# RULED SUBMANIFOLDS WITH HARMONIC GAUSS MAP 

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#### Abstract

Ruled submanifolds of Minkowski space with harmonic Gauss map are studied. Apart from ruled submanifolds in Euclidean space, ruled submanifolds with degenerate rulings in Minkowski space draw our attention. In particular, we completely classify ruled submanifolds with harmonic Gauss map and we also characterize minimal ruled submanifolds with degenerate rulings by means of harmonic Gauss map.


## 1. Introduction

In eighteenth century, the so-called minimal surfaces were introduced when the graph of a certain function minimizes the area among surfaces with the fixed boundary. Since then, the theory of minimal submanifolds has been one of the most interesting topics in differential geometry.

In 1966, T. Takahashi showed: Let $x: M \rightarrow \mathbb{E}^{m}$ be an isometric immersion of a Riemannian manifold $M$ into the Euclidean space $\mathbb{E}^{m}$ and $\Delta$ the Laplace operator defined on $M$. If $\Delta x=\lambda x(\lambda \neq 0)$ holds, then $M$ is a minimal submanifold in a hypersphere of Euclidean space ([17]). Extending this point of view, in the late 1970's B.-Y. Chen introduced the notion of finite type immersion of Riemannian manifolds into Euclidean space ([4, 5]). In particular, minimal submanifolds of Euclidean space can be considered as a spacial case of submanifolds of finite type or those with harmonic immersion. The notion of finite type immersion was extended to submanifolds in

[^0]pseudo-Euclidean space in 1980's: A pseudo-Riemannian submanifold $M$ of an $m$ dimensional pseudo-Euclidean space $\mathbb{E}_{s}^{m}$ with signature $(m-s, s)$ is said to be of finite type if its position vector field $x$ can be expressed as a finite sum of eigenvectors of the Laplacian $\Delta$ of $M$, that is, $x=x_{0}+\sum_{i=1}^{k} x_{i}$, where $x_{0}$ is a constant map, $x_{1}, \cdots, x_{k}$ non-constant maps such that $\Delta x_{i}=\lambda_{i} x_{i}, \lambda_{i} \in \mathbb{R}, i=1,2, \cdots, k([4,5])$.

Such a notion can be naturally extended to a smooth map defined on submanifolds of pseudo-Euclidean space. A smooth map $\phi$ on an $n$-dimensional pseudo-Riemannian submanifold $M$ of $\mathbb{E}_{s}^{m}$ is said to be of finite type if $\phi$ is a finite sum of $\mathbb{E}_{s}^{m}$-valued eigenfunctions of $\Delta$. We also similarly define a smooth map of $k$-type on $M$ as that of immersion $x$. A very typical and interesting smooth map on the submanifold $M$ of Euclidean space or pseudo-Euclidean space is the Gauss map. In particular, we say that a differentiable map $\phi$ is harmonic if $\Delta \phi=0$.

A ruled surface is one of the most natural geometric objects in the classical differential geometry and has been dealt with some geometric conditions ([1, 2, 6, 7, 11, 12, 13, 14, 15, 16]). Due to Beltrami equation, the submanifolds of Euclidean space or Minkowski space with harmonic immersion are the minimal ones.

We now have a question: Can we completely classify ruled submanifolds in Minkowski space with harmonic Gauss map?

In this article, we study ruled submanifolds in the Minkowski space $\mathbb{L}^{m}$ with harmonic Gauss map and we characterize minimal ruled submanifolds with degenerate rulings by means of harmonic Gauss map.

All of geometric objects under consideration are smooth and submanifolds are assumed to be connected unless otherwise stated.

## 2. Preliminaries

Let $\mathbb{E}_{s}^{m}$ be an $m$-dimensional pseudo-Euclidean space of signature $(m-s, s)$. In particular, for $m \geq 2, \mathbb{E}_{1}^{m}$ is called a Lorentz-Monkowski $m$-space or simply Minkowski $m$-space, which is denoted by $\mathbb{L}^{m}$. A curve in $\mathbb{L}^{m}$ is said to be space-like, time-like or null if its tangent vector field is space-like, time-like or null, respectively. Let $x: M \rightarrow \mathbb{E}_{s}^{m}$ be an isometric immersion of an $n$-dimensional pseudo-Riemannian manifold $M$ into $\mathbb{E}_{s}^{m}$. From now on, a submanifold in $\mathbb{E}_{s}^{m}$ always means pseudoRiemannian, that is, each tangent space of the submanifold is non-degenerate.

Let $\left(x_{1}, x_{2}, \cdots, x_{n}\right)$ be a local coordinate system of $M$ in $\mathbb{E}_{s}^{m}$. For the components $g_{i j}$ of the pseudo-Riemannian metric $\langle\cdot, \cdot\rangle$ on $M$ induced from that of $\mathbb{E}_{s}^{m}$, we denote by $\left(g^{i j}\right)$ (respectively, $\mathcal{G}$ ) the inverse matrix (respectively, the determinant) of the matrix $\left(g_{i j}\right)$. Then, the Laplacian $\Delta$ on $M$ is given by

$$
\Delta=-\frac{1}{\sqrt{|\mathcal{G}|}} \sum_{i, j} \frac{\partial}{\partial x_{i}}\left(\sqrt{|\mathcal{G}|} g^{i j} \frac{\partial}{\partial x_{j}}\right)
$$

We now choose an adapted local orthonormal frame $\left\{e_{1}, e_{2}, \cdots, e_{m}\right\}$ in $\mathbb{E}_{s}^{m}$ such that $e_{1}, e_{2}, \cdots, e_{n}$ are tangent to $M$ and $e_{n+1}, e_{n+2}, \cdots, e_{m}$ normal to $M$. The Gauss map $G: M \rightarrow G(n, m) \subset \mathbb{E}^{N}\left(N={ }_{m} C_{n}\right), G(p)=\left(e_{1} \wedge e_{2} \wedge \cdots \wedge e_{n}\right)(p)$, of $x$ is a smooth map which carries a point $p$ in $M$ to an oriented $n$-plane in $\mathbb{E}_{s}^{m}$ which is obtained from the parallel translation of the tangent space of $M$ at $p$ to an $n$ plane passing through the origin in $\mathbb{E}_{s}^{m}$, where $G(n, m)$ is the Grassmannian manifold consisting of all oriented $n$-planes through the origin of $\mathbb{E}_{s}^{m}$.

An indefinite scalar product $<\cdot \cdot \cdot \gg$ on $G(n, m) \subset \mathbb{E}^{N}$ is defined by

$$
\ll e_{i_{1}} \wedge \cdots \wedge e_{i_{n}}, e_{j_{1}} \wedge \cdots \wedge e_{j_{n}} \gg=\operatorname{det}\left(\left\langle e_{i_{l}}, e_{j_{k}}\right\rangle\right) .
$$

Then, $\left\{e_{i_{1}} \wedge e_{i_{2}} \wedge \cdots \wedge e_{i_{n}} \mid 1 \leq i_{1}<\cdots<i_{n} \leq m\right\}$ is an orthonormal basis of $\mathbb{E}_{k}^{N}$ for some positive integer $k$.

Now, we define a ruled submanifold $M$ in $\mathbb{L}^{m}$. A non-degenerate $(r+1)$ dimensional submanifold $M$ in $\mathbb{L}^{m}$ is called a ruled submanifold if $M$ is foliated by $r$-dimensional totally geodesic submanifolds $E(s, r)$ of $\mathbb{L}^{m}$ along a regular curve $\alpha=\alpha(s)$ on $M$ defined on an open interval $I$. Thus, a parametrization of a ruled submanifold $M$ in $\mathbb{L}^{m}$ can be given by

$$
x=x\left(s, t_{1}, t_{2}, \cdots, t_{r}\right)=\alpha(s)+\sum_{i=1}^{r} t_{i} e_{i}(s), \quad s \in I, t_{i} \in I_{i},
$$

where $I_{i}$ 's are some open intervals for $i=1,2, \cdots, r$. For each $s, E(s, r)$ is open in $\operatorname{Span}\left\{e_{1}(s), e_{2}(s), \cdots, e_{r}(s)\right\}$, which is the linear span of linearly independent vector fields $e_{1}(s), e_{2}(s), \cdots, e_{r}(s)$ along the curve $\alpha$. Here we assume $E(s, r)$ are either non-degenerate or degenerate for all $s$ along $\alpha$. We call $E(s, r)$ the rulings and $\alpha$ the base curve of the ruled submanifold $M$. In particular, the ruled submanifold $M$ is said to be cylindrical if $E(s, r)$ is parallel along $\alpha$, or non-cylindrical otherwise.

Remark 2.1. ([9]).
(1) If the rulings of $M$ are non-degenerate, then the base curve $\alpha$ can be chosen to be orthogonal to the rulings as follows: Let $V$ be a unit vector field on $M$ which is orthogonal to the rulings. Then $\alpha$ can be taken as an integral curve of $V$.
(2) If the rulings are degenerate, we can choose a null base curve which is transversal to the rulings: Let $V$ be a null vector field on $M$ which is not tangent to the rulings. An integral curve of $V$ can be the base curve.

By solving a system of ordinary differential equations similarly set up in relation to a frame along a curve in $\mathbb{L}^{m}$ as given in [3], we have

Lemma 2.2. ([10]). Let $V(s)$ be a smooth $l$-dimensional non-degenerate distribution in the Minkowski $m$-space $\mathbb{L}^{m}$ along a curve $\alpha=\alpha(s)$, where $l \geq 2$ and $m \geq 3$.

Then, we can choose orthonormal vector fields $e_{1}(s), \cdots, e_{m-l}(s)$ along $\alpha$ which generate the orthogonal complement $V^{\perp}(s)$ satisfying $e_{i}^{\prime}(s) \in V(s)$ for $1 \leq i \leq m-l$.

## 3. Non-degenerate Rulings

Let $M$ be an $(r+1)$-dimensional ruled submanifold in $\mathbb{L}^{m}$ generated by nondegenerate rulings. By Remark 2.1, the base curve $\alpha$ can be chosen to be orthogonal to the rulings. Without loss of generality, we may assume that $\alpha$ is a unit speed curve, that is, $\left\langle\alpha^{\prime}(s), \alpha^{\prime}(s)\right\rangle=\varepsilon(= \pm 1)$. From now on, the prime ' denotes $d / d s$ unless otherwise stated. By Lemma 2.2, we may choose orthonormal vector fields $e_{1}(s), \cdots, e_{r}(s)$ along $\alpha$ satisfying

$$
\begin{equation*}
\left\langle\alpha^{\prime}(s), e_{i}(s)\right\rangle=0,\left\langle e_{i}^{\prime}(s), e_{j}(s)\right\rangle=0, \quad i, j=1,2, \cdots, r . \tag{3.1}
\end{equation*}
$$

A parametrization of $M$ is given by

$$
\begin{equation*}
x=x\left(s, t_{1}, t_{2}, \cdots, t_{r}\right)=\alpha(s)+\sum_{i=1}^{r} t_{i} e_{i}(s) . \tag{3.2}
\end{equation*}
$$

In this section, we always assume that the parametrization (3.2) satisfies the condition (3.1). Then, $M$ has the Gauss map

$$
G=\frac{1}{\left\|x_{s}\right\|} x_{s} \wedge x_{t_{1}} \wedge \cdots \wedge x_{t_{r}}
$$

or, equivalently

$$
\begin{equation*}
G=\frac{1}{|q|^{1 / 2}}\left(\Phi+\sum_{i=1}^{r} t_{i} \Psi_{i}\right) \tag{3.3}
\end{equation*}
$$

where $q$ is the function of $s, t_{1}, t_{2}, \cdots, t_{r}$ defined by

$$
q=\left\langle x_{s}, x_{s}\right\rangle, \quad \Phi=\alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r} \quad \text { and } \quad \Psi_{i}=e_{i}^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r}
$$

Now, we separate the cases into two typical types of ruled submanifolds which are cylindrical or non-cylindrical.

Theorem 3.1. The cylindrical ruled submanifolds in $\mathbb{L}^{m}$ generated by nondegenerate rulings have harmonic Gauss map if and only if $M$ is part of an $(r+1)$ plane or a cylinder over the base curve $\alpha(s)$ which is a plane curve in a degenerate plane given by $\alpha(s)=s^{2} \boldsymbol{C}+s \boldsymbol{D}$ for some constant null vector $\boldsymbol{C}$ and a constant space-like unit vector $\boldsymbol{D}$ satisfying $\langle\boldsymbol{C}, \boldsymbol{D}\rangle=0$.

Proof. Let $M$ be a cylindrical $(r+1)$-dimensional ruled submanifold in $\mathbb{L}^{m}$ generated by non-degenerate rulings, which is parameterized by (3.2). We may assume that $e_{1}, e_{2}, \cdots, e_{r}$ generating the rulings are constant vectors.

The Laplacian $\Delta$ of $M$ is then naturally expressed by

$$
\Delta=-\varepsilon \frac{\partial^{2}}{\partial s^{2}}-\sum_{i=1}^{r} \varepsilon_{i} \frac{\partial^{2}}{\partial t_{i}^{2}}
$$

where $\varepsilon_{i}=\left\langle e_{i}(s), e_{i}(s)\right\rangle= \pm 1$ and the Gauss map $G$ of $M$ is given by

$$
G=\alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r}
$$

If we denote by $\Delta^{\prime}$ the Laplacian of $\alpha$, that is $\Delta^{\prime}=-\varepsilon \frac{\partial^{2}}{\partial s^{2}}$, we have the Laplacian $\Delta G$ of the Gauss map

$$
\begin{equation*}
\Delta G=\Delta^{\prime} \alpha^{\prime} \wedge e_{1} \wedge \cdots \wedge e_{r} \tag{3.4}
\end{equation*}
$$

We now suppose that the Gauss map $G$ is harmonic, that is $\Delta G=0$. From (3.4), we have

$$
\Delta^{\prime} \alpha^{\prime}=0
$$

The converse is straightforward.
We need the following lemmas for later use.
Lemma 3.2. Let $M$ be an $(r+1)$-dimensional non-cylindrical ruled submanifold parameterized by (3.2) in $\mathbb{L}^{m}$. Suppose that $e_{1}^{\prime}, e_{2}^{\prime}, \cdots, e_{r}^{\prime}$ are non-null and some of generators of rulings $e_{1}, \cdots, e_{k}$ are constant vector fields along $\alpha$. Then we have the Laplacian

$$
\Delta=\frac{1}{2 q^{2}} \frac{\partial q}{\partial s} \frac{\partial}{\partial s}-\frac{1}{q} \frac{\partial^{2}}{\partial s^{2}}-\frac{1}{2 q} \sum_{i=k+1}^{r} \varepsilon_{i} \frac{\partial q}{\partial t_{i}} \frac{\partial}{\partial t_{i}}-\sum_{i=1}^{r} \varepsilon_{i} \frac{\partial^{2}}{\partial t_{i}^{2}}
$$

Proof. The isometric immersion $x$ of $M$ can be put

$$
x\left(s, t_{1}, \cdots, t_{r}\right)=\alpha(s)+\sum_{i=1}^{k} t_{i} e_{i}(s)+\sum_{j=k+1}^{r} t_{j} e_{j}(s) .
$$

Then, we have

$$
x_{s}=\alpha^{\prime}(s)+\sum_{j=k+1}^{r} t_{j} e_{j}^{\prime}(s), \quad x_{t_{i}}=e_{i}(s)
$$

for $i=1,2, \cdots, r$. As we introduced in the beginning of this section, the function $q$ is given by

$$
\begin{equation*}
q=\left\langle x_{s}, x_{s}\right\rangle=\varepsilon+\sum_{i=k+1}^{r} 2 u_{i} t_{i}+\sum_{i, j=k+1}^{r} w_{i j} t_{i} t_{j} \tag{3.5}
\end{equation*}
$$

where $u_{i}=\left\langle\alpha^{\prime}, e_{i}^{\prime}\right\rangle, w_{i j}=\left\langle e_{i}^{\prime}, e_{j}^{\prime}\right\rangle, i, j=k+1, \cdots, r$. Note that $q$ is a polynomial in $t=\left(t_{k+1}, \cdots, t_{r}\right)$ with functions in $s$ as coefficients.

Then, the Laplacian $\Delta$ is easily obtained by

$$
\Delta=\frac{1}{2 q^{2}} \frac{\partial q}{\partial s} \frac{\partial}{\partial s}-\frac{1}{q} \frac{\partial^{2}}{\partial s^{2}}-\frac{1}{2 q} \sum_{i=k+1}^{r} \varepsilon_{i} \frac{\partial q}{\partial t_{i}} \frac{\partial}{\partial t_{i}}-\sum_{i=1}^{r} \varepsilon_{i} \frac{\partial^{2}}{\partial t_{i}^{2}}
$$

From now on, for a polynomial $F(t)$ in $t=\left(t_{1}, t_{2}, \cdots, t_{r}\right)$, $\operatorname{deg} F(t)$ denotes the degree of $F(t)$ in $t=\left(t_{1}, t_{2}, \cdots, t_{r}\right)$ unless otherwise stated.

By Lemma 3.2, $\Delta G=0$ is rewritten as

$$
\begin{array}{r}
\left(\frac{\partial q}{\partial s}\right)^{2}\left(\Phi+\sum_{j=k+1}^{r} \Psi_{j} t_{j}\right)-\frac{3}{2} q \frac{\partial q}{\partial s}\left(\Phi^{\prime}+\sum_{j=k+1}^{r} \Psi_{j}^{\prime} t_{j}\right)-\frac{1}{2} q \frac{\partial^{2} q}{\partial s^{2}}\left(\Phi+\sum_{j=k+1}^{r} \Psi_{j} t_{j}\right) \\
+q^{2}\left(\Phi^{\prime \prime}+\sum_{j=k+1}^{r} \Psi_{j}^{\prime \prime} t_{j}\right)+\frac{1}{2} q \sum_{i=k+1}^{r} \varepsilon_{i}\left(\frac{\partial q}{\partial t_{i}}\right)^{2}\left(\Phi+\sum_{j=k+1}^{r} \Psi_{j} t_{j}\right)  \tag{3.6}\\
-\frac{1}{2} q^{2} \sum_{i=k+1}^{r} \varepsilon_{i} \frac{\partial q}{\partial t_{i}} \Psi_{i}-\frac{1}{2} q^{2} \sum_{i=k+1}^{r} \varepsilon_{i} \frac{\partial^{2} q}{\partial t_{i}^{2}}\left(\Phi+\sum_{j=k+1}^{r} \Psi_{j} t_{j}\right)=0
\end{array}
$$

To deal with (3.6), we have two possible cases either $\frac{\partial q}{\partial s} \neq 0$ or $\frac{\partial q}{\partial s}=0$ on some open interval.

Lemma 3.3. Let $M$ be an $(r+1)$-dimensional non-cylindrical ruled submanifold parameterized by (3.2) in $\mathbb{L}^{m}$ with harmonic Gauss map. Let $e_{1}, e_{2}, \cdots, e_{r}$ be orthonormal generators of the rulings along $\alpha$ generating the rulings. If $e_{i}^{\prime}$ are non-null for $i=1,2, \cdots, r$ and some of generators of the rulings $e_{1}, \cdots, e_{k}$ are constant vector fields along $\alpha$, then we have

$$
e_{i}^{\prime}=\varepsilon u_{i} \alpha^{\prime}
$$

Proof. We will prove this according to the following steps.
Step 1. In this step, we show that $w_{i j}=\varepsilon u_{i} u_{j}$ for $i, j=k+1, \cdots, r$.
Case 1. Suppose that $\frac{\partial q}{\partial s} \neq 0$. We may assume $\frac{\partial q}{\partial s} \neq 0$ on an open interval $\mathcal{I}$ for this case. Then, each term of the left side of (3.6) involves $\left(\frac{\partial q}{\partial s}\right)^{2}$ or $q$ and thus we have

$$
\begin{equation*}
\left(\frac{\partial q}{\partial s}\right)^{2}=q(t) P(t) \tag{3.7}
\end{equation*}
$$

for some polynomial $P(t)$ in $t$ of degree 2 with functions in $s$ as coefficients. Comparing the both sides of (3.7), we can get

$$
P(t)=\sum_{i, j} \theta_{i j} t_{i} t_{j}
$$

for some symmetric functions $\theta_{i j}$ of $s$. Together with this equation and (3.7), we have

$$
\begin{equation*}
\varepsilon \theta_{i j}=4 u_{i}^{\prime} u_{j}^{\prime}, \tag{3.8}
\end{equation*}
$$

$$
\begin{equation*}
u_{i} \theta_{j h}+u_{j} \theta_{h i}+u_{h} \theta_{i j}=2\left(u_{i}^{\prime} w_{j h}^{\prime}+u_{j}^{\prime} w_{h i}^{\prime}+u_{h}^{\prime} w_{i j}^{\prime}\right) \tag{3.9}
\end{equation*}
$$

(3.10) $2\left(w_{i j}^{\prime} w_{h l}^{\prime}+w_{i h}^{\prime} w_{j l}^{\prime}+w_{l i}^{\prime} w_{j h}^{\prime}\right)=w_{i j} \theta_{h l}+w_{i h} \theta_{j l}+w_{j l} \theta_{i h}+w_{l i} \theta_{j h}+w_{h l} \theta_{i j}+w_{j h} \theta_{l i}$
for $i, j, h, l=k+1, \cdots, r$. From (3.8) and (3.10), we see that $\varepsilon=1$.
If $u_{i}^{\prime}=0$ for all $i=k+1, \cdots, r$, then $\left(\frac{\partial q}{\partial s}\right)=0$, which is a contradiction. Therefore, there exists $i_{0} \in\{k+1, \cdots, r\}$ such that $u_{i_{0}}^{\prime} \neq 0$. And, (3.10) yields $\left(w_{i_{0} i_{0}}^{\prime}\right)^{2}=4 w_{i_{0} i_{0}}\left(u_{i_{0}}^{\prime}\right)^{2}$. Also, (3.9) implies $\left(w_{i_{0} i_{0}}^{\prime}\right)^{2}=4 u_{i_{0}}^{2}\left(u_{i_{0}}^{\prime}\right)^{2}$. Thus we have

$$
\begin{equation*}
w_{i_{0} i_{0}}=u_{i_{0}}^{2} \tag{3.11}
\end{equation*}
$$

for $i_{0}$ satisfying $u_{i_{0}}^{\prime} \neq 0$.
By replacing $h, l$ with $i, j$, respectively in (3.9) and (3.10), we get

$$
\begin{gather*}
2 u_{i} \theta_{i j}+u_{j} \theta_{i i}=4 u_{i}^{\prime} w_{i j}^{\prime}+2 u_{j}^{\prime} w_{i i}^{\prime},  \tag{3.12}\\
4 w_{i j}^{\prime} w_{i j}^{\prime}+2 w_{i i}^{\prime} w_{j j}^{\prime}=4 w_{i j} \theta_{i j}+w_{i i} \theta_{j j}+w_{j j} \theta_{i i} . \tag{3.13}
\end{gather*}
$$

If $u_{i_{0}}^{\prime} \neq 0$ and $u_{j_{0}}^{\prime} \neq 0$ for some $i_{0}$ and $j_{0}$, then equation (3.12) with the aid of (3.8) and (3.11) implies

$$
\begin{equation*}
w_{i_{0} j_{0}}^{\prime}=u_{i_{0}} u_{j_{0}}^{\prime}+u_{i_{0}}^{\prime} u_{j_{0}} \tag{3.14}
\end{equation*}
$$

Substituting (3.14) into (3.13), we have

$$
\begin{equation*}
w_{i_{0} j_{0}}=u_{i_{0}} u_{j_{0}} \tag{3.15}
\end{equation*}
$$

for $i_{0}$ and $j_{0}$.
Suppose that $u_{i}^{\prime}\left(s_{0}\right)=0$ at some $s_{0}$ for $i=k+1, \cdots, r$. Let

$$
\Lambda_{0}=\left\{i \mid u_{i}^{\prime}\left(s_{0}\right)=0, k+1 \leq i \leq r\right\} .
$$

Then, at $s_{0}$, equation (3.7) can be written as

$$
\begin{align*}
& 4 \sum_{i, j \notin \Lambda_{0}} u_{i}^{\prime} u_{j}^{\prime} t_{i} t_{j}+4 \sum_{i, j, h \notin \Lambda_{0}} u_{i}^{\prime} w_{j h}^{\prime} t_{i} t_{j} t_{h}+\sum_{i, j, h, l \notin \Lambda_{0}} w_{i j}^{\prime} w_{h l}^{\prime} t_{i} t_{j} t_{h} t_{l} \\
& =\left(1+2 \sum_{i=k+1}^{r} u_{i} t_{i}+\sum_{i, j=k+1}^{r} w_{i j} t_{i} t_{j}\right)\left(\sum_{i, j \notin \Lambda_{0}} \theta_{i j} t_{i} t_{j}\right), \tag{3.16}
\end{align*}
$$

from which,

$$
\begin{equation*}
u_{i}\left(s_{0}\right)=0 \tag{3.17}
\end{equation*}
$$

and

$$
\begin{equation*}
w_{i j}\left(s_{0}\right)=0 \tag{3.18}
\end{equation*}
$$

for $i \in \Lambda_{0}$ and $j=k+1, \cdots, r$.
Therefore, by (3.15), (3.17) and (3.18), we have

$$
\begin{equation*}
w_{i j}(s)=u_{i}(s) u_{j}(s) \tag{3.19}
\end{equation*}
$$

for all $i, j=k+1, \cdots, r$ and for all $s$.
Case 2. Suppose $\frac{\partial q}{\partial s}=0$ on some open interval $U$.
On $U$, all functions $u_{i}$ and $w_{i j}$ are constants. So, from (3.6) we also have

$$
\begin{equation*}
\sum_{i=k+1}^{r} \varepsilon_{i}\left(\frac{\partial q}{\partial t_{i}}\right)^{2}=q(t) P_{1}(t) \tag{3.20}
\end{equation*}
$$

where $P_{1}(t)$ is a function of $t$.
Suppose that there exist $j_{1}, \cdots, j_{l} \in\{k+1, \cdots, r\}$ such that $\left(\frac{\partial q}{\partial t_{j_{i}}}\right)^{2}$ are not the multiples of $q(t)$ for $i=1, \cdots, l$. Because of (3.20), we get

$$
\begin{equation*}
\sum_{i=1}^{l} \varepsilon_{j_{i}}\left(\frac{\partial q}{\partial t_{j_{i}}}\right)^{2}=q(t) P_{2}(t) \tag{3.21}
\end{equation*}
$$

for some function $P_{2}(t)$. Since all of $\left(\frac{\partial q}{\partial t_{j_{1}}}\right)^{2}, \cdots,\left(\frac{\partial q}{\partial t_{j_{l}}}\right)^{2}$ are not the multiples of $q(t)$,

$$
\left(\frac{\partial q}{\partial t_{j_{i}}}\right)^{2}=a_{j_{i}} q(t)+r_{j_{i}}(t)
$$

for some constants $a_{j_{i}}$ and polynomials $r_{j_{i}}(t)$ in $t$ with $\operatorname{deg} r_{j_{i}}(t) \leq 1$ for $i=1, \cdots, l$. Then, $\sum_{i=1}^{l} \varepsilon_{j_{i}} r_{j_{i}}(t)$ must be a multiple of $q(t)$ because of (3.21), which is a contradiction. Thus, we have

$$
\left(\frac{\partial q}{\partial t_{i}}\right)^{2}=4 \varepsilon u_{i}^{2} q(t)
$$

for all $i=k+1, \cdots, r$. If we compare the both sides of the above equation, we easily see that

$$
\begin{equation*}
w_{i j}=\varepsilon u_{i} u_{j} \tag{3.22}
\end{equation*}
$$

for all $i, j=k+1, \cdots, r$.
Step 2. In this step, we prove that $e_{i}^{\prime}=\varepsilon u_{i} \alpha^{\prime}$ for $i=k+1, \cdots, r$.
Equations (3.19) and (3.22) imply

$$
\begin{equation*}
\left\langle e_{i}^{\prime}-\varepsilon u_{i} \alpha^{\prime}, e_{i}^{\prime}-\varepsilon u_{i} \alpha^{\prime}\right\rangle=0 \tag{3.23}
\end{equation*}
$$

for all $i=k+1, \cdots, r$.
Suppose that $M$ is Lorentzian. Then, the normal space of $M$ at each point is space-like. By (3.23), we see that the normal components of $e_{i}^{\prime}$ vanish and thus we get

$$
\begin{equation*}
e_{i}^{\prime}=\varepsilon u_{i} \alpha^{\prime} \tag{3.24}
\end{equation*}
$$

for all $i=k+1, \cdots, r$.
We now assume that $M$ is space-like. Then, $\varepsilon=1$. Since $w_{i i}=u_{i}^{2}$ for $i=$ $k+1, \cdots, r$, we can put

$$
e_{i}^{\prime}=u_{i} \alpha^{\prime}+\sum_{a=r+1}^{m-1} \lambda_{a}^{i} e_{a}
$$

where $\sum_{a=r+1}^{m-1} \lambda_{a}^{i} e_{a}$ is vanishing or a null vector field along $\alpha$.
Suppose $\sum_{a=r+1}^{m-1} \lambda_{a}^{i} e_{a}$ is a null vector field along $\alpha$ for some $i=k+1, \cdots, r$. By the hypothesis, $u_{i}$ is non-zero. In case of $u_{i}^{\prime}\left(s_{0}\right)=0$ at some $s_{0}$, if we follow the argument developed in Case 1 of Step 1 above, we see that $u_{i}\left(s_{0}\right)=0$, a contradiction. Therefore, $u_{i}^{\prime} \neq 0$ for all $s$. Then, we get

$$
\begin{aligned}
q=\left\langle x_{s}, x_{s}\right\rangle & =\left(1+\sum_{i=k+1}^{r} t_{i} u_{i}\right)^{2}+\sum_{a=r+1}^{m-1} \varepsilon_{a}\left(\sum_{i=k+1}^{r} \lambda_{a}^{i} t_{i}\right)^{2} \\
& =\left(1+\sum_{i=k+1}^{r} t_{i} u_{i}\right)^{2}
\end{aligned}
$$

Without loss of generality, we may assume that $1+\sum_{i=k+1}^{r} t_{i} u_{i}>0$. Hence we may put

$$
G=\Phi+\frac{1}{\tilde{q}} \sum_{a=r+1}^{m-1}\left(\sum_{i=k+1}^{r} \lambda_{a}^{i} t_{i}\right) \xi_{a}
$$

where $\xi_{a}=e_{a} \wedge e_{1} \wedge e_{2} \wedge \cdots \wedge e_{r}$ for $a=r+1, \cdots, m-1$ and $\tilde{q}^{2}=q$. By straightforward computation we have the Laplacian $\Delta G$ of the Gauss map

$$
\Delta G=\frac{\sum_{j=k+1}^{r} u_{j}^{\prime} t_{j}}{\tilde{q}^{3}} \Phi^{\prime}-\frac{1}{\tilde{q}^{2}} \Phi^{\prime \prime}+\left\{\frac{\sum_{j=k+1}^{r} u_{j}^{\prime \prime} t_{j}}{\tilde{q}^{4}}-\frac{3\left(\sum_{j=k+1}^{r} u_{j}^{\prime} t_{j}\right)^{2}}{\tilde{q}^{5}}\right\} \sum_{a=r+1}^{m-1}\left(\sum_{i=k+1}^{r} \lambda_{a}^{i} t_{i}\right) \xi_{a}
$$

$$
\begin{aligned}
& \quad 3 \sum_{j=k+1}^{r} u_{j}^{\prime} t_{j} \\
& +\frac{\tilde{q}^{4}}{m-1}\left(\sum_{a=r+1}^{r}\left(\lambda_{a}^{i}\right)^{\prime} t_{i}\right) \xi_{a}+\frac{3 \sum_{j=k+1}^{r} u_{j}^{\prime} t_{j}}{\tilde{q}^{4}} \sum_{a=r+1}^{m-1}\left(\sum_{i=k+1}^{r} \lambda_{a}^{i} t_{i}\right) \xi_{a}^{\prime} \\
& -\frac{1}{\tilde{q}^{3}} \sum_{a=r+1}^{m-1}\left(\sum_{i=k+1}^{r}\left(\lambda_{a}^{i}\right)^{\prime \prime} t_{i}\right) \xi_{a}-\frac{2}{\tilde{q}^{3}} \sum_{a=r+1}^{m-1}\left(\sum_{i=k+1}^{r}\left(\lambda_{a}^{i}\right)^{\prime} t_{i}\right) \xi_{a}^{\prime} \\
& -\frac{1}{\tilde{q}^{3}} \sum_{a=r+1}^{m-1}\left(\sum_{i=k+1}^{r} \lambda_{a}^{i} t_{i}\right) \xi_{a}^{\prime \prime}-\frac{1}{\tilde{q}^{3}} \sum_{i=k+1}^{r} u_{i}^{2}\left(\sum_{a=r+1}^{m-1}\left(\sum_{h=k+1}^{r} \lambda_{a}^{h} t_{h}\right) \xi_{a}\right) \\
& +\frac{1}{\tilde{q}^{2}} \sum_{i=k+1}^{r} u_{i}\left(\sum_{a=r+1}^{m-1} \lambda_{a}^{i} \xi_{a}\right) .
\end{aligned}
$$

Since the Gauss map is harmonic, $\Delta G=0$ and thus

$$
\begin{align*}
0= & \tilde{q}^{2}\left(\sum_{j=k+1}^{r} u_{j}^{\prime} t_{j}\right) \Phi^{\prime}-\tilde{q}^{3} \Phi^{\prime \prime} \\
& +\left\{\tilde{q}\left(\sum_{j=k+1}^{r} u_{j}^{\prime \prime} t_{j}\right)-3\left(\sum_{j=k+1}^{r} u_{j}^{\prime} t_{j}\right)^{2}\right\} \sum_{a=r+1}^{m-1}\left(\sum_{i=k+1}^{r} \lambda_{a}^{i} t_{i}\right) \xi_{a} \\
& +3 \tilde{q}\left(\sum_{j=k+1}^{r} u_{j}^{\prime} t_{j}\right) \sum_{a=r+1}^{m-1}\left\{\left(\sum_{i=k+1}^{r}\left(\lambda_{a}^{i}\right)^{\prime} t_{i}\right) \xi_{a}+\left(\sum_{i=k+1}^{r} \lambda_{a}^{i} t_{i}\right) \xi_{a}^{\prime}\right\}  \tag{3.25}\\
& -\tilde{q}^{2} \sum_{a=r+1}^{m-1}\left\{\left(\sum_{i=k+1}^{r}\left(\lambda_{a}^{i}\right)^{\prime \prime} t_{i}\right) \xi_{a}-2\left(\sum_{i=k+1}^{r}\left(\lambda_{a}^{i}\right)^{\prime} t_{i}\right) \xi_{a}^{\prime}-\left(\sum_{i=k+1}^{r} \lambda_{a}^{i} t_{i}\right) \xi_{a}^{\prime \prime}\right\} \\
& -\tilde{q}^{2} \sum_{i=k+1}^{r} u_{i}^{2}\left(\sum_{a=r+1}^{m-1}\left(\sum_{h=k+1}^{r} \lambda_{a}^{h} t_{h}\right) \xi_{a}\right)+\tilde{q}^{3} \sum_{i=k+1}^{r} u_{i}\left(\sum_{a=r+1}^{m-1} \lambda_{a}^{i} \xi_{a}\right) .
\end{align*}
$$

In equation (3.25) all the coefficients of terms in $t$ vanish. So, we can see easily that

$$
\begin{equation*}
\Phi^{\prime \prime}=\sum_{i=k+1}^{r} u_{i}\left(\sum_{a=r+1}^{m-1} \lambda_{a}^{i} \xi_{a}\right) . \tag{3.26}
\end{equation*}
$$

Considering the coefficient of $t_{i_{0}}$ for some $i_{0} \in\{k+1, \cdots, r\}$, we have

$$
\begin{equation*}
u_{i_{0}}^{\prime} \Phi^{\prime}=\sum_{a=r+1}^{m-1}\left\{\left(\lambda_{a}^{i_{0}}\right)^{\prime \prime} \xi_{a}-2\left(\lambda_{a}^{i_{0}}\right)^{\prime} \xi_{a}^{\prime}-\lambda_{a}^{i_{0}} \xi_{a}^{\prime \prime}+\left(\sum_{h=k+1}^{r} u_{h}^{2}\right) \lambda_{a}^{i_{0}} \xi_{a}\right\} . \tag{3.27}
\end{equation*}
$$

Using (3.26) and (3.27), (3.25) is rewritten as

$$
\begin{align*}
& \left.\tilde{q}\left(\sum_{j=k+1}^{r} u_{j}^{\prime \prime} t_{j}\right)-3\left(\sum_{j=k+1}^{r} u_{j}^{\prime} t_{j}\right)^{2}\right\} \sum_{a=r+1}^{m-1}\left(\sum_{i=k+1}^{r} \lambda_{a}^{i} t_{i}\right) \xi_{a} \\
& +3 \tilde{q}\left(\sum_{j=k+1}^{r} u_{j}^{\prime} t_{j}\right) \sum_{a=r+1}^{m-1}\left\{\left(\sum_{i=k+1}^{r}\left(\lambda_{a}^{i}\right)^{\prime} t_{i}\right) \xi_{a}+\left(\sum_{i=k+1}^{r} \lambda_{a}^{i} t_{i}\right) \xi_{a}^{\prime}\right\}=0 . \tag{3.28}
\end{align*}
$$

Comparing the coefficients of $t_{i_{0}}^{2}$ and $t_{i_{0}}^{3}$, we obtain

$$
\left(u_{i_{0}}^{\prime}\right)^{2} \sum_{a=r+1}^{m-1} \lambda_{a}^{i_{0}} \xi_{a}=0
$$

Since $u_{i_{0}}^{\prime}$ is non-zero, we know that $\lambda_{a}^{i_{0}}=0$ for all $a=r+1, \cdots, m-1$.
Therefore, we have

$$
\begin{equation*}
e_{i}^{\prime}=u_{i} \alpha^{\prime} \tag{3.29}
\end{equation*}
$$

for all $i=k+1, \cdots, r$.
Consequently, by (3.24) and (3.29), we obtain

$$
e_{i}^{\prime}=\varepsilon u_{i} \alpha^{\prime}
$$

for all $i=k+1, \cdots, r$.
We now prove that the non-cylindrical ruled submanifold $M$ in $\mathbb{L}^{m}$ satisfying the conditions of Lemma 3.3 is an $(r+1)$-plane.

Theorem 3.4. Let $M$ be an $(r+1)$-dimensional non-cylindrical ruled submanifold parameterized by (3.2) in $\mathbb{L}^{m}$ with harmonic Gauss map. Let $e_{1}, e_{2}, \cdots, e_{r}$ be orthonormal generators of the rulings along the base curve $\alpha$. If $e_{i}^{\prime}$ are non-null for $i=1,2, \cdots, r$ and some of generators of the rulings $e_{1}, \cdots, e_{k}$ are constant vector fields along $\alpha$, then $M$ is part of an $(r+1)$-plane in $\mathbb{L}^{m}$. Proof. By Lemma 3.3, we have

$$
\begin{aligned}
q & =\left\langle x_{s}, x_{s}\right\rangle=\left\langle\left(1+\varepsilon \sum_{i=k+1}^{r} t_{i} u_{i}\right) \alpha^{\prime},\left(1+\varepsilon \sum_{j=k+1}^{r} t_{j} u_{j}\right) \alpha^{\prime}\right\rangle \\
& =\left(1+\varepsilon \sum_{i=k+1}^{r} t_{i} u_{i}\right)^{2}\left\langle\alpha^{\prime}, \alpha^{\prime}\right\rangle=\varepsilon\left(1+\varepsilon \sum_{i=k+1}^{r} t_{i} u_{i}\right)^{2}
\end{aligned}
$$

and hence $G=\Phi$. From $\Delta G=0$ we obtain $\Phi^{\prime \prime}+\varepsilon \Phi^{\prime \prime} u_{i} t_{i}-\varepsilon \Phi^{\prime} u_{i}^{\prime} t_{i}=0$. Therefore, $\Phi^{\prime \prime}=0$ and hence $\Phi^{\prime} u_{i}^{\prime}=0$.

If $u_{i}^{\prime} \neq 0$ for some $i=k+1, \cdots, r$, we get $\Phi^{\prime}=0$.

Suppose that $u_{i}^{\prime} \equiv 0$ for all $i=k+1, \cdots, r$. But, $\Phi^{\prime \prime}=0$ implies

$$
\begin{equation*}
\alpha^{\prime \prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{r}+\sum_{i=k+1}^{r} \alpha^{\prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{i}^{\prime} \wedge \cdots \wedge e_{r}=0 \tag{3.30}
\end{equation*}
$$

This gives

$$
\begin{equation*}
\alpha^{\prime \prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{r} \wedge e_{j}^{\prime}=0 \tag{3.31}
\end{equation*}
$$

for all $j=k+1, \cdots, r$. By virtue of (3.31) and Lemma 3.3, we get $\alpha^{\prime \prime \prime} \wedge \alpha^{\prime}=0$. Together with this fact and (3.30), we have

$$
\alpha^{\prime \prime}=-\sum_{i=k+1}^{r} \varepsilon_{i} u_{i} e_{i}
$$

and hence $\Phi^{\prime}=0$.
Therefore, $M$ is part of an $(r+1)$-dimensional plane in $\mathbb{L}^{m}$.
We now deal with the case that some of generators of rulings have null derivatives.
Lemma 3.5. Let $M$ be an $(r+1)$-dimensional non-cylindrical ruled submanifold parameterized by (3.2) in $\mathbb{L}^{m}$ with harmonic Gauss map. If some generators $e_{j_{1}}, e_{j_{2}}$, $\cdots, e_{j_{k}}$ of the rulings have null derivatives along the base curve $\alpha$ for $j_{1}<j_{2}<$ $\cdots<j_{k} \in\{1,2, \cdots, r\}$, then the Gauss map $G$ has of the form

$$
\begin{equation*}
G=\Phi+\sum_{i=1}^{k} t_{j_{i}} \Psi_{j_{i}} \tag{3.32}
\end{equation*}
$$

for the harmonic vectors $\Phi$ and $\Psi_{j_{i}}$.
Proof. We can rewrite the parametrization (3.2) of $M$ as

$$
x\left(s, t_{1}, \cdots, t_{r}\right)=\alpha(s)+\sum_{i \neq j_{1}, j_{2}, \cdots, j_{k}} t_{i} e_{i}(s)+\sum_{i=1}^{k} t_{j_{i}} e_{j_{i}}(s)
$$

and its Laplace operator is given by

$$
\Delta=\frac{1}{2 q^{2}} \frac{\partial q}{\partial s} \frac{\partial}{\partial s}-\frac{1}{q} \frac{\partial^{2}}{\partial s^{2}}-\frac{1}{2 q} \sum_{i=1}^{r} \varepsilon_{i} \frac{\partial q}{\partial t_{i}} \frac{\partial}{\partial t_{i}}-\sum_{i=1}^{r} \varepsilon_{i} \frac{\partial^{2}}{\partial t_{i}^{2}}
$$

Then, there are possible two cases such that either $e_{j_{k+1}}, \cdots, e_{j_{r}}$ generating the rulings except $e_{j_{1}}(s), e_{j_{2}}(s), \cdots, e_{j_{k}}(s)$ are constant vector fields or $e_{i}^{\prime} \neq 0$ for some $i=j_{k+1}, \cdots, j_{r}$ if $k<r$.

Case 1. Suppose that $e_{j_{k+1}}, \cdots, e_{j_{r}}$ are constant vector fields.
Subcase 1.1. Let deg $q(t)=0$. In this case, $e_{j_{i}}^{\prime}$ are null with $e_{j_{i}}^{\prime}(s) \wedge e_{j_{l}}^{\prime}(s)=0$ for $i, l=1,2, \cdots, k$ and $\left\langle\alpha^{\prime}(s), e_{j}^{\prime}(s)\right\rangle=0$ for $j=j_{1}, j_{2}, \cdots, j_{k}$. Then $M$ has the Gauss map

$$
G=\Phi+\sum_{i=1}^{k} t_{j_{i}} \Psi_{j_{i}} .
$$

Thus, we have

$$
\Delta G=-\left(\Phi^{\prime \prime}(s)+\sum_{i=1}^{k} t_{j_{i}} \Psi_{j_{i}}^{\prime \prime}(s)\right)
$$

Hence, $\Phi$ and $\Psi_{j_{i}}$ are harmonic if the Gauss map $G$ is harmonic.
Subcase 1.2. Let $\operatorname{deg} q(t)=1$. In this case, $\left\langle\alpha^{\prime}(s), e_{j_{i}}^{\prime}(s)\right\rangle \neq 0$ for some $j_{i}$ $(1 \leq i \leq k)$ and the null vector fields $e_{j_{i}}^{\prime}$ satisfy $e_{j_{i}}^{\prime} \wedge e_{j_{l}}^{\prime}=0$ for $i, l=1,2, \cdots, k$. The Gauss map $G$ of $M$ has the form

$$
G=\frac{\tilde{G}(t)}{(\tilde{\varepsilon} q)^{1 / 2}},
$$

where $\operatorname{deg} \tilde{G}(t) \leq 1$. Computing $\Delta G$ and using $\Delta G=0$, we get

$$
\begin{align*}
& \left(\frac{\partial q}{\partial s}\right)^{2}\left(\Phi+\sum_{i=1}^{k} \Psi_{j_{i}} t_{j_{i}}\right)-\frac{3}{2} q \frac{\partial q}{\partial s}\left(\Phi^{\prime}+\sum_{i=1}^{k} \Psi_{j_{i}}^{\prime} t_{j_{i}}\right)-\frac{1}{2} q \frac{\partial^{2} q}{\partial s^{2}}\left(\Phi+\sum_{i=1}^{k} \Psi_{j_{i}} t_{j_{i}}\right) \\
& +q^{2}\left(\Phi^{\prime \prime}+\sum_{i=1}^{k} \Psi_{j_{i}}^{\prime \prime} t_{j_{i}}\right)+\frac{1}{2} q \sum_{i=1}^{k} \varepsilon_{j_{i}}\left(\frac{\partial q}{\partial t_{j_{i}}}\right)^{2}\left(\Phi+\sum_{l=1}^{k} \Psi_{j_{l}} t_{j_{l}}\right)  \tag{3.33}\\
& -\frac{1}{2} q^{2} \sum_{i=1}^{k} \varepsilon_{j_{i}} \frac{\partial q}{\partial t_{j_{i}}} \Psi_{j_{i}}=0
\end{align*}
$$

Suppose $u_{j_{i}}^{\prime} \equiv 0$ for all $i=1, \cdots, k$. Then, $\frac{\partial q}{\partial s}=0$ and $\frac{\partial^{2} q}{\partial s^{2}}=0$. Together with (3.33) and these facts, we have

$$
q\left(\Phi^{\prime \prime}+\sum_{i=1}^{k} \Psi_{j_{i}}^{\prime \prime} t_{j_{i}}\right)+\frac{1}{2} \sum_{i=1}^{k} \varepsilon_{j_{i}}\left(\frac{\partial q}{\partial t_{j_{i}}}\right)^{2}\left(\Phi+\sum_{l=1}^{k} \Psi_{j_{l}} t_{j_{l}}\right)-\frac{1}{2} q \sum_{i=1}^{k} \varepsilon_{j_{i}} \frac{\partial q}{\partial t_{j_{i}}} \Psi_{j_{i}}=0 .
$$

In this equation, since $u_{j_{0}} \neq 0$ for some $j_{0}$, we can easily see that $\Psi_{j_{0}}^{\prime \prime}$ vanishes. Then we obtain two equations as follows:

$$
\begin{gather*}
\Phi^{\prime \prime}+2 \varepsilon \sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}}^{2} \Phi-\sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}} \Psi_{j_{i}}=0  \tag{3.34}\\
\sum_{i=1}^{k} \Phi^{\prime \prime} u_{j_{i}} t_{j_{i}}+\sum_{l=1}^{k} \sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}}^{2} t_{j_{l}} \Psi_{j_{l}}-\sum_{i, l=1}^{k} \varepsilon_{j_{i}} u_{j_{l}} t_{j_{l}} u_{j_{i}} \Psi_{j_{i}}=0 . \tag{3.35}
\end{gather*}
$$

Substituting (3.34) into (3.35), we get

$$
\sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}}^{2}\left\{\sum_{l=1}^{k}\left(2 \varepsilon u_{j_{l}} \Phi-\Psi_{j_{l}}\right) t_{j_{l}}\right\}=0
$$

If $\sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}}^{2} \equiv 0$, then there exists $e_{h}$ for some $h=j_{1}, j_{2}, \cdots, j_{k}$ such that $\varepsilon_{h}=-1$ since deg $q(t)=1$. It is a contradiction because of the causal character of $e_{h}$. Therefore, $\sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}}^{2} \neq 0$ on an open interval $\mathcal{J}$. So, we have $2 \varepsilon u_{j_{i}} \Phi=\Psi_{j_{i}}$ on $\mathcal{J}$, that is,

$$
\left(2 \varepsilon u_{j_{i}} \alpha^{\prime}-e_{j_{i}}^{\prime}\right) \wedge e_{1} \wedge \cdots \wedge e_{r}=0
$$

Since $2 \varepsilon u_{j_{i}} \alpha^{\prime}-e_{j_{i}}^{\prime}$ is orthogonal to $e_{l}$ for each $l=1,2, \cdots, r, 2 \varepsilon u_{j_{i}} \alpha^{\prime}-e_{j_{i}}^{\prime}$ has to be vanishing. But, it is a contradiction because of the characters of $\alpha^{\prime}$ and $e_{j_{i}}^{\prime}$ for all $j_{1}, \cdots, j_{k}$.

Hence, there exists a non-zero function $u_{j_{0}}^{\prime}$ in some open interval $\mathcal{U}$ for some $j_{0}=j_{1}, j_{2}, \cdots, j_{k}$.

On the other hand, equation (3.33) shows that all the coefficients of terms in $t$ vanish. Especially, if we examine the coefficients of $t_{j_{0}}^{3}, t_{j_{0}}^{2}, t_{j_{0}}^{1}$ and $t_{j_{0}}^{0}$, then we have the following four equations:

$$
\begin{gather*}
4\left(u_{j_{0}}^{\prime}\right)^{2} \Psi_{j_{0}}-6 u_{j_{0}}^{\prime} u_{j_{0}} \Psi_{j_{0}}^{\prime}-2 u_{j_{0}}^{\prime \prime} u_{j_{0}} \Psi_{j_{0}}+4 u_{j_{0}}^{2} \Psi_{j_{0}}^{\prime \prime}=0  \tag{3.36}\\
4\left(u_{j_{0}}^{\prime}\right)^{2} \Phi-6 u_{j_{0}}^{\prime} u_{j_{0}} \Phi^{\prime}-2 u_{j_{0}}^{\prime \prime} u_{j_{0}} \Phi+4 u_{j_{0}}^{2} \Phi^{\prime \prime}-3 \varepsilon u_{j_{0}}^{\prime} \Psi_{j_{0}}^{\prime}-\varepsilon u_{j_{0}}^{\prime \prime} \Psi_{j_{0}}+4 \varepsilon u_{j_{0}} \Psi_{j_{0}}^{\prime \prime} \\
+4\left(\sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}}^{2}\right) u_{j_{0}} \Psi_{j_{0}}-4\left(\sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}} \Psi_{j_{i}}\right) u_{j_{0}}^{2}=0 \\
-3 \varepsilon u_{j_{0}}^{\prime} \Phi^{\prime}-\varepsilon u_{j_{0}}^{\prime \prime} \Phi+\Psi_{j_{0}}^{\prime \prime}+4 \varepsilon u_{j_{0}} \Phi^{\prime \prime}+4\left(\sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}}^{2}\right) u_{j_{0}} \Phi \\
+2 \varepsilon\left(\sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}}^{2}\right) \Psi_{j_{0}}-4 \varepsilon\left(\sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}} \Psi_{j_{i}}\right) u_{j_{0}}=0 \\
\Phi^{\prime \prime}+2 \varepsilon\left(\sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}}^{2}\right) \Phi-\left(\sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}} \Psi_{j_{i}}\right)=0 .
\end{gather*}
$$

Substituting (3.39) into (3.38), we get

$$
\begin{equation*}
-3 \varepsilon u_{j_{0}}^{\prime} \Phi^{\prime}-\varepsilon u_{j_{0}}^{\prime \prime} \Phi+\Psi_{j_{0}}^{\prime \prime}-4\left(\sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}}^{2}\right) u_{j_{0}} \Phi+2 \varepsilon\left(\sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}}^{2}\right) \Psi_{j_{0}}=0 \tag{3.40}
\end{equation*}
$$

Putting (3.36) and (3.39) into (3.37), we obtain

$$
\begin{array}{r}
4\left(u_{j_{0}}^{\prime}\right)^{2} \Phi-6 u_{j_{0}}^{\prime} u_{j_{0}} \Phi^{\prime}-2 u_{j_{0}}^{\prime \prime} u_{j_{0}} \Phi+3 \varepsilon u_{j_{0}}^{\prime} \Psi_{j_{0}}^{\prime}+\varepsilon u_{j_{0}}^{\prime \prime} \Psi_{j_{0}} \\
+4\left(\sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}}^{2}\right) u_{j_{0}} \Psi_{j_{0}}-8 \varepsilon u_{j_{0}}^{2}\left(\sum_{i=1}^{k} \varepsilon_{j_{i}} u_{j_{i}}^{2}\right) \Phi-4 \varepsilon \frac{\left(u_{j_{0}}^{\prime}\right)^{2}}{u_{j_{0}}} \Psi_{j_{0}}=0 \tag{3.41}
\end{array}
$$

Multiplying $2 \varepsilon u_{j_{0}}$ with (3.40) and substituting the equation obtained in such a way into (3.41), we get $2 \varepsilon u_{j_{0}} \Phi=\Psi_{j_{0}}$ because $u_{j_{0}}^{\prime}$ is non-zero.

Then, one can easily see that $\alpha^{\prime} \wedge e_{j_{0}}^{\prime}=0$, which is a contradiction. Therefore, we can conclude that no ruled submanifolds with $\operatorname{deg} q=1$ have harmonic Gauss map.

Subcase 1.3. Let deg $q(t)=2$. Using the similar argument developed in Lemma 3.3, we have

$$
\alpha^{\prime} \wedge e_{i}^{\prime}=0
$$

for $i=j_{1}, \cdots, j_{k}$, which is a contradiction. Therefore, no ruled submanifolds with $\operatorname{deg} q=2$ have harmonic Gauss map $G$.

Case 2. Suppose that $e_{i}^{\prime} \neq 0$ for some $i=j_{k+1}, \cdots, j_{r}$.
In this case, we may assume that $e_{i}^{\prime} \neq 0$ for all $i=j_{k+1}, \cdots, j_{r}$, otherwise the ruled submanifold $M$ is a cylinder built over the ruled submanifold parameterized by the base curve $\alpha$ and the rulings generated by $e_{i}$ 's except those constant vector fields. Then, $e_{i}^{\prime}$ are non-null for all $i=j_{k+1}, \cdots, j_{r}$ and $\operatorname{deg} q=2$.

If we again follow a similar argument in the proof of Lemma 3.3 , we have

$$
\alpha^{\prime} \wedge e_{i}^{\prime}=0
$$

for all $i=1,2, \cdots, r$. This is a contradiction.
This completes the proof.
It is easy to show that if the Gauss map $G$ of a ruled submanifold with nondegenerate rulings in $\mathbb{L}^{m}$ has of the form (3.32), $G$ is harmonic. Therefore, combining the results of Theorem 3.4 and Proposition 3.5, we conclude

Theorem 3.6. Let $M$ be an ( $r+1$ )-dimensional non-cylindrical ruled submanifold with non-degenerate rulings in the Minkowski $m$-space $\mathbb{L}^{m}$. Then, $M$ has harmonic Gauss map if and only if $M$ is part of either an $(r+1)$-plane or a ruled submanifold up to cylinders over a certain submanifold with the parametrization given by

$$
x\left(s, t_{1}, t_{2}, \cdots, t_{r}\right)=f(s) \boldsymbol{N}+s \boldsymbol{E}+\sum_{j=1}^{r} t_{j}\left(p_{j}(s) \boldsymbol{N}+\boldsymbol{F}_{j}\right)
$$

for some smooth functions $f$ and $p_{j}$, and some constant vector fields $\boldsymbol{N}, \boldsymbol{E}, \boldsymbol{F}_{j}$ with $\langle\boldsymbol{E}, \boldsymbol{E}\rangle=1,\langle\boldsymbol{N}, \boldsymbol{N}\rangle=\langle\boldsymbol{N}, \boldsymbol{E}\rangle=\left\langle\boldsymbol{N}, \boldsymbol{F}_{j}\right\rangle=\left\langle\boldsymbol{E}, \boldsymbol{F}_{j}\right\rangle=0$, and $\left\langle\boldsymbol{F}_{j}, \boldsymbol{F}_{i}\right\rangle=\delta_{j i}$ for $i, j=1,2, \cdots, r$.

Proof. Suppose that $M$ has harmonic Gauss map.
We now suppose that a non-cylindrical ruled submanifold $M$ with non-degenerate rulings is not part of an ( $r+1$ )-plane and it is parameterized by (3.2). We may also assume that the derivatives of the orthonormal vector fields $e_{1}, e_{2}, \cdots, e_{r}$ defining the rulings never vanish, otherwise $M$ is a cylinder built over those submanifolds. As we see in the proof of Lemma 3.5, only Subcase 1.1 can occur. Therefore, we have $q=1$ and $e_{j}^{\prime}$ are null vector fields with $e_{j}^{\prime}(s) \wedge e_{k}^{\prime}(s)=0$ and $\left\langle\alpha^{\prime}(s), e_{j}^{\prime}(s)\right\rangle=0$ for $j, k=1,2, \cdots, r$. Then $\Delta G=0$ implies that

$$
\begin{equation*}
\Phi^{\prime \prime}(s)=0 \tag{3.42}
\end{equation*}
$$

and

$$
\begin{equation*}
\Psi_{j}^{\prime \prime}(s)=0 \tag{3.43}
\end{equation*}
$$

for all $j=1,2, \cdots, r$. Since $e_{j}^{\prime} \wedge e_{k}^{\prime}=0$ for $j, k=1,2, \cdots, r$, we have

$$
\Psi_{j}^{\prime \prime}=e_{j}^{\prime \prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{r}+\sum_{i=1}^{r} e_{j}^{\prime \prime} \wedge \cdots \wedge e_{i}^{\prime} \wedge \cdots \wedge e_{r}=0
$$

for each $j=1,2, \cdots, r$. This implies

$$
\begin{equation*}
e_{j}^{\prime \prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{r} \wedge e_{l}^{\prime}=0 \tag{3.44}
\end{equation*}
$$

for $j, l=1,2, \cdots, r$. Therefore, the vector fields $e_{j}^{\prime \prime \prime}, e_{1}, \cdots, e_{r}, e_{l}^{\prime}$ are linearly dependent for all $s$. So, (3.1) and the fact that $e_{j}^{\prime} \wedge e_{i}^{\prime}=0$ for $i, j=1,2, \cdots r$ imply

$$
e_{j}^{\prime \prime \prime} \wedge e_{l}^{\prime}=0
$$

for $j, l=1,2, \cdots, r$.
Since $e_{j}^{\prime \prime \prime} \wedge e_{k}^{\prime}=0$ and $e_{j}^{\prime} \wedge e_{k}^{\prime}=0$ for all $j, k=1,2, \cdots r$, we get

$$
\begin{equation*}
e_{j}^{\prime \prime} \wedge e_{k}^{\prime}=0 \tag{3.45}
\end{equation*}
$$

for all $j, k=1,2, \cdots, r$.
On the other hand, (3.42) gives

$$
0=\alpha^{\prime \prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{r}+2 \sum_{i=1}^{r} \alpha^{\prime \prime} \wedge \cdots \wedge e_{i}^{\prime} \wedge \cdots \wedge e_{r}+\sum_{i=1}^{r} \alpha^{\prime} \wedge \cdots \wedge e_{i}^{\prime \prime} \wedge \cdots \wedge \cdots \wedge e_{r}
$$

Thus, we have

$$
\alpha^{\prime \prime \prime} \wedge e_{1} \wedge \cdots \wedge e_{r} \wedge e_{j}^{\prime}=0
$$

and hence

$$
\begin{equation*}
\alpha^{\prime \prime \prime} \wedge e_{j}^{\prime}=0 \tag{3.46}
\end{equation*}
$$

for all $j=1,2, \cdots, r$.
Since $\left\langle\alpha^{\prime}, \alpha^{\prime}\right\rangle=1,\left\langle\alpha^{\prime}, e_{j}^{\prime}\right\rangle=0$ and $e_{j}^{\prime \prime} \wedge e_{k}^{\prime}=0$ for $j, k=1,2, \cdots, r$, we see that

$$
\left\langle\alpha^{\prime \prime}, \alpha^{\prime \prime}\right\rangle=0
$$

along $\alpha$. From this, $\alpha^{\prime \prime}=0$ or $\alpha^{\prime \prime}$ is null and hence up to translation we may put

$$
\begin{equation*}
\alpha(s)=f(s) \mathbf{N}+s \mathbf{E} \tag{3.47}
\end{equation*}
$$

where $\mathbf{N}$ is a constant null vector, $\mathbf{E}$ a constant space-like unit vector satisfying $\langle\mathbf{N}, \mathbf{E}\rangle=$ 0 and $f$ a smooth function.

Since $e_{j}^{\prime \prime} \wedge e_{i}^{\prime}=0$ and $\alpha^{\prime \prime} \wedge e_{j}^{\prime}=0$ for all $i, j=1,2, \cdots, r$, we may have

$$
e_{j}=p_{j}(s) \mathbf{N}+\mathbf{F}_{j}
$$

for some non-zero smooth function $p_{j}$ and orthonormal space-like constant vector fields $\mathbf{F}_{j}$ along $\alpha$ satisfying $\left\langle\mathbf{N}, \mathbf{F}_{j}\right\rangle=0$ for $j=1,2, \cdots, r$.

Consequently, up to translation the parametrization (3.2) of $M$ can be put

$$
\begin{equation*}
x\left(s, t_{1}, t_{2}, \cdots, t_{r}\right)=f(s) \mathbf{N}+s \mathbf{E}+\sum_{j=1}^{r} t_{j}\left(p_{j}(s) \mathbf{N}+\mathbf{F}_{j}\right) . \tag{3.48}
\end{equation*}
$$

Conversely, for some smooth functions $f$ and $p_{j}$ defined along $\alpha$ and some constant vector fields $\mathbf{N}, \mathbf{E}, \mathbf{F}_{j}(j=1,2, \cdots, r)$ satisfying above conditions, it is easy to show that a non-cylindrical ruled submanifold parameterized by (3.48) satisfies

$$
\Delta G=0 .
$$

This completes the proof.
Remark. In Theorem 3.6, if the base curve $\alpha$ is a straight line and the generators $e_{i}$ satisfy $e_{i}^{\prime \prime}=0$ along $\alpha(i=1,2, \cdots, r)$, the ruled submanifold $M$ is minimal.

## 4. Degenerate Rulings

Let $M$ be an $(r+1)$-dimensional ruled submanifold in $\mathbb{L}^{m}$ with degenerate rulings $E(s, r)$ along a regular curve and let its parametrization be given by $\tilde{x}(s, t)$ where $t=\left(t_{1}, t_{2}, \cdots, t_{r}\right)$. Since $E(s, r)$ is degenerate, it can be spanned by a degenerate frame $\left\{B(s)=e_{1}(s), e_{2}(s), \cdots, e_{r}(s)\right\}$ such that

$$
\langle B(s), B(s)\rangle=\left\langle B(s), e_{i}(s)\right\rangle=0, \quad\left\langle e_{i}(s), e_{j}(s)\right\rangle=\delta_{i j}, \quad i, j=2,3, \cdots, r .
$$

Without loss of generality as Lemma 2.2, we may assume that

$$
\left\langle e_{i}^{\prime}(s), e_{j}(s)\right\rangle=0, \quad i, j=2,3, \cdots, r
$$

Since the tangent space of $M$ at $\tilde{x}(s, t)$ is a Minkowski $(r+1)$-space which contains the degenerate ruling $E(s, r)$, there exists a tangent vector field $A$ to $M$ which satisfies

$$
\langle A(s, t), A(s, t)\rangle=0, \quad\langle A(s, t), B(s)\rangle=-1, \quad\left\langle A(s, t), e_{i}(s)\right\rangle=0, \quad i=2,3, \cdots, r
$$

at $\tilde{x}(s, t)$.
Let $\alpha(s)$ be an integral curve of the vector field $A$ on $M$. Then we can define another parametrization $x$ of $M$ as follows:

$$
x\left(s, t_{1}, t_{2}, \cdots, t_{r}\right)=\alpha(s)+\sum_{i=1}^{r} t_{i} e_{i}(s)
$$

where $\alpha^{\prime}(s)=A(s)$.
Lemma 4.1. ([9]). We may assume that $\left\langle A(s), B^{\prime}(s)\right\rangle=0$ for all $s$.
Two of the present authors proved the following lemma.
Lemma 4.2. ([10]). Let $M$ be a ruled submanifold with degenerate rulings. Then, the following are equivalent.
(1) $M$ is minimal.
(2) $B^{\prime}(s)$ is tangent to $M$.

If we put $P=\left\langle x_{s}, x_{s}\right\rangle$ and $Q=-\left\langle x_{s}, x_{t_{1}}\right\rangle$, Lemma 4.1 implies

$$
\begin{gathered}
P(s, t)=2 \sum_{i=2}^{r} u_{i}(s) t_{i}+\sum_{i, j=1}^{r} w_{i j}(s) t_{i} t_{j} \\
Q(s, t)=1+\sum_{i=2}^{r,} v_{i}(s) t_{i}
\end{gathered}
$$

where $v_{i}(s)=\left\langle B^{\prime}(s), e_{i}(s)\right\rangle, u_{i}(s)=\left\langle A(s), e_{i}^{\prime}(s)\right\rangle, w_{i j}(s)=\left\langle e_{i}^{\prime}(s), e_{j}^{\prime}(s)\right\rangle$ for $i, j=$ $1,2, \cdots, r$. Note that $P$ and $Q$ are polynomials in $t=\left(t_{1}, t_{2}, \cdots, t_{r}\right)$ with functions in $s$ as coefficients. Then the Laplacian $\Delta$ of $M$ can be expressed as follows:

$$
\begin{gathered}
\Delta=\frac{1}{Q^{2}}\left\{\frac{\partial \bar{P}}{\partial t_{1}} \frac{\partial}{\partial t_{1}}-2 Q \sum_{i=2}^{r} v_{i} \frac{\partial}{\partial t_{i}}+2 Q \frac{\partial^{2}}{\partial s \partial t_{1}}+\bar{P} \frac{\partial^{2}}{\partial t_{1}^{2}}\right. \\
\left.-2 Q \sum_{i=2}^{r} v_{i} t_{1} \frac{\partial^{2}}{\partial t_{1} \partial t_{i}}-Q^{2} \sum_{i=2}^{r} \frac{\partial^{2}}{\partial t_{i}^{2}}\right\},
\end{gathered}
$$

where $\bar{P}=P-t_{1}^{2} \sum_{i=2}^{r} v_{i}^{2}$.

By definition of an indefinite scalar product $\ll, \gg$ on $G(r+1, m)$, we may put

$$
\ll x_{s} \wedge x_{t_{1}} \wedge x_{t_{2}} \wedge \cdots \wedge x_{t_{r}}, x_{s} \wedge x_{t_{1}} \wedge x_{t_{2}} \wedge \cdots \wedge x_{t_{r}} \gg=-Q^{2}
$$

Let $\bar{\varepsilon}=\operatorname{sign} Q(t)$. Then we have the Gauss map

$$
\begin{aligned}
G= & \frac{1}{\overline{\varepsilon Q}} x_{s} \wedge x_{t_{1}} \wedge x_{t_{2}} \wedge \cdots \wedge x_{t_{r}} \\
= & \frac{1}{\bar{\varepsilon} Q}\left\{A \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r}+t_{1} B^{\prime} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r}\right. \\
& \left.+\sum_{i=2}^{r} t_{i} e_{i}^{\prime} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r}\right\} .
\end{aligned}
$$

We now define a $G$-kind ruled submanifold in Minkouski $m$-space. For a null curve $\tilde{\alpha}(s)$ in $\mathbb{L}^{m}$, we consider a null frame $\left\{A(s), B(s)=e_{1}(s), e_{2}(s), \cdots, e_{m-1}(s)\right\}$ along $\tilde{\alpha}(s)$ satisfying

$$
\begin{array}{r}
\langle A(s), A(s)\rangle=\langle B(s), B(s)\rangle=\left\langle A(s), e_{i}(s)\right\rangle=\left\langle B(s), e_{i}(s)\right\rangle=0, \\
\langle A(s), B(s)\rangle=-1, \quad\left\langle e_{i}(s), e_{j}(s)\right\rangle=\delta_{i j}, \quad \tilde{\alpha}^{\prime}(s)=A(s)
\end{array}
$$

for $i, j=2,3, \cdots, m-1$.
Let $X(s)$ be the matrix $\left(A(s) B(s) e_{2}(s) \cdots e_{m-1}(s)\right)$ consisting of column vectors of $A(s), B(s), e_{2}(s), \cdots, e_{m-1}(s)$ with respect to the standard coordinate system in $\mathbb{L}^{m}$. Then we have

$$
X^{t}(s) E X(s)=T,
$$

where $X^{t}(s)$ denotes the transpose of $X(s), E=\operatorname{diag}(-1,1, \cdots, 1)$ and

$$
T=\left(\begin{array}{ccccc}
0 & -1 & 0 & \cdots & 0 \\
-1 & 0 & 0 & \cdots & 0 \\
0 & 0 & 1 & \cdots & 0 \\
\vdots & \vdots & \vdots & \ddots & \vdots \\
0 & 0 & 0 & \cdots & 1
\end{array}\right)
$$

Consider a system of ordinary differential equations

$$
\begin{equation*}
X^{\prime}(s)=X(s) M(s) \tag{4.1}
\end{equation*}
$$

where

$$
M(s)=\left(\begin{array}{cccccccccc}
0 & 0 & v_{2} & v_{3} & \cdots & v_{r} & 0 & 0 & \cdots & 0 \\
0 & 0 & u_{2} & u_{3} & \cdots & u_{r} & u_{r+1} & u_{r+2} & \cdots & u_{m-1} \\
u_{2} & v_{2} & 0 & 0 & \cdots & 0 & z_{2, r+1} & z_{2, r+2} & \cdots & z_{2, m-1} \\
u_{3} & v_{3} & 0 & 0 & \cdots & 0 & z_{3, r+1} & z_{3, r+2} & \cdots & z_{3, m-1} \\
\vdots & \vdots & \vdots & \vdots & & \vdots & \vdots & \vdots & & \vdots \\
u_{r} & v_{r} & 0 & 0 & \cdots & 0 & z_{r, r+1} & z_{r, r+2} & \cdots & z_{r, m-1} \\
u_{r+1} & 0 & -z_{2, r+1} & -z_{3, r+1} & \cdots & -z_{r, r+1} & 0 & 0 & \cdots & 0 \\
u_{r+2} & 0 & -z_{2, r+2} & -z_{3, r+2} & \cdots & -z_{r, r+2} & 0 & 0 & \cdots & 0 \\
\vdots & \vdots & \vdots & \vdots & & \vdots & \vdots & \vdots & & \vdots \\
u_{m-1} & 0 & -z_{2, m-1} & -z_{3, m-1} & \cdots & -z_{r, m-1} & 0 & 0 & \cdots & 0
\end{array}\right),
$$

where $v_{i}(2 \leq i \leq r), u_{j}(2 \leq j \leq m-1)$ and $z_{a, b}(2 \leq a \leq r, \quad r+1 \leq b \leq m-1)$ are some smooth functions of $s$.

For a given initial condition $X(0)=\left(A(0) B(0) e_{2}(0) \cdots e_{m-1}(0)\right)$ satisfying $X^{t}(0) E X(0)=T$, there exists a unique solution to $X^{\prime}(s)=X(s) M(s)$ on the whole domain $I$ of $\tilde{\alpha}(s)$ containing 0 . Since $T$ is symmetric and $M T$ is skew-symmetric, $\frac{d}{d s}\left(X^{t}(s) E X(s)\right)=0$ and hence we have

$$
X^{t}(s) E X(s)=T
$$

for all $s \in I$. Therefore, $A(s), B(s), e_{2}(s), \cdots, e_{m-1}(s)$ form a null frame along a null curve $\tilde{\alpha}(s)$ in $\mathbb{L}^{m}$ on $I$. Let $\alpha(s)=\int_{0}^{s} A(u) d u$.

Then, we can define a parametrization for a ruled submanifold $M$ by

$$
\begin{equation*}
x\left(s, t_{1}, t_{2}, \cdots, t_{r}\right)=\alpha(s)+t_{1} B(s)+\sum_{i=2}^{r} t_{i} e_{i}(s) \tag{4.2}
\end{equation*}
$$

Definition 4.3. A ruled submanifold $M$ with the parametrization

$$
x\left(s, t_{1}, t_{2}, \cdots, t_{r}\right)=\alpha(s)+t_{1} B(s)+\sum_{i=2}^{r} t_{i} e_{i}(s), \quad s \in J, \quad t_{i} \in I_{i}
$$

satisfying (4.1) and (4.2) is called a $G$-kind ruled submanifold.
Remark 4.4. The terminology $G$-kind ruled submanifolds is named by ruled submanifolds generated by the Gauss map.

We now prove
Theorem 4.5. Let $M$ be a ruled submanifold in $\mathbb{L}^{m}$ with degenerate rulings. $M$ has harmonic Gauss map if and only if $M$ is an open portion of a $G$-kind ruled submanifold.

Proof. We assume that the ruled submanifold $M$ is parameterized by

$$
x\left(s, t_{1}, t_{2}, \cdots, t_{r}\right)=\alpha(s)+t_{1} B(s)+\sum_{i=2}^{r} t_{i} e_{i}(s), \quad s \in J, \quad t_{i} \in I_{i}
$$

such that $\langle A(s), A(s)\rangle=\langle B(s), B(s)\rangle=\left\langle A(s), e_{i}(s)\right\rangle=\left\langle B(s), e_{i}(s)\right\rangle=0,\langle A(s)$, $B(s)\rangle=-1,\left\langle e_{i}(s), e_{j}(s)\right\rangle=\delta_{i j}$, and $\left\langle e_{i}^{\prime}(s), e_{j}(s)\right\rangle=0$ for $i, j=2,3, \cdots, r$, where $J$ and $I_{i}$ are some open intervals and $\alpha^{\prime}(s)=A(s)$. Furthermore, we assume that $\left\langle A(s), B^{\prime}(s)\right\rangle=0$ for all $s$.

Suppose that $M$ has harmonic Gauss map $G$. We then have two possible cases according to the degree of $Q$.

Case 1. Suppose that $\operatorname{deg} Q(t)=0$, that is, $Q=1$ and $v_{i}(s)=0$ for all $i=2,3, \cdots, r$. By definition, we get

$$
\begin{aligned}
\Delta G & =2 \sum_{i=1}^{r}\left\langle B^{\prime}, e_{i}^{\prime}\right\rangle t_{i} B^{\prime} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r} \\
& +2 B^{\prime \prime} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r}+2 \sum_{i=2}^{r} B^{\prime} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{i}^{\prime} \wedge \cdots \wedge e_{r}
\end{aligned}
$$

From $\Delta G=0$ we have

$$
\left\langle B^{\prime}, e_{i}^{\prime}\right\rangle B^{\prime} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r}=0
$$

for all $i=1,2, \cdots, r$. From Lemma 4.1, we see that $B^{\prime}(s)$ must be space-like. Thus, $B$ is a null constant vector field. By using Lemma 4.2, we see that $M$ is minimal.

Let $V(s)=\left\{A(s), B(s), e_{2}(s), \cdots, e_{r}(s)\right\}$ be a smooth distribution of index 1 along $\alpha$ satisfying $\langle A(s), A(s)\rangle=\langle B(s), B(s)\rangle=\left\langle A(s), e_{i}(s)\right\rangle=\left\langle B(s), e_{i}(s)\right\rangle=0$, $\langle A(s), B(s)\rangle=-1,\left\langle e_{i}(s), e_{j}(s)\right\rangle=\delta_{i j}$, and $\left\langle e_{i}^{\prime}(s), e_{j}(s)\right\rangle=0$ for all $s$ and $i, j=$ $2,3, \cdots, r$. Then, by Lemma 2.2, we can choose an orthonormal basis $\left\{e_{r+1}, \cdots\right.$, $\left.e_{m-1}\right\}$ for the orthogonal complement $V^{\perp}(s)$ satisfying $e_{h}^{\prime}(s) \in V(s)$ for all $h=$ $r+1, \cdots, m-1$. Thus we may put

$$
\begin{aligned}
& A^{\prime}(s)=\sum_{i=2}^{m-1} u_{i}(s) e_{i}(s), \\
& B^{\prime}(s)=0, \\
& e_{j}^{\prime}(s)=u_{j}(s) B(s)+\sum_{a=r+1}^{m-1}\left(-z_{j, a}(s)\right) e_{a}(s), \quad j=2, \cdots, r, \\
& e_{a}^{\prime}(s)=u_{h}(s) B(s)+\sum_{i=2}^{r} z_{i, a}(s) e_{i}(s), \quad a=r+1, \cdots, m-1 .
\end{aligned}
$$

For a certain initial condition the above system of linear ordinary differential equations has a unique solution to (4.1) with $v_{i}=0(i=2, \cdots, r)$ and $z_{a, b}=0(a, b=$ $r+1, r+2, \cdots, m-1$ ). The solution defines part of a $G$-kind ruled submanifold.

Case 2. Suppose that $\operatorname{deg} Q(t)=1$. Let $V(s)=\left\{A(s), B(s), e_{2}(s), \cdots, e_{r}(s)\right\}$ be a smooth distribution of index 1 along $\alpha$. Then we can choose an orthonormal basis $\left\{e_{r+1}, \cdots, e_{m-1}\right\}$ for the orthogonal complement $V^{\perp}(s)$ satisfying $e_{h}^{\prime}(s) \in V(s)$ for $h=r+1, \cdots, m-1$. Then we may put

$$
\begin{aligned}
A^{\prime}(s) & =\sum_{i=2}^{m-1} u_{i}(s) e_{i}(s) \\
B^{\prime}(s) & =\sum_{i=2}^{m-1} v_{i}(s) e_{i}(s) \\
e_{j}^{\prime}(s) & =v_{j}(s) A(s)+u_{j}(s) B(s)+\sum_{b=r+1}^{m-1}\left(-z_{j, b}(s)\right) e_{b}(s), \quad j=2, \cdots, r \\
e_{a}^{\prime}(s) & =v_{a}(s) A(s)+u_{a}(s) B(s)+\sum_{b=2}^{r} z_{b, a}(s) e_{b}(s), \quad a=r+1, \cdots, m-1
\end{aligned}
$$

The straightforward computation provides

$$
\begin{align*}
\Delta G= & \frac{2 \bar{\varepsilon}}{Q^{3}} \sum_{h=r+1}^{m-1}\left\{\left(\sum_{i=1}^{r}\left\langle B^{\prime}, e_{i}^{\prime}\right\rangle t_{i}-\sum_{i=2}^{r} v_{i}^{\prime} t_{i}\right) v_{h}+v_{h}^{\prime} Q\right\} e_{h} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r} \\
& +\frac{2 \bar{\varepsilon}}{Q^{2}} \sum_{h=r+1}^{m-1} v_{h}^{2} A \wedge B \wedge e_{2} \wedge \cdots \wedge e_{r}  \tag{4.3}\\
& +\frac{2 \bar{\varepsilon}}{Q^{2}} \sum_{i=2}^{r} \sum_{h=r+1}^{m-1} v_{i} v_{h} e_{h} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{i-1} \wedge A \wedge e_{i+1} \wedge \cdots \wedge e_{r} \\
& -\frac{2 \bar{\varepsilon}}{Q^{2}} \sum_{i=2}^{r} \sum_{h, l=r+1}^{m-1} v_{h} z_{i, l} e_{h} \wedge B \wedge e_{2} \wedge \cdots \wedge e_{i-1} \wedge e_{l} \wedge e_{i+1} \wedge \cdots \wedge e_{r}
\end{align*}
$$

Since the Gauss map is harmonic, $\Delta G=0$. Thus, the functions $v_{a}$ are vanishing for $a=r+1, \cdots, m-1$. Therefore, $M$ is part of a $G$-kind ruled submanifold.

Conversely, if $M$ is a $G$-kind ruled submanifold, then $\operatorname{deg} Q \leq 1$. From equation (4.3), we can see that $v_{r+1}=\cdots=v_{m-1}=0$ implies $\Delta G=0$. Therefore, a $G$-kind ruled submanifold $M$ has harmonic Gauss map $G$. This completes the proof.

In [10], two of the present authors set up a characterization of minimal ruled submanifolds in Minkowski space. Together with Theorem 4.5, we have a new characterization of minimal ruled submanifolds with degenerate rulings in Minkowski space.

Theorem 4.6. Let $M$ be a ruled submanifold in $\mathbb{L}^{m}$ with degenerate rulings. The following are equivalent:
(1) $M$ is minimal.
(2) $M$ has harmonic Gauss map.
(3) $M$ is an open portion of a G-kind ruled submanifold.

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