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# A BOUNDED TRANSFORM APPROACH TO SELF-ADJOINT OPERATORS: FUNCTIONAL CALCULUS AND AFFILIATED VON NEUMANN ALGEBRAS 

CHRISTIAN BUDDE and KLAAS LANDSMAN<br>Communicated by Q.-W. Wang


#### Abstract

Spectral theory and functional calculus for unbounded self-adjoint operators on a Hilbert space are usually treated through von Neumann's Cayley transform. Using ideas of Woronowicz, we redevelop this theory from the point of view of multiplier algebras and the so-called bounded transform (which establishes a bijective correspondence between closed operators and pure contractions). This also leads to a simple account of the affiliation relation between von Neumann algebras and self-adjoint operators.


## 1. Introductory overview

The theory of unbounded self-adjoint operators on a Hilbert space was initiated by von Neumann [7], partly motivated by mathematical problems of quantum mechanics. The monograph by Schmüdgen [10] presents an excellent survey of the present state of the art.

Von Neumann's approach was based on the Cayley transform, and in its subsequent development the notion of a spectral measure played an important role, especially in defining a functional calculus. We consider this route a bit indirect and will avoid both by first invoking the bounded transform instead of the Cayley transform; that is, the formal expressions

$$
\begin{align*}
& S=T{\sqrt{I+T^{2}}}^{-1},  \tag{1.1}\\
& T=S{\sqrt{I-S^{2}}}^{-1}, \tag{1.2}
\end{align*}
$$

[^0]Keywords. bounded transform, self-adjoint operators, von Neumann algebras.

Our aim here is to generalize these results to the case where $T$ is unbounded. This indeed turns out to be possible so that our main results are as follows. Throughout the remainder of this article, we assume that $T^{*}=T$ is possibly unbounded with bounded transform $S$.

Theorem 1.1. The spectra of $T$ and its bounded transform $S$ are related by

$$
\begin{align*}
\sigma(T) & =\left\{\mu\left(1-\mu^{2}\right)^{-1 / 2}: \mu \in \tilde{\sigma}(S)\right\}  \tag{1.11}\\
\sigma(S) & =\left\{\lambda\left(1+\lambda^{2}\right)^{-1 / 2}: \lambda \in \sigma(T)\right\}^{-} \tag{1.12}
\end{align*}
$$

where - denotes the closure in $\mathbb{R}$, and we abbreviate

$$
\begin{equation*}
\tilde{\sigma}(S)=\sigma(S) \cap(-1,1) \tag{1.13}
\end{equation*}
$$

Note that $\tilde{\sigma}(S)=\sigma(S)$ if and only if $T$ is bounded (in which case $\sigma(S)$ is a compact subset of $(-1,1)$ since $\pm 1 \in \sigma(S)$ if and only if $T$ is unbounded). We define the following operator algebras within $B(\mathcal{H})$ :

$$
\begin{equation*}
C_{\bullet}^{*}(S)=\left\{g(S): g \in C_{\bullet}(\tilde{\sigma}(S))\right\} \tag{1.14}
\end{equation*}
$$

where $\bullet$ is $b, c$, or 0 so that we have defined $C_{c}^{*}(S), C_{0}^{*}(S)$, and $C_{b}^{*}(S)$. Notice that $C(\sigma(S))$ consists of all $g \in C_{b}(\tilde{\sigma}(S))$ for which $\lim _{y \rightarrow \pm 1} g(y)$ exists, where this limit is 0 if and only if $g \in C_{0}(\tilde{\sigma}(S))$; hence, we have the inclusions (of which the first set implies the second)

$$
\begin{align*}
C_{c}(\tilde{\sigma}(S)) & \subseteq C_{0}(\tilde{\sigma}(S)) \subseteq C(\sigma(S)) \subseteq C_{b}(\tilde{\sigma}(S))  \tag{1.15}\\
C_{c}^{*}(S) & \subseteq C_{0}^{*}(S) \subseteq C^{*}(S) \subseteq C_{b}^{*}(S) \tag{1.16}
\end{align*}
$$

with equalities if and only if $T$ is bounded. This means that $g(S)$ is defined for $g \in C_{0}(\tilde{\sigma}(S))$, and hence a fortiori also for $g \in C_{c}(\tilde{\sigma}(S))$. Consequently, $f(T)$ may be defined by (1.8) whenever $f \in C_{0}(\sigma(T))$, including $f \in C_{c}(\sigma(T))$. To pass to the larger class $f \in C_{b}(\sigma(T))$, we define $C_{0}^{*}(S) \mathcal{H}$ as the linear span of all vectors of the form $g(S) \psi$, where $g \in C_{0}(\tilde{\sigma}(S))$ and $\psi \in \mathcal{H}$. Then $C_{0}^{*}(S) \mathcal{H}$ is dense in $\mathcal{H}$ (see Lemma 2.1). In the spirit of Woronowicz (see [5, Chapter 10], [12]), we then initially define $f(T)$ for $f \in C_{b}(\sigma(T))$ on the domain $C_{0}^{*}(S) \mathcal{H}$ by linear extension of the formula

$$
\begin{equation*}
f(T)_{0} h(T) \psi=(f h)(T) \psi, \tag{1.17}
\end{equation*}
$$

where $h \in C_{0}(\sigma(T))$, and hence also $f h \in C_{0}(\sigma(T))$ since $C_{b}(\sigma(T))$ is the mutiplier algebra of $C_{0}(\sigma(T))$. Then $f(T)_{0}$ is bounded (see Lemma 2.2), and we define $f(T)$ as its closure; that is,

$$
\begin{equation*}
f(T)=f(T)_{0}^{-} \tag{1.18}
\end{equation*}
$$

This also works for $f \in C(\sigma(T))$, in which case $f(T)_{0}$ may no longer be bounded, but remains closable (see Lemma 2.3) so that we may once again define $f(T)$ as its closure (cf. (1.18)). We have the following theorem (see also Theorem 1.4).

Theorem 1.2. If $f \in C(\sigma(T))$ is real-valued, then $f(T)$ is self-adjoint; that is, $f(T)_{0}^{-}=f(T)_{0}^{*}$, and, more generally, $f(T)^{*}=f^{*}(T)$. Furthermore, the continuous functional calculus $f \mapsto f(T)$ restricts to an isometric *-homomorphism from $C_{0}(\sigma(T))$ (with supremum-norm) to $C^{*}(S)$.

In addition, the map $f \mapsto f(T)$ has the reassuring special cases

$$
\begin{align*}
\mathbf{1}_{\sigma(T)}(T) & =I  \tag{1.19}\\
\operatorname{id}(T) & =T  \tag{1.20}\\
(\mathrm{id}-z)^{-1}(T) & =(T-z)^{-1}, \quad z \in \rho(T), \tag{1.21}
\end{align*}
$$

where $\mathbf{1}_{\sigma(T)}(x)=1$ and $\operatorname{id}(x)=x(x \in \sigma(T))$, and therefore it does what it is supposed to do.

Finding the right analogue of (1.10) for unbounded $T=T^{*}$ first requires a redefinition of $W^{*}(T)$, which is standard (see [8]). If $T$ is unbounded and $R \in$ $B(\mathcal{H})$, then we say that $R$ and $T$ commute, written $T R \subset R T$, if $R \psi \in \mathcal{D}(T)$ and $R T \psi=T R \psi$ for any $\psi \in \mathcal{D}(T)$. Let $\{T\}^{\prime}$ be the set of all bounded operators that commute with $T$. If $T^{*}=T$, then $\{T\}^{\prime}$ is a unital, strongly closed $*$-subalgebra of $B(\mathcal{H})$, and hence a von Neumann algebra (see [8]). Its commutant

$$
\begin{equation*}
W^{*}(T)=\{T\}^{\prime \prime} \tag{1.22}
\end{equation*}
$$

is a von Neumann algebra, too. If $T$ is bounded, then $W^{*}(T)$ is the von Neumann algebra generated by $T$, which coincides with $C^{*}(T)^{\prime \prime}$. As usual, we call a closed unbounded operator $X$ affiliated to a von Neumann algebra $A \subset B(H)$, written $X \eta A$, if and only if $X R \subset R X$ for each $R \in A^{\prime}$. For example, if $T^{*}=T$, then $T \eta W^{*}(T)$, and if $T \eta A$, then $W^{*}(T) \subseteq A$; in other words, $W^{*}(T)$ is the smallest von Neumann algebra such that $T$ is affiliated to it.

As a result of independent interest as well as a lemma for Theorem 1.4, we may then adapt [8, Lemma 5.2.8] to the bounded transform, as in this theorem.

Theorem 1.3. Let $A \subset B(H)$ be a von Neumann algebra. Then $T \eta A$ if and only if $S \in A$.

Denoting the (Banach) space of (bounded) Borel functions on $\sigma(T)$ (equipped with the supremum-norm) by $\mathcal{B}_{(b)}(\sigma(T))$, we may still define $f(T)$ by (1.8) and the usual Borel functional calculus for the bounded transform $S$.

Theorem 1.4. The map $f \mapsto f(T)$ is a norm-decreasing *-homomorphism from $\mathcal{B}_{b}(\sigma(T))$ to

$$
\begin{equation*}
W^{*}(T)=W^{*}(S) \tag{1.23}
\end{equation*}
$$

More generally, if $f \in \mathcal{B}(\sigma(T))$, then $f(T)$ is affiliated with $W^{*}(T)$.
The remainder of this paper simply consists of the proofs of these theorems.

## 2. Proofs

This section contains all proofs. We will not repeat the theorems.
2.1. Proof of Theorem 1.1. The operator $\sqrt{1-S^{2}}$ is a bijection from $\mathcal{H}$ to $\mathcal{R}\left(\sqrt{1-S^{2}}\right)=\mathcal{D}(T)$ (see [4], proof of Theorem 1). Let $\lambda \in \rho(T) \equiv \mathbb{C} \backslash \sigma(T)$ so that $T-\lambda I$ is a bijection from $\mathcal{D}(T)$ to $\mathcal{H}$. Thus, by composition, we have a bijection $\mathcal{H} \rightarrow \mathcal{H}$; equivalently, $(T-\lambda I)\left(\sqrt{I-S^{2}}\right)$ is invertible, which in turn is equivalent to invertibility of $S-\lambda \sqrt{I-S^{2}}$. Thus, $\lambda \in \rho(T) \Longleftrightarrow S-\lambda \sqrt{I-S^{2}}$
is a bijection, or, expressed contrapositively, $\lambda \in \sigma(T) \Longleftrightarrow S-\lambda \sqrt{I-S^{2}}$ is not invertible in $B(\mathcal{H})$. This is the case if and only if $S-\lambda \sqrt{I-S^{2}}$ is not invertible in $C^{*}(S)$, which, using the Gelfand isomorphism $C^{*}(S) \cong C(\sigma(S))$, in turn is true if and only if the function $k_{\lambda}(x)=x-\lambda \sqrt{1-x^{2}}$ is not invertible in $C(\sigma(S))$; that is, if and only if $0 \in \sigma\left(k_{\lambda}\right)$. Since in $C(X)$ we have $\sigma(f)=\mathcal{R}(f)$ (with $X$ a compact Hausdorff space), and since $\sigma(S)$ is indeed compact and Hausdorff because $S$ is bounded, we obtain $\lambda \in \sigma(T)$ if and only if $0 \in \mathcal{R}\left(k_{\lambda}\right)$. If $\pm 1$ lie in $\sigma(S)$, then they cannot give rise to $0 \in \mathcal{R}\left(k_{\lambda}\right)$ since $k_{\lambda}( \pm 1)= \pm 1$ for each $\lambda$; hence, $0 \in \mathcal{R}\left(k_{\lambda}\right)$ if and only if $\lambda=\mu\left(1-\mu^{2}\right)^{-1 / 2}$ for some $\mu \in \sigma(S) \cap(-1,1)$, which yields (1.11).

The same argument shows that $\mu \in \sigma(S) \cap(-1,1)$ comes from $\lambda \in \sigma(T)$. But since $\sigma(S)$ is compact and hence closed in $[-1,1]$, we obtain (1.12).
2.2. Proof of Theorem 1.2. This proof relies on three lemmas.

Lemma 2.1. Let $C_{c}^{*}(S) \mathcal{H}$ be the linear span of all vectors of the form $g(S) \psi$, where $g \in C_{c}(\tilde{\sigma}(S))$ and $\psi \in \mathcal{H}$. Then $C_{c}^{*}(S) \mathcal{H}$ is dense in $H$.
Proof. Define $g_{n}:(-1,1) \rightarrow[0,1]$ by putting $g_{n}(x)=0$ for $x \in\left(-1, \frac{1}{n}-1\right] \cup$ $\left[1-\frac{1}{n}, 1\right), g_{n}(x)=1$ if $x \in\left[\frac{2}{n}-1,1-\frac{2}{n}\right]$, and linear interpolation in between. The ensuing sequence converges pointwise to the unit 1 on $(-1,1)$. Restricting each $g_{n}$ to $\tilde{\sigma}(S)$, the continuous functional calculus gives $g_{n}(S) \rightarrow \mathbf{1}_{\tilde{\sigma}(S)}$ strongly. Therefore, for any $\psi \in \mathcal{H}$, we have a sequence $\psi_{n}=g_{n}(S) \psi$ in $C_{c}^{*}(S) \mathcal{H}$ such that $\psi_{n} \rightarrow \psi$.
Lemma 2.2. For $f \in C_{b}(\sigma(T))$, define an operator $f(T)_{0}$ on the domain $C_{0}^{*}(S) \mathcal{H}$ by (1.17). Then $f(T)_{0}$ is bounded with bound

$$
\begin{equation*}
\left\|f(T)_{0}\right\| \leq\|f\|_{\infty} \tag{2.1}
\end{equation*}
$$

Proof. Let $\varepsilon>0$. If $h \in C_{0}(\sigma(T))$, then $f h \in C_{0}(\sigma(T))$ so that we can find a compact subset $K \subset \sigma(T)$ such that $|h(x) f(x)|<\varepsilon$ for each $x \notin K$. Let $\tilde{h}=h \circ u$ (see (1.4)). Then $\tilde{h} \in C_{0}(\tilde{\sigma}(S))$ whenever $h \in C_{0}(\sigma(T))$; in fact, we have an isometric isomorphism

$$
\begin{equation*}
C_{0}(\sigma(T)) \xlongequal{\rightrightarrows} C_{0}(\tilde{\sigma}(S)), \quad h \mapsto h \circ u \tag{2.2}
\end{equation*}
$$

Contractivity of the Borel functional calculus for bounded operators on $\mathcal{H}$ gives

$$
\left\|\left(\widetilde{\mathbf{1}_{K^{c}} f h}\right)(S) \psi\right\| \leq\left\|\left(\widetilde{\left(\mathbf{1}_{K^{c}} f h\right.}\right)(S)\right\|\|\psi\| \leq\left\|\widetilde{\boldsymbol{1}_{K^{c}} f h}\right\|_{\infty}\|\psi\|<\varepsilon\|\psi\|
$$

Using also the homomorphism property of the Borel functional calculus, we then find that

$$
\begin{aligned}
\|(f h)(T) \psi\| & =\|(\widetilde{f h})(S) \psi\| \\
& \left.=\| \widetilde{\left(\mathbf{1}_{K} f h\right.}\right)(S)+\left(\widetilde{f h}-\widetilde{\mathbf{1}_{K} f h}\right)(S) \psi \| \\
& \leq\left\|\left(\widetilde{\left(\mathbf{1}_{K} f h\right.}\right)(S) \psi\right\|+\left\|\left(\widetilde{\mathbf{1}_{K^{c}} f h}\right)(S) \psi\right\| \\
& =\left\|\widetilde{\left(\mathbf{1}_{K} f\right)}(S) \tilde{h}(S) \psi\right\|+\left\|\left(\widetilde{\left(\mathbf{1}_{K^{c}} f h\right.}\right)(S) \psi\right\|
\end{aligned}
$$

$$
\begin{aligned}
& <\left\|\widetilde{\left(\mathbf{1}_{K} f\right)}\right\|_{\infty}\|h(T) \psi\|+\varepsilon\|\psi\| \\
& \leq\|f\|_{\infty}\|h(T) \psi\|+\varepsilon\|\psi\|
\end{aligned}
$$

since $\left\|\widetilde{\left(\mathbf{1}_{K} f\right)}\right\|_{\infty} \leq\|\tilde{f}\|_{\infty}=\|f\|_{\infty}$. Since the last expression above is independent of $K$, we may let $\varepsilon \rightarrow 0$, obtaining boundedness of $f(T)$ as well as (2.1).

The last claim in Theorem 1.2 now follows from the continuous functional calculus for $S$ and the isometric isomorphism (2.2). Although isometry may be lost if we go from $C_{0}(\sigma(T))$ to $C_{b}(\sigma(T))$, it easily follows from (1.17)-(1.18) that the map $f \mapsto f(T)$ at least defines a ${ }^{*}$-homomorphism $C_{b}(\sigma(T)) \rightarrow B(H)$. This property will be used after Lemma 2.4 below.

Lemma 2.3. For $f \in C(\sigma(T))$, define an operator $f(T)_{0}$ on the domain $C_{c}^{*}(S) \mathcal{H}$ by (1.17). Then $f(T)_{0}$ is closable. Moreover, if $f$ is real-valued $\left(f^{*}=f\right)$, then $f(T)_{0}$ is symmetric.

Proof. Suppose that $h_{1}(T) \psi_{1}$ and $h_{2}(T) \psi_{2}$ lie in $\mathcal{D}\left(f(T)_{0}\right)$. Then we may compute

$$
\begin{align*}
\left\langle h_{2}(T) \psi_{2}, f(T)_{0} h_{1}(T) \psi_{1}\right\rangle & =\left\langle\psi_{2}, h_{2}(T)^{*}\left(f h_{1}\right)(T) \psi_{1}\right\rangle \\
& =\left\langle\psi_{2},\left(\overline{h_{2}} f h_{1}\right)(T) \psi_{1}\right\rangle ;  \tag{2.3}\\
\left\langle\left(h_{2} \bar{f}\right)(T) \psi_{2}, h_{1}(T) \psi_{1}\right\rangle & \left.=\left\langle\psi_{2}, h_{2} \bar{f}\right)(T)^{*} h_{1}(T) \psi_{1}\right\rangle \\
& =\left\langle\psi_{2},\left(\overline{h_{2}} f h_{1}\right)(T) \psi_{1}\right\rangle, \tag{2.4}
\end{align*}
$$

where in the first equality in (2.3) we have $h_{2} \in C_{0}(\sigma(T))$ so that the operator $h_{2}(T)=h_{2} \circ u(S)$ is defined by (1.8), and hence is bounded (see Section 1). The continuous functional calculus for $S$ then gives $h_{2}(T)^{*}=\overline{h_{2}}(T)$ as well as $\overline{h_{2}}(T)\left(f h_{1}\right)(T)=\left(\overline{h_{2}} f h_{1}\right)(T)$, and similarly in (2.4).

This implies that $\mathcal{D}\left(f(T)_{0}\right) \subseteq \mathcal{D}\left(f(T)_{0}^{*}\right)$. Since $\mathcal{D}\left(f(T)_{0}\right)$ is dense, so is $\mathcal{D}\left(f(T)_{0}^{*}\right)$, which implies that $f(T)_{0}$ is closable. The second claim is obvious from (2.3)-(2.4).

Proof of Theorem 1.2. To prove Theorem 1.2, we use a well-known result of Nelson [6] (see also [9]) (this step was suggested to us by Nigel Higson). For convenience we recall this result (without proof).

Lemma 2.4. Let $\{U(t)\}_{t \in \mathbb{R}}$ be a strongly continuous unitary group of operators on a Hilbert space $\mathcal{H}$. Let $R: \mathcal{D}(R) \rightarrow \mathcal{H}$ be densely defined and symmetric. Assume that $\mathcal{D}(R)$ is invariant under $\{U(t)\}_{t \in \mathbb{R}}$; that is, assume that $U(t): \mathcal{D}(R) \rightarrow$ $\mathcal{D}(R)$ for each $t$, and also that $\{U(t)\}_{t \in \mathbb{R}}$ is strongly differentiable on $\mathcal{D}(R)$. Then -idU $(t) / d t$ is essentially self-adjoint on $\mathcal{D}(R)$, and its closure is the self-adjoint generator of $\{U(t)\}_{t \in \mathbb{R}}$ (given by Stone's theorem). In particular, if $(d U(t) / d t) \psi=$ $i R U(t) \psi$ for each $\psi \in \mathcal{D}(R)$, then $R$ is essentially self-adjoint.

Set $R=f(T)_{0}$ for $f \in C(\sigma(T))$ so that

$$
\begin{equation*}
\mathcal{D}(R)=C_{c}^{*}(S) \mathcal{H} \tag{2.5}
\end{equation*}
$$

and for each $t \in \mathbb{R}$ define $U(t)$ via the (bounded) function $x \mapsto \exp (i t f(x))$ on $\sigma(T)$; that is, for $h \in C_{c}(\sigma(T))$ and $\psi \in \mathcal{H}$, we initially define

$$
\begin{equation*}
U(t)_{0} h(T) \psi=\left(e^{i t f} h\right)(T) \psi \tag{2.6}
\end{equation*}
$$

Then $U(t)_{0}$ is bounded by Lemma 2.2, and we define $U(t)$ as the closure of $U(t)_{0}$. The remark before Lemma 2.3 then implies that $t \mapsto U(t)$ defines a unitary representation of $\mathbb{R}$ on $\mathcal{H}$. Strong continuity of this representation follows from an $\varepsilon / 3$ argument. First, for

$$
\begin{equation*}
\varphi=h(T) \psi \tag{2.7}
\end{equation*}
$$

and assuming that $\|\psi\|=1$ for simplicity, equations (2.6) and (2.1) give

$$
\begin{equation*}
\|U(t) \varphi-\varphi\| \leq\left\|e^{i t f} h-h\right\|_{\infty} \leq\|h\|_{\infty}\left\|e^{i t f}-\mathbf{1}\right\|_{\infty}^{(K)} \tag{2.8}
\end{equation*}
$$

where $K$ is the (compact) support of $h$ in $\sigma(T)$. Since the exponential function is uniformly convergent on any compact set, this gives $\lim _{t \rightarrow 0}\|U(t) \varphi-\varphi\|=0$ for $\varphi$ of the form (2.7); taking finite linear combinations thereof gives the same result for any $\varphi \in C_{c}^{*}(S) \mathcal{H}$. Thus, for any $\varepsilon>0$, we can find $\delta>0$ so that $\|U(t) \varphi-\varphi\|<\varepsilon / 3$ whenever $|t|<\delta$. For general $\psi^{\prime} \in H$, we find $\varphi \in C_{c}^{*}(S) H$ such that $\left\|\varphi-\psi^{\prime}\right\|<\varepsilon / 3$, and we estimate

$$
\begin{aligned}
\left\|U(t) \psi^{\prime}-\psi^{\prime}\right\| & \leq\left\|U(t) \psi^{\prime}-U(t) \varphi\right\|+\|U(t) \varphi-\varphi\|+\left\|\varphi-\psi^{\prime}\right\| \\
& \leq \varepsilon / 3+\varepsilon / 3+\varepsilon / 3=\varepsilon
\end{aligned}
$$

since $\left\|U(t) \psi^{\prime}-U(t) \varphi\right\|=\left\|\psi^{\prime}-\varphi\right\|$ by unitarity of $U(t)$. Thus, $\lim _{t \rightarrow 0} \| U(t) \psi-$ $\psi \|=0$ for any $\psi \in \mathcal{H}$ so that the unitary representation $t \mapsto U(t)$ is strongly continuous. Similarly,

$$
\begin{equation*}
\left\|\frac{U(t+s) \varphi-U(t) \varphi}{s}-i R U(t) \varphi\right\| \leq\left\|\frac{e^{i s f} h-h}{s}-i f h\right\|_{\infty} \tag{2.9}
\end{equation*}
$$

assuming (2.7), so that by the same argument as in (2.8) we obtain

$$
\begin{equation*}
\frac{d U(t)}{d t} \varphi=i R U(t) \varphi \tag{2.10}
\end{equation*}
$$

initially for any $\varphi$ of the form (2.7), and hence, taking finite sums, for any $\varphi \in$ $\mathcal{D}(R)$ (see (2.5)). The final part of Lemma 2.4 then shows that $f(T)_{0}$ is essentially self-adjoint on its domain $C_{c}^{*}(S) \mathcal{H}$. Its closure $f(T)$ is therefore self-adjoint, and Theorem 1.2 is proved.

We now prove the examples (1.19)-(1.21), of which the first is trivial. Writing $T_{0}$ for the operator $\operatorname{id}(T)_{0}$, the definition (1.17) gives

$$
T_{0} \varphi=T \varphi
$$

for $\varphi \in \mathcal{D}\left(T_{0}\right)=C_{c}^{*}(S) \mathcal{H}$. Let $\psi \in \mathcal{D}\left(T_{0}^{-}\right)$so that there is a sequence $\left(\varphi_{n}\right)$ in $\mathcal{D}\left(T_{0}\right)$ such that $\varphi_{n} \rightarrow \varphi$ and $\left(T_{0} \varphi_{n}\right)$ converges. Since $T$ is closed, it follows that $T_{0} \varphi_{n}=T \varphi_{n} \rightarrow T \varphi$ so that $\varphi \in \mathcal{D}(T)$; hence, $T_{0}^{-} \subset T$. Since both operators are self-adjoint, this implies that $T_{0}^{-}=T$, which proves (1.20).

The proof of (1.21) is easier since $(T-z)^{-1}$ is bounded: writing

$$
f(x)=(x-z)^{-1},
$$

where $z \notin \sigma(T)$ is fixed and $x \in \sigma(T)$, we have

$$
f(T)_{0} h(T) \psi=(f h)(T) \psi=(T-z)^{-1} h(T) \psi
$$

and hence

$$
f(T)_{0} \varphi=(T-z)^{-1} \varphi
$$

for any $\varphi \in \mathcal{D}\left(f(T)_{0}\right)=C_{c}^{*}(S) \mathcal{H}$. If $\varphi_{n} \rightarrow \varphi$ for $\varphi \in \mathcal{H}$ and $\varphi_{n} \in \mathcal{D}\left(f(T)_{0}\right)$, boundedness and hence continuity of the resolvent implies that

$$
f(T) \varphi=\lim _{n \rightarrow \infty} f(T)_{0} \varphi_{n}=\lim _{n \rightarrow \infty}(T-z)^{-1} \varphi_{n}=(T-z)^{-1} \varphi
$$

2.3. Proof of Theorem 1.3. The first step consists in the observation that $T \eta A$ if and only if $U T=T U$ (or, equivalently, $U T U^{*}=T$ ) for each unitary $U \in A^{\prime}$ [11, Proposition 5.3.4].

The second step is to show that $U T=T U$ if and only if $S U=U S$ for any unitary $U$. This is a simple computation. First, suppose that $U T U^{*}=T$. Then

$$
\begin{aligned}
U\left(1+T^{2}\right)^{-1} U^{*} & =\left(U\left(1+T^{2}\right) U^{*}\right)^{-1}=\left(\left(U+U T^{2}\right) U^{*}\right)^{-1} \\
& =\left(U U^{*}+U T^{2} U^{*}\right)^{-1}=\left(1+U T U^{*} U T U^{*}\right)^{-1} \\
& =\left(1+T^{2}\right)^{-1}
\end{aligned}
$$

If $R$ is bounded and positive, then $U R=R U$ if and only if $U \in C^{*}(R)^{\prime}$, and since $\sqrt{R} \in C^{*}(R)$ by the continuous functional calculus, we also have $U \sqrt{R}=\sqrt{R} U$. Consequently,

$$
\begin{aligned}
U S U^{*} & =U\left(T \sqrt{\left(1+T^{2}\right)^{-1}}\right) U^{*} \\
& =\left(U T U^{*}\right)\left(U \sqrt{\left(1+T^{2}\right)^{-1}} U^{*}\right)=T \sqrt{\left(1+T^{2}\right)^{-1}}=S
\end{aligned}
$$

Similarly, if $S U=U S$, then

$$
\begin{aligned}
U T U^{*} & =U S \sqrt{1-S^{2}} \\
& =S U \sqrt{1-S^{2}} U^{*} U^{*}=S\left(U \sqrt{1-S^{2}} U^{*}\right)^{-1}=S \sqrt{1-S^{2}}
\end{aligned}
$$

Third, as in the first step, $S U=U S$ for any unitary $U \in A^{\prime}$ if and only if $S \in A^{\prime \prime}=A$.
2.4. Proof of Theorem 1.4. Equation (1.23) in Theorem 1.4 follows from Theorem 1.3: taking $A=W^{*}(T)$ so that $T \eta A$ yields $S \in W^{*}(T)$, and hence $W^{*}(S) \subseteq W^{*}(T)$. On the other hand, taking $A=W^{*}(S)$, in which case $S \in A$, gives $T \eta W^{*}(S)$, and hence $W^{*}(T) \subseteq W^{*}(S)$.

Similarly to (2.2), we have an isometric isomorphism

$$
\begin{equation*}
\mathcal{B}_{b}(\sigma(T)) \xlongequal{\leftrightharpoons} \mathcal{B}_{b}(\tilde{\sigma}(S)), \quad h \mapsto h \circ u \tag{2.11}
\end{equation*}
$$

so that the first claim of Theorem 1.4 follows from the Borel functional calculus for the bounded operator $S$ (see [8]). The proof of the last one is, mutatis mutandis, practically the same as in [8, Theorem 5.3.8], so we omit the details (see [2]).

As explained in [8, Section 5.3], there exists a Borel measure $\mu$ on $\sigma(T)$ such that the map $f \mapsto f(T)$ may also be seen as a so-called essential ${ }^{*}$-homomorphism from $\mathcal{B}(\sigma(T)) / \mathcal{N}(\sigma(T))$ into the *-algebra of normal operators affiliated with $W^{*}(T)$,
where $\mathcal{N}(\sigma(T))$ is the set of $\mu$-null functions on $\sigma(T)$. This remains true in our approach with the same proof (see [2]).

## 3. Epilogue

Let us finally note that, although the present article was inspired by the work of Woronowicz, the $C^{*}$-algebraic affiliation relation he defines in [12, Definition 1.1] (as did, independently, also Baaj and Julg in [1]) has not been used here. If we call his relation $\eta^{\prime}$ to avoid confusion with the $W^{*}$-algebraic relation $\eta$ we do use, if $A \subset B(\mathcal{H})$, then we have $T \eta^{\prime} A \Rightarrow T \in A$ (and hence $T$ is bounded) (cf. [12, Proposition 1.3]). Woronowicz does not define a $C^{*}$-algebraic counterpart of the von Neumann algebra $W^{*}(T)$, but it might be reasonable to define $C^{*}(T)$ as the smallest $C^{*}$-algebra $A$ in $B(\mathcal{H})$ such that $T \eta^{\prime} A$. It follows from [12, Example 4] that this would give $C^{*}(T)=C_{0}^{*}(S)$ as defined in (1.14). This $C^{*}$-algebra contains $S$ (and hence $T$ ) if and only if $T$ is bounded, in which case $C_{0}^{*}(S)=C^{*}(S)$ and hence $C^{*}(T)=C^{*}(S)$, as in our approach (see (1.5)). Also, in general (i.e., if $T$ is possibly unbounded), the bicommutant $C^{*}(T)^{\prime \prime}$ coincides with $W^{*}(T)$ as defined in the usual way (1.22). This follows from $C_{0}^{*}(S)^{\prime \prime}=C^{*}(S)^{\prime \prime}=W^{*}(S)$ and (1.10).

Of course, we could also redefine $\eta^{\prime}$, now calling it $\eta^{\prime \prime}$, by stipulating that $T \eta^{\prime \prime} A$ whenever $S \in A$, and redefine $C^{*}(T)$ accordingly (i.e., as the smallest $C^{*}$-algebra $A$ in $B(\mathcal{H})$ such that $T \eta^{\prime \prime} A$ ). This would give (1.5) even if $T$ is unbounded, though in a somewhat empty way.

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Radboud University Nijmegen, Institute for Mathematics, Astrophysics and Particle Physics, Heyendaalseweg 135, 6525 AJ Nijmegen, The Netherlands.

E-mail address: Christian.Budde@vub.ac.be; landsman@math.ru.nl


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