

# ON WALKER DIFFERENTIATION IN DIFFERENTIABLE MANIFOLDS WITH $r$ - $\pi$ -STRUCTURE

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After Walker [1]<sup>1)</sup> has introduced a new covariant differentiation in almost complex spaces under which the almost complex structure  $F_i^h$  is covariant constant and applied such differentiation to the Nijenhuis tensor  $N_{ji}^h$  to obtain a new differential invariant constructed only from the tensor  $F_i^h$ , Yano [2] showed that there exist an infinitely many Walker differentiations (the one obtained by Walker in the case of almost complex space is also contained) in almost product as well as almost complex spaces which leave their structure tensor  $F_i^h$  covariant constant.

In this note we intend to give formal extensions of some results by Walker and Yano to a more general differentiable manifold with an  $r$ - $\pi$ -structure [3]. Such spaces contain as special cases almost complex spaces, almost product spaces and also almost contact manifolds studied by Boothby-Wang, Gray, Sasaki and others.

As preparation we treat in §1 the hybrid and purity of tensors of differentiable manifold with  $r$ - $\pi$ -structure. In §2 we prove that the torsion tensor of an  $r$ - $\pi$ -structure is hybrid. In §3 we determine all Walker differentiations which leave the fundamental tensor  $F_i^h$  of the  $r$ - $\pi$ -structure covariant constant. In §4 the one corresponding to that of Walker is considered and a differential invariant constructed only from  $F_i^h$  is obtained. In §5 specialization to almost contact manifold is given.

1. If in a differentiable manifold  $V_n$  there exist  $r$  ( $2 \leq r \leq n$ ) differentiable distributions  $T_1, \dots, T_r$  which assign  $r$  complemented subspaces of dimension  $\geq 1$  in the complexified tangent space  $T_x^c$  ( $T_x^c = T_1 + \dots + T_r$ : direct sum) at each point  $x \in V_n$ , then  $V_n$  is said to have an  $r$ - $\pi$ -structure [3].

It is known that for a manifold  $V_n$  to have an  $r$ - $\pi$ -structure, it is necessary and sufficient that the manifold has a non-degenerate tensor field  $F_i^h$  satisfying

$$(1.1) \quad \overset{r}{F}_j^i = \lambda^r \delta_j^i,$$

where  $\lambda$  is a non-zero fixed complex constant, and we have put

$$(1.2) \quad \overset{s}{F}_j^i = F_{h_1}^i F_{h_1}^{h_2} \dots F_j^{h_{s-1}}, \quad \overset{1}{F}_j^i = F_j^i \quad \text{and} \quad \overset{0}{F}_j^i = \delta_j^i.$$

The tensor  $F_j^i$  satisfying (1.1) is called the fundamental tensor associated with

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1) Number in bracket refers to the reference at the end of the paper.

the  $r$ - $\pi$ -structure. It is evident that

$$(1.3) \quad \bar{F}_j^i = \frac{1}{\lambda^{ar}} \bar{F}_j^{ar-s}, \quad (a, s: \text{positive integers}, r > ar - s \geq 0).$$

Now let  $p_\alpha^i$  be the projection tensor to the distribution  $T_\alpha (\alpha = 1, \dots, r)$ , then we have

$$(1.4) \quad p_\alpha^i = \frac{1}{r} \sum_{s=0}^{r-1} \frac{1}{(\lambda w_\alpha)^s} \bar{F}_\alpha^i,$$

where  $w_\alpha (\alpha = 1, \dots, r)$  are the  $r$ -th power roots of unity.

Define

$$(1.5) \quad \Phi_{aj}^{ib} = \sum_{\alpha=1}^r p_\alpha^i p_j^b,$$

then by use of (1.4) and the following identities

$$\sum_{\alpha=1}^r \frac{1}{w_\alpha} = \sum_{\alpha=1}^r \frac{1}{w_\alpha^2} = \dots = \sum_{\alpha=1}^r \frac{1}{w_\alpha^{r-1}} = 0,$$

we have the following expression:

$$(1.6) \quad \Phi_{aj}^{ib} = \frac{1}{r} \left( \delta_a^i \delta_j^b + \frac{1}{\lambda^r} \sum_{t=1}^{r-1} \bar{F}_a^i \bar{F}_j^b \right).$$

Furthermore, we define

$$(1.7) \quad \Phi_{aj}^{ib} = \delta_a^i \delta_j^b - \Phi_{aj}^{ib} = \frac{1}{r} \left\{ (r-1) \delta_a^i \delta_j^b - \frac{1}{\lambda^r} \sum_{t=1}^{r-1} \bar{F}_a^i \bar{F}_j^b \right\}.$$

Then for the two tensors in (1.6) and (1.7), we have

$$(1.8) \quad \left\{ \begin{array}{l} \Phi_{aj}^{ib} + \Phi_{aj}^{ib} = \delta_a^i \delta_j^b, \\ \Phi_{hj}^{ik} \Phi_{ak}^{hb} = \Phi_{aj}^{ib}, \quad \Phi_{hj}^{ik} \Phi_{ak}^{hb} = \Phi_{aj}^{ib}, \\ \Phi_{hj}^{ik} \Phi_{ak}^{hb} = \Phi_{hj}^{ik} \Phi_{ak}^{hb} = 0. \end{array} \right.$$

For, by use of

$$p_\alpha^i p_\beta^h = \begin{cases} p_\alpha^i & \text{if } \alpha = \beta \\ 0 & \text{if } \alpha \neq \beta, \end{cases}$$

we have for example

$$\Phi_{hj}^{ik} \Phi_{ak}^{hb} = \left( \sum_{\alpha=1}^r p_\alpha^i p_j^k \right) \left( \sum_{\beta=1}^r p_\alpha^h p_\beta^b \right) = \sum_{\alpha=1}^r \sum_{\beta=1}^r (p_\alpha^i p_\alpha^h) (p_j^k p_\beta^b) = \sum_{\alpha=1}^r p_\alpha^i p_\alpha^b = \Phi_{aj}^{ib}.$$

Using this relation and (1.7) other relations in (1.8) can be easily seen.



$C^\infty$  of the principal fibre space  $E_\pi(V_n)$  which consists of all adapted bases relative to all points of  $V_n$ . Then at each point of the neighborhood  $U$  an adapted basis  $(e_i)$  is assigned. Let  $(\theta^i)$  be the dual cobasis of  $(e_i)$ , then we have

$$(2.5) \quad d\theta^i = \frac{1}{2} C_{jk}^i \theta^j \wedge \theta^k, \text{ where } C_{jk}^i + C_{kj}^i = 0.$$

Then by definition the torsion tensor  $t_{jk}^i$  of an  $r$ - $\pi$ -structure is the one with the following components with respect to the adapted basis  $(e_i)$ :

$$(2.6) \quad t_{b_\alpha c_\alpha}^{a_\alpha} = C_{b_\alpha c_\alpha}^{a_\alpha}, \quad t_{jk}^i = 0 \text{ for other indices.}$$

Hence

$$(2.7) \quad t_{b_\alpha c_\alpha}^{a_\alpha} = t_{b_\alpha c_\alpha}^{a_\alpha} = 0.$$

Therefore,  $t_{jk}^i$  is hybrid both with respect to  $i, j$  and to  $i, k$ . Q.E.D.

In a previous paper [3] of the present author it is shown that

$$(2.8) \quad t_{jk}^m = \frac{2(r-1)^2}{r^2} \Phi \Phi' S_{jk}^m,$$

where  $S_{jk}^m$  is the torsion tensor of a  $\pi$ -connection  $\Gamma_{jk}^m$  and the operation  $\Phi$  and  $\Phi'$  are defined as follows:

$$(2.9) \quad \Phi S_{jk}^m = S_{jk}^m - \frac{1}{r-1} \frac{1}{\lambda^r} \sum_{s=1}^{r-1} F_k^{k_1} S_{jk_1}^h F_h^{r-s} = \frac{r}{r-1} \Phi_{hk}^{mk} S_{jk_1}^h,$$

$$(2.10) \quad \Phi' S_{jk}^m = S_{jk}^m - \frac{1}{r-1} \frac{1}{\lambda^r} \sum_{s=1}^{r-1} F_j^{j_1} S_{j_1 k}^h F_h^{r-s} = \frac{r}{r-1} \Phi_{hj}^{mj} S_{j_1 k}^h.$$

From these relations we have the following by virtue of (1.8)<sub>3</sub>:

$$(2.11) \quad \Phi t_{jk}^m = 0.$$

Since  $\Phi$  and  $\Phi'$  commute, it follows that  $t_{jk}^m$  is hybrid both with respect to  $m, j$  and  $m, k$ .

3. Hereafter we assume that the torsion tensor of the  $r$ - $\pi$ -structure  $t_{jk}^i$  does not vanish.

Let  $U$  and  $V$  be two local coordinate neighborhoods of  $V_n$  such that  $U \cap V \neq \emptyset$  and let  $(x^i)$  and  $(x^\alpha)$  be respectively the local coordinates in  $U$  and  $V$ .

Suppose  $\Gamma_{jkl}^i$  and  $\Gamma_{\beta\gamma\delta}^\alpha$  be respectively the components of a geometric object having the following transformation law which is associative and reversible:

$$(3.1) \quad \Gamma_{jkl}^i = X_\alpha^i X_j^\beta X_k^\gamma X_l^\delta \Gamma_{\beta\gamma\delta}^\alpha - X_{\alpha\beta}^i X_j^\beta X_k^\gamma X_l^\delta t_{\gamma\delta}^\alpha,$$

where  $t_{jk}^i$  is the torsion tensor of the considered  $r$ - $\pi$ -structure of  $V_n$  and

$$(3.2) \quad X_\alpha^i = \frac{\partial x^i}{\partial x^\alpha}, \quad X_{\alpha\beta}^i = \frac{\partial^2 x^i}{\partial x^\alpha \partial x^\beta} \text{ and } X_i^\alpha = \frac{\partial x^\alpha}{\partial x^i}$$

are formed from the coordinate transformation  $x^i = x^i(x^\alpha)$ .

Then the following are tensors provided that  $v^i$ ,  $w_j$  and  $\phi_j^i$  are tensors:

$$(3.3) \quad \left\{ \begin{array}{l} \nabla_{kl} f = \frac{\partial f}{\partial x^m} t_{kl}^m, \quad (f: \text{scalar}), \\ \nabla_{kl} v^i = \frac{\partial v^i}{\partial x^m} t_{kl}^m + \Gamma_{jkl}^i v^j, \\ \nabla_{kl} w_j = \frac{\partial w_j}{\partial x^m} t_{kl}^m - \Gamma_{jkl}^i w_i, \\ \nabla_{kl} \phi_j^i = \frac{\partial \phi_j^i}{\partial x^m} t_{kl}^m + \Gamma_{mkl}^i \phi_j^m - \Gamma_{jkl}^m \phi_m^i, \text{ etc.} \end{array} \right.$$

This sort of covariant differentiation is a formal extension of the one first introduced by Walker and following Yano we shall call the left hand side terms of these formulas as *Walker derivatives*. The same calculation rules as those of ordinary covariant differentiation hold also in the present case.

We shall now determine all Walker differentiations which satisfy<sup>2)</sup>

$$(3.4) \quad \nabla_{kl} F_j^i = \frac{\partial F_j^i}{\partial x^m} t_{kl}^m + \Gamma_{mkl}^i F_j^m - \Gamma_{jkl}^m F_m^i = 0,$$

where  $F_j^i$  is the fundamental tensor of the considered  $r$ - $\pi$ -structure.

Let  $\nabla_m$  denote the covariant derivative with respect to a fixed affine connection  $\Gamma_{jk}^i$ , then (3.4) may be written as

$$(3.5) \quad (\nabla_m F_j^i) t_{kl}^m + T_{mkl}^i F_j^m - T_{jkl}^m F_m^i = 0,$$

where we put

$$(3.6) \quad T_{jkl}^i = \Gamma_{jkl}^i - \Gamma_{jh}^i t_{kl}^h.$$

From (1.1) and (3.5) we have

$$(3.7) \quad T_{jkl}^n - \frac{1}{\lambda^r} F_j^m T_{mkl}^i \bar{F}_i^{r-1n} = \frac{1}{\lambda^r} t_{kl}^m (\nabla_m F_j^i) \bar{F}_i^{r-1m}.$$

Since from  $\nabla_{kl} F_j^i = 0$  it follows that

$$(3.8) \quad \nabla_{kl} \bar{F}_j^i = 0 \quad (s=1, \dots, r-1),$$

so we have the following:

$$(3.9) \quad T_{jkl}^n - \frac{1}{\lambda^r} \bar{F}_j^m T_{mkl}^i \bar{F}_i^{r-sn} = \frac{1}{\lambda^r} t_{kl}^m (\nabla_m \bar{F}_j^i) \bar{F}_i^{r-sm}, \quad (s=1, \dots, r-1).$$

2) This method is quite the same as that of determination of all  $\pi$ -connections of an  $r$ - $\pi$ -structure which was reported (unpublished) by the present author in a seminar conducted by Prof. S. Sasaki at Tôhoku University, Sendai, Japan in the spring of 1960.

Adding these relations, we have by virtue of (1. 7)

$$(3.10) \quad \Phi_{aj}^{nb} T_{bkl}^a = t_{kl}^m P_{mj}^n,$$

where

$$(3.11) \quad P_{mj}^n = \frac{1}{r} \frac{1}{\lambda^r} \sum_{t=1}^{r-1} (\nabla_m F_j^t) F_i^{r-t}.$$

Thus (3.10) is a necessary condition to be satisfied by  $T_{jkl}^i$  so that the Walker differentiation with coefficients  $\Gamma_{jkl}^i$  to satisfy (3. 4).

First of all, we shall solve the system of equations (3.10) for unknown tensor  $T_{jkl}^i$ .

It is evident that

$$(3.12) \quad \Phi_{ai}^{nb} (t_{kl}^m P_{mb}^a) = t_{kl}^m \Phi_{ai}^{nb} P_{mb}^a.$$

However we show that

$$(3.13) \quad \Phi_{ai}^{nb} P_{mb}^a = 0.$$

For, each term of  $P_{mb}^a$  is written as

$$(3.14) \quad (\nabla_j F_i^{h_t}) F_{h_t}^{r-t} = \sum_{s=0}^{r-1} F_i^{h_s} (\nabla_j F_{h_s}^{h_{s+1}}) F_{h_{s+1}}^{r-s-1}$$

and we have

$$(3.15) \quad \begin{aligned} & \Phi_1 [F_i^{h_s} (\nabla_j F_{h_s}^{h_{s+1}}) F_{h_{s+1}}^{r-s-1}] \\ &= \frac{1}{r} \left\{ F_i^{h_s} (\nabla_j F_{h_s}^{h_{s+1}}) F_{h_{s+1}}^{r-s-1} + \frac{1}{\lambda^r} \sum_{t=1}^{r-1} F_i^{h_s} (\nabla_j F_{h_s}^{h_{s+1}}) F_{h_{s+1}}^{r+s-1} \right\}. \end{aligned}$$

However,

$$(3.16) \quad \begin{aligned} & \frac{1}{\lambda^r} \sum_{t=1}^{r-1} F_i^{h_s} (\nabla_j F_{h_s}^{h_{s+1}}) F_{h_{s+1}}^{r+t-s-1} \\ &= \sum_{t=1}^s F_i^{h_s} (\nabla_j F_{h_s}^{h_{s+1}}) F_{h_{s+1}}^{r+t-s-1} + \sum_{t=s+1}^{r-1} F_i^{h_s} (\nabla_j F_{h_s}^{h_{s+1}}) F_{h_{s+1}}^{t-s-1} \\ &= \sum_{u=0}^{s-1} F_i^{h_s} (\nabla_j F_{h_s}^{h_{s+1}}) F_{h_{s+1}}^{r-u-1} + \sum_{v=s+1}^{r-1} F_i^{h_s} (\nabla_j F_{h_s}^{h_{s+1}}) F_{h_{s+1}}^{v-s-1}, \end{aligned}$$

so we have

$$(3.17) \quad \Phi_1 [F_i^{h_s} (\nabla_j F_{h_s}^{h_{s+1}}) F_{h_{s+1}}^{r-s-1}] = \frac{1}{r} \sum_{s=0}^{r-1} F_i^{h_s} (\nabla_j F_{h_s}^{h_{s+1}}) F_{h_{s+1}}^{r-s-1} = \frac{1}{r} \nabla_j F_i^r = 0.$$

Thus we have (3.13) and consequently from (3.12) and (1. 8)<sub>1</sub>

$$(3.18) \quad \Phi_{\alpha j}^{hb}(t_{kl}^m P_{mb}^a) = t_{kl}^m P_{mj}^h,$$

that is  $T_{jkl}^h = t_{kl}^m P_{mj}^h$  is a special solution of (3.10).

Let  $T_{jkl}^h$  be any solution of (3.10) and put

$$(3.19) \quad W_{jkl}^h = T_{jkl}^h - t_{kl}^m P_{mj}^h,$$

then from (3.10) and (3.18) we have

$$(3.20) \quad \Phi W_{jkl}^h = 0 \quad \text{or} \quad \Phi W_{jkl}^h = W_{jkl}^h.$$

Conversely for any tensor  $W_{jkl}^h$ ,  $t_{kl}^m P_{mj}^h + \Phi W_{jkl}^h$  is a solution of (3.10).

Thus *the most general solution of (3.10) are expressed as*

$$(3.21) \quad T_{jkl}^i = t_{kl}^m P_{mj}^h + \Phi_{\alpha j}^{ib} W_{bkl}^a,$$

where  $W_{jkl}^h$  is an arbitrary tensor.

In the next place, we shall show that the solutions (3.21) of (3.10) all satisfy (3. 7) which is equivalent to (3. 4). Since

$$(3.22) \quad t_{kl}^p P_{pj}^n - \frac{1}{\lambda^r} F_j^m (t_{kl}^p P_{pm}^i) F_i^{r-1} = t_{kl}^p (P_{pj}^n - \frac{1}{\lambda^r} F_j^n P_{pm}^i F_i^{r-1}),$$

it needs only to show that

$$(3.23) \quad P_{pj}^n - \frac{1}{\lambda^r} F_j^m P_{pm}^i F_i^{r-1} = \frac{1}{\lambda^r} (\nabla_p F_j^i) F_i^{r-1}$$

and

$$(3.24) \quad \frac{1}{\lambda^r} F_j^m (\Phi W_{mkl}^i) F_i^{r-1} = \Phi W_{jkl}^n.$$

For the proof of (3.23), if we take account of

$$(3.25) \quad (\nabla_j F_i^{s-1}) F_{h_1}^{r-s} = (\nabla_j F_i^{m_1}) F_{m_1}^{r-1} + F_i^{m_1} (\nabla_j F_{m_1}^{s-1}) F_{h_1}^{r-s},$$

then by virtue of

$$(3.26) \quad (\nabla_j F_i^{m_1}) F_{m_1}^{r-1} = -F_i^{m_1} (\nabla_j F_{m_1}^{r-1}),$$

$P_{ji}^h$  can be written as follows:

$$(3.27) \quad P_{ji}^h = \frac{1}{\lambda^r} \left\{ (\nabla_j F_i^{h_1}) F_{h_1}^{r-1} + \frac{1}{r} \sum_{s=1}^{r-1} F_i^{m_1} (\nabla_j F_{m_1}^s) F_{h_1}^{r-s-1} \right\}.$$

Furthermore, from (3.11) we have

$$(3.28) \quad F_i^{t_1} P_{ji}^{r-1} F_{h_1}^h = \frac{1}{r} \sum_{s=1}^{r-1} F_i^{t_1} (\nabla_j F_{i_1}^s) F_{h_1}^{r-s-1}.$$

Hence, from (3.27) and (3.28) we have (3.23).

For the proof of (3.24), we have by the definition of  $\Phi$ :

$$\begin{aligned} & \frac{1}{\lambda^r} F_i^{i_1} (\Phi W_{i_{kl} h_1} F_{h_1}^{r-1}) \\ &= \frac{1}{r} \frac{1}{\lambda^r} \left\{ F_i^{i_1} W_{i_{kl} h_1} F_{h_1}^{r-1} + \frac{1}{\lambda^r} \sum_{t=1}^{r-1} F_a^{r+t-1} W_{bkl}^a F_i^{r-t+1} \right\} \\ &= \frac{1}{r} \frac{1}{\lambda^r} \left\{ F_{h_1}^{r-1} W_{i_{kl} h_1} F_i^{i_1} + \sum_{s=0}^{r-2} F_a^{r-s} W_{bkl}^a F_i^{r-s} \right\}; \quad s = t - 1 \\ &= \frac{1}{r} \left\{ W_{ikl}^h + \frac{1}{\lambda^r} \sum_{s=1}^{r-1} F_a^{r-s} W_{bkl}^a F_i^{r-s} \right\} = \Phi W_{ikl}^h. \end{aligned}$$

Thus the most general  $T_{jkl}^i$  which gives rise to a Walker differentiation satisfying (3.4) is given by (3.21).

4. We consider now a special case with the following

$$(4.1) \quad W_{jkl}^i = -\Phi_{aj}^{ib} (\nabla_b t_{kl}^a) + \frac{1}{r} \frac{1}{\lambda^r} \left\{ \sum_{t=1}^{r-1} F_a^{t-1} t_{hk}^a (\nabla_l F_j^{r-t}) - \sum_{t=1}^{r-1} F_a^{t-1} t_{hl}^a (\nabla_k F_j^{r-t}) \right\}.$$

For such  $W_{jkl}^i$ , it can be shown that

$$(4.2) \quad \Phi_{aj}^{ib} W_{bkl}^a = W_{jkl}^i.$$

For the proof of (4.2) if we take account of (1.8)<sub>2</sub>, it needs only to show that

$$(4.3) \quad \Phi_{aj}^{ib} \left\{ \sum_{t=1}^{j-1} F_a^{t-1} t_{\beta k}^a (\nabla_l F_b^{r-t}) \right\} = \sum_{t=1}^{r-1} F_a^{t-1} t_{\beta k}^a (\nabla_l F_b^{r-t}).$$

Let the left hand side of (4.3) be denoted by  $A$ , then we have

$$(4.4) \quad A = \sum_{t=1}^{r-1} F_a^{t-1} t_{\beta k}^a (\nabla_l F_b^{r-t}) + \frac{1}{\lambda^r} \sum_{s=1}^{r-1} \sum_{t=1}^{r-1} F_a^{s+t-1} t_{\beta k}^a (\nabla_l F_b^{r-t}) F_j^{r-t}.$$

Let the second block of the right hand side of (4.4) be denoted as  $B$ , then we have by virtue of  $F_j^{2r-t-s} = F_b^{r-t} F_j^{r-s}$

$$(4.5) \quad B = -\frac{1}{\lambda^r} \sum_{s=1}^{r-1} \sum_{t=1}^{r-1} F_a^{s+t-1} t_{\beta k}^a F_b^{r-t} (\nabla_l F_j^{r-s}) + \frac{1}{\lambda^r} \sum_{s=1}^{r-1} \sum_{t=1}^{r-1} F_a^{s+t-1} t_{\beta k}^a (\nabla_l F_j^{2r-t-s}).$$

Now denote the first and second block of the right hand side of (4.5) respectively as  $C$  and  $D$ . Then by making use of the fact that  $t_{jk}^i$  is hybrid with respect to  $i$  and  $j$ , that is

$$(4.6) \quad \Phi_{ab}^{ih} t_{hl}^a = t_{bl}^a + \frac{1}{\lambda^r} \sum_{t=1}^{r-1} F_a^{t-1} t_{hl}^a F_b^{r-t} = 0,$$



we have

$$(4.7) \quad C = -\frac{1}{\lambda^r} \sum_{s=1}^{r-1} F_\gamma^i \left( \sum_{t=1}^{r-1} F_\alpha^\gamma t_{\beta k}^\alpha F_b^{r-t} \right) (\nabla_l F_j^b) = \sum_{s=1}^{r-1} F_\gamma^i t_{bk}^\gamma (\nabla_l F_j^b).$$

However, by virtue of (1.2) and  $\nabla_l F_j^\beta = 0$  we have

$$(4.8) \quad \begin{aligned} D &= \frac{1}{\lambda^r} \sum_{s=1}^{r-1} \sum_{u=s+1}^{r+s-1} F_\alpha^i t_{\beta k}^\alpha (\nabla_l F_j^\beta) \quad (u = s + t) \\ &= \frac{1}{\lambda^r} \sum_{s=1}^{r-1} \left( \sum_{u=s+1}^{r-1} + \sum_{u=r+1}^{r+s-1} \right) F_\alpha^i t_{\beta k}^\alpha (\nabla_l F_j^\beta) \\ &= \sum_{s=1}^{r-1} \left\{ \sum_{u=s+1}^{r-1} F_\alpha^i t_{\beta k}^\alpha (\nabla_l F_j^\beta) + \sum_{v=1}^{s-1} F_\alpha^i t_{\beta k}^\alpha (\nabla_l F_j^\beta) \right\}, (v = u - r). \end{aligned}$$

Substitute (4.7) and (4.8) into (4.5), we have

$$(4.9) \quad B = \sum_{s=1}^{r-1} \sum_{u=1}^{r-1} F_\alpha^i t_{\beta k}^\alpha (\nabla_l F_j^\beta) = (r-1) \sum_{u=1}^{r-1} F_\alpha^i t_{\beta k}^\alpha (\nabla_l F_j^\beta).$$

Put (4.9) into (4.4) we get (4.3).

Finally, under the assumption that the affine connection  $\Gamma_{jk}^i$  is symmetric, we shall calculate the following  $\Gamma_{jkl}^i$  obtained by putting (3.21) and (4.1) into (3.6):

$$(4.10) \quad \begin{aligned} \Gamma_{jkl}^i &= \Gamma_{jh}^i t_{kl}^h + P_{mj}^i t_{kl}^m \\ &\quad - \Phi_{aj}^{ib} (\nabla_b t_{kl}^a) + \frac{1}{r} \frac{1}{\lambda^r} \left\{ \sum_{t=1}^{r-1} F_\alpha^i t_{hk}^a (\nabla_l F_j^{r-t}) - \sum_{u=1}^{r-1} F_\alpha^i t_{kl}^a (\nabla_k F_j^{r-u}) \right\}. \end{aligned}$$

First of all, we have

$$(4.11) \quad \begin{aligned} P_{mj}^i t_{kl}^m &= \frac{1}{r} \frac{1}{\lambda^r} t_{kl}^m \sum_{t=1}^{r-1} \left( \frac{\partial F_j^a}{\partial x^m} \right)^{r-t} F_a^i \\ &\quad - \frac{r-1}{r} \Gamma_{jm}^i t_{kl}^m + \frac{1}{r} \frac{1}{\lambda^r} t_{hl}^m \sum_{t=1}^{r-1} F_j^h \Gamma_{hm}^a F_a^{r-t} F_a^i, \end{aligned}$$

and

$$(4.12) \quad \begin{aligned} -\Phi_{aj}^{ib} (\nabla_b t_{kl}^a) &= -\Phi_{aj}^{ib} \left( \frac{\partial t_{kl}^a}{\partial x^b} \right) \\ &\quad - \frac{1}{r} \Gamma_{hj}^i t_{kl}^h - \frac{1}{r} \frac{1}{\lambda^r} \sum_{t=1}^{r-1} F_\alpha^i F_j^{r-t} \Gamma_{mb}^a t_{kl}^m \\ &\quad + \frac{1}{r} \frac{1}{\lambda^r} \sum_{t=1}^{r-1} F_\alpha^i F_j^{r-t} (\Gamma_{kb}^h t_{hl}^a) + \frac{1}{r} \frac{1}{\lambda^r} \sum_{t=1}^{r-1} F_\alpha^i F_j^{r-t} (\Gamma_{lb}^h t_{kh}^a) \\ &\quad + \frac{1}{r} \Gamma_{kj}^h t_{hl}^i + \frac{1}{r} \Gamma_{lj}^h t_{kh}^i. \end{aligned}$$

The last two terms are equal to the following by the fact that  $t_{jk}^i$  is hybrid with respect to  $i$  and  $j$ :

$$(4.13) \quad -\frac{1}{r} \frac{1}{\lambda^r} \sum_{t=1}^{r-1} F_a^i t_{hl}^a \bar{F}_b^{r-t} \Gamma_{kj}^b - \frac{1}{r} \frac{1}{\lambda^r} \sum_{t=1}^{r-1} F_a^i t_{kh}^a \bar{F}_b^{r-t} \Gamma_{lj}^b.$$

Moreover, we have

$$(4.14) \quad -\frac{1}{r} \frac{1}{\lambda^r} \sum_{t=1}^{r-1} F_a^i t_{hl}^a \nabla_k \bar{F}_j^{r-t} \\ = -\frac{1}{r} \frac{1}{\lambda^r} \sum_{t=1}^{r-1} F_a^i t_{hl}^a \left( \frac{\partial \bar{F}_j^{r-t}}{\partial x^k} + \Gamma_{bk}^h \bar{F}_j^{r-t} - \Gamma_{jk}^b \bar{F}_b^{r-t} \right).$$

Thus from (4.10), (4.11), (4.12), (4.13) and (4.14) we have

$$(4.15) \quad \Gamma_{jkl}^i = \frac{1}{r} \frac{1}{\lambda^r} t_{kl}^m \sum_{t=1}^{r-1} \left( \frac{\partial F_j^a}{\partial x^m} \right)^{r-t} \bar{F}_a^i - \Phi_{aj}^{ib} \left( \frac{\partial t_{kl}^a}{\partial x^b} \right) \\ - \frac{1}{r} \frac{1}{\lambda^r} \left\{ \sum_{t=1}^{r-1} F_a^i t_{hl}^a \frac{\partial \bar{F}_j^{r-t}}{\partial x^k} - \sum_{t=1}^{r-1} F_a^i t_{hk}^a \frac{\partial \bar{F}_j^{r-t}}{\partial x^l} \right\}.$$

Since  $t_{jk}^i$  is constructed only from  $\bar{F}_j^s$  and their first order partial derivatives, it is seen that  $\Gamma_{jkl}^i$  is expressed by  $\bar{F}_j^s$  and their partial derivatives up to the second order.

Therefore, the tensor  $\nabla_{lm} t_{jk}^i$  is constructed only from  $\bar{F}_j^s$  ( $s = 1, \dots, r-1$ ) and their first and second order partial derivatives.

5. As an example we consider now an almost contact manifold [4] (or a manifold with a  $(\phi, \xi, \eta)$ -structure), that is one over which there exists a tensor field  $\phi_j^i$ , a contravariant vector field  $\xi^i$  and a covariant vector field  $\eta_j$  such that the following conditions are satisfied:

$$(5.1) \quad \text{rank } |\phi_j^i| = 2n, \\ (5.2) \quad \phi_j^i \xi^j = 0, \quad \phi_j^i \eta_i = 0, \\ (5.3) \quad \xi^i \eta_i = 1, \\ (5.4) \quad \phi_j^i \phi_k^j = -\delta_k^i + \xi^i \eta_k.$$

It is known [5] that such manifold has a  $3\pi$ -structure whose fundamental tensor is given as follows:

$$(5.5) \quad \begin{cases} F_j^i = \frac{1}{2} (-\delta_j^i + 3\xi^i \eta_j - \sqrt{3} \phi_j^i) \\ \bar{F}_j^i = \frac{1}{2} (-\delta_j^i + 3\xi^i \eta_j + \sqrt{3} \phi_j^i) \\ \bar{\bar{F}}_j^i = \delta_j^i. \end{cases}$$

Therefore  $\nabla_{kl} F_j^i = 0$  is equivalent to  $\nabla_{kl} \phi_j^i = 0$ .

In this case we have

$$(5.6) \quad \begin{cases} \Phi_{1aj}^{ib} = \frac{1}{2} (\delta_a^i \delta_j^b - \delta_a^i \xi^b \eta_j - \xi^i \eta_a \delta_j^b + 3\xi^i \eta_a \xi^b \eta_j - \phi_a^i \phi_j^b) \\ \Phi_{2aj}^{ib} = \frac{1}{2} (\delta_a^i \delta_j^b + \delta_a^i \xi^b \eta_j + \xi^i \eta_a \delta_j^b - 3\xi^i \eta_a \xi^b \eta_j + \phi_a^i \phi_j^b) \end{cases}$$

and

$$(5.7) \quad P_{ji}^l = \frac{1}{2} \{ -(\nabla_j \phi_i^h) \phi_h^l - (\nabla_j \xi^l) \eta_i + 2(\nabla_j \eta_i) \xi^l + 3\xi^l \eta_a (\nabla_j \xi^a) \eta_i \}.$$

The torsion tensor of the 3- $\pi$ -structure is given by

$$(5.8) \quad t_{jk}^i = \frac{1}{4} \{ -N_{jk}^i - 3\xi^i (\partial_k \eta_j - \partial_j \eta_k) + 5(N_j \eta_k - N_k \eta_j) \\ + \xi^i \phi_j^p \phi_k^q (\partial_p \eta_q - \partial_q \eta_p) - N_p^i (\phi_j^p \eta_k - \phi_k^p \eta_j) \},$$

where  $\partial_j \eta_k = \frac{\partial \eta_k}{\partial x^j}$  and

$$(5.9) \quad \begin{cases} N_{jk}^i = \phi_k^q (\partial_q \phi_j^i - \partial_j \phi_q^i) - \phi_j^p (\partial_p \phi_k^i - \partial_k \phi_p^i) - \eta_j \partial_k \xi^i + \eta_k \partial_j \xi^i, \\ N_j^i = \xi^q (\partial_q \phi_j^i - \partial_j \phi_q^i) - \phi_j^q \partial_q \xi^i, \\ N_j = \xi^p (\partial_p \eta_j - \partial_j \eta_p). \end{cases}$$

It is easily shown that

$$(5.10) \quad \eta_a t_{hk}^{a\xi^h} = 0.$$

Then  $W_{jkl}^i$  in (4.1) turns out to be the following by virtue of (5.10):

$$(5.11) \quad W_{jkl}^i = -\Phi_{1aj}^{ib} (\nabla_b t_{kl}^a) + \frac{1}{2} \phi_a^i \{ (\nabla_k \phi_j^b) t_{bl}^a - (\nabla_l \phi_j^b) t_{bk}^a \} \\ + \frac{1}{2} \eta_j \{ t_{bl}^i (\nabla_k \xi^b) - t_{bk}^i (\nabla_l \xi^b) \} - \frac{3}{2} \xi^i \eta_j \eta_a \{ t_{bl}^a (\nabla_k \xi^b) \\ - t_{bk}^a (\nabla_l \xi^b) \} + \frac{1}{2} \xi^b \{ t_{bl}^i (\nabla_k \eta_j) - t_{bk}^i (\nabla_l \eta_j) \}.$$

Finally  $\Gamma_{jkl}^i$  in (4.15) is written as follows in this case:

$$(5.12) \quad \Gamma_{jkl}^i = -\Phi_{1aj}^{ib} \frac{\partial t_{kl}^a}{\partial x^b} - \frac{1}{2} \phi_a^i \frac{\partial \phi_j^a}{\partial x^m} t_{kl}^m - \frac{1}{2} \eta_j \frac{\partial \xi^i}{\partial x^m} t_{kl}^m + \xi^i \frac{\partial \eta_j}{\partial x^m} t_{kl}^m \\ + \frac{3}{2} \xi^i \eta_a \eta_j \frac{\partial \xi^a}{\partial x^m} t_{kl}^m + \frac{1}{2} \phi_a^i \frac{\partial \phi_j^b}{\partial x^k} t_{bl}^a - \frac{1}{2} \phi_a^i \frac{\partial \phi_j^b}{\partial x^l} t_{bk}^a \\ + \frac{1}{2} \eta_j t_{hl}^i \frac{\partial \xi^h}{\partial x^k} - \frac{1}{2} \eta_j t_{hk}^i \frac{\partial \xi^h}{\partial x^l} - \frac{3}{2} \xi^i \eta_j \eta_a t_{hl}^a \frac{\partial \xi^h}{\partial x^k}$$

$$+ \frac{3}{2} \xi^i \eta_j \eta_a t_{hk}^a \frac{\partial \xi^h}{\partial x^l} + \frac{1}{2} \xi^b t_{bl}^i \frac{\partial \eta_j}{\partial x^k} - \frac{1}{2} \xi^b t_{bk}^i \frac{\partial \eta_j}{\partial x^l}.$$

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