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# COMPUTATION OF RIEMANN MATRICES FOR THE HYPERBOLIC CURVES OF DETERMINANTAL POLYNOMIALS 

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#### Abstract

The numerical range of a matrix, according to Kippenhahn, is determined by a hyperbolic determinantal form of linear Hermitian matrices associated to the matrix. On the other hand, using Riemann theta functions, Helton and Vinnikov confirmed that a hyperbolic form always admits a determinantal representation of linear real symmetric matrices. The Riemann matrix of the hyperbolic curve plays the main role in the existence of real symmetric matrices. In this article, we implement computations of the Riemann matrix and the Abel-Jacobi variety of the hyperbolic curve associated to a determinantal polynomial of a matrix. Further, we prove that the lattice of the AbelJacobi variety is decomposed into the direct sum of two orthogonal lattices for some $4 \times 4$ Toeplitz matrices.


## 1. Introduction

Let $T$ be an $n \times n$ complex matrix. The real ternary form $F_{T}(x, y, z)$ associated to $T$ is defined by

$$
F_{T}(x, y, z)=\operatorname{det}\left(z I_{n}+x \Re(T)+y \Im(T)\right),
$$

where the two Hermitian matrices $\Re(T)=\left(T+T^{*}\right) / 2, \Im(T)=\left(T-T^{*}\right) /(2 i)$ correspond to the Cartesian decomposition $T=\Re(T)+i \Im(T)$. The form $F_{T}(x, y, z)$ is deeply related to the numerical range $W(T)$ of $T$ defined by

$$
W(T)=\left\{\xi^{*} T \xi: \xi \in \mathbf{C}^{n}, \xi^{*} \xi=1\right\} .
$$

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$F(x, y, 1)=0$ and $F_{y}(x, y, 1)=0$ at $P$. The $\delta$-number of the curve $C_{F}$ at $P$ is defined as

$$
\delta(P)=\frac{1}{2}\left(\operatorname{dim} \mathbb{C}[[X, Y]] /\left(f, f_{Y}\right)-m+s\right)
$$

Let $Q_{1}, Q_{2}, \ldots, Q_{\ell}$ be the set of all singular points of the algebraic curve $C_{F}$. The genus $g\left(C_{F}\right)$ of the curve $C_{F}$ is given by

$$
g\left(C_{F}\right)=\frac{1}{2}(n-1)(n-2)-\sum_{j=1}^{\ell} \delta\left(Q_{j}\right)
$$

An irreducible curve is called a rational (resp., elliptic) curve if its genus $g=0$ (resp., $g=1$ ).

We assume that the form $F(x, y, z)$ is irreducible. By using the resolution of singular points of the curve, we can construct a compact Riemann surface $\Gamma$ which is a parameter space of the curve $C_{F}$. This Riemann surface is called the nonsingular model $\tilde{C}_{F}$ of the curve $C_{F}$. If the curve $C_{F}$ has no singular points, the curve itself is a Riemann surface. We consider the canonical projection $\pi$ of $\tilde{C}_{F}$ onto $C_{F}: F(x, y, z)=0$. For a sufficiently small open neighborhood $V$ of a singular point $P$ of $C_{F}$, the set $\pi^{-1}(V)$ is composed of a finite number of connected open sets $V_{1}, V_{2}, \ldots, V_{q}$ of $\tilde{C}_{F}$ which form a multilevel crossing around $P$. The complex affine curve $F(x, y, 1)=0$ is parameterized by two rational functions $L_{1}$, $L_{2}$ on $\tilde{C}_{F}: x=L_{1}(s), y=L_{2}(s)$. The nonsingular model $\tilde{C}_{F}$ is homeomorphic to one of the compact surfaces: the Riemann sphere $\mathbb{C} \cup\{\infty\}$, the torus $\mathbb{C} /(\mathbb{Z}+i \mathbb{Z})$, or a $g$-hole torus with $g \geq 2$. The genus $g$ of the curve $C_{F}$ coincides with the number of holes of the compact Riemann surface $\tilde{C}_{F}$. (For references on Riemann surfaces, see, e.g., [5], [4], [10].)

In this article, we use the complex analytic method to study invariants of the numerical ranges of matrices. The complex analytic technique is often involved in studying operators on a Hilbert space (see [6], [15]). We investigate some invariants of the Riemann surface $\tilde{C}_{F}$ of the algebraic curve associated to a matrix $T$. In particular, we compute the Riemann matrix of $\tilde{C}_{F}$. The Riemann matrix is a common invariant for the birationally equivalent class of algebraic curves, and hence the invariant for the curve $C_{F}$ and its dual curve. The Riemann matrix generates a lattice of the Abel-Jacobi variety of $\tilde{C}_{F}$. We prove that the lattice is decomposed into the direct sum of two orthogonal sublattices for some $4 \times 4$ Toeplitz matrices.

## 2. Riemann matrix

Let $\Gamma$ be a compact Riemann surface with genus $g$. We recall the definition of a Riemann matrix which adopts the standard notation used in [5]. For an arbitrary base point $P_{0}$ of $\Gamma$, define a cycle to be a continuous map $f:[0,1] \rightarrow \Gamma$ with $f(0)=f(1)=P_{0}$. Two such maps $f_{1}, f_{2}$ produce a new map $f_{2} \circ f_{1}$ by the rule

$$
\left(f_{2} \circ f_{1}\right)(t)=f_{1}(2 t), \quad\left(f_{2} \circ f_{1}\right)(1 / 2+t)=f_{2}(2 t),
$$

$0 \leq t \leq 1 / 2$. The inverse $f^{-1}$ of $f$ is given by $f^{-1}(t)=f(1-t)$. Two maps $f, g$ are identified if they are homotopic, that is, there is a continuous map
$F:[0,1] \times[0,1] \rightarrow \Gamma$ satisfying

$$
F(0, t)=f(t), \quad F(1, t)=g(t), \quad F(s, 0)=F(s, 1)=P_{0}
$$

The fundamental group $\Pi_{1}(\Gamma)$ of $\Gamma$ is the collection of cycles of $\Gamma$ under the homotopic relation. We identify an element of $\Pi_{1}(\Gamma)$ with a closed, oriented, piecewise smooth curve on $\Gamma$. Let $\omega$ be an arbitrary $C^{(1)}$-differential 1-form on $\Gamma$. On a local coordinate $z=x+i y$ in an open neighborhood $V$ of a point $P$ of $\Gamma$, the differential form $\omega$ is expressed as $\omega=Q(z, \bar{z}) d z+R(z, \bar{z}) d \bar{z}$, by $C^{(1)}$-differentiable complex-valued functions $Q, R$. We define a differential 2 -form $d \omega$ by

$$
d \omega=\left(\frac{\partial R}{\partial z}-\frac{\partial Q}{\partial \bar{z}}\right) d z \wedge d \bar{z}
$$

If $d \omega=0$, the differential form $\omega$ is closed. If there exists a holomorphic function $U$ on an open neighborhood of an arbitrary point $P$ of $\Gamma$ satisfying $\omega=U_{z} d z$, then $\omega$ is called a holomorphic 1-form or an abelian differential form of the first kind. A holomorphic 1-form is a closed form.

Given two cycles $a, b$, the equation $\int_{a^{-1} \circ b^{-1} \text { oaob }} \omega=0$ holds for any holomorphic differential 1-form $\omega$. Let $N$ be the normal subgroup of $\Pi_{1}(\Gamma)$ generated by $\left\{a^{-1} \circ\right.$ $\left.b^{-1} \circ a \circ b: a, b \in \Pi_{1}(\Gamma)\right\}$. The abelian group $H_{1}(\Gamma, \mathbb{Z})=\Pi_{1}(\Gamma) / N$ is called the 1-dimensional homology group (see [4]). A cycle $c$ of $\Gamma$ is homologous to zero if $\int_{c} \omega=0$ for any holomorphic differential 1-form $\omega$. For two cycles $a$ and $b$ with different base points, the intersection index $a \cdot b$ of two cycles counts the number of intersections, taking the orientation of the cycles into account. If two cycles $a$ and $b$ do not intersect transversally, their intersection index $a \cdot b$ is zero. If they intersect once, $a \cdot b$ is 1 if the outer product of the tangent vectors $t_{a} \times t_{b}$ points out from the surface, and -1 if the outer product points into the surface. The homology group $H_{1}(\Gamma) \cong \mathbb{Z}^{2 g}$ has a canonical basis $\left\{a_{1}, \ldots, a_{g}, b_{1}, \ldots, b_{g}\right\}$ satisfying the symplectic relations

$$
a_{i} \cdot a_{j}=0, \quad b_{i} \cdot b_{j}=0, \quad a_{i} \cdot b_{j}=\delta_{i j}
$$

where $\delta_{i j}$ is the Kronecker delta (see [10]). Any element $c$ of $H_{1}(\Gamma)$ is expressed as

$$
c=m_{1} a_{1}+\cdots+m_{g} a_{g}+m_{g+1} b_{1}+\cdots+m_{2 g} b_{g}
$$

for some integers $m_{1}, \ldots, m_{2 g}$. The set of all holomorphic differential 1-forms is a $g$-dimensional complex vector space. Let $\left\{\omega_{1}, \ldots, \omega_{g}\right\}$ be a basis for the holomorphic differential 1-form vector space. The period matrix $\Omega$ of $\Gamma$ is a $g \times 2 g$ complex matrix defined as $\Omega=[A B]$, where $A=\left(A_{i j}\right)_{i, j=1}^{g}$ and $B=\left(B_{i j}\right)_{i, j=1}^{g}$ are given by

$$
A_{i j}=\oint_{a_{j}} \omega_{i}, \quad B_{i j}=\oint_{b_{j}} \omega_{i}
$$

The matrix $A$ is known to be invertible. The Riemann matrix of $\Gamma$ is defined as $R=A^{-1} B$. The imaginary part matrix $\left(\Im R_{i j}\right)_{i, j \leq g}$ is positive definite. The Riemann matrix is closely related to the geometric and analytic structure of the Riemann surface under consideration (see [10]).

The period matrix $\Omega$ generates a nondegenerate lattice

$$
\Lambda=\left\{m_{1} A_{1}+\cdots+m_{g} A_{g}+m_{g+1} B_{1}+\cdots+m_{2 g} B_{g}: m_{1}, m_{2}, \ldots, m_{2 g} \in \mathbb{Z}\right\}
$$

where $A_{1}, \ldots, A_{g}$ and $B_{1}, \ldots, B_{g}$ are, respectively, the columns of $A$ and $B$. This lattice is the period lattice associated to the basis $\left\{\omega_{1}, \ldots, \omega_{g}\right\}$ of the space of holomorphic 1-forms on $\Gamma$. Given a point $z_{0} \in \Gamma$, the holomorphic mapping $J$ : $\Gamma \rightarrow \mathbb{C}^{g} / \Lambda$ defined by

$$
J(z)=\left(\int_{z_{0}}^{z} \omega_{1}, \int_{z_{0}}^{z} \omega_{2}, \ldots, \int_{z_{0}}^{z} \omega_{g}\right) \bmod \Lambda
$$

is a holomorphic embedding of $\Gamma$ onto $J(\Gamma)=\mathbb{C}^{g} / \Lambda$. The $g$-dimensional compact space $\mathbb{C}^{g} / \Lambda$ is called the Abel-Jacobi variety of $\Gamma$.

Suppose that $\Gamma$ is a Riemann surface with genus $g$ associated to an irreducible hyperbolic curve $F_{T}(x, y, 1)=0$ defined by a complex $n \times n$ matrix $T$. It would be interesting to see whether there exists a canonical basis $\left\{a_{1}, a_{2}, \ldots, a_{g}, b_{1}, \ldots, b_{g}\right\}$ for the homology group $H_{1}(\Gamma)$ for which the Riemann matrix $R=A^{-1} B=$ $\left(e_{g+1}, \ldots, e_{2 g}\right)^{T}$ and the basis $e_{1}=(1,0, \ldots, 0)^{T}, \ldots, e_{g}=(0, \ldots, 0,1)^{T}$ span the lattice

$$
\Sigma=\left\{m_{1} e_{1}+\cdots+m_{g} e_{g}+m_{g+1} e_{g+1}+\cdots+m_{2 g} e_{2 g}: m_{1}, m_{2}, \ldots, m_{2 g} \in \mathbb{Z}\right\}
$$

so that $\Sigma$ is decomposed into the direct sum of two orthogonal lattices $\Sigma_{1}, \Sigma_{2}$, where each of $\Sigma_{1}$ and $\Sigma_{2}$ is spanned by $g$ elements of the column vectors $\left\{e_{1}, \ldots\right.$, $\left.e_{g}, e_{g+1}, \ldots, e_{2 g}\right\}$ and $\Sigma_{1} \cong \Sigma_{2} \cong \mathbb{Z}^{g}$ as free abelian groups. In the case $g=1$, the Riemann matrix $R$ is none other than the $\tau$-invariant of the elliptic curve $F_{T}(x, y, 1)=0$. In [3], the authors showed that the $\tau$-invariants of the elliptic curves for some $4 \times 4$ matrices are pure imaginary. In this article, we generalize the problem for $g \geq 2$, and we verify the orthogonal lattice decomposition in Sections 3 and 4 for some $4 \times 4$ Toeplitz matrices.

## 3. Quartic curve of genus 2

Using a graph-theoretic technique, C. Tretkoff and M. Tretkoff [21] provided an explicit method of constructing a canonical basis $\left\{a_{j}, b_{j}: j=1, \ldots, g\right\}$ for the homology group $H_{1}(\Gamma)=H_{1}(\Gamma, \mathbb{Z})$ of a Riemann surface $\Gamma$. Deconinck and van Hoeji [5] gave an algorithm to compute the period matrix of the Riemann surface using the Tretkoff method, where the Maple algcurves package, a collection of Maple programs for computations with algebraic curves, is performed (see also [4]). In this article, we follow the algorithm in [5], and implement the algcurves package to find a canonical basis for the group $H_{1}(\Gamma)$ associated to $4 \times 4$ Toeplitz matrices. We provide some numerical values to explain the process of computing the Riemann matrix.

Consider a Toeplitz matrix

$$
T=\left(\begin{array}{cccc}
0 & 2 & 2 a & 2 k  \tag{3.1}\\
0 & 0 & 2 & 2 a \\
0 & 0 & 0 & 2 \\
0 & 0 & 0 & 0
\end{array}\right)
$$

for $a \in \mathbb{R}, k=a^{2}-1$. Then

$$
\begin{aligned}
F_{T}(x, y, z)= & \left(\left(4-4 a^{2}\right) x^{2}+4 a^{3} x z-\left(4+a^{4}\right) z^{2}\right) y^{2} \\
& +\left(\left(4-4 a^{2}\right) x^{4}+4 a^{3} x^{3} z-\left(4+a^{4}\right) x^{2} z^{2}+z^{4}\right)
\end{aligned}
$$

If $a=0$, then the curve $C_{F}$ is a rational curve. We assume that $a>0$. If $a=1$, then the polynomial $F_{T}(x, y, z: a)=z\left(4 x^{3}+4 x y^{2}-5 x^{2} z-5 y^{2} z+z^{3}\right)$ is reducible. Since

$$
\begin{aligned}
F(x, 1, z)= & \left(-(2+2 a) x+\left(2+2 a+a^{2}\right) z\right)\left(-(2-2 a) x-\left(2-2 a+a^{2}\right) z\right) \\
& +\left(4-4 a^{2}\right) x^{4}+4 a^{3} x^{3} z-\left(4+a^{4}\right) x^{2} z^{2}+z^{4}
\end{aligned}
$$

the point $(x, y, z)=(0,1,0)$ is an ordinary double point under the condition $0<a<1$ or $a>1$. The roots of the polynomial

$$
\begin{aligned}
p(x)= & \left((2 a+2) x-\left(a^{2}+2 a+2\right)\right)\left((2 a-2) x-\left(a^{2}-2 a+2\right)\right) \\
& \times\left((2 a+2) x^{2}-a^{2} x-1\right)\left((2 a-2) x^{2}-a^{2} x+1\right)
\end{aligned}
$$

are branch points of the Riemann surface. For $a=\sqrt{2}$, two of these points coincide. We treat the case $a>\sqrt{2}$. In this case, the curve $F_{T}(x, y, z)=0$ is a hyperelliptic curve of genus 2 . The roots of $p(x)$ are given by

$$
\begin{aligned}
r_{1}=\frac{a^{2}-\sqrt{a^{4}+8 a+8}}{4(a+1)} & <r_{2}=\frac{a^{2}-\sqrt{a^{4}-8 a+8}}{4(a-1)}<r_{3}=\frac{a^{2}+\sqrt{a^{4}+8 a+8}}{4(a+1)} \\
& <r_{4}=\frac{a^{2}-2 a+2}{2(a-1)}<r_{5}=\frac{a^{2}+2 a+2}{2(a+1)} \\
& <r_{6}=\frac{a^{2}+\sqrt{a^{4}-8 a+8}}{4(a-1)}
\end{aligned}
$$

Theorem 3.1. Let $T$ be the $4 \times 4$ matrix defined in (3.1) with $a>\sqrt{2}$. Suppose that $\Gamma$ is the Riemann surface associated to the irreducible hyperbolic form $F_{T}(x, y, z)$. Then there exists a canonical basis $\left\{\tilde{a}_{1}, \tilde{a}_{2}, \tilde{b}_{1}, \tilde{b}_{2}\right\}$ for the homology group $H_{1}(\Gamma)$ such that the lattice

$$
\Sigma=\left\{m_{1} e_{1}+m_{2} e_{2}+m_{3} e_{3}+m_{4} e_{4}: m_{1}, m_{2}, m_{3}, m_{4} \in \mathbb{Z}\right\}
$$

spanned by the Riemann matrix $R=A^{-1} B=\left(e_{3}, e_{4}\right)^{T}$ and the basis $\left\{e_{1}=\right.$ $\left.(1,0), e_{2}=(0,1)\right\}$, is the direct sum of two mutually orthogonal lattices $\Sigma_{1}=$ $\left\{m_{1} e_{1}+m_{2} e_{2}: m_{1}, m_{2} \in \mathbb{Z}\right\}$ and $\Sigma_{2}=\left\{m_{3} e_{3}+m_{4} e_{4}: m_{3}, m_{4} \in \mathbb{Z}\right\}$.
Proof. Consider the curve $f(x, y)=F_{T}(x, y, 1)=0$ which is described by

$$
\begin{aligned}
& \left((2 a+2) x-\left(a^{2}+2 a+2\right)\right)\left((2 a-2) x-\left(a^{2}-2 a+2\right)\right) y^{2} \\
& \quad=\left((2 a-2) x^{2}-a^{2} x+1\right)\left(-(2 a+2) x^{2}+a^{2} x+1\right) .
\end{aligned}
$$

The basis $\left\{\omega_{1}, \omega_{2}\right\}$ for the space of holomorphic differential 1-forms is given by

$$
\omega_{1}=-\frac{1}{2} \frac{d x}{\left(\left(4 a^{2}-4\right) x^{2}-4 a^{3} x+\left(a^{4}+4\right)\right) y}
$$

and

$$
\omega_{2}=-\frac{1}{2} \frac{x d x}{\left(\left(4 a^{2}-4\right) x^{2}-4 a^{3} x+\left(a^{4}+4\right)\right) y} .
$$

The base point $x_{0}$ is taken on the real line with $x_{0}<r_{1}$, and hence the function $y=y(x)$ with $f(x, y(x))=0$ takes pure imaginary values at $x=x_{0}$. We define the two sheets of the surface $\Gamma$ as $\Im\left(y_{1}\left(x_{0}\right)\right)<0$ and $\Im\left(y_{2}\left(x_{0}\right)\right)>0$.

Suppose that $\left\{a_{1}, a_{2}, b_{1}, b_{2}\right\}$ is a canonical basis for the homology group $H_{1}(\Gamma)$ produced by the algorithm in [5] and the algcurves implementation. We construct a new basis depending on the canonical basis as follows.

1. The cycle $\tilde{a}_{1}=a_{1}$ starts on sheet 1 , encircles branch point $r_{1}=\left(a^{2}-\right.$ $\left.\sqrt{a^{4}+8 a+8}\right) /(4(a+1))$ to arrive at sheet 2 , and encircles branch point $r_{2}=\left(a^{2}-\sqrt{a^{4}-8 a+8}\right) /(4(a-1))$ to arrive at sheet 1 .
2. The cycle $\tilde{a}_{2}=-a_{2}$ starts on sheet 1 , encircles branch point $r_{3}=\left(a^{2}+\right.$ $\left.\sqrt{a^{4}+8 a+8}\right) /(4(a+1))$ to arrive at sheet 2 , and encircles branch point $r_{4}=\left(a^{2}-2 a+2\right) /(2(a-1))$ to arrive at sheet 1 .
3. The cycle $\tilde{b}_{2}=-b_{2}-a_{1}$ starts on sheet 1 , encircles branch point $r_{4}=$ $\left(a^{2}-2 a+2\right) /(2(a-1))$ to arrive at sheet 2 , and encircles branch point $r_{5}=\left(a^{2}+2 a+2\right) /(2(a+1))$ to arrive at sheet 1 .
4. The cycle $\tilde{b}_{3}=\tilde{b}_{1}-\tilde{b}_{2}=b_{1}+b_{2}+a_{2}$ starts on sheet 1 , encircles branch point $r_{2}=\left(a^{2}-\sqrt{a^{4}-8 a+8}\right) /(4(a-1))$ to arrive at sheet 2 , and encircles branch point $r_{3}=\left(a^{2}+\sqrt{a^{4}+8 a+8}\right) /(4(a+1))$ to arrive at sheet 1 .
Consider the set $\left\{\tilde{a}_{1}, \tilde{a}_{2}, \tilde{b}_{1}, \tilde{b}_{2}\right\}$ in the group $H_{1}(\Gamma)$ given by

$$
\begin{aligned}
& \tilde{a}_{1}=a_{1}, \quad \tilde{a}_{2}=-a_{2}, \\
& \tilde{b}_{1}=b_{1}-a_{1}+a_{2}, \quad \tilde{b}_{2}=-b_{2}-a_{1} .
\end{aligned}
$$

Then, the new basis $\left\{\tilde{a}_{1}, \tilde{a}_{2}, \tilde{b}_{1}, \tilde{b}_{2}\right\}$ of $H_{1}(\Gamma)$ satisfies the symplectic structure:

$$
\begin{aligned}
& \tilde{b}_{1} \cdot \tilde{b}_{2}=\left(b_{1}-a_{1}+a_{2}\right) \cdot\left(-\left(b_{2}+a_{1}\right)\right)=-\left(b_{1} \cdot a_{1}+a_{2} \cdot b_{2}\right)=0, \\
& \tilde{b}_{1} \cdot \tilde{b}_{1}=\left(b_{1}-a_{1}+a_{2}\right) \cdot\left(b_{1}-a_{1}+a_{2}\right)=-b_{1} \cdot a_{1}-a_{1} \cdot b_{1}=0, \\
& \tilde{b}_{2} \cdot \tilde{b}_{2}=\left(b_{2}+a_{1}\right) \cdot\left(b_{2}+a_{1}\right)=b_{2} \cdot b_{2}+a_{1} \cdot a_{1}+b_{2} \cdot a_{1}+a_{1} \cdot b_{2}=0, \\
& \tilde{a}_{1} \cdot \tilde{b}_{1}=a_{1} \cdot\left(b_{1}-a_{1}+a_{2}\right)=1, \\
& \tilde{a}_{2} \cdot \tilde{b}_{1}=0 \\
& \tilde{a}_{1} \cdot \tilde{b}_{2}=a_{1} \cdot\left[-\left(b_{2}+a_{1}\right)\right]=0, \\
& \tilde{a}_{2} \cdot \tilde{b}_{2}=-a_{2} \cdot\left[-\left(b_{2}+a_{2}\right)\right]=1, \\
& \tilde{a}_{1} \cdot \tilde{a}_{1}=\tilde{a}_{2} \cdot \tilde{a}_{2}=\tilde{a}_{1} \cdot \tilde{a}_{2}=\tilde{a}_{2} \cdot \tilde{a}_{1}=0 .
\end{aligned}
$$

Taking a base point $x_{0}<r_{1}$, we define the branch $y_{2}(x)$ of the function $y=y(x)$ with $F_{T}(x, y(x), 1)=0$ on sheet 2 in each interval $r_{j}<x<r_{j+1}, j=1,2,3,4$, according to the algorithm given in [5] on pages 33-34. We characterize $y_{2}(x)$ and the function $U(x)=-1 /\left(\left(\left(4 a^{2}-4\right) x^{2}-4 a^{3} x+\left(a^{4}+4\right)\right) y_{2}(x)\right)$ on each interval $r_{j}<x<r_{j+1}$ as follows.

On the interval $r_{1}<x<r_{2}$, the functions $y_{2}(x)$ and $U(x)$ are real valued, $y_{2}(x)>0$ and $U(x)<0$. Moreover, $U(x)=-\frac{1}{\sqrt{V(x)}}$, where

$$
\begin{aligned}
V(x)= & \left((2 a+2) x-\left(a^{2}+2 a+2\right)\right)\left((2 a-2) x-\left(a^{2}-2 a+2\right)\right) \\
& \times\left((2 a-2) x^{2}-a^{2} x+1\right)\left(-(2 a+2) x^{2}+a^{2} x+1\right) .
\end{aligned}
$$

On the interval $r_{3}<x<r_{4}$, the functions $y_{2}(x)$ and $U(x)$ are real valued, $y_{2}(x)<0$ and $U(x)>0$.

On the interval $r_{2}<x<r_{3}$, the functions $y_{2}(x)$ and $U(x)$ take pure imaginary values, $\Im(U(x))<0$ and $\Im\left(y_{2}(x)\right)<0$.

On the interval $r_{4}<x<r_{5}$, the functions $y_{2}(x)$ and $U(x)$ take pure imaginary values, $\Im(U(x))>0$ and $\Im\left(y_{2}(x)\right)<0$.

According to the labeling of the branches of $y(x)$ in [5], we compute that

$$
\tilde{a}_{11}=-\int_{a_{1}} \frac{1}{2} \frac{d x}{\left(\left(4 a^{2}-4\right) x^{2}-4 a^{3} x+\left(a^{4}+4\right)\right) y(x)} d x=-\int_{r_{1}}^{r_{2}} \frac{d x}{\sqrt{V(x)}}<0
$$

and

$$
\tilde{a}_{12}=-\int_{\tilde{a}_{2}} \frac{1}{2} \frac{d x}{\left(\left(4 a^{2}-4\right) x^{2}-4 a^{3} x+\left(a^{4}+4\right)\right) y(x)} d x=\int_{r_{3}}^{r_{4}} \frac{d x}{\sqrt{V(x)}}>0
$$

Similarly, we have

$$
\tilde{b}_{12}=-\int_{b_{2}} \frac{1}{2} \frac{d x}{\left(\left(4 a^{2}-4\right) x^{2}-4 a^{3} x+\left(a^{4}+4\right)\right) y(x)}=i \int_{r_{4}}^{r_{5}} \frac{d x}{\sqrt{-V(x)}}
$$

which is a pure imaginary number with $\Im\left(\tilde{b}_{12}\right)>0$, and

$$
\begin{aligned}
b_{11}+b_{12}+a_{12} & =-\int_{b_{1}+b_{2}+a_{2}} \frac{1}{2} \frac{d x}{\left(\left(4 a^{2}-4\right) x^{2}-4 a^{3} x+\left(a^{4}+4\right)\right) y(x)} \\
& =-i \int_{r_{2}}^{r_{3}} \frac{d x}{\sqrt{-V(x)}} d x
\end{aligned}
$$

which is also a pure imaginary number with $\Im\left(b_{11}+b_{12}+a_{12}\right)<0$.
We proceed to compute the entries of the periodic matrix. The values of the integrals

$$
\oint_{\tilde{a}_{1}} \omega_{2}=-\frac{1}{2} \oint_{a_{1}} \frac{x d x}{\left[\left(4 a^{2}-4\right) x^{2}-4 a^{3} x+\left(a^{4}+4\right)\right] y(x)} d x=-\int_{r_{1}}^{r_{2}} \frac{x d x}{\sqrt{V(x)}}
$$

and

$$
\oint_{\tilde{a}_{2}} \omega_{2}=-\frac{1}{2} \oint_{\tilde{a}_{2}} \frac{x d x}{\left[\left(4 a^{2}-4\right) x^{2}-4 a^{3} x+\left(a^{4}+4\right)\right] y(x)} d x=\int_{r_{3}}^{r_{4}} \frac{x d x}{\sqrt{V(x)}}
$$

are real numbers. The values of the integrals

$$
\oint_{\tilde{b}_{2}} \omega_{2}=-\frac{1}{2} \oint_{\tilde{b}_{2}} \frac{x d x}{\left[\left(4 a^{2}-4\right) x^{2}-4 a^{3} x+\left(a^{4}+4\right)\right] y(x)} d x=i \int_{r_{4}}^{r_{5}} \frac{x d x}{\sqrt{-V(x)}}
$$

and

$$
\oint_{\tilde{b}_{1}-\tilde{b}_{2}} \omega_{2}=-\frac{1}{2} \oint_{\tilde{b}_{1}-\tilde{b}_{2}} \frac{x d x}{\left[\left(4 a^{2}-4\right) x^{2}-4 a^{3} x+\left(a^{4}+4\right)\right] y(x)} d x=-i \int_{r_{2}}^{r_{3}} \frac{x d x}{\sqrt{-V(x)}}
$$

are pure imaginary numbers, and hence $\oint_{\tilde{b}_{1}} \omega_{2}$ is also pure imaginary.
Let $\tilde{\Omega}=[\tilde{A} \tilde{B}]$ be the periodic matrix of $\Gamma$ with respect to the new basis $\left\{\tilde{a}_{1}, \tilde{a}_{2}, \tilde{b}_{1}, \tilde{b}_{2}\right\}$. Then the matrix $\tilde{A}=\left(\begin{array}{ccc}\tilde{a}_{11} & \tilde{a}_{12} \\ \tilde{a}_{21} & \tilde{a}_{22}\end{array}\right)$ is invertible and its entries are real. Further, the entries of the matrix $\tilde{B}=\left(\begin{array}{lll}\tilde{b}_{11} & \tilde{b}_{12} \\ \tilde{b}_{21} & \tilde{b}_{22}\end{array}\right)$ are pure imaginary, and hence the Riemann matrix $R=\tilde{A}^{-1} \tilde{B}$ is a complex symmetric matrix whose entries are pure imaginary.

Let $R=\left(\begin{array}{c}r_{11} \\ r_{21} \\ r_{22}\end{array}\right)$, and set $e_{1}=(1,0), e_{2}=(0,1), e_{3}=\left(r_{11}, r_{12}\right), e_{4}=$ $\left(r_{21}, r_{22}\right)=\left(r_{12}, r_{22}\right)$. The space $\mathbb{C}^{g} \cong \mathbb{R}^{2 g}$ is equipped with the real part of the standard inner product. Then the lattice $\Sigma$ spanned by the vectors $e_{1}, e_{2}, e_{3}, e_{4}$ is the direct sum of two orthogonal lattices $\Sigma_{1}=\left\{n e_{1}+m e_{2}: n, m \in \mathbb{Z}\right\}$ and $\Sigma_{2}=\left\{n e_{3}+m e_{4}: n, m \in \mathbb{Z}\right\}$.

We run the algcurves package for the case $a=2$; the period matrix $[\tilde{A} \tilde{B}]$ constructed by Theorem 3.1 is approximately evaluated by

$$
\tilde{A} \approx\left(\begin{array}{cc}
-0.195915 & 0.410645 \\
-0.0199325 & 0.382071
\end{array}\right), \quad \tilde{B} \approx\left(\begin{array}{ll}
0.170162 i & 0.655061 i \\
0.553593 i & 0.865335 i
\end{array}\right) .
$$

## 4. Quartic curve of genus 3

C. Tretkoff and M. Tretkoff [21, p. 501] mentioned that the Riemann matrix of the Fermat curve $x^{4}+y^{4}=1$ is $\operatorname{diag}(i, 2 i, i)$. Hence, the lattice $\Sigma$ is spanned by six mutually orthogonal bases of the lattice. They also found that the Riemann matrix of the Klein-Hurwitz curve $-y^{7}+x(1-x)^{2}=0$ is $(-1+\sqrt{7} i) / 2 I_{3}$. In this case, the lattice $\Sigma$ for this real Riemann surface cannot be decomposed into the direct sum of two orthogonal lattices $\Sigma_{1}, \Sigma_{2}$; each one has three generators. However, we may find a suitable basis for the homology group $H_{1}(\Gamma)$ so that the Riemann surface generates a lattice which is the direct sum of two orthogonal lattices for some $4 \times 4$ Toeplitz matrices. (For the computation of the Riemann matrix and the split of the lattice $\Sigma$, we refer the reader to [7], [12], [19].)

Theorem 4.1. Let $T$ be the $4 \times 4$ nilpotent matrix given by

$$
T=\left(\begin{array}{llll}
0 & 2 & 0 & 6 \\
0 & 0 & 2 & 0 \\
0 & 0 & 0 & 2 \\
0 & 0 & 0 & 0
\end{array}\right)
$$

Then the homology group $H_{1}(\Gamma)$ of the Riemann surface of the algebraic curve $C_{F}$ of the polynomial $F_{T}(x, y, z)$ has a canonical basis $\left\{\tilde{a}_{1}, \tilde{a}_{2}, \tilde{a}_{3}, \tilde{b}_{1}, \tilde{b}_{2}, \tilde{b}_{3}\right\}$ for the homology group $H_{1}(\Gamma)$ such that the Riemann matrix $R=A^{-1} B$ of the period
matrix $[A B]$ is the form

$$
R=\left(\begin{array}{lll}
i r_{11} & i r_{12} & r_{13} \\
i r_{12} & i r_{22} & r_{23} \\
r_{13} & r_{23} & i r_{33}
\end{array}\right)
$$

for some real numbers $r_{i j}, 1 \leq i \leq j \leq 3$. Furthermore, the lattice $\Sigma$ generated by the six vectors $e_{1}=(1,0,0), e_{2}=(0,1,0), e_{3}=(0,0,1)$ and $e_{4}=$ $\left(i r_{11}, i r_{12}, r_{13}\right), e_{5}=\left(i r_{12}, i r_{22}, r_{23}\right), e_{6}=\left(r_{13}, r_{23}, i r_{33}\right)$ is decomposed into the direct sum of two mutually orthogonal lattices generated, respectively, by $\left\{e_{1}, e_{2}, e_{6}\right\}$ and $\left\{e_{3}, e_{4}, e_{5}\right\}$.

Proof. It is easy to find that

$$
F_{T}(x, y, z)=16 y^{4}+\left(20 x^{2}-12 z^{2}\right) y^{2}+4 x^{4}-12 x^{2} z^{2}+z^{4} .
$$

The curve $C_{F}$ of the polynomial $F_{T}(x, y, z)$ has no singular points with genus 3 . From the partial derivative $F_{y}(x, y, 1)=8 y\left(5 x^{2}+8 y^{2}-3\right)$, a basis for the space of holomorphic differential 1-forms on the Riemann surface $\Gamma$ associated to $C_{F}$ consists of

$$
\begin{aligned}
& \omega_{1}=\frac{1}{8} \frac{d x}{y\left(5 x^{2}+8 y^{2}-3\right)}, \quad \omega_{2}=\frac{1}{8} \frac{d x}{5 x^{2}+8 y^{2}-3} \\
& \omega_{3}=\frac{1}{8} \frac{x d x}{y\left(5 x^{2}+8 y^{2}-3\right)}
\end{aligned}
$$

(see [4]). Our strategy is to find a basis $\left\{\tilde{a}_{1}, \tilde{a}_{2}, \tilde{a}_{3}, \tilde{b}_{1}, \tilde{b}_{2}, \tilde{b}_{3}\right\}$ of the homology group $H_{1}(\Gamma)$ so that the period matrix $[\tilde{A} \tilde{B}]$ is the form

$$
\tilde{A}=\left(\begin{array}{ccc}
\alpha_{11} & \alpha_{12} & 0 \\
\alpha_{21} & 0 & i \alpha_{23} \\
i \alpha_{31} & 0 & \alpha_{33}
\end{array}\right), \quad \tilde{B}=\left(\begin{array}{ccc}
0 & i \beta_{12} & \beta_{13} \\
i \beta_{21} & 0 & \beta_{23} \\
\beta_{31} & \beta_{32} & i \beta_{33}
\end{array}\right)
$$

where $\alpha_{i j}$ and $\beta_{i j}$ are real numbers. Then the form $\tilde{A}^{-1} \tilde{B}$ is the kind required.
The curve $C_{F}$ has the following eight branch points:

$$
\begin{aligned}
& P_{1}: x=-i \sqrt{5 / 3} \approx-1.29099 i, \quad P_{2}: x=-\frac{i}{\sqrt{3}} \approx-0.577350 i, \\
& P_{3}: x=\frac{-2-\sqrt{2}}{2} \approx-1.70711, \quad P_{4}: x=\frac{-2+\sqrt{2}}{2} \approx-0.292893, \\
& P_{5}: x=\frac{2-\sqrt{2}}{2} \approx 0.292893, \quad P_{6}=\frac{2+\sqrt{2}}{2} \approx 1.70711, \\
& P_{7}: x=\frac{i}{\sqrt{3}} \approx 0.577350 i, \quad P_{8}: x=i \sqrt{5 / 3} \approx 1.29099 i .
\end{aligned}
$$

We take a base point $x_{0}=-2.27279$. The labels of the four branches of the function are based on the following labeling:

$$
\begin{aligned}
& y_{1}(-2.27279) \approx-2.26981 i, \quad y_{2}(-2.27279) \approx-0.744951 i \\
& y_{3}(-2.27279) \approx 0.744951 i, \quad y_{4}(-2.27279) \approx 2.26981 i
\end{aligned}
$$

The values of $y_{j}\left(x_{0}\right)$ are pure imaginary, and labeled as

$$
\Im\left(y_{1}\left(x_{0}\right)\right)<\Im\left(y_{2}\left(x_{0}\right)\right)<\Im\left(y_{3}\left(x_{0}\right)\right)<\Im\left(y_{4}\left(x_{0}\right)\right) .
$$

First, we define the cycles $\left\{a_{1}, a_{2}, a_{3}, b_{1}, c_{6}-a_{1}, c_{7}\right\}$ in the following way:

1. The cycle $a_{1}$ starts on sheet 1 , encircles branch point $P_{1}$ to arrive at sheet 2 , and encircles branch point $P_{2}$ to arrive at sheet 1 .
2. The cycle $a_{2}$ starts on sheet 1 , encircles branch point $P_{4}$ to arrive at sheet 4 , and encircles branch point $P_{5}$ to arrive at sheet 1 .
3. The cycle $a_{3}$ starts on sheet 3 , encircles branch point $P_{2}$ to arrive at sheet 4 , and encircles branch point $P_{7}$ to arrive at sheet 3 .
4. The cycle $b_{1}$ starts on sheet 1 , encircles branch point $P_{1}$ to arrive at sheet 2, and encircles point $P_{8}$ to arrive sheet 1 .
5. The cycle $c_{6}-a_{1}$ starts on sheet 1 , encircles branch point $P_{2}$ to arrive at sheet 2 , encircles branch point $P_{3}$ to arrive at sheet 3 , encircles branch point $P_{7}$ to arrive at sheet 4 , and encircles branch point $P_{4}$ to arrive at sheet 1.
6. The cycle $c_{7}$ starts on sheet 1 , encircles branch point $P_{1}$ to arrive at sheet 2 , encircles branch point $P_{3}$ to arrive at sheet 3 , encircles branch point $P_{8}$ to arrive at sheet 4 , and encircles branch point $P_{4}$ to arrive at sheet 1 .

The cycles $b_{2}, b_{3}$ are defined by $b_{2}=b_{1}-a_{1}-2\left(c_{6}-a_{1}\right)+c_{7}$ and $b_{3}=a_{2}+$ $\left(c_{6}-a_{1}\right)+a_{1}-c_{7}$. Then the set $\left\{a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}\right\}$ is a canonical basis for the homology group $H_{1}(\Gamma)$. We use another canonical basis $\left\{\tilde{a}_{1}=a_{1}, \tilde{a}_{2}=a_{2}, \tilde{a}_{3}=\right.$ $\left.a_{3}-a_{2}, \tilde{b}_{1}=b_{1}, \tilde{b}_{2}=b_{2}+b_{3}-a_{2}, \tilde{b}_{3}=b_{3}\right\}$. Then the cycles $\tilde{b}_{2}, b_{3}$ satisfy $\tilde{b}_{2}=$ $b_{1}-\left(c_{6}-c_{1}\right), b_{3}=a_{2}+\left(c_{6}-a_{1}\right)+a_{1}-c_{7}$.

We outline the method for computing the period matrix $[\tilde{A} \tilde{B}]$ by computing some specified entries.

The integrals over the cycle $\tilde{a}_{1}=a_{1}$ give $\tilde{a}_{i 1}$, the column one entries of $\tilde{A}$. On the line segment $\{i X:-\sqrt{5 / 3}<X<-1 / \sqrt{3}\}$, the function $y(x)$ takes imaginary values. The branches $y_{1}(x), y_{2}(x)$ of the function on this line segment satisfy $y_{2}(x)=\overline{y_{1}(x)}$ and $\Im\left(y_{1}(x)\right)<0<\Im\left(y_{2}(x)\right)$, and hence the four branches $y_{j}(x)$ are vertices of a rectangle with edges parallel to the real or imaginary axis and satisfy $y_{1}(x)+y_{2}(x)+y_{3}(x)+y_{4}(x)=0$. The function $W(x)=1 /\left(\left(5 x^{2}+\right.\right.$ $\left.8 y^{2}-3\right) y$ ) satisfies $W_{2}(x)=\overline{W_{1}(x)}$ for $y=y_{j}(x)$, and $\Im\left(W_{2}(x)\right)<0<\Im\left(W_{1}(x)\right)$. Then

$$
\begin{aligned}
\tilde{a}_{11} & =\frac{1}{8} \oint_{\tilde{a}_{1}} \frac{d x}{\left(5 x^{2}+8 y^{2}-3\right) y} \\
& =\frac{1}{8} \int_{-\sqrt{5 / 3 i} i}^{-i / \sqrt{3}} \frac{d x}{\left(5 x^{2}+8 y_{2}^{2}-3\right) y_{2}}+\frac{1}{8} \int_{-i / \sqrt{3}}^{-\sqrt{5 / 3} i} \frac{d x}{\left(5 x^{2}+8 y_{1}^{2}-3\right) y_{1}} \\
& \approx(0.0778069-0.0066436 i)+(0.0778069+0.0066436 i) \\
& \approx 0.155614 .
\end{aligned}
$$

The function $U(x)=1 /\left(5 x^{2}+8 y^{2}-3\right)$ on the line segment takes pure imaginary values, and $\Im\left(U_{2}(x)\right)<0<\Im\left(U_{1}(x)\right)$. Then

$$
\begin{aligned}
\tilde{a}_{21} & =\frac{1}{8} \oint_{\tilde{a}_{1}} \frac{d x}{\left(5 x^{2}+8 y^{2}-3\right)} \\
& =\frac{1}{8} \int_{-\sqrt{5 / 3} i}^{-i / \sqrt{3}} \frac{d x}{\left(5 x^{2}+8 y_{2}^{2}-3\right)}+\frac{1}{8} \int_{-i / \sqrt{3}}^{-\sqrt{5 / 3} i} \frac{d x}{\left(5 x^{2}+8 y_{1}^{2}-3\right)} \\
& \approx 0.072851+0.072851 \\
& \approx 0.145702 .
\end{aligned}
$$

The function $V(x)=x /\left(\left(5 x^{2}+8 y^{2}-3\right)\right)$ on the line segment takes imaginary values, and $V_{2}(x)=-\overline{V_{1}(x)}$. These branches satisfy $\Re\left(V_{2}(x)\right)<0<\Re\left(V_{1}(x)\right)$. Then

$$
\begin{aligned}
\tilde{a}_{31} & =\frac{1}{8} \oint_{\tilde{a}_{1}} \frac{x d x}{\left(5 x^{2}+8 y^{2}-3\right)} \\
& =\frac{1}{8} \int_{-\sqrt{5 / 3} i}^{-i / \sqrt{3}} \frac{x d x}{\left(5 x^{2}+8 y_{2}^{2}-3\right) y_{2}}+\frac{1}{8} \int_{-i / \sqrt{3}}^{-\sqrt{5 / 3} i} \frac{x d x}{\left(5 x^{2}+8 y_{1}^{2}-3\right) y_{1}} \\
& \approx(-0.00568944-0.0669032 i)+(0.00568944-0.0669032 i) \\
& \approx-0.133806 i .
\end{aligned}
$$

Similarly, the integrals over the cycle $\tilde{b}_{1}=b_{1}$ give $\tilde{b}_{i 1}$, the column one entries of $\tilde{B}$. We compute that

$$
\begin{aligned}
\tilde{b}_{11}= & \frac{1}{8} \int_{\tilde{b}_{1}} \frac{d x}{\left(5 x^{2}+8 y^{2}-3\right) y} \\
= & \frac{1}{8} \int_{-2.7279}^{-i / \sqrt{3}}\left(\frac{1}{\left(5 x^{2}+8 y_{1}^{2}-3\right) y_{1}}-\frac{1}{\left(5 x^{2}+8 y_{2}^{2}-3\right) y_{2}}\right) d x \\
& -\frac{1}{8} \int_{-2.27279}^{i / \sqrt{3}}\left(\frac{1}{\left(5 x^{2}+8 y_{1}^{2}-3\right) y_{1}}-\frac{1}{\left(5 x^{2}+8 y_{1}^{2}-3\right) y_{2}}\right) d x \\
& +\frac{1}{8} \int_{-\sqrt{5 / 3} i}^{-i / \sqrt{3}}\left(\frac{1}{\left(5 x^{2}+8 y_{2}^{2}-3\right) y_{2}}-\frac{1}{\left(5 x^{2}+8 y_{1}^{2}-3\right) y_{1}}\right) d x \\
& -\frac{1}{8} \int_{\sqrt{5 / 3 i}}^{i / \sqrt{3}}\left(\frac{1}{\left(5 x^{2}+8 y_{2}^{2}-3\right) y_{2}}-\frac{1}{\left(5 x^{2}+8 y_{1}^{2}-3\right) y_{1}}\right) d x \\
= & \frac{1}{4} \Re\left(\int_{-2.27279}^{-i / \sqrt{3}}\left(\frac{1}{\left(5 x^{2}+8 y_{1}^{2}-3\right) y_{1}}-\frac{1}{\left(5 x^{2}+8 y_{2}^{2}-3\right) y_{2}}\right) d x\right) \\
& -\frac{1}{4} \Re\left(\int_{-\sqrt{5 / 3} i}^{-i / \sqrt{3}}\left(\frac{1}{\left(5 x^{2}+8 y_{2}^{2}-3\right) y_{1}}-\frac{1}{\left(5 x^{2}+8 y_{1}^{2}-3\right) y_{2}}\right) d x\right) \\
\approx & 2 \Re(-0.155614+i 0.0721547)-2 \Re(-0.155614) \\
= & 0 .
\end{aligned}
$$

Next, we compute

$$
\begin{aligned}
\tilde{b}_{21}= & \frac{1}{8} \oint_{\tilde{b}_{1}} \frac{d x}{5 x^{2}+8 y^{2}-3} \\
= & \frac{1}{8} \int_{-2.27279}^{-i / \sqrt{3}} \frac{d x}{5 x^{2}+8 y_{1}^{2}-3}-\frac{1}{8} \int_{-2.27279}^{-i / \sqrt{3}} \frac{d x}{5 x^{2}+8 y_{2}^{2}-3} \\
& +\frac{1}{8} \int_{-2.27279}^{i / \sqrt{3}} \frac{d x}{5 x^{2}+8 y_{2}^{2}-3}-\frac{1}{8} \int_{-2.27279}^{i / \sqrt{3}} \frac{d x}{5 x^{2}+8 y_{1}^{2}-3} \\
& +\frac{1}{8} \int_{-\sqrt{5 / 3 i}}^{-i / \sqrt{3}} \frac{d x}{5 x^{2}+8 y_{2}^{2}-3}-\frac{1}{8} \int_{-\sqrt{5 / 3 i}}^{-i / \sqrt{3}} \frac{d x}{5 x^{2}+8 y_{1}^{2}-3} \\
& -\frac{1}{8} \int_{\sqrt{5 / 3 i}}^{i / \sqrt{3}} \frac{d x}{5 x^{2}+8 y_{2}^{2}-3}+\frac{1}{8} \int_{\sqrt{5 / 3 i}}^{i / \sqrt{3}} \frac{d x}{5 x^{2}+8 y_{1}^{2}-3} \\
= & \frac{4 i}{8} \Im\left(\int_{-2.27279}^{-i / \sqrt{3}} \frac{d x}{5 x^{2}+8 y^{2}-3}\right)-\frac{4 i}{8} \Im\left(\int_{-\sqrt{5 / 3} i}^{-i / \sqrt{3}} \frac{d x}{5 x^{2}+8 y_{1}^{2}-3}\right) \\
\approx & 4 i \Im(-0.0555593+0.053564 i)-2 i \Im(-0.145702) \\
= & 0.214256 i .
\end{aligned}
$$

From the definition, we have

$$
\begin{aligned}
\tilde{b}_{31}= & \frac{1}{8} \oint_{\tilde{b}_{1}} \frac{x d x}{\left(5 x^{2}+8 y^{2}-3\right) y} d x \\
= & \left(\frac{1}{8} \int_{-2.27279}^{-\sqrt{5 / 3} i} \frac{x d x}{\left(5 x^{2}+8 y_{1}^{2}-3\right) y_{1}}-\frac{1}{8} \int_{-2.27279}^{-\sqrt{5 / 3} i} \frac{x d x}{\left(5 x^{2}+8 y_{2}^{2}-3\right) y_{2}}\right) \\
& +\left(\frac{1}{8} \int_{-2.27279}^{\sqrt{5 / 3} i} \frac{x d x}{\left(5 x^{2}+8 y_{2}^{2}-3\right) y_{2}}-\frac{1}{8} \int_{-2.27279}^{\sqrt{5 / 3} i} \frac{x d x}{\left(5 x^{2}+8 y_{1}^{2}-3\right) y_{1}}\right) .
\end{aligned}
$$

On the line segment $\{x(t)=(-2.27279)(1-t)+i \sqrt{5 / 3} i t: 0<t<1\}$, we consider the branches $y_{1}(x)$ and $y_{2}(x)$ of the function $y=y(x)$ satisfying $F_{T}(x, y(x), 1)=0$. The curve $9 x(t)^{4}+18 x(t)^{2}+5,0<t<1$, passes through the half-line $\{x: x<0\}$ at some point $t_{0} \in(0.8302473,0.8302474)$. Denote the branch of square roots a complex number

$$
\sqrt{\lambda}=\sqrt{|\lambda|} \exp (i \operatorname{Arg} \lambda / 2)
$$

for $-\pi<\operatorname{Arg} \lambda \leq \pi$. We define

$$
\begin{aligned}
& Y_{2}(x)=\frac{1}{2 \sqrt{2}} \sqrt{-5 x^{2}+3+\sqrt{9 x^{4}+18 x^{2}+5}} \\
& Y_{1}(x)=\frac{1}{2 \sqrt{2}} \sqrt{-5 x^{2}+3-\sqrt{9 x^{4}+18 x^{2}+5}}
\end{aligned}
$$

On the line $(-2.27279,-\sqrt{5 / 3} i)$, the functions $y_{2}(x), y_{1}(x)$ for $0<t<t_{0}$ are expressed, respectively, as

$$
\frac{x}{\left(5 x^{2}+8 Y_{2}^{2}-3\right) Y_{2}}, \quad \frac{x}{\left(5 x^{2}+8 Y_{1}^{2}-3\right) Y_{1}} .
$$

The functions $y_{1}(x), y_{2}(x)$ are expressed, respectively, as

$$
\frac{x}{\left(5 x^{2}+8 Y_{2}^{2}-3\right) Y_{2}}, \quad \frac{x}{\left(5 x^{2}+8 Y_{1}^{2}-3\right) Y_{1}}
$$

for $t_{0}<t<1$. Direct computations show that

$$
\begin{aligned}
& \frac{1}{8} \int_{-2.27279}^{-\sqrt{5 / 3} i} \frac{x d x}{\left(5 x^{2}+8 y_{1}^{2}-3\right) y_{1}}-\frac{1}{8} \int_{-2.27279}^{-\sqrt{5 / 3} i} \frac{x d x}{\left(5 x^{2}+8 y_{2}^{2}-3\right) y_{2}} \\
& \quad \approx 0.133806-0.0562865 i,
\end{aligned}
$$

and

$$
\begin{aligned}
& \frac{1}{8} \int_{-2.27279}^{\sqrt{5 / 3} i} \frac{x d x}{\left(5 x^{2}+8 y_{2}^{2}-3\right) y_{2}}-\frac{1}{8} \int_{-2.27279}^{\sqrt{5 / 3} i} \frac{x d x}{\left(5 x^{2}+8 y_{1}^{2}-3\right) y_{1}} \\
& \quad \approx 0.133806+0.0562865 i .
\end{aligned}
$$

Therefore,

$$
\begin{aligned}
\tilde{b}_{31} & =\frac{1}{8} \oint_{b_{1}} \frac{x d x}{\left(5 x^{2}+8 y^{2}-3\right) y} \\
& \approx(0.133806-0.0562865 i)+(0.133806+0.0562865 i) \approx 0.267613
\end{aligned}
$$

The remaining entries of the period matrix $[\tilde{A} \tilde{B}]$ can be evaluated in a similar way. We obtain numerically the matrix

$$
\tilde{A}=\left(\begin{array}{ccc}
0.155614 & -0.311228 & 0 \\
0.145702 & 0 & -0.214256 i \\
-0.133806 i & 0 & 0.267613
\end{array}\right)
$$

and

$$
\tilde{B}=\left(\begin{array}{ccc}
0 & -0.349582 i & -0.155614 \\
0.214256 i & 0 & 0.145702 \\
0.267613 & 0.267613 & 0.133806 i
\end{array}\right)
$$

Thus the Riemann matrix is numerically given by

$$
R=\tilde{A}^{-1} \tilde{B}=\left(\begin{array}{ccc}
1.69486 i & 0.847432 i & 0.152568 \\
0.847432 i & 1.54696 i & 0.576284 \\
0.152568 & 0.576284 & 0.576284 i
\end{array}\right)
$$

We equip the space $\mathbb{C}^{g} \cong \mathbb{R}^{2 g}$ with the real part of the standard inner product. Then it is easy to verify that

$$
\begin{aligned}
& \left\langle e_{1}, e_{3}\right\rangle=\left\langle e_{1}, e_{4}\right\rangle=\left\langle e_{1}, e_{5}\right\rangle=0, \\
& \left\langle e_{2}, e_{3}\right\rangle=\left\langle e_{2}, e_{4}\right\rangle=\left\langle e_{2}, e_{5}\right\rangle=0, \\
& \left\langle e_{6}, e_{3}\right\rangle=\left\langle e_{6}, e_{4}\right\rangle=\left\langle e_{6}, e_{4}\right\rangle=0 .
\end{aligned}
$$

Hence the two lattices

$$
\Sigma_{1}=\left\{n e_{1}+m e_{2}+\ell e_{6}: n, m, \ell \in \mathbb{Z}\right\}, \quad \Sigma_{2}=\left\{n e_{3}+m e_{4}+\ell e_{5}: n, m, \ell \in \mathbb{Z}\right\}
$$

are mutually orthogonal, and $\Sigma=\Sigma_{1} \oplus \Sigma_{2}$.

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