \rightarrow \infty} \frac{1}{m^{1+\frac{1}{H}}} \log \left(\hat{\mathcal{A}}_{m,M}^1(\theta) \right) \leq t \frac{\lambda_{\rho,M}}{\rho}.$$

This bound is valid for any $\rho \in \mathbb{N}$, and a small variation of Proposition 3.5 (replacing the distribution γ by its smoothed version γ_M^1 and summing over diagonal terms) asserts that for all $M > 0$ there exists a quantity $\lambda_M(\theta)$ verifying:

$$\lim_{\rho \rightarrow \infty} \frac{\lambda_{\rho,M}}{\rho} = \lambda_M(\theta), \quad \text{and} \quad \lim_{M \rightarrow \infty} \lambda_M(\theta) = \mathcal{E}_\theta.$$

This yields our upper bound taking into account (4.15), the fact that we consider $\mathcal{A}_{m,M}^1(\theta)$ with $\theta = \frac{pcH}{2}$ and p arbitrarily close to 1, plus relation (4.22). \square

Using Corollary 1.2.5 in [3] we deduce the following corollary.

Corollary 4.5. Consider the solution u of equation (1.1). Then the following tail estimate holds true

$$\lim_{a \rightarrow \infty} a^{-(1+H)} \log (\mathbb{P} (\log (u(t, x)) \geq a)) = -\hat{c}_{H, t}, \tag{4.27}$$

where the constant $\hat{c}_{H, t}$ is defined by

$$\hat{c}_{H, t} = \left[\frac{(1+H)c_H}{2} \left(\left(1 + \frac{1}{H} \right) t\mathcal{E} \right)^H \right]^{-1}.$$

Remark 4.6. Denote by $q_{H, t}$ the quantity showing up in the right hand side of equation (1.3). Then we have $\hat{c}_{H, t} = q_{H, t}^{-(1+H)}$.

Proof of Corollary 4.5. Recall that we have proved Theorem 4.4. Furthermore, some simple arguments based on Hölder’s inequality allow us to extend this result to real valued powers. Namely, the limit in (4.8) can be taken along positive real numbers instead of integers.

Let us thus consider a sequence of real numbers $(a_n)_{n \geq 1}$ converging to ∞ . Related to this sequence, we also introduce the sequence of random variables $Y_n = a_n^{-1/H} \log(u(t, x))$, and the sequence $(\rho_n)_{n \geq 1}$ with $\rho_n = a_n^{1+\frac{1}{H}}$. Then a direct application of (4.8) yields, for an additional parameter $\beta \geq 0$

$$\lim_{n \rightarrow \infty} \frac{1}{\rho_n} \log (\mathbb{E} [\exp (\beta \rho_n Y_n)]) = \Lambda(\beta), \quad \text{where} \quad \Lambda(\beta) \equiv \left(\frac{c_H}{2} \right)^{\frac{1}{H}} t\mathcal{E} \beta^{1+\frac{1}{H}}. \tag{4.28}$$

Notice that the above asymptotic result does not enable a direct application of Ellis-Gartner’s theorem, since the limit in (4.28) is only obtained for $\beta \geq 0$ (while Ellis-Gartner would require limits for $\beta < 0$ too). We will thus apply a large deviation theorem for positive random variables (Theorem 1.2.3 in [3]), which can be summarized as follows. Assume relation (4.28) holds true for $\beta \geq 0$, and define Λ^* as

$$\Lambda^*(\lambda) = \sup \{ \lambda\beta - \Lambda(\beta) \}. \tag{4.29}$$

If Λ^* is smooth and convex on $(0, \infty)$, and if $\lim_{\lambda \rightarrow \infty} \Lambda^*(\lambda) = \infty$, then the following tail estimate holds true for $\lambda > 0$

$$\lim_{n \rightarrow \infty} \frac{1}{\rho_n} \log (\mathbb{P} (Y_n \geq \lambda)) = -\Lambda^*(\lambda). \tag{4.30}$$

The application of the latter result raises 2 additional questions that we address now, similarly to what is done in [5]:

(i) Our random variable Y_n is not positive. However, observe that

$$\text{sgn}(Y_n) = \text{sgn}(Y_1) = \text{sgn}(\log(u(t, x))), \quad \text{and} \quad \mathbb{P}(Y_1 > 0) > 0.$$

We thus decompose the exponential moments of Y_n into

$$\begin{aligned} & \mathbb{E} [\exp (\beta \rho_n Y_n)] \\ &= \mathbb{E} [\exp (\beta \rho_n Y_n) \mid Y_1 \leq 0] \mathbb{P} (Y_1 \leq 0) + \mathbb{E} [\exp (\beta \rho_n Y_n) \mid Y_1 > 0] \mathbb{P} (Y_1 > 0) \\ &\leq 1 + \mathbb{E} [\exp (\beta \rho_n Y_n) \mid Y_1 > 0]. \end{aligned}$$

Owing to this relation, plus our limiting result (4.28), it is readily checked that

$$\lim_{n \rightarrow \infty} \frac{1}{\rho_n} \log (\mathbb{E} [\exp (\beta \rho_n Y_n) \mid Y_1 > 0]) = \Lambda(\beta).$$

Furthermore, conditioned to the event $(Y_1 > 0)$, the random variable Y_n can be considered as positive.

(ii) One can easily compute Λ^* thanks to relation (4.29), and the reader can check that $\Lambda^*(\lambda) = \hat{c}_{H,t} \lambda^{1+H}$. In particular, Λ^* is smooth and convex on $(0, \infty)$, and we also have $\lim_{\lambda \rightarrow \infty} \Lambda^*(\lambda) = \infty$.

Taking into account the last considerations, we have thus obtained a conditioned version of (4.30), that is

$$\lim_{n \rightarrow \infty} \frac{1}{\rho_n} \log (\mathbb{P} (Y_n \geq \lambda | Y_1 > 0)) = -\Lambda^*(\lambda).$$

One can then transform this relation into a nonconditioned one, owing to the fact that $\mathbb{P}(Y_1 > 0)$ is a fixed strictly positive quantity. Recalling the notation for Y_n and ρ_n , we end up with

$$\lim_{n \rightarrow \infty} \frac{1}{a_n^{1+\frac{1}{H}}} \log \left(\mathbb{P} \left(\log(u(t, x)) \geq a_n^{\frac{1}{H}} \lambda \right) \right) = -\hat{c}_{H,t} \lambda^{1+H}.$$

Some elementary changes of variable now yield our claim (4.27). □

Remark 4.7. Notice that the constant $\hat{c}_{H,t}$ appearing in (4.27) is precisely $c_0(H)^{-\frac{1}{1+H}} (t\mathcal{E})^{-H}$, where $c_0(H)$ is defined in (1.4).

5 Proof of Theorem 1.1

In this section we start from the tail behavior for the random variable $\log(u(t, x))$ provided by (4.27). We carry out a localization and discretization procedure which will allow us to evaluate the growth of $x \mapsto u(t, x)$.

5.1 Proof of the lower bound by localization

Our approach is based on a method (introduced in [8] and [9]) involving localizations of the driving noise \dot{W} in the space-time domain. Let us start by some elementary preliminaries (whose proofs are left to the reader) concerning the localizing function.

Lemma 5.1. *Let ℓ be the function defined by:*

$$\ell(x) = \frac{1 - \cos x}{\pi x^2}, \quad x \in \mathbb{R}, \quad \text{or equivalently } \mathcal{F}\ell(\xi) = (1 - |\xi|) I_{\{|\xi| \leq 1\}}, \quad \xi \in \mathbb{R}.$$

For any $\beta > 0$ define $\ell_\beta(x) = \beta \ell(\beta x)$. Then the Fourier transform of $\ell_\beta(x)$ is given by:

$$\mathcal{F}\ell_\beta(\xi) = \left(1 - \frac{|\xi|}{\beta} \right) I_{\{|\xi| \leq \beta\}}, \quad \xi \in \mathbb{R}. \tag{5.1}$$

We now give a representation of the noise \dot{W} as a convolution of a certain kernel with respect to a space-time white noise.

Lemma 5.2. *Let $\tilde{\gamma}$ be the distribution defined by*

$$\tilde{\gamma}(x) = \left(\frac{c_H}{2\pi} \right)^{1/2} \int_{\mathbb{R}} e^{i\xi x} |\xi|^{\frac{1}{2}-H} d\xi = \left(\frac{c_H}{2\pi} \right)^{1/2} \mathcal{F} \left(|\xi|^{\frac{1}{2}-H} \right).$$

Then the Gaussian field $\{W(t, \phi), t \geq 0, \phi \in \mathcal{S}(\mathbb{R})\}$ introduced in (2.3) can be represented as

$$W(t, \phi) = \int_0^t \int_{\mathbb{R}} (\phi * \tilde{\gamma})(x) \hat{W}(ds, dx), \tag{5.2}$$

where \hat{W} is a standard space-time white noise on \mathbb{R}^2 .

Proof. It is easily checked that $\tilde{\gamma} * \tilde{\gamma} = c_H \gamma$ is the spatial covariance of the fractional Brownian sheet W . In fact,

$$\tilde{\gamma} * \tilde{\gamma} = \frac{c_H}{2\pi} \mathcal{F}(|\xi|^{\frac{1}{2}-H}) * \mathcal{F}(|\xi|^{\frac{1}{2}-H}) = c_H \mathcal{F}\mu = c_H \gamma.$$

Let now W be the Gaussian field given by (5.2). Then for any $s, t \geq 0$ and $\phi, \psi \in \mathcal{S}(\mathbb{R})$, we can write:

$$\begin{aligned} \mathbb{E}[W(t, \phi)W(s, \psi)] &= (s \wedge t) \langle \phi * \tilde{\gamma}, \psi * \tilde{\gamma} \rangle_{L^2(\mathbb{R})} = (2\pi)^{-1} (s \wedge t) \langle \mathcal{F}(\phi * \tilde{\gamma}), \mathcal{F}(\psi * \tilde{\gamma}) \rangle_{L^2(\mathbb{R})} \\ &= (2\pi)^{-1} (s \wedge t) \langle \mathcal{F}\phi \mathcal{F}\tilde{\gamma}, \mathcal{F}\psi \mathcal{F}\tilde{\gamma} \rangle_{L^2(\mathbb{R})} \\ &= c_H (s \wedge t) \int_{\mathbb{R}} \mathcal{F}\phi(\xi) \overline{\mathcal{F}\psi(\xi)} \mu(d\xi), \end{aligned} \tag{5.3}$$

which corresponds to expression (2.1). □

We now turn to a description of the localized approximation of u which will be used in the sequel. For this step, we fix a parameter $\beta \geq 1$ and consider the approximation $\{W_\beta(t, \phi), \phi \in \mathcal{S}(\mathbb{R})\}$ of the fractional Brownian field W defined by

$$W_\beta(t, \phi) = \int_0^t \int_{\mathbb{R}} ([(\mathcal{F}\ell_\beta)\tilde{\gamma}] * \phi)(x) \hat{W}(ds, dx). \tag{5.4}$$

From (5.4) we obtain the following expression for the covariance function of the random field $\{W_\beta(t, \phi), t \geq 0, \phi \in \mathcal{S}(\mathbb{R})\}$:

$$\mathbb{E}[W_\beta(t, \phi)W_\beta(t, \psi)] = t c_H \int_{\mathbb{R}} (\ell_\beta * |\cdot|^{\frac{1}{2}-H})^2(\xi) \mathcal{F}\phi(\xi) \overline{\mathcal{F}\psi(\xi)} d\xi. \tag{5.5}$$

Notice that ℓ_β is an approximation of the identity as β tends to infinity, so that W_β has to be seen as an approximation of W . On the other hand, the spatial covariance of the noise W_β , given by $(\mathcal{F}\ell_\beta)\tilde{\gamma} * (\mathcal{F}\ell_\beta)\tilde{\gamma}$, has compact support. In this sense we call it *localized*.

Remark 5.3. We have followed the notation of [8] for our localization step. However let us stress the fact that, though the localization is made through the Fourier transform of ℓ_β , it is a localization in *direct* spatial coordinates.

Having the approximation (5.4) in hand, we can now define the following Picard approximation of the solution u to equation (1.1). Namely, we set $U_{\beta,0}(t, x) = 1$ and for $n \geq 1$, we define

$$U_{\beta,n+1}(t, x) = 1 + \int_0^t \int_{x-\beta\sqrt{t}}^{x+\beta\sqrt{t}} p_{t-s}(y-x) U_{\beta,n}(s, y) W_\beta(ds, dy). \tag{5.6}$$

The next result states the independence of the random variables $\{U_{\beta,n}(t, x_i), i \geq 1\}$ if the points x_i 's are far enough from each other. This property is crucial in order to establish our almost sure spatial behavior.

Lemma 5.4. *Choose any fixed $\beta \geq 1$ and let $n = \lceil \log \beta \rceil + 1$. Then for any sequence of points $\{x_1 < x_2 < \dots\}$ such that $x_{i+1} - x_i > 2n\beta(1 + \sqrt{t})$ for each $i \geq 1$, the random variables $\{U_{\beta,n}(t, x_i), i \geq 1\}$, defined by (5.6), are independent.*

Proof. The proof is exactly the same as the proof of Lemma 5.4 of [9], and is omitted for sake of conciseness. □

Next let us recall the following elementary result borrowed from [14], which will help us to bound moments of iterated integrals.

Lemma 5.5. For $m \geq 1$ let $\alpha \in (-1+\varepsilon, 1)^m$ with $\varepsilon > 0$ and set $|\alpha| = \sum_{i=1}^m \alpha_i$. For $t \in [0, T]$, the m -th dimensional simplex over $[0, t]$ is denoted by $S_m(t) = \{(r_1, r_2, \dots, r_m) \in \mathbb{R}^m : 0 < r_1 < \dots < r_m < t\}$. Then there is a constant $c > 0$ such that

$$J_m(t, \alpha) := \int_{S_m(t)} \prod_{i=1}^m (r_i - r_{i-1})^{\alpha_i} dr \leq \frac{c^m t^{|\alpha|+m}}{\Gamma(|\alpha| + m + 1)},$$

where by convention, $r_0 = 0$.

The random field $u(t, x)$ admits a chaos expansion (see, for instance, formula (5.9) in [14]). Namely, we have

$$u(t, x) = \sum_{n=0}^{\infty} I_n(f_n(\cdot, t, x)), \tag{5.7}$$

where $f_0(t, x) = 1$ and for any $n \geq 1$,

$$f_n(s_1, x_1, \dots, s_n, x_n, t, x) = \frac{1}{n!} p_{t-s_{\sigma(n)}}(x - x_{\sigma(n)}) \cdots p_{s_{\sigma(2)}-s_{\sigma(1)}}(x_{\sigma(2)} - x_{\sigma(1)}). \tag{5.8}$$

Here σ denotes the permutation of $\{1, 2, \dots, n\}$ such that $0 < s_{\sigma(1)} < \dots < s_{\sigma(n)} < t$ and I_n is the multiple Itô-Wiener integral with respect to the fractional Brownian field W .

The same kind of formula holds for $U_{\beta,n}(t, x)$ defined by (5.6). Namely, denote by $p^{(\beta)}$ the kernel defined by

$$p^{(\beta)}(s, t; y) = p_{t-s}(y) \mathbf{1}_{\{|y| \leq \beta \sqrt{t}\}}.$$

Then for $n \geq 1$, one can recast formula (5.6) as

$$U_{\beta,n+1}(t, x) = 1 + \int_0^t \int_{\mathbb{R}} p^{(\beta)}(s, t; y - x) U_{\beta,n}(s, y) W_{\beta}(ds, dy).$$

By iteration, similarly to [14] and (5.7), we have

$$U_{\beta,n}(t, x) = \sum_{k=0}^n I_{\beta,k}(f_{\beta,k}(\cdot, t, x)), \tag{5.9}$$

where $f_{\beta,0}(t, x) = 1$ and for $k \geq 1$,

$$f_{\beta,k}(s_1, x_1, \dots, s_k, x_k, t, x) = \frac{1}{k!} p^{(\beta)}(s_{\sigma(k)}, t; x - x_{\sigma(k)}) \cdots p^{(\beta)}(s_{\sigma(1)}, s_{\sigma(2)}; x_{\sigma(2)} - x_{\sigma(1)}). \tag{5.10}$$

In the above expression, $I_{\beta,k}$ is the multiple Itô-Wiener integral of order k with respect to the Gaussian process $W_{\beta}(t, x)$. In the next proposition we are going to show that the sequence $U_{\beta,n}(t, x)$ converges in L^2 , and defines a random field $U_{\beta}(t, x)$ given by

$$U_{\beta}(t, x) = \sum_{n=0}^{\infty} I_{\beta,n}(f_{\beta,n}(\cdot, t, x)). \tag{5.11}$$

On the other hand, we will also see that $U_{\beta}(t, x)$ converges in L^2 to $u(t, x)$ as β tends to infinity.

Proposition 5.6. For any $(t, x) \in \mathbb{R}_+ \times \mathbb{R}$ and for any $p \geq 1$, the sequence $U_{\beta,n}(t, x)$ defined by (5.6) converges in L^p to the random variable $U_{\beta}(t, x)$ defined by (5.11), as n tends to infinity. Furthermore, we have the following estimates for the differences of the solutions: there is a finite constant C (dependent on t but independent of β and p) such that for any $\beta \geq 1$ and $p \geq 2$,

$$\|u(t, x) - U_{\beta}(t, x)\|_{L^p(\Omega)} \leq \beta^{-(\frac{1}{2}-H)} \exp\left(C p^{\frac{1}{H}}\right) \tag{5.12}$$

and for two constants $c_1, c_2 > 0$:

$$\|U_{\beta,n}(t, x) - U_{\beta}(t, x)\|_{L^p(\Omega)} \leq \frac{(c_1 p^{1/2})^n}{\Gamma\left(\frac{nH}{2} + 1\right)} \exp\left(c_2 p^{\frac{1}{H}}\right). \tag{5.13}$$

Proof. In order to simplify the notation we will omit the dependence on (t, x) in some of the terms of our computations. The proof will be done in several steps.

Step 1: Estimates on a fixed chaos. In this step we consider a fixed chaos n . Having expressions (5.7) and (5.11) in mind, we shall estimate the expectation $\mathbb{E}[|I_n(f_n) - I_{\beta,n}(f_{\beta,n})|^2]$. Let us start with the following decomposition:

$$\begin{aligned} \mathbb{E}[|I_n(f_n) - I_{\beta,n}(f_{\beta,n})|^2] &\leq 2\mathbb{E}[|I_n(f_n) - I_n(f_{\beta,n})|^2] + 2\mathbb{E}[|I_n(f_{\beta,n}) - I_{\beta,n}(f_{\beta,n})|^2] \\ &=: 2(A_1 + A_2), \end{aligned} \tag{5.14}$$

and let us estimate A_1 and A_2 separately.

Consider first the term A_1 . Some elementary computation reveals that the Fourier transform of f_n is given by

$$\mathcal{F}f_n(s_1, \xi_1, \dots, s_n, \xi_n) = \frac{1}{n!} e^{-ix(\xi_{\sigma(n)} + \dots + \xi_{\sigma(1)})} \prod_{i=1}^n e^{-\frac{1}{2}(s_{\sigma(i+1)} - s_{\sigma(i)})|\xi_{\sigma(i)} + \dots + \xi_{\sigma(1)}|^2},$$

where we have used the convention $s_{\sigma(n+1)} = t$. The same kind of computation can be performed for $f_{\beta,n}$. Specifically, let $\rho^{(\beta)}(s, t; \xi)$ be the Fourier transform of $p^{(\beta)}(s, t; x)$, namely,

$$\rho^{(\beta)}(s, t; \xi) = \frac{1}{\sqrt{2\pi(t-s)}} \int_{-\beta\sqrt{t}}^{\beta\sqrt{t}} e^{-ix\xi - \frac{x^2}{2(t-s)}} dx. \tag{5.15}$$

Then the Fourier transform of $f_{\beta,n}$ is given by

$$\mathcal{F}f_{\beta,n}(s_1, \xi_1, \dots, s_n, \xi_n) = \frac{1}{n!} \prod_{i=1}^n \rho^{(\beta)}(s_{\sigma(i)}, s_{\sigma(i+1)}; \xi_{\sigma(i)} + \dots + \xi_{\sigma(1)}) e^{-ix(\xi_{\sigma(n)} + \dots + \xi_{\sigma(1)})}. \tag{5.16}$$

In the sequel, we also make use of the notation

$$\tilde{\rho}_j^{(\beta)} = \rho^{(\beta)}(s_{\sigma(j)}, s_{\sigma(j+1)}; \xi_{\sigma(j)} + \dots + \xi_{\sigma(1)}) - e^{-\frac{1}{2}(s_{\sigma(j+1)} - s_{\sigma(j)})|\xi_{\sigma(j)} + \dots + \xi_{\sigma(1)}|^2}. \tag{5.17}$$

Now a straightforward application of Parseval's identity yields

$$\begin{aligned} A_1 &= n! \|\mathcal{F}f_n - \mathcal{F}f_{\beta,n}\|_{L^2([0,t]^n \times \mathbb{R}^n, \lambda^n \times \mu^n)}^2 \\ &= \frac{1}{n!} \int_{[0,t]^n} \int_{\mathbb{R}^n} \left| e^{-ix(\xi_{\sigma(n)} + \dots + \xi_{\sigma(1)})} \prod_{i=1}^n \rho^{(\beta)}(s_{\sigma(i)}, s_{\sigma(i+1)}; \xi_{\sigma(i)} + \dots + \xi_{\sigma(1)}) \right. \\ &\quad \left. - e^{-ix(\xi_{\sigma(n)} + \dots + \xi_{\sigma(1)})} \prod_{i=1}^n e^{-\frac{1}{2}(s_{\sigma(i+1)} - s_{\sigma(i)})|\xi_{\sigma(i)} + \dots + \xi_{\sigma(1)}|^2} \right|^2 \prod_{i=1}^n |\xi_i|^{1-2H} d\xi ds, \end{aligned}$$

where we have set $d\xi = d\xi_1 \dots d\xi_n$, $ds = ds_1 \dots ds_n$ and where λ denotes the Lebesgue measure. Thus, using a telescoping sum argument, we can write

$$A_1 \leq \frac{n}{n!} \sum_{j=1}^n A_{j,n}, \tag{5.18}$$

where we recall that we have defined $\tilde{\rho}_j^{(\beta)}$ in (5.17), and where we set:

$$A_{j,n} = \int_{[0,t]^n} \int_{\mathbb{R}^n} \left| \tilde{\rho}_j^{(\beta)} \right|^2 \prod_{i=j+1}^n e^{-\frac{1}{2}(s_{\sigma(i+1)} - s_{\sigma(i)})|\xi_{\sigma(i)} + \dots + \xi_{\sigma(1)}|^2} \\ \times \prod_{i=1}^{j-1} \left| \rho^{(\beta)}(s_{\sigma(i)}, s_{\sigma(i+1)}; \xi_{\sigma(i)} + \dots + \xi_{\sigma(1)}) \right|^2 \prod_{i=1}^n |\xi_i|^{1-2H} d\xi ds. \quad (5.19)$$

For the time being, denote $t = s_{\sigma(j+1)}$, $s = s_{\sigma(j)}$, $\eta = \xi_{\sigma(j)} + \dots + \xi_{\sigma(1)}$, let C be a generic constant (possibly depending on H and t), and let us estimate $\tilde{\rho}_j^{(\beta)}$. It is easy to see that

$$\rho^{(\beta)}(s, t; \eta) - e^{-\frac{1}{2}(t-s)\eta^2} = \frac{1}{\sqrt{2\pi(t-s)}} \int_{|x| \geq \beta\sqrt{t}} e^{ix\eta - \frac{x^2}{2(t-s)}} dx \\ = \frac{1}{\sqrt{2\pi(t-s)}} \int_{|x| \geq \beta\sqrt{t}} \cos(\eta x) e^{-\frac{x^2}{2(t-s)}} dx \\ = \frac{-2}{\eta\sqrt{2\pi(t-s)}} e^{-\frac{\beta^2 t}{2(t-s)}} \sin(\eta\beta\sqrt{t}) - \frac{1}{\eta\sqrt{2\pi(t-s)}} \int_{|x| \geq \beta\sqrt{t}} \sin(\eta x) de^{-\frac{x^2}{2(t-s)}}.$$

Therefore, trivially bounding $|\sin(\eta x)|$ by 1 in the integral above, we get

$$|\rho^{(\beta)}(s, t; \eta) - e^{-\frac{1}{2}(t-s)\eta^2}| \leq \frac{C}{|\eta|\sqrt{t-s}} e^{-\frac{\beta^2}{2}}.$$

On the other hand, it is readily checked from (5.17) that $|\rho^{(\beta)}(s, t; \eta)|$ is bounded by a constant C . Thus, for any $\theta \in [0, 1]$ (possibly depending on η), we have

$$|\rho^{(\beta)}(s, t; \eta) - e^{-\frac{1}{2}(t-s)\eta^2}| \leq \frac{C}{|\eta|^\theta (t-s)^{\theta/2}} e^{-\frac{\beta^2 \theta}{2}}. \quad (5.20)$$

Substituting this bound into (5.19) yields

$$A_{j,n} \leq C^2 e^{-\beta^2 \theta} \int_{[0,t]^n} \int_{\mathbb{R}^n} \frac{(s_{\sigma(j+1)} - s_{\sigma(j)})^{-\theta}}{|\xi_{\sigma(j)} + \dots + \xi_{\sigma(1)}|^{2\theta}} \prod_{i=j+1}^n e^{-(s_{\sigma(i+1)} - s_{\sigma(i)})|\xi_{\sigma(i)} + \dots + \xi_{\sigma(1)}|^2} \\ \times \prod_{i=1}^{j-1} \left| \rho^{(\beta)}(s_{\sigma(i)}, s_{\sigma(i+1)}; \xi_{\sigma(i)} + \dots + \xi_{\sigma(1)}) \right|^2 \prod_{i=1}^n |\xi_i|^{1-2H} d\xi ds. \quad (5.21)$$

Making the change of variable $\xi_{\sigma(i)} + \dots + \xi_{\sigma(1)} = \eta_i$, for all $i = 1, 2, \dots, n$, we obtain

$$A_{j,n} \leq C^2 e^{-\beta^2 \theta} \int_{[0,t]^n} \int_{\mathbb{R}^n} (s_{\sigma(j+1)} - s_{\sigma(j)})^{-\theta} |\eta_j|^{-2\theta} \prod_{i=j+1}^n e^{-(s_{\sigma(i+1)} - s_{\sigma(i)})\eta_i^2} \\ \times \prod_{i=1}^{j-1} |\rho^{(\beta)}(s_{\sigma(i)}, s_{\sigma(i+1)}; \eta_i)|^2 \prod_{i=1}^n |\eta_i - \eta_{i-1}|^{1-2H} d\eta ds,$$

where we have set $\eta_0 = 0$. We can now invoke the elementary bound $|\eta_i - \eta_{i-1}|^{1-2H} \leq |\eta_i|^{1-2H} + |\eta_{i-1}|^{1-2H}$, and we obtain

$$A_{j,n} \leq C^2 e^{-\beta^2 \theta} \sum_{\alpha \in D_n} \int_{[0,t]^n} \int_{\mathbb{R}^n} (s_{\sigma(j+1)} - s_{\sigma(j)})^{-\theta} |\eta_j|^{-2\theta} \prod_{i=j+1}^n e^{-(s_{\sigma(i+1)} - s_{\sigma(i)})\eta_i^2} \\ \times \prod_{i=1}^{j-1} |\rho^{(\beta)}(s_{\sigma(i)}, s_{\sigma(i+1)}; \eta_i)|^2 \prod_{i=1}^n |\eta_i|^{\alpha_i} d\eta ds, \quad (5.22)$$

where D_n is a subset of multi-indices of length n satisfying the following rules: $\text{Card}(D_n) = 2^n$ and for any $\alpha \in D_n$ we have

$$|\alpha| \equiv \sum_{i=1}^n \alpha_i = n(1 - 2H), \quad \text{and} \quad \alpha_i \in \{0, 1 - 2H, 2(1 - 2H)\}, \quad i = 1, \dots, n. \quad (5.23)$$

Now we perform the integration on each variable η_i in relation (5.22). If $i \geq j + 1$, it is readily checked that

$$\int_{\mathbb{R}} |\eta_i|^{\alpha_i} e^{-(s_{\sigma(i+1)} - s_{\sigma(i)})\eta_i^2} d\eta_i = C(s_{\sigma(i+1)} - s_{\sigma(i)})^{-\frac{\alpha_i+1}{2}}. \quad (5.24)$$

In the case $i \leq j - 1$, we also claim that

$$\int_{\mathbb{R}} |\rho^{(\beta)}(s_{\sigma(i)}, s_{\sigma(i+1)}; \eta_i)|^2 |\eta_i|^{\alpha_i} d\eta_i \leq C(s_{\sigma(i+1)} - s_{\sigma(i)})^{-\frac{\alpha_i+1}{2}}. \quad (5.25)$$

In fact, recalling our definition (5.15) and thanks to an elementary change of variable, this integral can be written as

$$\frac{1}{2\pi} (s_{\sigma(i+1)} - s_{\sigma(i)})^{-\frac{\alpha_i+1}{2}} \int_{\mathbb{R}} |\xi|^{\alpha_i} \left| \int_{-\beta\sqrt{s_{\sigma(i+1)}}/\sqrt{s_{\sigma(i+1)} - s_{\sigma(i)}}}^{\beta\sqrt{s_{\sigma(i+1)}}/\sqrt{s_{\sigma(i+1)} - s_{\sigma(i)}}} e^{ix\xi - \frac{x^2}{2}} dx \right|^2 d\xi,$$

which is bounded by a constant times $(s_{\sigma(i+1)} - s_{\sigma(i)})^{-\frac{\alpha_i+1}{2}}$ by Lemma 6.1 in the Appendix below. It remains to consider the integral over the variable η_j in (5.22), which is given by

$$(s_{\sigma(j+1)} - s_{\sigma(j)})^{-\theta} \int_{\mathbb{R}} |\eta_j|^{-2\theta + \alpha_j} d\eta_j.$$

We decompose the integral $\int_{\mathbb{R}} |\eta_j|^{-2\theta + \alpha_j} d\eta_j$ into two parts: on the region $|\eta_j| \leq 1$ we take $\theta = \frac{\alpha_j+1}{2} - \frac{\delta}{2}$ and on the region $|\eta_j| > 1$ we take $\theta = \frac{\alpha_j+1}{2} + \frac{\delta}{2}$, for some $\delta > 0$ to be fixed later. In this way, we obtain

$$\begin{aligned} & (s_{\sigma(j+1)} - s_{\sigma(j)})^{-\theta} \int_{\mathbb{R}} |\eta_j|^{-2\theta + \alpha_j} d\eta_j \\ & \leq C \left((s_{\sigma(j+1)} - s_{\sigma(j)})^{-\frac{1}{2}(\alpha_j+1-\delta)} + (s_{\sigma(j+1)} - s_{\sigma(j)})^{-\frac{1}{2}(\alpha_j+1+\delta)} \right) \\ & \leq C (s_{\sigma(j+1)} - s_{\sigma(j)})^{-\frac{1}{2}(\alpha_j+1+\delta)}. \end{aligned} \quad (5.26)$$

Therefore, plugging (5.24), (5.25) and (5.26) into (5.22), we end up with the following relation for any $1 \leq j \leq n$

$$\begin{aligned} A_{j,n} & \leq C e^{-\beta^2\theta} \sum_{\alpha} \int_{[0,t]^n} \prod_{i \neq j} (s_{\sigma(i+1)} - s_{\sigma(i)})^{-\frac{\alpha_i+1}{2}} (s_{\sigma(j+1)} - s_{\sigma(j)})^{-\frac{\alpha_j+1}{2} - \frac{\delta}{2}} ds \\ & = C e^{-\beta^2\theta} n! \sum_{\alpha} \int_{S_n(t)} \prod_{i \neq j} (s_{i+1} - s_i)^{-\frac{\alpha_i+1}{2}} (s_{j+1} - s_j)^{-\frac{\alpha_j+1}{2} - \frac{\delta}{2}} ds, \end{aligned} \quad (5.27)$$

where $\theta = \frac{\alpha_j+1}{2} - \frac{\delta}{2}$. We now wish to apply Lemma 5.5 in order to bound the right hand side of (5.27), and we first discuss the nature of the exponents involved: (i) First we have to ensure that each exponent in the integral showing up in (5.27) is lower bounded by -1 . These exponents are $\geq -\frac{\max_{i \leq n} \alpha_i + 1}{2} - \frac{\delta}{2}$, and recall that $\max_{i \leq n} \alpha_i \leq 2(1 - 2H)$ according to (5.23). We can thus ensure that each exponent is greater than -1 , provided that $0 < \delta < 4H - 1$. (ii) Invoking relation (5.23) again, it is readily checked that the sum of the exponents in (5.27) is $-n + nH - \frac{\delta}{2}$. (iii) We wish to choose θ as large as

possible in order to ensure the maximal exponential decay for $A_{j,n}$. According to our previous considerations, we have taken $\theta = \frac{\alpha_j+1}{2} - \frac{\delta}{2}$ with δ of the form $4H - 1 - \varepsilon$ for an arbitrarily small ε . Referring once more to (5.27), we can just ensure $\alpha_j \geq 0$, which yields $\theta \geq \frac{1-\delta}{2} = 1 - 2H + \frac{\varepsilon}{2} \geq 1 - 2H$. With those considerations in mind, we can now apply Lemma 5.5 to relation (5.27) in order to conclude that

$$A_{j,n} \leq \frac{n!C^n}{\Gamma(nH + \frac{1}{2})} e^{-\beta^2(1-2H)}, \tag{5.28}$$

where C is a constant depending on H and t . Substituting (5.28) into (5.18) yields (recall that the constant C might change from line to line)

$$A_1 \leq \frac{C^n}{\Gamma(nH + \frac{1}{2})} e^{-\beta^2(1-2H)}.$$

Using the inequality $\Gamma(nH + \frac{1}{2}) \geq \frac{\Gamma(nH+1)}{nH+\frac{1}{2}}$, we obtain

$$A_1 \leq \frac{C^n}{\Gamma(nH + 1)} e^{-\beta^2(1-2H)}. \tag{5.29}$$

Going back to our decomposition (5.14), let us now deal with the term A_2 . The spectral measure of the noise W_β has a density equal to $c_H(\ell_\beta * |\cdot|^{\frac{1}{2}-H})^2(\xi)$. Therefore, thanks to another telescoping sum argument, we get

$$\begin{aligned} A_2 &= c_H^n n! \int_{[0,t]^n} \int_{\mathbb{R}^n} |\mathcal{F}f_{\beta,n}(s_1, \xi_1, \dots, s_n, \xi_n)|^2 \left| \prod_{i=1}^n |\xi_i|^{\frac{1}{2}-H} - \prod_{i=1}^n (\ell_\beta * |\cdot|^{\frac{1}{2}-H})(\xi_i) \right|^2 d\xi ds \\ &\leq c_H^n n n! \sum_{j=1}^n \int_{[0,t]^n} \int_{\mathbb{R}^n} |\mathcal{F}f_{\beta,n}(s_1, \xi_1, \dots, s_n, \xi_n)|^2 \prod_{i=1}^{j-1} |\xi_i|^{1-2H} \\ &\quad \times \left| |\xi_j|^{\frac{1}{2}-H} - (\ell_\beta * |\cdot|^{\frac{1}{2}-H})(\xi_j) \right|^2 \prod_{i=j+1}^n (\ell_\beta * |\cdot|^{\frac{1}{2}-H})^2(\xi_i) d\xi ds. \end{aligned}$$

In addition, notice that

$$\left| |\xi_j|^{\frac{1}{2}-H} - (\ell_\beta * |\cdot|^{\frac{1}{2}-H})(\xi_j) \right| \leq c_{1,H} \beta^{-(\frac{1}{2}-H)},$$

where $c_{1,H} = \int_{\mathbb{R}} |\eta|^{\frac{1}{2}-H} \ell(\eta) d\eta$. This follows easily from

$$\left| |\xi_j|^{\frac{1}{2}-H} - (\ell_\beta * |\cdot|^{\frac{1}{2}-H})(\xi_j) \right| \leq \beta^{-(\frac{1}{2}-H)} \int_{\mathbb{R}} \ell(\eta) \left| |\beta\xi_j|^{\frac{1}{2}-H} - |\beta\xi_j - \eta|^{\frac{1}{2}-H} \right| d\eta,$$

and the inequality $\left| |\beta\xi_j|^{\frac{1}{2}-H} - |\beta\xi_j - \eta|^{\frac{1}{2}-H} \right| \leq |\eta|^{\frac{1}{2}-H}$. In the same way we can show that

$$(\ell_\beta * |\cdot|^{\frac{1}{2}-H})(\xi_i) \leq |\xi_i|^{\frac{1}{2}-H} + c_{1,H} \beta^{-(\frac{1}{2}-H)}.$$

Taking into account that $\beta \geq 1$, this leads to the estimate

$$A_2 \leq C^n n! \beta^{-(1-2H)} \int_{[0,t]^n} \int_{\mathbb{R}^n} |\mathcal{F}f_{\beta,n}(s_1, \xi_1, \dots, s_n, \xi_n)|^2 \prod_{i=1}^n (|\xi_i|^{1-2H} \vee 1) d\xi ds.$$

We now start from the expression (5.16) for $\mathcal{F}f_{\beta,n}(s_1, \xi_1, \dots, s_n, \xi_n)$, we make the change of variable $\xi_{\sigma(i)} + \dots + \xi_{\sigma(1)} = \eta_i$, for all $i = 1, 2, \dots, n$, and we bound $|\eta_i - \eta_{i-1}|^{1-2H}$ by

$|\eta_i|^{1-2H} + |\eta_{i-1}|^{1-2H}$ as in the case of our term A_1 . This yields

$$\begin{aligned} A_2 &\leq C^n \beta^{-(1-2H)} \int_{S_n(t)} \int_{\mathbb{R}^n} |\rho^{(\beta)}(s_i, s_{i+1}; \eta_i)|^2 \prod_{i=1}^n (|\eta_i - \eta_{i-1}|^{1-2H} \vee 1) d\eta ds \\ &\leq C^n \beta^{-(1-2H)} \sum_{\alpha \in D_n} \int_{S_n(t)} \int_{\mathbb{R}^n} |\rho^{(\beta)}(s_i, s_{i+1}; \eta_i)|^2 \prod_{i=1}^n (|\eta_i|^{\alpha_i} \vee 1) d\eta ds \\ &\leq C^n \beta^{-(1-2H)} \sum_{\alpha \in D_n} \int_{S_n(t)} (s_{i+1} - s_i)^{-\frac{\alpha_i+1}{2}} ds \leq \frac{C^n}{\Gamma(nH + 1)} \beta^{-(1-2H)}, \end{aligned} \tag{5.30}$$

where we recall that $S_n(t)$ denotes the n -dimensional simplex of $[0, t]^n$. We now conclude, putting together (5.29) and (5.30), that

$$\mathbb{E}[|I_n(f_n) - I_{\beta,n}(f_{\beta,n})|^2] \leq \frac{C^n}{\Gamma(nH + 1)} \beta^{-(1-2H)}. \tag{5.31}$$

In a similar way, we can also obtain the following estimate (whose proof is left to the patient reader), where the constant C is independent of β

$$\mathbb{E}[|I_{\beta,n}(f_{\beta,n})|^2] \leq \frac{C^n}{\Gamma(nH + 1)}. \tag{5.32}$$

Step 2: L^p -estimates. Recall that $U_{\beta,n}(t, x)$ is defined by the finite sum (5.9). Let us first get the convergence of this finite sum to a random variable $U_{\beta}(t, x)$ formally defined by the series (5.11). To this aim, recall that for a functional F_n which belongs to the n -th chaos of a Wiener space and $p \geq 2$, we have the hypercontractivity inequality $\|F_n\|_{L^p(\Omega)} \leq p^{\frac{n}{2}} \|F_n\|_{L^2(\Omega)}$. We thus get

$$\begin{aligned} \|U_{\beta,n}(t, x) - U_{\beta}(t, x)\|_{L^p(\Omega)} &\leq \sum_{k=n+1}^{\infty} \|I_{\beta,k}(f_{\beta,k})\|_{L^p(\Omega)} \leq \sum_{k=n+1}^{\infty} p^{\frac{k}{2}} \|I_{\beta,k}(f_{\beta,k})\|_{L^2(\Omega)} \\ &\leq \sum_{k=n+1}^{\infty} \frac{(Cp^{1/2})^k}{\Gamma(\frac{kH}{2} + 1)}, \end{aligned}$$

where the last inequality is due to (5.32). Furthermore, the following inequality, valid for $z \geq 0$ and $a > 0$, is an easy consequence of estimates on Mittag-Leffler functions which can be found in [12]

$$\sum_{k=n+1}^{\infty} \frac{z^k}{\Gamma(ak + 1)} \leq \frac{c_1 z^n}{\Gamma(an + 1)} e^{c_2 z^{\frac{1}{a}}},$$

where c_1, c_2 are two universal constants. Plugging this bound into our previous estimate, we end up with

$$\|U_{\beta,n}(t, x) - U_{\beta}(t, x)\|_{L^p(\Omega)} \leq \frac{(c_3 p^{1/2})^n}{\Gamma(\frac{nH}{2} + 1)} \exp\left(c_4 p^{\frac{1}{H}}\right),$$

which shows our claim (5.13).

The same kind of consideration also allows to derive inequality (5.12). Namely, write

$$\|u(t, x) - U_{\beta}(t, x)\|_{L^p(\Omega)} \leq \sum_{k=0}^{\infty} \|I_k(f_k) - I_{\beta,k}(f_{\beta,k})\|_{L^p(\Omega)} \leq \beta^{-(1-2H)} \sum_{k=0}^{\infty} \frac{(Cp^{1/2})^k}{\Gamma(\frac{kH}{2} + 1)},$$

where we resort to hypercontractivity and (5.31) for the last step. This easily yields (5.12) by the same kind of argument as before. \square

Corollary 5.7. Consider $p \geq 1$. Under the same assumptions as in Proposition 5.6, suppose that $\beta = \exp(Mp^{\frac{1}{H}})$ and $n = \lceil \log \beta \rceil + 1$ for some constant $M > 0$. Then, for any $\nu > 0$, there exists $M > 0$, such that

$$\|u(t, x) - U_{\beta,n}(t, x)\|_{L^p(\Omega)} \leq \exp \left\{ -\nu p^{\frac{1}{H}} \right\}. \tag{5.33}$$

Proof. From (5.12) and (5.13) we obtain

$$\|u(t, x) - U_{\beta,n}(t, x)\|_{L^p(\Omega)} \leq e^{Cp^{\frac{1}{H}}} \left(\beta^{-(\frac{1}{2}-H)} + \frac{C^n p^{n/2}}{\Gamma(\frac{nH}{2} + \frac{1}{2})} \right),$$

for some constant C depending on H and t . Using the asymptotic properties of the Gamma function, this is bounded by

$$e^{Cp^{\frac{1}{H}}} \left(\beta^{-(\frac{1}{2}-H)} + C^n p^{\frac{n}{2}} n^{-\frac{nH}{2}} \right). \tag{5.34}$$

To bound the above right-hand side, we have first (recall that $\beta = \exp(Mp^{\frac{1}{H}})$)

$$e^{Cp^{\frac{1}{H}}} \beta^{-(\frac{1}{2}-H)} = \exp \left\{ p^{\frac{1}{H}} \left[C - \left(\frac{1}{2} - H \right) M \right] \right\},$$

which is less than $\frac{1}{2} \exp(-\nu p^{\frac{1}{H}})$ for M large enough. For the second summand in the right-hand side of (5.34), we obtain (provided $n = \lceil \log \beta \rceil + 1$) the upper bound

$$\exp \left\{ p^{\frac{1}{H}} \left(C + M \log C - \frac{M}{2} \log M \right) + \frac{1}{2} \log p \right\},$$

which again is less than $\frac{1}{2} \exp(-\nu p^{\frac{1}{H}})$ for M large enough. This completes the proof of the corollary. \square

We are now ready to give the proof of our lower bound.

Proposition 5.8. Under the assumptions of Theorem 1.1, for all $t > 0$ we have

$$\liminf_{R \rightarrow \infty} (\log R)^{-\frac{1}{1+H}} \log \left(\max_{|x| \leq R} u(t, x) \right) \geq [\hat{c}_{H,t}]^{-\frac{1}{1+H}} \quad \text{a.s.}, \tag{5.35}$$

where $\hat{c}_{H,t}$ is defined in Corollary 4.5 and is related to (1.3) by Remark 4.6.

Proof. We divide this proof in two steps: first we determine a main contribution to the maximum, given by our approximations $U_{\beta,n}$ suitably discretized. Then we will evaluate the main contribution.

Step 1: Fluctuation results. Fix $R > 0$ and consider a given $\nu > 0$. Referring to the notation of Corollary 5.7, we wish to choose p in inequality (5.33) such that we obtain

$$\mathbb{E} [|u(t, x) - U_{\beta,n}(t, x)|^p] \leq R^{-\nu}. \tag{5.36}$$

It is readily checked that this is achieved for $p = p(R) = (\log R)^{\frac{H}{1+H}}$, which is greater than 1 if $R \geq e$. We thus choose this p , the corresponding β and n in Corollary 5.7 being then given by

$$\beta = \beta(R) = \exp \left\{ M (\log R)^{\frac{1}{1+H}} \right\} \quad \text{and} \quad n = \left\lceil M (\log(R))^{\frac{1}{1+H}} \right\rceil + 1. \tag{5.37}$$

For a fixed $t > 0$, we now wish to produce some independent random variables $U_{\beta,n}(t, x_j)$, with $x_j \in [-R, R]$. For this we need $x_{j+1} - x_j > 2n\beta(1 + t^{1/2})$ for all j . Set $N = 2n\beta(1 + t^{1/2})$. We choose the set of points

$$\mathcal{N}_R = \left\{ 2kN : k \in \mathbb{Z}, -\left\lfloor \frac{R}{2N} \right\rfloor \leq k \leq \left\lfloor \frac{R}{2N} \right\rfloor \right\}.$$

If $|\mathcal{N}_R| = (2 \lfloor \frac{R}{2N} \rfloor + 1)$ denotes the cardinality of \mathcal{N}_R , one can check, using the expressions of β and n given in (5.37), that for any $\varepsilon > 0$ the following inequalities hold

$$c_{H,M,\varepsilon} R^{1-\varepsilon} \leq |\mathcal{N}_R| \leq R, \tag{5.38}$$

for R large enough, where $c_{H,M,\varepsilon}$ is a positive constant.

We can now study $\max_{z \in \mathcal{N}_R} |u(t, z) - U_{\beta,n}(t, z)|$. For any $\eta > 0$ we have

$$\begin{aligned} \mathbb{P} \left(\log \max_{z \in \mathcal{N}_R} |u(t, z) - U_{\beta,n}(t, z)| \geq \eta(\log R)^{\frac{1}{1+H}} \right) \\ \leq |\mathcal{N}_R| \mathbb{P} \left(\log |u(t, 0) - U_{\beta,n}(t, 0)| \geq \eta(\log R)^{\frac{1}{1+H}} \right) \end{aligned} \tag{5.39}$$

Furthermore, a simple application of Markov's inequality yields, for an arbitrary $p \geq 1$

$$\mathbb{P} \left(\log |u(t, 0) - U_{\beta,n}(t, 0)| \geq \eta(\log R)^{\frac{1}{1+H}} \right) \leq \frac{\mathbb{E} [|u(t, x) - U_{\beta,n}(t, x)|^p]}{\exp \left(\eta p (\log R)^{\frac{1}{1+H}} \right)},$$

so that choosing $p = p(R)$ and invoking relation (5.36), we can recast this relation as

$$\mathbb{P} \left(\log |u(t, 0) - U_{\beta,n}(t, 0)| \geq \eta(\log R)^{\frac{1}{1+H}} \right) \leq R^{-(\nu+\eta)}. \tag{5.40}$$

Going back to inequality (5.39) and choosing $\nu = 3$, we have obtained the following inequality for R large enough

$$\mathbb{P} \left(\log \max_{z \in \mathcal{N}_R} |u(t, z) - U_{\beta,n}(t, z)| \geq \eta(\log R)^{\frac{1}{1+H}} \right) \leq R^{-2}. \tag{5.41}$$

Notice that this decay in R is sufficient to apply Borel-Cantelli's lemma. Considering for instance a sequence $R = m$, we get

$$\lim_{m \rightarrow \infty} (\log m)^{-\frac{1}{1+H}} \log \max_{z \in \mathcal{N}_m} |u(t, z) - U_{\beta(m),n(m)}(t, z)| = 0 \quad \text{a.s.} \tag{5.42}$$

which is enough to assert that

$$\liminf_{R \rightarrow \infty} (\log R)^{-\frac{1}{1+H}} \log \left(\max_{|x| \leq R} u(t, x) \right) \geq \liminf_{R \rightarrow \infty} (\log R)^{-\frac{1}{1+H}} \log \left(\max_{x \in \mathcal{N}_R} U_{\beta,n}(t, x) \right), \tag{5.43}$$

where we recall that $\beta = \beta(R)$ and $n = n(R)$ in the right-hand side of (5.43) are given by (5.37). We will now evaluate the right-hand side of (5.43), identified with our main contribution.

Step 2: Evaluation of the main term. Fix $\lambda > 0$ and $\delta > 0$ arbitrarily small, satisfying the following condition

$$\hat{c}_{H,t} (\lambda + \delta)^{1+H} < 1 - \delta, \tag{5.44}$$

where $\hat{c}_{H,t}$ is the constant introduced in Corollary 4.5 (observe that λ is arbitrarily close to $\hat{c}_{H,t}^{-1/(1+H)}$). Using the independence property established in Lemma 5.4, we can write

$$\begin{aligned} \mathbb{P} \left(\log \max_{z \in \mathcal{N}_R} |U_{\beta,n}(t, z)| \leq \lambda(\log R)^{\frac{1}{1+H}} \right) \\ = \left(1 - \mathbb{P} \left(\log |U_{\beta,n}(t, 0)| > \lambda(\log R)^{\frac{1}{1+H}} \right) \right)^{|\mathcal{N}_R|}. \end{aligned} \tag{5.45}$$

In the following lines, we will write $u = u(t, 0)$ and $U_{\beta,n} = U_{\beta,n}(t, 0)$ to alleviate notations. Then recall that for $a > 0$ and $b \in \mathbb{R}$, the following elementary relation holds true

$$\log a \leq \log 2 + \max \{ \log(|a - b|), \log(|b|) \}. \tag{5.46}$$

Applying this inequality to $a = u$ and $b = |U_{\beta,n}|$, we obtain

$$\log u \leq \log 2 + \max \{ \log (|u - U_{\beta,n}|), \log (|U_{\beta,n}|) \}.$$

Hence, if we assume that R is large enough, so that $\log 2 \leq \delta(\log R)^{\frac{1}{1+H}}$, we have

$$\begin{aligned} \left\{ \log u > (\lambda + \delta)(\log R)^{\frac{1}{1+H}} \right\} &\subset \left\{ \max \{ \log (|u - U_{\beta,n}|), \log (|U_{\beta,n}|) \} > \lambda(\log R)^{\frac{1}{1+H}} \right\} \\ &= \left\{ \log (|u - U_{\beta,n}|) > \lambda(\log R)^{\frac{1}{1+H}} \right\} \cup \left\{ \log (|U_{\beta,n}|) > \lambda(\log R)^{\frac{1}{1+H}} \right\} \end{aligned}$$

Owing to simple additivity properties of \mathbb{P} , we thus get

$$\begin{aligned} &\mathbb{P} \left(\log (|U_{\beta,n}|) > \lambda(\log R)^{\frac{1}{1+H}} \right) \\ &\geq \mathbb{P} \left(\log u > (\lambda + \delta)(\log R)^{\frac{1}{1+H}} \right) - \mathbb{P} \left(\log (|u - U_{\beta,n}|) > \lambda(\log R)^{\frac{1}{1+H}} \right) \\ &\geq \mathbb{P} \left(\log u > (\lambda + \delta)(\log R)^{\frac{1}{1+H}} \right) - \frac{1}{R^{3+\lambda}}, \end{aligned} \tag{5.47}$$

where the last inequality is a direct consequence of (5.40). Now recall that we have chosen λ fulfilling condition (5.44). Applying Corollary 4.5 in this context yields

$$\lim_{R \rightarrow \infty} \frac{1}{\log R} \log \mathbb{P} \left(\log u > (\lambda + \delta)(\log R)^{\frac{1}{1+H}} \right) > 1 - \delta,$$

and thus, for R large enough the following holds true

$$\mathbb{P} \left(\log u > (\lambda + \delta)(\log R)^{\frac{1}{1+H}} \right) > \frac{1}{R^{1-\delta}}.$$

Plugging this relation into (5.47) gives, for R large enough

$$\mathbb{P} \left(\log (|U_{\beta,n}|) > \lambda(\log R)^{\frac{1}{1+H}} \right) > \frac{1}{R^{1-\frac{\delta}{2}}} \tag{5.48}$$

We now gather (5.45), (5.38) and (5.48) in order to get

$$\mathbb{P} \left(\log \max_{z \in \mathcal{N}_R} |U_{\beta,n}(t, z)| \leq \lambda(\log R)^{\frac{1}{1+H}} \right) \leq \left(1 - \frac{1}{R^{1-\frac{\delta}{2}}} \right)^{|\mathcal{N}_R|} \leq \exp \left(-c_{H,M,\varepsilon} R^{\frac{\delta}{2}-\varepsilon} \right).$$

In conclusion, since we can choose $\varepsilon < \frac{\delta}{2}$, we have established the bound

$$\mathbb{P} \left\{ \log \max_{z \in \mathcal{N}_R} |U_{\beta,n}(t, z)| \leq \lambda(\log R)^{\frac{1}{1+H}} \right\} \leq \exp \{ -R^v \}$$

for some $v > 0$ and for R large enough. Resorting again to Borel-Cantelli's lemma, this implies

$$\liminf_{m \rightarrow \infty} (\log m)^{-\frac{1}{1+H}} \log \max_{z \in \mathcal{N}_m} |U_{\beta(m),n}(t, z)| \geq \lambda \quad \text{a.s.} \tag{5.49}$$

Step 3: Conclusion. Combining the above inequality (5.49) with (5.42), we have obtained:

$$\liminf_{m \rightarrow \infty} (\log m)^{-\frac{1}{1+H}} \log \max_{z \in \mathcal{N}_m} u(t, z) \geq \lambda \quad \text{a.s.}$$

By the fact that

$$\max_{z \in \mathcal{N}_R} u(t, z) \leq \max_{|x| \leq R} u(t, x)$$

and by the monotonicity of $\max_{|x| \leq R} u(t, x)$ in R , we can now easily deduce that:

$$\liminf_{R \rightarrow \infty} (\log R)^{-\frac{1}{1+H}} \log \max_{|x| \leq R} u(t, x) \geq \lambda \quad \text{a.s.}$$

Finally, recall that λ satisfies condition (5.44), and can thus be chosen arbitrarily close to $\hat{c}_{H,t}^{-1/(1+H)}$. Our proof is thus easily concluded. \square

5.2 Proof of the upper bound

The proof of the upper bound is based on a quantification of the fluctuations of u in boxes around the points $x_j \in \mathcal{N}_R$, defined in the proof of Proposition 5.8. We will first need an evaluation of the modulus of continuity of u in the space variable.

Proposition 5.9. *For any $\beta \in (0, 2H - 1/2)$, there exists a constant C depending on α , H and t , such that for any $x, y \in \mathbb{R}$ and any $p \geq 2$,*

$$\mathbb{E} [|u(t, x) - u(t, y)|^p] \leq |x - y|^{p\beta} \exp\left(Cp^{1+\frac{1}{H}}\right). \tag{5.50}$$

Proof. First we estimate the L^2 norm using the Wiener chaos expansion of the solution and the notation used in the proof of Proposition 5.6. In this way we can write

$$\mathbb{E} (|I_n(f_n(\cdot, t, x)) - I_n(f_n(\cdot, t, y))|^2) = n! \|f_n(\cdot, t, x) - f_n(\cdot, t, y)\|_{\mathcal{H}^{\otimes n}}^2 = n! c_H^n L_n(x, y), \tag{5.51}$$

where we have set

$$L_n(x, y) = \|\mathcal{F}f_n(\cdot, t, x) - \mathcal{F}f_n(\cdot, t, y)\|_{L^2([0,t]^n \times \mathbb{R}^n, \lambda^n \times \mu^n)}^2.$$

Furthermore, it is readily checked that

$$L_n(x, y) = \frac{1}{(n!)^2} \int_{[0,t]^n} \int_{\mathbb{R}^n} K_{xy}(\xi) \prod_{i=1}^n e^{-\frac{1}{2}(s_{\sigma(i+1)} - s_{\sigma(i)})|\xi_{\sigma(i)} + \dots + \xi_{\sigma(1)}|^2} |\xi_{\sigma(i)}|^{1-2H} d\xi ds, \tag{5.52}$$

where

$$K_{xy}(\xi) \equiv \left| e^{-ix(\xi_{\sigma(n)} + \dots + \xi_{\sigma(1)})} - e^{-iy(\xi_{\sigma(n)} + \dots + \xi_{\sigma(1)})} \right|^2 = \left| e^{-ix(\xi_n + \dots + \xi_1)} - e^{-iy(\xi_n + \dots + \xi_1)} \right|^2.$$

Notice that one can recast identity (5.52) as

$$L_n(x, y) = \frac{1}{n!} \int_{S_n(t)} \int_{\mathbb{R}^n} K_{xy}(\xi) \prod_{i=1}^n e^{-\frac{1}{2}(s_{i+1} - s_i)|\xi_i + \dots + \xi_1|^2} |\xi_i|^{1-2H} d\xi ds,$$

where we recall that $S_n(t)$ is defined in Lemma 5.5. Furthermore, the kernel K_{xy} can be bounded as follows, for all $\beta \in (0, 1)$

$$|K_{xy}(\xi)| \leq |x - y|^{2\beta} |\xi_1 + \dots + \xi_n|^{2\beta}.$$

Making the change of variable $\xi_i + \dots + \xi_1 = \eta_i$, for all $i = 1, 2, \dots, n$, with the convention $\eta_0 = 0$ and using the bound $|\eta_i - \eta_{i-1}|^{1-2H} \leq |\eta_i|^{1-2H} + |\eta_{i-1}|^{1-2H}$ as in the proof of Proposition 5.6, we obtain

$$L_n(x, y) \leq \frac{|x - y|^{2\beta}}{n!} \sum_{\alpha \in D_n} \int_{S_n(t)} \int_{\mathbb{R}^n} \left(\prod_{i=1}^{n-1} e^{-\frac{1}{2}(s_{i+1} - s_i)\eta_i^2} |\eta_i|^{\alpha_i} \right) e^{-\frac{1}{2}(t - s_n)\eta_n^2} |\eta_n|^{\alpha_n + 2\beta} d\eta ds,$$

where we recall that D_n has been introduced in (5.23). Integrating with respect to the variables η_i and using Lemma 5.5 thus yields

$$L_n(x, y) \leq c^n \frac{|x - y|^{2\beta}}{n!} \sum_{\alpha \in D_n} \int_{S_n(t)} \left(\prod_{i=1}^{n-1} (s_{i+1} - s_i)^{-\frac{\alpha_i + 1}{2}} \right) (t - s_n)^{-\frac{\alpha_n + 2\beta + 1}{2}} ds,$$

where

$$c = \max \left\{ \int_{\mathbb{R}} e^{-\frac{1}{2}x^2} |x|^\delta dx : \delta \in [0, 4] \right\}.$$

At this point we can repeat the discussion following inequality (5.27). We find that, in order to ensure the convergence of the integral above, we have to choose $\beta < 2H - \frac{1}{2}$. Then applying Lemma 5.5, we end up with

$$L_n(x, y) \leq \frac{C^n |x - y|^{2\beta}}{n! \Gamma(nH + \frac{1}{2})}.$$

Plugging this relation into (5.51) and applying the hypercontractivity property on a fixed chaos, we have thus obtained

$$\|I_n(f_n(\cdot, t, x)) - I_n(f_n(\cdot, t, y))\|_{L^p(\Omega)} \leq \frac{C^n p^{\frac{n}{2}} |x - y|^\beta}{\Gamma(\frac{nH}{2} + 1)},$$

from which (5.50) is obtained exactly as in Proposition 5.6. □

We are now ready to prove the upper bound part of Theorem 1.1.

Proposition 5.10. *Under the assumptions of Theorem 1.1, for all $t > 0$ we have*

$$\liminf_{R \rightarrow \infty} (\log R)^{-\frac{1}{1+H}} \log \left(\max_{|x| \leq R} u(t, x) \right) \leq [\hat{c}_{H,t}]^{-\frac{1}{1+H}} \quad \text{a.s.}, \quad (5.53)$$

where $\hat{c}_{H,t}$ is defined in Corollary 4.5 and is related to (1.3) by Remark 4.6.

Proof. We shall use the same kind of notation as in the proof of Proposition 5.8, sometimes with a slightly different meaning (which should be clear from the context). Fix $R > 0$ and divide the interval $[-R, R]$ into subintervals I_j with the same length, for $j = 1, \dots, \mathcal{N}_R$ (notice that \mathcal{N}_R is now a cardinal instead of being a set as in Proposition 5.8), of length less than or equal to ℓ , for some $\ell > 0$ to be chosen later. Pick one point x_j of each interval I_j . By convention, we assume that I_1 contains 0, and we choose $x_1 = 0$. For any $x \in I_j$ we can write

$$u(t, x) \leq u(t, x_j) + |u(t, x) - u(t, x_j)|,$$

and hence:

$$\max_{|x| \leq R} u(t, x) \leq \max_j u(t, x_j) + \max_j \sup_{x, y \in I_j, |x-y| \leq \ell} |u(t, x) - u(t, y)|.$$

Therefore a simple elaboration of (5.46) yields

$$\log \max_{|x| \leq R} u(t, x) \leq \log 2 + \log \max_j \left(u(t, x_j), \sup_{x, y \in I_j, |x-y| \leq \ell} |u(t, x) - u(t, y)| \right). \quad (5.54)$$

Consider now $\lambda, \delta > 0$. Choose \mathcal{N}_R and ℓ large enough, so that $\cup_{j \leq \mathcal{N}_R} I_j$ covers $[-R, R]$. Owing to the stationarity of u , we can write

$$\begin{aligned} \mathbb{P} \left(\log \max_{|x| \leq R} u(t, x) \geq (\lambda + \delta)(\log R)^{\frac{1}{1+H}} \right) &\leq \sum_{j=1}^{\mathcal{N}_R} \mathbb{P} \left(\log \max_{x \in I_j} u(t, x) \geq (\lambda + \delta)(\log R)^{\frac{1}{1+H}} \right) \\ &\leq \mathcal{N}_R \mathbb{P} \left(\log \max_{x \in I_1} u(t, x) \geq (\lambda + \delta)(\log R)^{\frac{1}{1+H}} \right). \end{aligned}$$

Hence, inequality (5.54) enables to get, for R large enough

$$\mathbb{P} \left(\log \max_{|x| \leq R} u(t, x) \geq (\lambda + \delta)(\log R)^{\frac{1}{1+H}} \right) \leq \mathcal{N}_R \mathbb{P} \left(\log u(t, 0) \geq \lambda(\log R)^{\frac{1}{1+H}} \right) + F_R, \quad (5.55)$$

where F_R is a fluctuation term given by

$$F_R = \mathcal{N}_R \mathbb{P} \left(\log \sup \{ |u(t, x) - u(t, y)| : x, y \in I_1, |x - y| \leq \ell \} \geq \lambda (\log R)^{\frac{1}{1+H}} \right). \quad (5.56)$$

Furthermore, according to Corollary 4.5, for any $\rho > 0$ arbitrarily small and R large enough we have

$$\mathbb{P} \left\{ \log u(t, 0) \geq \lambda (\log R)^{\frac{1}{1+H}} \right\} \leq R^{-(\hat{c}_{H,t} - \rho) \lambda^{1+H}}.$$

Let us now specify our parameters: we assume that ρ, λ satisfy $(\hat{c}_{H,t} - \rho) \lambda^{1+H} > 1$, that is λ^{1+H} is arbitrarily close to $\hat{c}_{H,t}^{-1/(1+H)}$, as in condition (5.44). Then we can choose $\mathcal{N}_R = R^\eta$, with $1 < \eta < (\hat{c}_{H,t} - \rho) \lambda^{1+H}$, and $\ell = CR^{1-\eta}$ so that $\cup_{j \leq \mathcal{N}_R} I_j$ covers $[-R, R]$. We get

$$\mathcal{N}_R \mathbb{P} \left\{ \log u(t, 0) \geq \lambda (\log R)^{\frac{1}{1+H}} \right\} \leq R^{-\nu}, \quad \text{with } \nu = (\hat{c}_{H,t} - \rho) \lambda^{1+H} - \eta, \quad (5.57)$$

an notice that $\nu > 0$.

The fluctuation term F_R defined by (5.56) can be handled as follows: by Chebychev's inequality, we have

$$F_R \leq R^\eta e^{-p \lambda (\log R)^{\frac{1}{1+H}}} \mathbb{E} [\sup \{ |u(t, x) - u(t, y)|^p : x, y \in I_1, |x - y| \leq \ell \}]. \quad (5.58)$$

Consider now $0 < \gamma < \beta < 2H - \frac{1}{2}$ and p such that $\beta - \gamma > p^{-1}$. According to Garsia's lemma [13], we have

$$\mathbb{E} \left[\|u(t, \cdot)\|_{\gamma, I_1}^p \right] \leq c_{\gamma, p} \int_{I_1^2} \frac{\mathbb{E} (|u(t, x) - u(t, y)|^p)}{|x - y|^{\gamma p + 2}} dx dy,$$

where $\|f\|_{\gamma, I_1}$ stands for the γ -Hölder norm of f on the interval I_1 . Plugging the result of Proposition 5.9 into this inequality, we obtain

$$\mathbb{E} \left[\|u(t, \cdot)\|_{\gamma, I_1}^p \right] \leq c_{\gamma, \beta, p} |I_1|^{p(\beta - \gamma)} e^{Cp^{1+\frac{1}{H}}} = c_{\gamma, \beta, p} \ell^{p(\beta - \gamma)} e^{Cp^{1+\frac{1}{H}}},$$

and going back to (5.58), we end up with

$$F_R \leq c_{\gamma, \beta, p} R^\eta \ell^{p(\beta - \gamma)} \exp \left(-p \lambda (\log R)^{\frac{1}{1+H}} + Cp^{1+\frac{1}{H}} \right).$$

Recall that ℓ is of the form $CR^{-(\eta-1)}$, and we also take p such that the dominant term in the exponential above is $p \lambda (\log R)^{\frac{1}{1+H}}$. This is achieved for instance by taking

$$p = \left(\frac{\lambda}{2C} \right)^H (\log(R))^{\frac{H}{1+H}},$$

which implies $Cp^{1+\frac{1}{H}} = \frac{1}{2} p \lambda (\log R)^{\frac{1}{1+H}}$. In this way we obtain

$$F_R \leq c_{\gamma, \beta, p} R^{\eta - p(\beta - \gamma)(\eta - 1)} \exp \left(-\frac{1}{2} p \lambda (\log R)^{\frac{1}{1+H}} \right). \quad (5.59)$$

With the values of the parameters we have considered so far, observe that we can pick p such that $p(\beta - \gamma) > \frac{\eta}{\eta - 1}$. In this case we have $\eta - p(\beta - \gamma)(\eta - 1) = -\kappa$ with $\kappa > 0$, and we can recast (5.59) as

$$F_R \leq c_{\gamma, \beta, p} R^{-\kappa}. \quad (5.60)$$

Gathering our bounds (5.57) and (5.60) into (5.55), we thus have

$$\mathbb{P} \left(\log \max_{|x| \leq R} u(t, x) \geq (\lambda + \delta) (\log R)^{\frac{1}{1+H}} \right) \leq R^{-\nu} + R^{-\kappa}. \quad (5.61)$$

We can now conclude in the following way: relation (5.61) asserts that

$$\sum_{m \geq 1} \mathbb{P} \left\{ \log \max_{|x| \leq 2^m} u(t, x) \geq (\lambda + \delta)(\log 2^m)^{\frac{1}{1+H}} \right\} < \infty.$$

Hence Borel-Cantelli's lemma applies, and we deduce, almost surely

$$\limsup_{m \rightarrow \infty} (\log 2^m)^{-\frac{1}{1+H}} \log \left(\max_{|x| \leq 2^m} u(t, x) \right) \geq \lambda + \delta \geq (\hat{c}_{H,t} + \rho)^{-\frac{1}{1+H}} + \delta.$$

Because $\delta > 0$ and $\rho > 0$ are arbitrary, we thus get that almost surely

$$\limsup_{m \rightarrow \infty} (\log 2^m)^{-\frac{1}{1+H}} \log \max_{|x| \leq 2^m} u(t, x) \leq \hat{c}_{H,t}^{-\frac{1}{1+H}},$$

which implies

$$\limsup_{R \rightarrow \infty} (\log R)^{-\frac{1}{1+H}} \log \max_{|x| \leq R} u(t, x) \leq \hat{c}_{H,t}^{-\frac{1}{1+H}}.$$

This completes the proof of the upper bound. □

6 Appendix

Lemma 6.1. *For any $\alpha \in (0, 1)$ and we have*

$$\sup_{\beta \geq 1} \int_{\mathbb{R}} |\xi|^\alpha \left| \int_{[-\beta, \beta]} e^{ix\xi - \frac{x^2}{2}} dx \right|^2 d\xi < \infty.$$

Proof. Clearly the integral over $\{|\xi| \leq 1\}$ is uniformly bounded. So, it suffices to show that

$$\sup_{\beta \geq 1} \int_{|\xi| > 1} |\xi|^\alpha \left| \int_0^\beta \cos(x\xi) e^{-\frac{x^2}{2}} dx \right|^2 d\xi < \infty. \tag{6.1}$$

Integrating by parts, we obtain

$$\int_0^\beta \cos(x\xi) e^{-\frac{x^2}{2}} dx = e^{-\frac{\beta^2}{2}} \frac{\sin(\beta\xi)}{\xi} + \int_0^\beta x e^{-\frac{x^2}{2}} \frac{\sin(x\xi)}{\xi} dx.$$

Therefore,

$$\sup_{\beta \geq 1} \left| \int_0^\beta \cos(x\xi) e^{-\frac{x^2}{2}} dx \right| \leq |\xi|^{-1},$$

which implies (6.1). □

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