# CLASSICAL SOLUTIONS OF THE STEFAN PROBLEM 

Ei-Ichi Hanzawa

(Received August 22, 1979, revised October 13, 1980)

Introduction. The purpose of the present paper is to prove the local (in-time) existence of the classical solutions for the initial value problem of the one-phase multidimensional Stefan problem, by using Nash's implicit function theorem.

The Stefan problem is a mathematical model for melting of a body of ice in contact with water. The initial value problem for the one-phase Stefan problem is formulated as follows:
(S-1) The unknowns are the thermal distribution in water and the shape of ice. The initial data are prescribed.
(S-2) The temperature of ice is maintained at $0^{\circ} \mathrm{C}$. (The problem in which one considers the thermal distribution in ice is called the twophase problem, which is not discussed in this paper.)
(S-3) The thermal distribution $u$ satisfies the heat equation $\left(\partial_{t}-\Delta\right) u=0$, where $t$ is the time variable and $\Delta$ stands for the Laplacian with respect to the space variables $x=\left(x_{1}, \cdots, x_{n}\right)$.
(S-4) The body of ice melts, at each point of the interface, with velocity in proportion to the normal gradient of $u$. The locus of the interface in the ( $x, t$ )-space is the free boundary to be determined.
(S-5) The region occupied by water has possibly another connected component of the boundary. This component is fixed as $t$ varies, and the heat may be supplied through it. The temperature is always nonnegative.
This is a naive and typical free boundary problem, posed by Stefan [36][39].

In the one dimensional case, this problem (and also the two-phase problem) has been extensively studied. The problem of existence and uniqueness of the classical solution was settled by Rubinstein [27], [28]. It has also been proved that the classical solutions exist globally in time, for the initial and boundary data in various classes of function spaces (Rubinstein [29]; Friedman [10]; Cannon and Hill [5]; Cannon, Hill and Primicerio [6]; Cannon and Primicerio [7]). An excellent historical survey for the result before 1967 is provided by a monograph by Rubinstein [30]. See also Nogi [25] and Yamaguti and Nogi [40].

On the other hand, in the multidimensional case, as the ice melts, it may possibly break into two or more pieces in a finite time. This means that the classical solutions may not exist for all time in general, even if the given data are sufficiently smooth.

Let us briefly refer to studies in the multidimensional one-phase Stefan problem. The classical solutions are not expected to exist for all time. This fact motivates the study of the solutions in a generalized sense, i.e., the weak solutions. In [17], Kamenomostskaja introduced the notion of the weak solution of this problem, and proved its global existense and uniqueness. Her work was generalized by Oleinik [26] and Friedman [11]. The formulation of the problem as a parabolic variational inequality was initiated by Duvaut [9]. This method was developed by Friedman and Kinderlehrer [12], Caffarelli [1]-[3], Caffarelli and Friedman [4] and Kinderlehrer and Nirenberg [18], [19]. In [1]-[4], the Lipschitz continuity of the free boundary and the continuity of the thermal distribution up to the free boundary have been proved. Since the free boundary of the classical solution should be of $C^{1}$-class and the thermal distribution $u$ of the classical solution should have the derivatives $\partial_{x_{i}} u$ which are continuous up to the free boundary, the type of such conclusions as in [1]-[4] is slightly weaker than the required one. In [12] and [19], a case in which we can obtain the $C^{\infty}$ solution is posed. To formulate the problem as a variational inequality, one needs the positivity of the initial and boundary data. Further, in order to obtain the smoothness result in [12] and [19], they need a restrictive geometrical assumption on the initial and boundary data which assures that the melting is rapid and free from the breaking (see [12] and [19]). With assumptions of such kind, one may get around the difficulty, explained later in this introduction, in the multidimensional Stefan problem. In the case in [12] and [19], however, the smoothness up to the initial time was not proved.

What we do in the present paper is to construct the classical solutions in a sufficiently small time interval in general. Our proof has an advantage in revealing the character of the difficulty in the multidimensional problem.

In order to state our result, we introduce the following notations. Let $\Omega_{0}$ be a bounded domain in $R^{n}, n \geqq 2$, with $C^{\infty}$ boundary. The domain $\Omega_{0}$ is regarded as a region occupied by water. Suppose the boundary $\partial \Omega_{0}$ has two connected components $\Gamma_{0}$ and $J_{0}$, where the exterior boundary $\Gamma_{0}$ is in contact with ice, and the heat is supplied through the interior boundary $J_{0}$. Given $0<T_{0}<\infty$, we set, for $0<T<T_{0}, \Omega_{T}=\Omega_{0} \times[0, T]$,
$J_{T}=J_{0} \times[0, T], \quad \Gamma_{T}=\Gamma_{0} \times[0, T]$. As the ice melts, the interface $\Gamma_{0}$ varies and forms a free boundary, which will be diffeomorphic to $\Gamma_{T}$, as long as $T$ is small enough. We shall parametrize this free boundary by the distance function $\rho$ from $\Gamma_{0}$ (in $\boldsymbol{R}^{n}$ ), and denote it by $\Gamma_{\rho, T}$. The corresponding space-time domain in $\boldsymbol{R}^{n} \times \boldsymbol{R}$, which will be diffeomorphic to $\Omega_{T}$, is denoted by $\Omega_{\rho, T}$. The one-phase Stefan problem is a problem determining the free boundary $\Gamma_{\rho, T}$ and the thermal distribution $u$ in the region $\Omega_{\rho, T}$ occupied by water. By the preceding formulation (S-1)-(S-5), we are led to the following equations for $(\rho, u)$ :

$$
\begin{equation*}
\left(\partial_{t}-\Delta\right) u=0 \quad \text { in } \quad \Omega_{\rho, T} \tag{u}
\end{equation*}
$$

$$
\begin{gather*}
\left.u\right|_{t=0}=a_{0} .  \tag{u}\\
u=b_{0} \quad \text { on } \quad J_{T} .  \tag{u}\\
u=0 \quad \text { on } \quad \Gamma_{\rho, T} . \tag{u}
\end{gather*}
$$

$$
\begin{equation*}
\partial_{t} \Phi_{\rho}-c_{0}\left\langle\operatorname{grad} \Phi_{\rho}, \operatorname{grad} u\right\rangle=0 \quad \text { on } \quad \Gamma_{\rho, T} . \tag{u}
\end{equation*}
$$

Here $a_{0}$ (resp. $b_{0}$ ) is a nonnegative function on $\Omega_{0}$ (resp. $J_{T_{0}}$ ), $\Phi_{\rho}$ is a defining function of $\Gamma_{\rho, T}$ and $c_{0}$ is a positive (because of its physical meaning) constant. The equation $\left(5_{u}\right)$ is the so-called Stefan condition which makes the problem complicated. Now our result is as follows (for the precise statement, see Section 1).

Theorem. If $a_{0}$ and $b_{0}$ are sufficiently smooth and satisfy some compatibility conditions, then, for a sufficiently small T, there exists a classical solution ( $\rho, u$ ) of $\left(1_{u}\right)-\left(5_{u}\right)$.

The above theorem is obtained by using Nash's implicit function theorem. Recently, several articles have appeared on the applications of Nash's implicit function theorem (Guillemin [13], in differential geometry; Hamilton [15], in complex analysis; Hörmander [16] and Schaeffer [31]-[33], in free boundary problems; Klainerman [20], in the theory of nonlinear wave equations; Zehnder [41], in Hamiltonian mechanics). However, Nash's theorem and its character are not so popular yet. Hence, before explaining why and how this theorem is used in our proof, let us sketch the idea of this theorem.

The classical implicit function theorem or the inverse function theorem in the finite dimensional Euclidean space asserts the existence of the local inverse of a smooth mapping with the nonzero Jacobian determinant. This theorem extends to a Fréchet differentiable, i.e., a linearizable, mapping between infinine dimensional Banach spaces, provided the Fréchet derivative, i.e., the linearized operator, has the bounded inverse, namely,

Banach's implicit function theorem. However, it is sometimes possible that the linearized operator has merely an unbounded inverse. In the case that the Banach spaces considered are certain function spaces, e.g., the set of functions of $C^{k}$-class, and that the given mapping corresponds to a nonlinear differential equation, such a phenomenon happens when the operator solving the linearized problem does not gain the differentiability of the same order as that of the equation. In such a case, the original problem is said to have a derivative loss. Nash encountered such a problem with derivative loss in his work [23] on the isometric embedding problem for Riemannian manifolds and overcame this difficulty by establishing a new method, Nash's implicit function theorem.

The statement of Nash's implicit function theorem is somewhat complicated. The following is the simplified version due to Moser [22]. Consider two finite scales of Banach spaces $E_{0} \supset E_{1} \supset \cdots \supset E_{k}$ and $F_{1} \supset F_{2} \supset \cdots \supset F_{k}$, e.g., $E_{i}=C^{m i+p}, F_{i}=C^{m i+q}$ with $p>q$ and $m>0$. Let $\mathscr{F}$ be a nonlinear operator defined in a neighborhood $V$ of 0 of $E_{1}$ into $F_{1}$ such that $\mathscr{F}(\rho) \in F_{i}$ for $\rho \in V \cap E_{i}$. Nash's theorem asserts that the equation $\mathscr{F}(\rho)=0$ has a solution $\rho \in V$, provided the following assumptions are satisfied (for the exact statement, see Section 5):
(N-1) An operator $\mathscr{S}_{\theta}: E_{0} \rightarrow E_{k}$ with parameter $\theta \geqq 1$ is defined so that for $0 \leqq i \leqq j \leqq k$,

$$
\left|\mathscr{S}_{\theta} f\right|_{E_{j}} \leqq C \theta^{j-i}|f|_{E_{i}}, \quad\left|f-\mathscr{S}_{\theta} f\right|_{E_{i}} \leqq C \theta^{-(j-i)}|f|_{E_{j}}
$$

This $\mathscr{S}_{\theta}$ is called a smoothing operator.
(N-2) The operator $\mathscr{F}$ is Fréchet differentiable, i.e., $\mathscr{F}$ is linearizable.
(N-3) The linearized equation $D \mathscr{F}(\rho) \delta \rho=\delta G$ can be solved possibly with derivative loss for each $\rho \in V$, i.e., there exists a linear operator $\mathcal{J}(\rho): F_{i} \rightarrow E_{i-1}$ for $i \geqq 1$ and $\rho \in V \cap E_{i}$, which is a right inverse to $D \mathscr{F}(\rho)$.
(N-4) For $i \geqq 1$ and $\rho \in V \cap E_{i}$,

$$
|\mathscr{J}(\rho) \mathscr{F}(\rho)|_{E_{i-1}} \leqq C\left(1+|\rho|_{E_{i}}\right) .
$$

(N-5) The value $|\mathscr{F}(0)|_{F_{1}}$ is sufficiently small.
The assumptions ( $\mathrm{N}-2,3,5$ ) will be in no need of explanation. In contrast to these, it seems that the assumptions (N-1) and (N-4) are artificial and indistinct, though it is obvious from the proof of Nash's theorem ([16]; [22]; [23]; Schwartz [34]; Sergeraert [35]) that they are essential and indispensable. Later in this introduction, we will give a short account for ( $\mathrm{N}-1$ ) and ( $\mathrm{N}-4$ ). Further, we add that the present paper has two mathematical cores, which are to solve the linearized Stefan
problem and to introduce a logical frame of a general character in which the assumptions ( $\mathrm{N}-1$ ) and ( $\mathrm{N}-4$ ) can be verified.

Returning to our Stefan problem, let us sketch briefly how the problem is rewritten as a nonlinear operator equation $\mathscr{F}(\rho)=0$ and the essential feature of solving the linearized Stefan problem $D \mathscr{F}(\rho) \delta \rho=\delta G$. Regarding the thermal distribution $u$ as an auxiliary unknown determined from the distance function $\rho$ with $\left(1_{u}\right)-\left(4_{u}\right)$ and introducing an operator $\mathscr{F}$ which transforms $\rho$ to the pull back of the function $\partial_{t} \Phi_{\rho}-c_{0}\left\langle\operatorname{grad} \Phi_{\rho}\right.$, grad $u\rangle$ on the free boundary $\Gamma_{\rho, T}$ to the function on the flattened boundary $\Gamma_{T}$, we are led to a nonlinear equation $\mathscr{F}(\rho)=0$ which is equivalent to $\left(1_{u}\right)-\left(5_{u}\right)$. This equation can be linearized and the concrete form of the linearized problem $D \mathscr{F}(\rho) \delta \rho=\delta G$ consists of two parts, as follows. The first part, corresponding to $\left(1_{u}\right)-\left(4_{u}\right)$, is an initial boundary value problem of the Dirichlet type for a linear second order parabolic equation in $\Omega_{T}$, where the unknown is an auxiliary one corresponding to the formal Fréchet derivative of $u$, and given data are $\rho$ and $\delta \rho$. The second part, corresponding to the Stefan condition $\left(5_{u}\right)$, is a linear first order equation of hyperbolic type for $\delta \rho$ on the flattend boundary $\Gamma_{T}$, with the data containing the normal derivative of the unknown of the first part. To solve the above linearized problem, we eliminate $\delta \rho$ from the system, by substituting the solution $\delta \rho$ of the latter hyperbolic equation for the data $\delta \rho$ in the former initial boundary value problem of parabolic type. Then, introducing a new unknown $\delta X$ for convenience, we find that the essential point in solving the above system without $\delta \rho$ is to solve a linear parabolic initial boundary value problem for $\delta X$ in $\Omega_{T}$ whose boundary condition on $\Gamma_{T}$ is given by a linear hyperbolic first order equation. We extend this first order operator to one in the domain $\Omega_{T}$. Then, to invert this first order operator, we need the nonnegativity of the coefficients of the normal derivation in this operator on the boundary $\Gamma_{T}$. In fact, this condition assures that the characteristic curves starting from points on the domain $\Omega_{0}$ at the initial time cover the cylinder $\Omega_{T}$. In our case, by using the maximum principle of the heat equation, we can prove the required nonnegativity because the temperature on $\Omega_{0}$ at the initial time and on the fixed boundary $J_{T}$ are nonnegative (namely, this physical requirement has also a natural mathematical meaning). Then, after a technical deformation, the above problem for $\delta X$ is solved by decomposing it to a parabolic initial boundary value problem of Dirichlet type and an initial value problem for the above first order operator in $\Omega_{T}$. Consequently, the solution $\delta \rho$ of the linearized Stefan problem $D \mathscr{F}(\rho) \delta \rho=\delta G$ is obtained as a linear combination containing the normal
derivative of $\delta X$.
Since, except in the one-dimensional case, the operator $\delta G \mapsto \delta X$ gains only the same regularity in the weighted Hölder spaces, standard in the theory of parabolic equations, the linearized Stefan problem can be solved, but actually with an essential derivative loss. This is the reason why we need Nash's implicit function theorem.

From the above consideration on solving the linearized Stefan problem, we find that the reason why derivative loss occurs in the multidimensional Stefan problem consists in the following fact: Since the mapping solving the linear first order equation of hyperbolic type which corresponds to the Stefan condition $\left(5_{u}\right)$ gains the regularity only along the characteristic curves, it cannot cover the bad influence of the diffusion effect, i.e., the normal derivative of $\delta X$. This seems to be a new observation showing why the multidimensional Stefan problem is difficult.

In the one-dimensional case, since direction of the diffusion can be covered by the characteristic curve, we can solve the linearized problem without derivative loss. Therefore, the local existence theorem for classical solutions is obtained by usual Banach's implicit function theorem, because, in an appropriate setting, the norm $|\mathscr{F}(0)|$ can be taken small if $T$ is sufficiently small.

Now, we account for the remaining assumptions (N-1) and (N-4). The general principle to verify the assumption (N-4), i.e., the estimate $|\mathscr{J}(\rho) \mathscr{F}(\rho)|_{i-1} \leqq C\left(1+|\rho|_{i}\right)$, is as follows: This estimate is automatically obtained if we prove that each of all mathematical operations constructing $\mathscr{F}$ and $\mathscr{F}$ is a "balanced" operator. Here the meaning of the term "balanced" will be illustrated by the following examples. In the weighted Hölder spaces, addition, multiplication, division, composition and to take the solution $u$ of a parabolic initial boundary value problem

$$
\left(\partial_{t}-\sum A_{i j} \partial_{x_{i}} \partial_{x_{j}}-\sum A_{i} \partial_{x_{i}}-A_{0}\right) u=f \text { in } \quad \Omega_{T},\left.\quad u\right|_{t=0}=a,\left.\quad u\right|_{\partial \Omega_{T}}=b,
$$

are balanced operations, because we have estimates of the following type:

$$
\begin{aligned}
& |f+g|_{(r)} \leqq|f|_{(r)}+|g|_{(r)}, \quad|f g|_{(r)} \leqq C\left(|f|_{(r)}|g|_{(0)}+|f|_{(0)}|g|_{(r)}\right), \\
& |f / g|_{(r)} \leqq C\left[|f|_{(r)}+\left(1+|g|_{(r)}\right)|f|_{(0)}\right], \quad \text { if } \quad \text { inf } g \geqq B>0, \\
& |f \circ g|_{(r)} \leqq C\left(|f|_{(r)}+|f|_{(1)}|g|_{(r)}\right), \quad \text { if } \quad r \geqq 1 \quad \text { and }|g|_{(1)} \leqq B, \\
& |u|_{(r+2)} \leqq C\left[\left(\sum\left|A_{i j}\right|_{(r)}+\sum\left|A_{i}\right|_{(r)}+\left|A_{0}\right|_{(r)}\right)\left(|f|_{(\varepsilon)}+|a|_{(\varepsilon+2)}+|b|_{(\varepsilon+2)}\right)\right. \\
& \left.+|f|_{(r)}+|a|_{(r+2)}+|b|_{(r+2)}\right], \\
& \quad \text { if } r \geqq \varepsilon>0, \quad r \text { and } \varepsilon \text { are not integers and } \\
& \quad\left|A_{i j}\right|_{(\varepsilon)},\left|A_{i}\right|_{(\varepsilon)}, \quad\left|A_{0}\right|_{(\varepsilon)} \leqq B,
\end{aligned}
$$

where $\left|\left.\right|_{(r)}\right.$ is the weighted Hölder norm of $r$-th order. Now the preceding "general principle" is an ovbious fact. This "principle" is consciously used, e.g., in [16] and [31]-[33]. Further, e.g., in [16], in place of the parabolic boundary value problem and the weighted Hölder norm, with an elliptic boundary value problem and the usual Hölder norm, $\mathscr{F}^{-}$and $\mathscr{F}$ are constructed only by operations listed in the above examples. Then, with the fact that to construct smoothing orerators on a usual Hölder space is not so difficult, the above consideration is sufficient to apply Nash's theorem to the problem in [16]. Moreover, it seems that the above consideration is also sufficient for the other applications of Nash's theorem listed above in this introduction.

Our new trouble on the verification of ( $\mathrm{N}-4$ ) in the Stefan problem lies in the fact that, to invert the linearized problem, we had to take the Neumann series to solve linear integral equations of Volterra type several times. If an operator norm is less than 1 and we can complete the proof by taking Neumann series only once, then to prove (N-4) is not so difficult (see [33, Lemma 8.1]); however, both cannot be expected in our proof. The method adopted in the present paper is as follows. Instead of the usual weighted Hölder spaces $C^{(r)}$, ( $r$ denotes the order of regularity), we can use $C_{\ddagger}^{(r)}$ spaces each of which consists of all $C^{(r)}$ functions whose derivatives up to $r$-th order are vanishing at the initial time, because the considered problem is a linearized one, i.e., a variational one. In $C_{\#}^{(r)}$ spaces, the integral operator $\int_{0}^{t}$ commutes with the norm, i.e., $\left|\int_{0}^{t} f\right|_{(r), T} \leqq \int_{0}^{T}|f|_{(r), t} d t$, where the subscripts $T$ and $t$ denote the width of the time intervals. It should be noted that this commutativity does not hold in $C^{(r)}$ spaces. Now, in $C_{\sharp}^{(r)}$ spaces, by using the iterated estimate $\left|\int_{0}^{t} \cdots \int_{0}^{t} f\right|_{(r)} \leqq\left(t^{n} / n!\right)|f|_{(r)}$, it is obvious from the standard argument that the operator solving a linear integral equation of Volterra type is a balanced one.

Then, we account for the assumption (N-1) on the existence of smoothing operators. Since we are working in the setting of $C_{*}^{(r)}$ spaces, we have to construct smoothing operators on $C_{\#}^{(r)}$ spaces. It is known that for $C^{m}, m$ is an integer, Sobolev and usual Hölder spaces, smoothing operators can be constructed as integral operators whose kernels are defined by using Fourier transform (see, e.g., [16] and [23]). In the same method, we can construct smoothing operators on $C^{(r)}$ spaces; however, these operators do not preserve the $C_{\sharp}^{(r)}$ property, i.e., the image of a $C_{\#}^{(r)}$ function is not necessarily a $C_{\#}^{(r)}$ function. Therefore the above method cannot be used in our setting. On the other hand, in his work on the
isometric embedding problem for analytic Riemannian manifolds [24], in order to get smoothing operators which do not shrink the radii of convergence of real analytic functions at each point in the domain, Nash introduced a new method to construct smoothing operators. In [24], smoothing operators are constructed as integral operators, but their kernels are defined not by using Fourier transform, but by taking a linear combination of functions each of which is obtained from the heat kernel with an appropriate coordinate transform. (See also Gromov [14, Section 3] whose exposition may be more precise than that in [24].) Then, though the appearance of the problem is different from that in [24] and we must take a $C_{\#}^{(\infty)}$ function instead of the heat kernel, we can construct smoothing operators on $C_{\sharp}^{(r)}$ spaces by using the method in [24].

The finite scale with $C_{\sharp}^{(r)}$ spaces and the smoothing operators constructed in the above method constitute the logical frame, announced in this introduction, in which the operator solving a linear integral equation of Volterra type becomes a balanced one. This frame seems to be natural and useful in applying Nash's theorem to initial value problems.

The outline of this paper is as follows. In Section 1, we give the exact statement of our result and define the weighted Hölder spaces with which our theorem is formulated. In Section 2, we add a technical assumption to our theorem to simplify the account of the proof and restate our theorem by using some notations which are also introduced in this section. The elimination of this technical assumption will be carried out in Section 10. In Section 3, we collect together some fundamental lemmas which assure that multiplication, division, composition and to take the solution of a parabolic initial boundary value problem are balanced operators, in the weighted Hölder spaces. In Section 4, we introduce a nonlinear operator $\mathscr{F}$ by which Theorem' in Section 2 is restated as the equation $\mathscr{F}(\rho)=0$. We apply Nash's implicit function theorem to this operator equation. In Section 5, we state Nash's theorem and pose two finite scales of Banach spaces on which Nash's theorem is applied. In Section 6, smoothing operators in our setting are constructed; in other words, the assumption (N-1) is verified. In Section 7, the Fréchet differentiability of $\mathscr{F}$, i.e., the fact that the problem is linearizable, is proved; in other words, the assumption (N-2) is verified. In Section 8, we again collect together some technical lemmas, which include two key lemmas to our proof. One of the two lemmas assures that in $C_{\ddot{\#}}^{(r)}$ spaces, the operator solving a linear integral equation of Volterra type is a balanced one (see Lemma 8.A.4). The other lemma asserts that we can
solve a linear parabolic initial boundary value problem whose boundary condition is given by a first order equation such that the coefficient of the normal derivative is nonnegative. This lemma also asserts that the operator solving the problem is a balanced one (see Lemma 8.B.9). In Section 9, by using the lemmas in Section 8, we solve the linearized problem $D \mathscr{F}(\rho) \delta \rho=\delta G$ and observe that the solution $\delta \rho$ satisfies the assumptions ( $\mathrm{N}-3$ ) and ( $\mathrm{N}-4$ ). Section 9 contains the key fact which enables us to prove the result of this paper, that is, the fact that the coefficient corresponding to the one in the lemma in Section 8 on the solvability of a parabolic initial boundary value problem is, in fact, nonnegative (see (9.16)). At the end of Section 9, the assumption (N-5) is verified, so that the proof of Theorem' in Section 2 is completed. In Section 10, we sketch how our proof is modified when the technical assumption in Section 2 is replaced by the general compatibility condition in Theorem in Section 1.

The author would like to thank the referee for his advice on the style of this paper and on the literature on the Stefan problem.

1. The result. In this section, we state our result. At the beginning, we prepare some notations.
(A) Let $n$ be an integer with $n \geqq 2$. Let $\Omega_{0}$ be a bounded domain in $\boldsymbol{R}^{n}$ whose boundary consists of the outside component $\Gamma_{0}$ and the inside one $J_{0}$. Suppose $\Gamma_{0}$ and $J_{0}$ are $C^{\infty}$. Let $T_{0}$ be a positive constant. For $T \in\left(0, T_{0}\right]$, let $\Omega_{T}=\Omega_{0} \times[0, T], \Gamma_{T}=\Gamma_{0} \times[0, T]$, and $J_{T}=J_{0} \times[0, T]$.
(B) We define the Hölder spaces.

Definition 1.B.1. Let $d$ be a positive integer. Let $D$ be a domain with $C^{\infty}$ boundary in $\boldsymbol{R}^{d}$ or $D=\boldsymbol{R}^{d}$. Let $r \geqq 0$. Then $C^{r}(\bar{D})$ is the set of real-valued functions $f$ on $\bar{D}$ such that:
(i) The derivatives $\partial_{x}^{\alpha} f$ with $|\alpha| \leqq[r]$ are continuously extended to $\bar{D}$, where $\alpha=\left(\alpha_{1}, \cdots, \alpha_{n}\right)$ denotes a $d$-tuple of non-negative integers, $\partial_{x}=\partial_{x_{1}}^{\alpha_{1}} \cdots \partial_{x_{n}}^{\alpha_{n}},|\alpha|=\sum_{i=1}^{d} \alpha_{i}$, and $[r]$ is the greatest integer not greater than $r$.
(ii) The norm $|f|_{r}$ is finite. Here $|f|_{r}$ is defined by:
(i) $|f|_{r}=\sum_{|\alpha| \leq r} \sup _{x \in \bar{D}}\left|\partial_{x}^{\alpha} f(x)\right|$ if $r$ is an integer.
(ii) $|f|_{r}=\sum_{|\alpha| \leqslant[r]} \sup _{x \in \bar{D}}\left|\partial_{x}^{\alpha} f(x)\right|+\sum_{|\alpha|=[r]} \sup _{x, y \in \bar{D}}\left|\partial_{x}^{\alpha} f(x)-\partial_{x}^{\alpha} f(y)\right| / \mid x-$ $\left.y\right|^{r-[r]}$ if $r$ is not an integer where $|x-y|$ denotes $\left(\sum_{i=1}^{d}\left|x_{i}-y_{i}\right|^{2}\right)^{1 / 2}$.

Definition 1.B.2. Let $d$ be a positive integer. Let $I$ be an interval in $\boldsymbol{R}$. Let $D$ be a domain with piecewise- $C^{1}$ boundary in $\boldsymbol{R}^{d} \times I$. Let $r \geqq 0$. Then $C^{(r)}(\bar{D})$ is the set of real-valued functions $f$ on $\bar{D}$ such that:
(i) The derivatives $\partial_{x}^{\alpha} \partial_{t}^{a} f(x, t)$ with $|\alpha|+2 a \leqq[r]$ are continuously extended to $\bar{D}$.
(ii) The norm $|f|_{(r)}$ is finite.

The norm $|f|_{(r)}$ is defined by:
(i) $|f|_{(r)}=\sum_{i=0}^{r}\langle f\rangle_{(i)}$ if $r$ is an integer.
(ii) $|f|_{(r)}=\sum_{i=0}^{[r]}\langle f\rangle_{(i)}+\langle f\rangle_{(r)}$ if $r$ is not an integer.

Here $\langle f\rangle_{(i)}$ and $\langle f\rangle_{(r)}$ are defined by:
(i) $\langle f\rangle_{(i)}=\sum_{|\alpha|+2 a=i} \sup _{(x, t) \in \bar{D}}\left|\partial_{x}^{\alpha} \partial_{t}^{a} f(x, t)\right|+\sum_{|\alpha|+2 a=i-1} \sup _{(x, t),(x, s) \in \bar{D}}$ $\left|\partial_{x}^{\alpha} \partial_{t}^{a} f(x, t)-\partial_{x}^{\alpha} \partial_{t}^{a} f(x, s)\right| /|t-s|^{1 / 2}$ if $i$ is an integer.
(ii) $\langle f\rangle_{(r)}=\sum_{|\alpha|+2 a=[r]} \sup _{(x, t),(y, t) \in \bar{D}}\left|\partial_{x}^{\alpha} \partial_{t}^{z} f(x, t)-\partial_{x}^{\alpha} \partial_{t}^{a} f(y, t)\right| /|x-y|^{r-[r]}+$ $\sum_{|\alpha|+2 a=[r]} \sup _{(x, t),(x, s) \in \bar{D}}\left|\partial_{x}^{\alpha} \partial_{t}^{a} f(x, t)-\partial_{x}^{\alpha} \partial_{t}^{a} f(x, s)\right| /|t-s|^{(r-[r]) / 2}+\sum_{|\alpha|+2 a=[r]-1}$ $\sup _{(x, t),(x, s) \in \bar{D}}\left|\partial_{x}^{\alpha} \partial_{t}^{a} f(x, t)-\partial_{x}^{\alpha} \partial_{t}^{a} f(x, s)\right| /|t-s|^{(r-[r]+1) / 2}$ if $r$ is not an integer.

Definition 1.B.3. Let $M$ be a compact $C^{\infty}$ manifold (with or without boundary). Let $I$ be an interval in $\boldsymbol{R}$. Let $r \geqq 0$. Then the set $C^{(r)}(M \times I)$ and the norm $\left|\left.\right|_{(r)}\right.$ in it are defined by using a finite covering of $M$ by coordinate neighborhoods and a $C^{\infty}$ partition of unity subordinate to it.

Remark 1.B.4. All the normed spaces in Definitions 1.B.1-1.B. 3 are Banach spaces.
(C) Let $n_{\omega}$ be the outward unit normal at $\omega \in \Gamma_{0}$. Let $\gamma_{0}$ be a positive constant so small that a mapping $x: \Gamma_{0} \times\left[-\gamma_{0}, \gamma_{0}\right] \rightarrow \boldsymbol{R}^{n}$ defined by $(\omega, \lambda) \mapsto \omega+\lambda n_{\omega}$ is regular and one-to-one. Let

$$
N_{0}=\left\{x(\omega, \lambda) ;(\omega, \lambda) \in \Gamma_{0} \times\left[-\gamma_{0}, \gamma_{0}\right]\right\}
$$

and

$$
N_{0}^{-}=\left\{x(\omega, \lambda) ;(\omega, \lambda) \in \Gamma_{0} \times\left[-\gamma_{0}, 0\right]\right\}
$$

We denote the inverse mapping of $x$ on $N_{0}$ onto $\Gamma_{0} \times\left[-\gamma_{0}, \gamma_{0}\right]$ by $x \mapsto$ $(\omega(x), \lambda(x))$. Clearly the mappings $x(\omega, \lambda), \omega(x)$ and $\lambda(x)$ are $C^{\infty}$. We often use $\left(\omega_{1}(x), \cdots, \omega_{n-1}(x)\right)$ as local coordinates of $x \in N_{0}$, where $\omega_{i}$ is the $i$-th component of $\omega$ with respect to local coordinates in $\Gamma_{0}$.
(D) For $\rho \in C^{0}\left(\Gamma_{T}\right)$ with $|\rho|_{0}<\gamma_{0}$, let

$$
\Gamma_{\rho, T}=\left\{(x(\omega, \rho(\omega, t)), t) ;(\omega, t) \in \Gamma_{r}\right\}
$$

and let $\Omega_{\rho, T}$ be the domain in $R^{n} \times[0, T]$ bounded by $\Gamma_{\rho, T}$ and $J_{T}$. For $\rho \in C^{0}\left(\Gamma_{T}\right)$ with $|\rho|_{0}<\gamma_{0}$ and for $(x, t) \in N_{0} \times[0, T]$, let

$$
\begin{equation*}
\Phi_{\rho}(x, t)=\lambda(x)-\rho(\omega(x), t) \tag{1.1}
\end{equation*}
$$

Note that

$$
\Gamma_{\rho, T}=\left\{(x, t) \in N_{0} \times[0, T] ; \Phi_{\rho}(x, t)=0\right\}
$$

(E) Now our result can be stated as follows.

Theorem. Let $r_{0}=n_{0}+\varepsilon_{0}$ where $n_{0}$ is an integer with $n_{0} \geqq 7$ and $0<\varepsilon_{0}<1$. Suppose that:
(A.1) $\quad a_{0} \in C^{r_{0}+43}\left(\bar{\Omega}_{0}\right)$ and $b_{0} \in C^{\left(r_{0}+39\right)}\left(J_{T_{0}}\right)$.
(A.2) The pair $\left\{a_{0}, b_{0}\right\}$ satisfies the compatibility condition up to order $\left[\left(r_{0}+39\right) / 2\right]$ of the Dirichlet problem for the heat equation, on $J_{0}$.
(A.3) $\quad a_{0} \geqq 0$ on $\bar{\Omega}_{0}, b_{0} \geqq 0$ on $J_{T_{0}}$ and $c_{0}$ is a positive constant.
(A.4) The function $a_{0}$ satisfies the compatibility condition up to order $\left[\left(r_{0}+39\right) / 2\right]$ of the Stefan problem (see Remark 1.E. 4 in the following).
Then, for a sufficiently small $T \in\left(0, T_{0}\right]$, there exist $\rho \in C^{\left(r_{0}\right)}\left(\Gamma_{T}\right)$ with $\left.\rho\right|_{t=0}=0$ and $|\rho|_{0}<\gamma_{0}$ and $u \in C^{\left(r_{0}\right)}\left(\bar{\Omega}_{\rho, T}\right)$ which satisfy:

$$
\begin{equation*}
\left(\partial_{t}-\Delta\right) u=0 \text { in } \Omega_{\rho, T} . \tag{u}
\end{equation*}
$$

$$
\begin{equation*}
u=b_{0} \quad \text { on } \quad J_{T} . \tag{u}
\end{equation*}
$$

$$
\begin{equation*}
\partial_{t} \Phi_{\rho}-c_{0}\left\langle\operatorname{grad} \Phi_{\rho}, \operatorname{grad} u\right\rangle=0 \quad \text { on } \quad \Gamma_{\rho, T} . \tag{u}
\end{equation*}
$$

Here $\Delta=\sum_{i=1}^{n} \partial_{x_{i}}^{2}$, $\operatorname{grad}=\left(\partial_{x_{1}}, \cdots, \partial_{x_{n}}\right)$ and $\langle x, y\rangle=\sum_{i=1}^{n} x_{i} y_{i}$ for $x, y \in \boldsymbol{R}^{n}$.
Remark 1.E.1. The numbers 43 and 39 appearing in our theorem have no specific meaning. They come out because we employ Moser's version of Nash's implicit function theorem [22], which is the most popular one. Nash's original version [23] or Hörmander's version [16] provides smaller numbers, but such a refinement may not be essential.

Remark 1.E.2. Consider the mixed problem of the heat equation $\left(\partial_{t}-\Delta\right) u=0$ with the initial condition $u=a$ and the boundary condition $u=b$ (resp. $\partial_{n} u=b$ where $n$ is the outward unit normal of the boundary). The compatibility condition up to order $i$ of the above problem requires that $\partial_{t}^{j} b=\Delta^{j} a$ (resp. $\partial_{t}^{j} b=\partial_{n} \Delta^{j} a$ ) on the boundary at $t=0$ for $j=0, \cdots, i$. For other mixed problems of parabolic type, analogous definitions are adopted. (For details, see, e.g., Ladyzenskaja, Solonnikov and Uralceva [21, p. 319-320].)

Remark 1.E.3. The assumption (A.3), which is natural in the physical sense, enables us to solve the linearized Stefan problem (see (9.16)).

Remark 1.E.4. Let $i$ be a nonnegative integer and let $(\rho, u)$ be a $C^{(2 i)}$ solution of the above system $\left(1_{u}\right)-\left(5_{u}\right)$ with $\left.\rho\right|_{t=0}=0$. Because of $\left(4_{u}\right)$, on $\Gamma_{0}$ at $t=0$, the time variable derivatives up to order $i$ of $u$
along $\Gamma_{\rho, T}$ vanish. Therefore, we have algebraic relations of $u, \partial_{t} u, \cdots$, $\partial_{t}^{i} u$ and $\partial_{t} \rho, \cdots, \partial_{t}^{i} \rho$ on $\Gamma_{0}$ at $t=0$. On the other hand, in view of $\left(1_{u}\right)$, $\left(2_{u}\right)$ and $\left(5_{u}\right)$, all of the above derivatives are determined only by $a_{0}$. Then, we have algebraic relations of the derivatives $\partial_{x}^{\alpha} a_{0}$ with $|\alpha| \leqq 2 i$ on $\Gamma_{0}$. These relations constitute the compatibility condition up to order $i$ of the Stefan problem.

Remark 1.E.5. As is mentioned in Introduction, the uniqueness of the solution has been established even in the class of weak solutions.

Remark 1.E.6. The author does not know whether we can take $r_{0}=\infty$ in the above theorem. However, when the melting is rapid, our solution ( $\rho, u$ ) is of $C^{\infty}$-class at least for $0<t<T$, according to the regularity theorem for classical solutions by Kinderlehrer and Nirenberg [18], [19].

Remark 1.E.7. The width $T$ of the existence interval is determined only by $\left|a_{0}\right|_{r_{0}+43},\left|b_{0}\right|_{\left(r_{0}+39\right)}$ and the shape of $\Gamma_{0}$.
2. Technical assumption. We shall first prove our theorem, assuming the technical assumption

$$
\begin{equation*}
\partial_{x}^{\alpha} a_{0}=0 \quad \text { on } \quad \Gamma_{0} \quad \text { if } \quad|\alpha| \leqq\left[r_{0}+39\right] . \tag{T}
\end{equation*}
$$

The elimination of this assumption will be carried out in Section 10. In this section, assuming ( $T$ ) and introducing a few new notations, we restate our result in somewhat different form from that in Section 1.
(A) We introduce a new class of Hölder spaces.

Definition 2.A.1. Let $d$ be a positive integer. Let $r \geqq 0$. For $T \in\left(0, T_{0}\right]$ and for a domain $D$ with piecewise- $C^{1}$ boundary in $R^{d} \times[0, T]$, we set

$$
C_{\#}^{(r)}(\bar{D})=\left\{f \in C^{(r)}(\bar{D}) ;\left.\left(\partial_{t}^{a} f\right)\right|_{t=0}=0 \text { for } a=0, \cdots,[r / 2]\right\} .
$$

Further, for $T \in\left(0, T_{0}\right]$ and for a compact $C^{\infty}$ manifold $M$ (with or without boundary), $C_{\sharp}^{(r)}(M \times[0, T])$ is similarly defined. In these cases, when we emphasize that the interval of the "time variable" $t$ is $[0, T]$, we denote the norm $|f|_{(r)}$ by $|f|_{(r), T}$.

Remark 2.A.2. The above normed spaces are Banach spaces.
(B) For $T \in\left(0, T_{0}\right]$, we set

$$
V_{T}=\left\{\rho \in C_{\sharp}^{\left(r_{0}\right)}\left(\Gamma_{T}\right) ;|\rho|_{\left(r_{0}\right)}<\delta_{0}\right\},
$$

where $\delta_{0}$ is a positive constant so small that:
(i) $4 \delta_{0} \leqq \gamma_{0}$.
(ii) There is a positive constant $\sigma_{0}$ such that

$$
\sigma_{0}^{-1} \sum_{i=1}^{n} \xi_{i}^{2} \leqq \sum_{i, j=1}^{n} A_{\rho, i j}(x, t) \xi_{i} \xi_{j} \leqq \sigma_{0} \sum_{i=1}^{n} \xi_{i}^{2}
$$

for $\rho \in V_{T},(x, t) \in \Omega_{T}$, and $\xi \in \boldsymbol{R}^{n}$. Here $A_{\rho, i j}$ are the coefficients of the second order derivations in the operator $\mathscr{L}_{\rho}$ defined in (D) in this section.
(iii) There is a positive constant $B_{0}$ such that

$$
S_{\rho}(\omega, t) \geqq B_{0}^{-1}
$$

for $\rho \in V_{T}$ and $(\omega, t) \in \Gamma_{T}$. Here $S_{\rho}$ is the function defined in (D) in this section.
This constant $\delta_{0}$ is determined only by the shape of $\Gamma_{0}$.
(C) Choose a function $\chi_{0} \in C_{0}^{\infty}(R)$ so that:
(i) $\chi_{0}(\lambda)=1$ if $|\lambda| \leqq \delta_{0}$.
(ii) $\quad \chi_{0}(\lambda)=0$ if $|\lambda| \geqq 3 \delta_{0}$.
(iii) $\left|\partial_{\lambda} \chi_{0}(\lambda)\right| \leqq 3 / 4 \delta_{0}$ for $\lambda \in \boldsymbol{R}$.

For $\rho \in V_{T}$, define a diffeomorphism $e_{\rho}: \boldsymbol{R}^{n} \times[0, T] \rightarrow \boldsymbol{R}^{n} \times[0, T]$ by:

$$
\begin{align*}
& e_{\rho}(x(\omega, \lambda), t)=\left(x\left(\omega, \lambda+\chi_{0}(\lambda) \rho(\omega, t)\right), t\right)  \tag{2.1.i}\\
& \text { for } \quad(x, t)=(x(\omega, \lambda), t) \in N_{0} \times[0, T] . \\
& e_{\rho}(x, t)=(x, t) \quad \text { for } \quad(x, t) \in\left(\boldsymbol{R}^{n}-N_{0}\right) \times[0, T] . \tag{2.1.ii}
\end{align*}
$$

Note that $e_{\rho}\left(\Omega_{T}\right)=\Omega_{\rho, T}, e_{\rho}\left(\Gamma_{T}\right)=\Gamma_{\rho, T}$, and $\left.e_{\rho}\right|_{t=0}$ is the identity mapping.
Define a function $\eta: \boldsymbol{R} \times\left[-\delta_{0}, \delta_{0}\right] \rightarrow \boldsymbol{R}$ by $\eta(\lambda, \mu)+\chi_{0}(\eta(\lambda, \mu)) \mu=\lambda$ for $(\lambda, \mu) \in \boldsymbol{R} \times\left[-\delta_{0}, \delta_{0}\right]$. Since $\left|\partial_{\lambda} \chi_{0}(\lambda)\right| \leqq 3 /\left(4 \delta_{0}\right)$, the function $\eta$ is welldefined and $C^{\infty}$. We easily observe that:

$$
\begin{align*}
& e_{\rho}^{-1}(x(\omega, \lambda), t)=(x(\omega, \eta(\lambda, \rho(\omega, t))), t)  \tag{2.2.i}\\
& \quad \text { for } \quad(x, t)=(x(\omega, \lambda), t) \in N_{0} \times[0, T] . \\
& e_{\rho}^{-1}(x, t)=(x, t) \quad \text { for } \quad(x, t) \in\left(\boldsymbol{R}^{n}-N_{0}\right) \times[0, T] . \tag{2.2.ii}
\end{align*}
$$

(D) We define the operator $\mathscr{L}_{\rho}$ and the function $S_{\rho}$ by:

$$
\begin{equation*}
\mathscr{L}_{\rho} V=\left[\left(\partial_{t}-\Delta\right)\left(V \circ e_{\rho}^{-1}\right)\right] \circ e_{\rho} \quad \text { for } \quad V \in C^{(2)}\left(\bar{\Omega}_{T}\right) \tag{2.3.i}
\end{equation*}
$$

(2.3.ii) $\left[\left\langle\operatorname{grad} \Phi_{\rho}, \operatorname{grad}\left(V \circ e_{\rho}^{-1}\right)\right\rangle\right] \circ e_{\rho}=\left(\partial_{k} V\right) S_{\rho} \quad$ on $\quad \Gamma_{T}$

$$
\text { for } \quad V \in C^{1}\left(\bar{\Omega}_{T}\right) \text { with }\left.\quad V\right|_{\Gamma_{T}}=0 .
$$

For the explicit expressions, see Section 4.
(E) Now our restricted and modified theorem is as follows.

THEOREM'. Assume the assumptions (A.1)-(A.3) and (T). Then, there exist $\rho \in V_{T}$ and $u \in C^{\left(r_{0}\right)}\left(\bar{\Omega}_{\rho, T}\right)$ which satisfy $\left(1_{u}\right)-\left(5_{u}\right)$.

Remark 2.E.1. When we assume (T), we can replace the assumption
$a_{0} \in C^{r_{0}+43}\left(\bar{\Omega}_{0}\right)$ by a weaker one $a_{0} \in C^{r_{0}+39}\left(\bar{\Omega}_{0}\right)$.
3. Technical lemmas I. We collect together some lemmas which are used in the following sections. Throughout this paper, " $X$ is bounded with $Y^{\prime \prime}$ means that $X$ is bounded when $Y$ is.
(A) We begin with the following fact. By the term "a manifold", we mean both one with boundary and one without boundary.

Lemma 3.A.1. Let $M$ be a compact $C^{\infty}$ manifold. Let $\gamma<\delta$ and $0 \leqq p \leqq q \leqq r$. Let $f$ belong to $C^{(r)}(M \times[\gamma, \delta])$. Then

$$
|f|_{(q)} \leqq C|f|_{(p)}^{(r-q) /(r-p)}|f|_{(r)}^{(q-p) /(r-p)} .
$$

Here $C$ is a constant bounded with $r$ and $(\delta-\gamma)^{-1}$.
Lemma 3.A. 1 can be proved in the same way as [16, Theorem A.5] or [32, Corollary 1.3] (with [21, Lemma 3.2, p. 80]), which are analogues for $C^{r}$ spaces.

Corollary 3.A.2. Let $g$ and $h$ belong to $C^{(q+r)}(M \times[\gamma, \delta])$. Then

$$
|g|_{(q)}|h|_{(r)} \leqq C\left(|g|_{(q+r)}|h|_{0}+|g|_{0}|h|_{(q+r)}\right) .
$$

Here $C$ is a constant bounded with $q+r$ and $(\delta-\gamma)^{-1}$.
Proof. From Lemma 3.A. 1 and the obvious inequality $a^{1-\mu} b^{\mu} \leqq$ $(1-\mu) a+\mu b$ for $a, b \geqq 0$ and $\mu \in[0,1]$, the corollary follows immediately.

Corollary 3.A.3. Let $g$ and $h$ belong to $C^{(r)}(M \times[\gamma, \delta])$. Then their product fg belongs to $C^{(r)}(M \times[\gamma, \delta])$ and satisfies

$$
|f g|_{(r)} \leqq C\left(|f|_{(r)}|g|_{0}+|f|_{0}|g|_{(r)}\right) .
$$

Here $C$ is a constant bounded with $r$ and $(\delta-\gamma)^{-1}$.
Proof. From Definitions 1.B.2, 1.B.3, Leibnitz's formula and Corollary 3.A.2, the corollary follows immediately.

Remark 3.A.4. Corollary 3.A. 3 implies that $C^{(r)}(M \times[\gamma, \delta])$ is a ring and $C_{\#}^{(r)}(M \times[0, T])$ is an ideal in $C^{(r)}(M \times[0, T])$.

We introduce a special class of functions.
Definition 3.A.5. Let $M$ be a compact $C^{\infty}$ manifold. Let $I$ be a subset of nonnegative real numbers having the maximal element. Let $G \geqq 1$. Then, for $T \in\left(0, T_{0}\right], E_{G}^{I}(M \times[0, T])$ is the set of functions $f \in$ $C^{(\max I)}(M \times[0, T])$ having an extension $\tilde{f} \in C^{(\max I)}\left(M \times\left[-T_{0}, T\right]\right)$ such that $|\widetilde{f}|_{(q)} \leqq G|f|_{(q)}$ for $q \in I$. We call the above $\tilde{f}$ an $E_{G}^{\tau}$-extension of $f$.

Lemma 3.A.6. Let $M$ be a compact $C^{\infty}$ manifold. Let $r \geqq 0$. Then
we have:
(i) $C_{\sharp}^{(r)}(M \times[0, T]) \subset E_{1}^{[0, r]}(M \times[0, T])$.
(ii) If $r<2$, then $C^{(r)}(M \times[0, T])=E_{1}^{[0, r]}(M \times[0, T])$.
(iii) If $r \geqq 2$, then

$$
\left\{f \in C^{(r)}(M \times[0, T]) ; \partial_{t} f \in C_{\sharp}^{(r-2)}(M \times[0, T])\right\} \subset E_{1}^{[0, r]}(M \times[0, T])
$$

Proof. In the above cases, we can construct an $E_{1}^{[0, r]}$-extension $\tilde{f}$ of $f$ by putting $\tilde{f}(x, t)=f(x, 0)$ for $(x, t) \in M \times\left[-T_{0}, 0\right]$. This proves the lemma.

Lemma 3.A.7. Let $M$ be a compact $C^{\infty}$ manifold. Let $0 \leqq p, q \leqq r$ and $G \geqq 1$. Let $\alpha$ and $\beta$ be $n$-tuples of nonnegative integers. Suppose that:
(i) $f \in E_{G}^{\{p+|\alpha|+a, r+|\alpha|+a\}}(M \times[0, T])$.
(ii) $g \in E_{G}^{\prime q+|\beta|+b, r+|\beta|+b \mid}(M \times[0, T])$.

Then the product $\left(\partial_{x}^{\alpha} \partial_{t}^{\alpha} f\right)\left(\partial_{x}^{s} \partial_{t}^{b} g\right)$ belongs to $C^{(r)}(M \times[0, T])$ and satisfies

$$
\left|\left(\partial_{x}^{\alpha} \partial_{t}^{a} f\right)\left(\partial_{x}^{\beta} \partial_{t}^{b} g\right)\right|_{(r)} \leqq C\left(|f|_{(r+|\alpha|+a)}|g|_{(q+|\beta|+b)}+|f|_{(p+|\alpha|+a)}|g|_{(r+|\beta|+b)}\right) .
$$

Here $C$ is a constant bounded with $r,|\alpha|+a,|\beta|+b$, and $G$.
Proof. From Corollary 3.A. 3 and Definition 3.A.5, the lemma follows immediately.

Remark 3.A.8. In Lemma 3.A.7, the constant $C$ is independent of $T \in\left(0, T_{0}\right]$, when $G$ is.

Lemma 3.A.9. Let $d_{1}$ and $d_{2}$ be positive integers. Let $D_{1}\left(\right.$ resp. $\left.D_{2}\right)$ be a bounded domain with $C^{\infty}$ boundary in $\boldsymbol{R}^{d_{1}}\left(\right.$ resp. $\boldsymbol{R}^{d_{2}}$ ). Let $r \geqq 1$ and $G \geqq 1$. Suppose that:
(i) $f \in E_{G}^{[1, r)}\left(\bar{D}_{1} \times[0, T]\right)$.
(ii) The components $g_{1}, \cdots, g_{d_{1}}$ of a mapping $g: \bar{D}_{2} \times[0, T] \rightarrow \bar{D}_{1}$ belong to $E_{G}^{(1, r)}\left(\bar{D}_{2} \times[0, T]\right)$.
(iii) $\quad \sum_{i=1}^{d_{1}}\left|g_{i}\right|_{(1)} \leqq B$.

Then the composed function $(x, t) \mapsto f(g(x, t), t)$ belongs to $C^{(r)}\left(\bar{D}_{2} \times[0, T]\right)$ and satisfies

$$
|f(g(x, t), t)|_{(r)} \leqq C\left(|f|_{(r)}+|f|_{(1)} \sum_{i=1}^{d_{1}}\left|g_{i}\right|_{(r)}\right)
$$

Here $C$ is a constant bounded with $r, G$, and $B$.
Lemma 3.A. 9 can be proved in the same way as [16, Theorem A.8] or [32, Lemma 1.6], which are analogues for $C^{r}$ spaces.

Finally we state the following fact, which follows immediately from Definitions 1.B. 2 and 1.B.3.

Lemma 3.A.10. Let $M$ be a compact $C^{\infty}$ manifold. Let $\gamma<\delta$ and $r \geqq 0$. Let $B>0$. Suppose that:
(i) $f \in C^{(r)}(M \times[\gamma, \delta])$.
(ii) $|f(x, t)| \geqq B^{-1}$ for $(x, t) \in M \times[\gamma, \delta]$.

Then the function $(x, t) \mapsto 1 / f(x, t)$ belongs to $C^{(r)}(M \times[\gamma, \delta])$ and satisfies

$$
|1 / f|_{(r)} \leqq C\left(1+|f|_{(r)}\right)
$$

Here $C$ is a constant bounded with $r$ and $B$.
(B) Throughout (B), let $\mathscr{L}$ be a differential operator on $\bar{\Omega}_{T}$ of the form

$$
\mathscr{L}=\partial_{t}-\sum_{i, j=1}^{n} A_{i j}(x, t) \partial_{x_{i}} \partial_{x_{j}}-\sum_{i=1}^{n} A_{i}(x, t) \partial_{x_{i}}-A_{0}(x, t) \quad \text { for } \quad(x, t) \in \bar{\Omega}_{T}
$$

Definition 3.B.1. Let $0 \leqq k \leqq r$. Let $\sigma$ and $B$ be positive constants. We say that $\mathscr{L}$ is a ( $k, r, \sigma, B$ )-\#-parabolic operator if:
(i) $A_{i j}, A_{i}, A_{0} \in C^{(r)}\left(\bar{\Omega}_{T}\right)$ for $i, j=1, \cdots, n$.
(ii) If $r \geqq 2$, then $\partial_{t} A_{i j}, \partial_{t} A_{i}, \partial_{t} A_{0} \in C_{\sharp}^{(r-2)}\left(\bar{\Omega}_{T}\right)$ for $i, j=1, \cdots, n$.
(iii) $\sum_{i, j=1}^{n}\left|A_{i j}\right|_{(k)}+\sum_{i=1}^{n}\left|A_{i}\right|_{(k)}+\left|A_{0}\right|_{(k)} \leqq B$.
(iv) $\sum_{i, j=1}^{n}\left|\left(\left.A_{i j}\right|_{t=0}\right)\right|_{r}+\sum_{i=1}^{n}\left|\left(\left.A_{i}\right|_{t=0}\right)\right|_{r}+\left|\left(\left.A_{0}\right|_{t=0}\right)\right|_{r} \leqq B$.
(v) $\sigma^{-1} \sum_{i=1}^{n} \xi_{i}^{2} \leqq \sum_{i, j=1}^{n} A_{i j}(x, t) \xi_{i} \xi_{j} \leqq \sigma \sum_{i=1}^{n} \xi_{i}^{2} \quad$ for $\quad(x, t) \in \bar{\Omega}_{T} \quad$ and $\xi \in \boldsymbol{R}^{n}$.

Lemma 3.B.2. Let $i$ be a nonnegative integer. Let $\sigma$ and $B$ be positive constants. Suppose that:
(i) The operator $\mathscr{L}$ is $\left(\varepsilon_{0}, \varepsilon_{0}+i, \sigma, B\right)-\#$-parabolic.
(ii) $f \in C_{\#}^{\left(\varepsilon_{0}+i\right)}\left(\bar{\Omega}_{T}\right)$.
(iii) $u \in C^{\left(\varepsilon_{0}+i+2\right)}\left(\bar{\Omega}_{T}\right)$.
(iv) $\mathscr{L} u=f$ in $\Omega_{r}$.

Then $u$ belongs to $E_{G}^{\left(0, \varepsilon_{0}+2, \varepsilon_{0}+3, \cdots, \varepsilon_{0}+i+2\right)}\left(\bar{\Omega}_{T}\right)$. Here $G$ is a constant bounded with $i, \sigma$, and $B$.

Proof. Let $\widetilde{A}_{g h}, \widetilde{A}_{g}$, and $\widetilde{A}_{0}$ be extensions of $\left.A_{g h}\right|_{t=0},\left.A_{g}\right|_{t=0}$, and $\left.A_{0}\right|_{t=0}$ to $\boldsymbol{R}^{n}$, respectively, such that:
(i) $\quad \widetilde{A}_{g h}, \widetilde{A}_{g}, \widetilde{A}_{0} \in C^{\varepsilon_{0}+i}\left(\boldsymbol{R}^{n}\right)$.
(ii) $\left|\widetilde{A}_{g h}\right|_{0_{0}+j} \leqq C_{1}\left|\left(\left.A_{g h}\right|_{t=0}\right)\right|_{\varepsilon_{0}+j}, \quad\left|\widetilde{A}_{g}\right|_{\varepsilon_{0}+j} \leqq C_{1}\left|\left(\left.A_{g}\right|_{t=0}\right)\right|_{\varepsilon_{0}+j}, \quad\left|\widetilde{A}_{0}\right|_{\varepsilon_{0}+j} \leqq$ $C_{1}\left|\left(\left.A_{0}\right|_{t=0}\right)\right|_{\varepsilon_{0}+j}$, for $j=0, \cdots, i$.
Here $g, h=1, \cdots, n$ and $C_{1}$ is a constant bounded with $i$. (This is possible by the Hestenes-Whitney technique. See, e.g., [21, p. 296-297].) Let $\Omega_{1}$ be a domain with $C^{\infty}$ boundary in $\boldsymbol{R}^{n}$ such that:
(i) $\bar{\Omega}_{0} \subset \Omega_{1}$.
(ii) $(2 \sigma)^{-1} \sum_{g=1}^{n} \xi_{g}^{2} \leqq \sum_{g, h=1}^{n} \widetilde{A}_{g h}(x) \xi_{g} \xi_{h} \leqq 2 \sigma \sum_{g=1}^{n} \xi_{g}^{2}$ for $x \in \bar{\Omega}_{1}$ and $\xi \in \boldsymbol{R}^{n}$. The domain $\Omega_{1}$ is determined only by $\Omega_{0}, \sigma$, and $B$. Let $\widetilde{a}$ be an extension
of $\left.u\right|_{t=0}$ to $\bar{\Omega}_{1}$ such that:
(i) $\operatorname{supp} \widetilde{a} \subset \Omega_{1}$.
(ii) $\tilde{a} \in C^{\varepsilon_{0}+i+2}\left(\bar{\Omega}_{1}\right)$.
(iii) $|\widetilde{a}|_{\varepsilon_{0}+j+2} \leqq C_{2}\left|\left(\left.u\right|_{t=0}\right)\right|_{\varepsilon_{0}+j+2}$ for $j=0, \cdots, i$.

Here $C_{2}$ is a constant bounded with $i, \sigma$, and $B$. Let $\partial \Omega_{1}$ be the boundary of $\Omega_{1}$. Let $w$ be the solution of the mixed problem:

$$
\begin{gather*}
\left(\partial_{t}-\sum_{i, j=1}^{n} \widetilde{A}_{i j}(x) \partial_{x_{i}} \partial_{x_{j}}-\sum_{i=1}^{n} \widetilde{A}_{i}(x) \partial_{x_{i}}-\widetilde{A}_{0}(x)\right) w=0 \quad \text { in } \quad \Omega_{1} \times\left[0, T_{0}\right]  \tag{1}\\
\left.w\right|_{t=0}=\widetilde{a} \tag{2}
\end{gather*}
$$

$$
\begin{equation*}
w=0 \quad \text { on } \quad \partial \Omega_{1} \times\left[0, T_{0}\right] \tag{3}
\end{equation*}
$$

It is known that $w$ exists in $C^{\left(\varepsilon_{0}+i+2\right)}\left(\bar{\Omega}_{1} \times\left[0, T_{0}\right]\right)$ and satisfies $|w|_{(q)} \leqq$ $C_{3}|\widetilde{a}|_{q}$ for $q=0, \varepsilon_{0}+2, \varepsilon_{0}+3, \cdots, \varepsilon_{0}+i+2$. Here $C_{3}$ is a constant bounded with $i$. (See [21, Theorem 2.3, p. 16-17] and [21, Theorem 5.2, p. 320].) We easily observe that:
(i ) $\left.\left(\partial_{t}^{j} w\right)\right|_{t=0}=\left.\left(\partial_{t}^{j} u\right)\right|_{t=0}$ on $\bar{\Omega}_{0}$ for $j=0, \cdots,[(i+2) / 2]$.
(ii) $|w|_{(q), T_{0}} \leqq C_{2} C_{3}|u|_{(q), T}$ for $q=0, \varepsilon_{0}+2, \varepsilon_{0}+3, \cdots, \varepsilon_{0}+i+2$.

Let $\widetilde{w}$ be the Hestenes-Whitney extension of $w$ to $\bar{\Omega}_{1} \times\left[-T_{0}, T_{0}\right]$. Now we can construct an $E_{G}^{\left\{0, \varepsilon_{0}+2, \varepsilon_{0}+3, \cdots, \varepsilon_{0}+i+2\right)}$-extension $\tilde{u}$ of $u$ by setting $\tilde{u}(x, t)=$ $\widetilde{w}(x, t)$ for $(x, t) \in \bar{\Omega}_{0} \times\left[-T_{0}, 0\right]$. This proves the lemma.

Lemma 3.B.3. Let $i$ be a nonnegative integer. Let $\sigma$ and $B$ be positive constants. Suppose that:
(i) The operator $\mathscr{L}$ is $\left(\varepsilon_{0}, \varepsilon_{0}+i, \sigma, B\right)-\#$-parabolic.
(ii) $f \in C_{\#}^{\left(\varepsilon_{0}+i\right)}\left(\bar{\Omega}_{T}\right)$.
(iii) $a \in C^{\varepsilon_{0}+i+2}\left(\bar{\Omega}_{0}\right)$.
(iv) $b_{1} \in C^{\left(\varepsilon_{0}+i+2\right)}\left(J_{T}\right)$.
(v) $b_{2} \in C^{\left(\varepsilon_{0}+i+2\right)}\left(\Gamma_{T}\right)\left(\right.$ resp. $\left.b_{2} \in C^{\left(\varepsilon_{0}+i+1\right)}\left(\Gamma_{T}\right)\right)$.

Consider the following mixed problem:

$$
\begin{equation*}
\mathscr{L} u=f \quad \text { in } \quad \Omega_{T} \tag{1}
\end{equation*}
$$

$$
\begin{equation*}
\left.u\right|_{t=0}=a \tag{2}
\end{equation*}
$$

$$
\begin{equation*}
u=b_{1} \quad \text { on } \quad J_{T} \tag{3}
\end{equation*}
$$

(Here $\partial_{\lambda}$ is the partial differentiation with respect to $\lambda$ in the $\left(\omega_{1}, \cdots\right.$, $\left.\omega_{n-1}, \lambda\right)$ coordinates.) Suppose that:
(i) The set $\left\{a, b_{1}\right\}$ satisfies the compatibility condition of order $[(i+2) / 2]$ of the Dirichlet problem for the equation $\mathscr{L} u=f$ on $J_{0}$ at $t=0$.
(ii) The set $\left\{a, b_{2}\right\}$ satisfies the compactibility condition of order. $[(i+2) / 2]$ (resp. $[(i+1) / 2])$ of the Dirichlet (resp. Neumann) problem for the equation $\mathscr{L} u=f$ on $\Gamma_{0}$ at $t=0$.
Then the problem (1)-(4) has a unique solution $u$ in $C^{\left(\varepsilon_{0}+i+2\right)}\left(\bar{\Omega}_{T}\right)$ which satisfies

$$
\begin{aligned}
|u|_{\left(\varepsilon_{0}+i+2\right)} \leqq & C\left[\left(\sum_{g, h=1}^{n}\left|A_{g h}\right|_{\left(\varepsilon_{0}+i\right)}+\sum_{g=1}^{n}\left|A_{g}\right|_{\left(\varepsilon_{0}+i\right)}+\left|A_{0}\right|_{\left(\varepsilon_{0}+i\right)}\right)\right. \\
& \times\left(|f|_{\left(\varepsilon_{0}\right)}+|a|_{\varepsilon_{0}+2}+\left|b_{1}\right|_{\left(\varepsilon_{0}+2\right)}+\left|b_{2}\right|_{\left(\varepsilon_{0}+2\right)}\right) \\
& \left.+\left(|f|_{\left(\varepsilon_{0}+i\right)}+|a|_{\varepsilon_{0}+i+2}+\left|b_{1}\right|_{\left(\varepsilon_{0}+i+2\right)}+\left|b_{2}\right|_{\left(\varepsilon_{0}+i+2\right.}\right)\right]
\end{aligned}
$$

(resp. $|u|_{\left(\varepsilon_{0}+i+2\right)} \leqq C\left[\left(\sum_{g, h=1}^{n}\left|A_{g h}\right|_{\left(\varepsilon_{0}+i\right)}+\sum_{g=1}^{n}\left|A_{g}\right|_{\left(\varepsilon_{0}+i\right)}+\left|A_{0}\right|_{\left(\varepsilon_{0}+i\right)}\right)\left(|f|_{\left(\varepsilon_{0}\right)}+|a|_{\varepsilon_{0}+2}+\right.\right.$ $\left.\left.\left.\left|b_{1}\right|_{\left(\varepsilon_{0}+2\right)}+\left|b_{2}\right|_{\left(\varepsilon_{0}+1\right)}\right)+\left(|f|_{\left(\varepsilon_{0}+i\right)}+|a|_{\varepsilon_{0}+i+2}+\left|b_{1}\right|_{\left(\varepsilon_{0}+i+2\right)}+\left|b_{2}\right|_{\left(\varepsilon_{0}+i+1\right)}\right)\right]\right)$. Here C is a constant bounded with $i, \sigma$, and $B$.

For the proof, refer to [16, Theorem A.14] which is an analogue for the elliptic boundary value problem. With the aid of [21, Theorems 5.2, 5.3, p. 320-321], [21, Theorem 2.3, p. 16-17], Lemmas 3.A.6, 3.A.7 and Lemma 3.B.2, we can prove Lemma 3.B. 3 in a manner similar to the proof of [16, Theorem A.14].
4. Reformulation of the problem. In order to apply Nash's implicit function theorem, we reformulate Theorem' (in Section 2).
(A) For $\rho \in V_{T}$, by using the notations $\mathscr{L}_{\rho}$ and $S_{\rho}$ introduced in Section 2 and by setting $U=u \circ e_{\rho}$, we can reformulate $\left(1_{u}\right)-\left(5_{u}\right)$ as follows:

$$
\begin{equation*}
\mathscr{L}_{\rho} U=0 \quad \text { in } \quad \Omega_{T} . \tag{U}
\end{equation*}
$$

$$
\begin{equation*}
\left.U\right|_{t=0}=a_{0} . \tag{U}
\end{equation*}
$$

$$
\begin{array}{ll}
U=b_{0} & \text { on } \quad J_{T} . \\
U=0 & \text { on } \quad \Gamma_{T} . \tag{U}
\end{array}
$$

$$
\begin{equation*}
\partial_{t} \rho+c_{0}\left(\partial_{\lambda} U\right) S_{\rho}=0 \quad \text { on } \quad \Gamma_{T} . \tag{U}
\end{equation*}
$$

Let us express $\mathscr{L}_{\rho}$ as

$$
\mathscr{L}_{\rho}=\partial_{t}-\sum_{i, j=1}^{n} A_{\rho, i j}(x, t) \partial_{x_{i}} \partial_{x_{j}}-\sum_{i=1}^{n} A_{\rho, i}(x, t) \partial_{x_{i}}
$$

for $(x, t) \in \bar{\Omega}_{T}$ with $A_{\rho, i j}=A_{\rho, j i}$. By (2.1)-(2.3), a routine calculation gives us:

$$
\begin{equation*}
A_{\rho, i j}(x, t)=A_{i j}\left(x, \rho(\omega(x), t), \partial_{\omega_{1}} \rho(\omega(x), t), \cdots, \partial_{\omega_{n-1}} \rho(\omega(x), t)\right) \tag{4.1.i}
\end{equation*}
$$

for $i, j=1, \cdots, n$ and $(x, t) \in N_{0}^{-} \times[0, T]$, where $\left(\omega_{1}, \cdots, \omega_{n-1}\right)$ are local coordinates in $\Gamma_{0}$.

$$
\begin{equation*}
A_{\rho, i j}(x, t)=\delta_{i j} \tag{4.1.ii}
\end{equation*}
$$

for $i, j=1, \cdots, n$ and $(x, t) \in\left(\bar{\Omega}_{0}-N_{0}^{-}\right) \times[0, T]$, where $\delta_{i j}$ is Kronecker's delta.

$$
\begin{align*}
A_{\rho, i}(x, t)= & A_{i}\left(x, \rho(\omega(x), t), \partial_{\omega_{1}} \rho(\omega(x), t), \cdots, \partial_{\omega_{n-1}} \rho(\omega(x), t),\right.  \tag{4.1.iii}\\
& \partial_{t} \rho(\omega(x), t), \partial_{\omega_{1}}^{2} \rho(\omega(x), t), \partial_{\omega_{1}} \partial_{\omega_{2}} \rho(\omega(x), t), \cdots, \\
& \left.\partial_{\omega_{n-1}}^{2} \rho(\omega(x), t)\right)
\end{align*}
$$

for $i=1, \cdots, n$ and $(x, t) \in N_{0}^{-} \times[0, T]$.
(4.1.iv) $A_{\rho, i}(x, t)=0$ for $i=1, \cdots, n$ and $(x, t) \in\left(\bar{\Omega}_{0}-N_{0}^{-}\right) \times[0, T]$.

Here $A_{i j}\left(\right.$ resp. $\left.A_{i}\right)$ are $C^{\infty}$ functions on $N_{0}^{-} \times\left[-\delta_{0}, \delta_{0}\right] \times \boldsymbol{R}^{n-1}$ (resp. $\left.N_{0}^{-} \times\left[-\delta_{0}, \delta_{0}\right] \times \boldsymbol{R}^{\left(n^{2}+n\right) / 2}\right)$. We easily observe that if $\rho \in V_{T} \cap C_{\sharp}^{(r)}\left(\Gamma_{T}\right)$ for $r \geqq r_{0}$, then $\mathscr{L}_{\rho}$ is an $\left(r_{0}-2, r-2, \sigma_{0}, B\right)$-\#-parabolic operator (see (B) in Section 2). Here $B$ is a positive constant.

On the other hand, express $S_{\rho}$ as

$$
\begin{equation*}
S_{\rho}(\omega, t)=S\left(\omega, \rho(\omega, t), \partial_{\omega_{1}} \rho(\omega, t), \cdots, \partial_{\omega_{n-1}} \rho(\omega, t)\right) \tag{4.2}
\end{equation*}
$$

for $(\omega, t) \in \Gamma_{T}$. By (1.1), a routine calculation gives us

$$
\begin{aligned}
& S\left(\omega, \rho(\omega, t), \partial_{\omega_{1}} \rho(\omega, t), \cdots, \partial_{\omega_{n-1}} \rho(\omega, t)\right) \\
& \quad=\sum_{i=1}^{n}\left\{\partial_{x_{i}} \lambda(x(\omega, \rho(\omega, t)))-\sum_{j=1}^{n-1}\left[\partial_{\omega_{j}} \rho(\omega, t)\right]\left[\partial_{x_{i}} \omega_{j}(x(\omega, \rho(\omega, t)))\right]\right\}^{2}
\end{aligned}
$$

for $(\omega, t) \in \Gamma_{T}$. As is assumed in Section 1, the function $S_{\rho}$ has a positive constant $B_{0}^{-1}$ as a lower bound. Note that $\operatorname{grad} \lambda \neq 0$ on $\Gamma_{T}$ because $\Gamma_{0}$ is embedded in $\boldsymbol{R}^{n}$ without critical points. We easily observe that if $\rho \in V_{T} \cap C_{\#}^{(r)}\left(\Gamma_{T}\right)$ for $r \geqq r_{0}$, then $\partial_{t}\left(S_{\rho}\right) \in C_{\#}^{(r-3)}\left(\Gamma_{T}\right)$.
(B) Define a mapping $\mathscr{F}: V_{T} \rightarrow C_{\#}^{\left(r_{0}-2\right)}\left(\Gamma_{T}\right)$ by

$$
\begin{equation*}
\mathscr{F}(\rho)=\partial_{t} \rho+c_{0}\left[\left.\left(\partial_{\lambda} U\right)\right|_{\lambda=0}\right] S_{\rho} \quad \text { for } \quad \rho \in V_{T} . \tag{4.3}
\end{equation*}
$$

Here $U_{\rho}$ is the solution $U$ of $\left(1_{U}\right)-\left(4_{U}\right)$. This is possible by Lemma 3.B. 3 and the assumptions (A.2) and (T) in theorem.

Clearly Theorem' (in Section 2) is reformulated as:
Theorem' (reformulated form). Under the same assumptions as in Theorem' in Section 2, for a sufficiently small $T \in\left(0, T_{0}\right]$, there exists $\rho \in V_{T}$ with $\mathscr{F}(\rho)=0$.

Here we suppose $\rho$ is the only unknown of the problem because $u$ is obtained once $\rho$ is determined.
(C) By Lemmas 3.A. 9 and 3.B.3, we have

$$
\begin{align*}
\left|U_{\rho}\right|_{\left(r_{0}+i\right)} \leqq C\left(1+|\rho|_{\left(r_{0}+i\right)}\right) & \text { for }  \tag{4.4}\\
& i=0, \cdots, 39 \\
& \text { and } \rho \in V_{T} \cap C_{\ddagger}^{\left(r_{0}+i\right)}\left(\Gamma_{T}\right) .
\end{align*}
$$

Here $C$ is a constant bounded with $\left|a_{0}\right|_{r_{0}+39}$ and $\left|b_{0}\right|_{\left.\text {ro }_{0}+39\right)}$. By (4.3), (4.4) and Lemma 3.A.7, we have:

$$
\begin{align*}
\mathscr{F}\left(V_{T} \cap C_{\sharp}^{\left(r_{0}+i\right)}\left(\Gamma_{T}\right)\right) \subseteq C_{\sharp}^{\left(r_{0}+i-2\right)}\left(\Gamma_{T}\right) & \text { for } \quad i=0, \cdots, 40 .  \tag{4.5.i}\\
|\mathscr{F}(\rho)|_{\left(r_{0}+i-2\right)} \leqq C\left(1+|\rho|_{\left(r_{0}+i\right)}\right) & \text { for } \\
& i=0, \cdots, 40 \\
& \text { and } \rho \in V_{T} \cap C_{\sharp}^{\left(r_{0}+i\right)}\left(\Gamma_{T}\right) .
\end{align*}
$$

Here $C$ is a constant bounded with $\left|a_{0}\right|_{r_{0}+39}$ and $\left|b_{0}\right|_{\left(r_{0}+39\right)}$.
5. Nash's implicit function theorem. (A) We recall the well-known version of Nash's implicit function theorem due to Moser [22]. For the proof, refer to [22], [34, Chapter 2] and [35].

NASH'S IMPLICIT FUNCTION THEOREM. Let $E_{0}, E_{1}, \cdots$, and $E_{11}$ (resp. $F_{1}, F_{2}, \cdots$, and $F_{11}$ ) be real Banach spaces such that:
(i) $E_{0} \supset E_{1} \supset \cdots \supset E_{11}\left(\right.$ resp.$\left.F_{1} \supset F_{2} \supset \cdots \supset F_{11}\right)$
(ii) $|x|_{i} \leqq|x|_{i+1}$ for $i=0, \cdots, 11$ and $x \in E_{i+1}$ (resp. $|y|_{i} \leqq|y|_{i+1}$ for $i=1, \cdots, 11$ and $\left.y \in F_{i+1}\right)$.
Here $\left|\left.\right|_{i}\right.$ denotes the norm in $E_{i}\left(\right.$ resp. $\left.F_{i}\right)$. Let $\delta$ be a positive constant and let $V=\left\{x \in E_{1} ;\left|x_{1}\right|<\delta\right\}$. Let $\mathscr{F}: V \rightarrow F_{1}$ be a mapping such that $\mathscr{F}\left(V \cap E_{i}\right) \subseteq F_{i}$ for $i=1, \cdots, 11$.

Suppose that:
( I ) There exists a "smoothing" linear operator $\mathscr{S}_{0}: E_{0} \rightarrow E_{11}$ for $\theta \geqq 1$ which satisfies:
(I.i) $\left|\mathscr{S}_{v} x\right|_{j} \leqq C_{1} \theta^{j-i}|x|_{i} \quad$ for $\quad i, j$ with $0 \leqq i \leqq j \leqq 11$ and $x \in E_{i}$.
(I.ii) $\quad\left|x-\mathscr{S}_{\theta} x\right|_{i} \leqq C_{1} \theta^{-(j-i)}|x|_{j} \quad$ for $\quad i, j \quad$ with $\quad 0 \leqq i \leqq j \leqq 11$ and $x \in E_{j}$.
(II) There exists a "Fréchet derivative" linear operator $D \mathscr{F}(x)$ : $E_{1} \rightarrow F_{1}$ for $x \in V$ which satisfies:

$$
\begin{equation*}
|D \mathscr{F}(x) h|_{1} \leqq C_{2}|h|_{1} \quad \text { for } \quad h \in E_{1} . \tag{II.i}
\end{equation*}
$$

(II.ii) $|\mathscr{F}(x+h)-\mathscr{F}(x)-D \mathscr{F}(x) h|_{1} \leqq C_{2}|h|_{1}^{2} \quad$ for $\quad h \in E_{1}$

$$
\text { with } \quad x+h \in V
$$

(III) There exists a "right inverse" linear operator $\mathscr{I}(x): F_{1} \rightarrow E_{0}$ for $x \in V$ which satisfies:

$$
\begin{equation*}
\mathscr{J}(x) F_{i} \subseteq E_{i-1} \quad \text { for } \quad i=1, \cdots, 11 \text { and } x \in V \cap E_{i} . \tag{III.i}
\end{equation*}
$$

$$
\begin{equation*}
D \mathscr{F}(x) \mathscr{J}(x) y=y \quad \text { for } \quad x \in V \cap E_{2} \quad \text { and } \quad y \in F_{2} . \tag{III.ii}
\end{equation*}
$$

(III.iii) $\quad|\mathscr{J}(x) y|_{0} \leqq C_{3}|y|_{1} \quad$ for $\quad y \in F_{1}$.
(III.iv) $|\mathscr{J}(x) \mathscr{F}(x)|_{i-1} \leqq C_{3}\left(1+|x|_{i}\right)$ for $i=1, \cdots, 11$ and $x \in V \cap E_{i}$.

Here $C_{1}, C_{2}$, and $C_{3}$ are constants.
Then there is a positive constant $\varepsilon$ determined by $C_{1}, C_{2}, C_{3}$ and $\delta$ such that: If $|\mathscr{F}(0)|_{1}<\varepsilon$, then there exists $x \in V$ with $\mathscr{F}(x)=0$.
(B) We prove Theorem' (in Section 4) by using Nash's implicit function theorem. The setting in which we apply Nash's implicit function theorem is:
(i) $E_{i}=C_{\#}^{\left(r_{0}-4+4 i\right)}\left(\Gamma_{T}\right)$ for $i=0, \cdots, 11$.
(ii) $F_{i}=C_{\#}^{\left(r_{0}-6-\left(\varepsilon_{0} / 2\right)+4 i\right)}\left(\Gamma_{T}\right)$ for $i=1, \cdots, 11$.
(iii) $\delta=\delta_{0}$ (so that $V=V_{T}$ ).
(iv) The mapping $\mathscr{F}$ is the one defined by (4.3) (see (4.5.i)).

In the following sections, we show that the mapping $\mathscr{F}$ and the spaces $E_{i}, F_{i}$ together with $\delta$ satisfy the conditions in Nash's implicit function theorem.
6. Smoothing operators. We verify the condition (I) in Nash's theorem. Choose a function $\zeta_{0}$ on $\boldsymbol{R}^{n-1} \times \boldsymbol{R}$ so that:
(i) $\zeta_{0} \in C_{0}^{\infty}\left(\boldsymbol{R}^{n-1} \times \boldsymbol{R}\right)$.
(ii) $\zeta_{0}(x, t)=0$ for $x \in \boldsymbol{R}^{n-1}$ and $t \leqq 0$.
(iii) $\int_{\boldsymbol{R}^{n-1} \times \boldsymbol{R}} \zeta_{0}(x, t) d x d t=1$, where $d x=\prod_{i=1}^{n-1} d x_{i}$.

Define a sequence of functions $\left\{\zeta_{0}, \zeta_{1}, \cdots\right\}$ by

$$
\zeta_{i}(x, t)=\left(1-2^{i}\right)^{-1}\left[\zeta_{i-1}(x, t)-2^{i+n+1} \zeta_{i-1}(2 x, 4 t)\right]
$$

for $i=1,2, \cdots$ and $(x, t) \in \boldsymbol{R}^{n-1} \times \boldsymbol{R}$. For $i=0,1, \cdots$, we can easily verify that:
(i) $\zeta_{i} \in C_{0}^{\infty}\left(\boldsymbol{R}^{n-1} \times \boldsymbol{R}\right)$.
(ii) $\zeta_{i}(x, t)=0$ for $x \in \boldsymbol{R}^{n-1}$ and $t \leqq 0$.
(iii) $\int_{\mathbf{R}^{n-1} \times \mathbf{R}} \zeta_{i}(x, t) d x d t=1$.
(iv) $\int_{\boldsymbol{R}^{n-1} \times \boldsymbol{R}}^{\boldsymbol{R}^{n-1} \times \boldsymbol{R}} x^{\alpha} t^{a} \zeta_{i}(x, t) d x d t=0$ if $0<|\alpha|+2 a \leqq i$.

Let $i_{0}=\left[r_{0}\right]+40$. For $\bar{\theta} \geqq 1$, define a linear operator

$$
\overline{\mathscr{S}}_{\vec{\theta}}: C_{\#}^{(0)}\left(\boldsymbol{R}^{n-1} \times[0, T]\right) \rightarrow C_{\#}^{\left(i_{0}\right)}\left(\boldsymbol{R}^{n-1} \times[0, T]\right)
$$

by

$$
\begin{aligned}
{\left[\mathscr{S}_{\bar{\theta}} f\right](x, t)=} & \int_{0}^{T}\left\{\int_{\boldsymbol{R}^{n-1}}\left[\bar{\theta}^{n+1} \zeta_{i_{0}}\left(\bar{\theta}(x-\xi), \bar{\theta}^{2}(t-\tau)\right) f(\xi, \tau)\right] d \xi\right\} d \tau \\
& \text { for } \quad f \in C_{\ddagger}^{(0)}\left(\boldsymbol{R}^{n-1} \times[0, T]\right) \quad \text { and } \quad(x, t) \in \boldsymbol{R}^{n-1} \times[0, T]
\end{aligned}
$$

The fact that $\overline{\mathscr{S}}_{\bar{\theta}} f$ is in $C_{\ddagger}^{\left(i_{0}\right)}$ follows from the property (ii) of $\zeta_{i}$. We can verify that:
(i) $\left|\mathscr{S}_{\vec{\theta}} f\right|_{(r)} \leqq C \bar{\theta}^{r-q}|f|_{(q)}$ for real numbers $q, r$ with $0 \leqq q \leqq r \leqq$ $r_{0}+40$ and for $f \in C_{\#}^{(q)}\left(\boldsymbol{R}^{n-1} \times[0, T]\right)$.
(ii) $\left|f-\mathscr{S}_{\mathscr{\theta}} f\right|_{(q)} \leqq C \bar{\theta}^{-(r-q)}|f|_{(r)}$ for real numbers $q, r$ with $0 \leqq q \leqq$ $r \leqq r_{0}+40$ and for $f \in C_{\sharp}^{(r)}\left(\boldsymbol{R}^{n-1} \times[0, T]\right)$.
Here $C$ is a constant determined by $\zeta_{i_{0}}$. For the proof, refer to [16, Theorem A.10] where analogous inequalities are proved for $C^{r}$ spaces. The above inequalities (i) and (ii) can be proved similarly.

Set $\theta=\bar{\theta}^{1 / 4}$. Consider a finite covering of $\Gamma_{0}$ by coordinate neighborhoods and a $C^{\infty}$ partition of unity subordinate to it. Then, for $\theta \geqq 1$, we can construct a linear operator $\mathscr{S}_{\theta}: C_{\neq}^{\left(r_{0}-4\right)}\left(\Gamma_{T}\right) \rightarrow C_{\neq}^{\left(r_{0}+40\right)}\left(\Gamma_{T}\right)$ which satisfies:
(i) $\left|\mathscr{S}_{\theta} f\right|_{\left(r_{0}-4+4 j\right)} \leqq C \theta^{j-i}|f|_{\left(r_{0}-4+4 i\right)}$ for $i, j$ with $0 \leqq i \leqq j \leqq 11$ and $f \in C_{\#}^{\left(r_{0}-4+4 i\right)}\left(\Gamma_{T}\right)$.
(ii) $\left|f-\mathscr{S}_{\theta} f\right|_{\left(r_{0}-4+4 i\right)} \leqq C \theta^{-(j-i)}|f|_{\left(r_{0}-4+4 j\right)}$ for $i, j$ with $0 \leqq i \leqq j \leqq 11$ and $f \in C_{\ddagger}^{\left(r_{0}-4+4 j\right)}\left(\Gamma_{T}\right)$.
Here $C$ is a constant. This proves (I).
7. Fréchet differentiability of $\mathscr{F}$. We verify the condition (II) in Nash's implicit function theorem.
(A) Let $\mathscr{E}_{\rho}=\partial_{t}-\mathscr{L}_{\rho}$. For $\rho \in V_{T}$ and $\delta \rho \in C_{\#}^{\left(r_{0}\right)}\left(\Gamma_{T}\right)$, let $\delta \mathscr{E}_{\rho}=$ $\sum_{i, j=1}^{n}\left(\delta A_{\rho, i j}\right) \partial_{x_{i}} \partial_{x_{j}}+\sum_{i=1}^{n}\left(\delta A_{\rho, i}\right) \partial_{x_{i}}$ on $\Omega_{T}$. Here $\delta A_{\rho, i j}$ and $\delta A_{\rho, i}$ are defined by:
(7.1.i) $\delta A_{\rho, i j}=\left[(\partial / \partial \rho) A_{i j}\right] \delta \rho+\sum_{h=1}^{n}\left\{\left[\partial / \partial\left(\partial_{\omega_{h}} \rho\right)\right] A_{i j}\right\} \partial_{\omega_{h}} \delta \rho \quad$ on $\quad N_{0}^{-} \times[0, T]$.
(7.1.iii) $\delta A_{\rho, i}=\left[(\partial / \partial \rho) A_{i}\right] \delta \rho+\sum_{n=1}^{n-1}\left\{\left[\partial\left(\partial_{\omega_{h}} \rho\right)\right] A_{i}\right\} \partial_{\omega_{h}} \delta \rho+\left\{\left[\partial / \partial\left(\partial_{t} \rho\right)\right] A_{i}\right\} \partial_{t} \delta \rho$

$$
+\sum_{1 \leqq g \leqq n \leq n-1}\left\{\left[\partial / \partial\left(\partial_{\omega_{g}} \partial_{\omega_{h}} \rho\right)\right] A_{i}\right\} \partial_{\omega_{g}} \partial_{\omega_{h}} \delta \rho \quad \text { on } \quad N_{0}^{-} \times[0, T] .
$$

$$
\begin{equation*}
\delta A_{\rho, i}=0 \quad \text { on } \quad N_{0}^{-} \times[0, T] \tag{7.1.iv}
\end{equation*}
$$

(See (4.1).)
Let $\delta U_{\rho}$ be the solution of the problem:

$$
\mathscr{L}_{\rho}\left(\delta U_{\rho}\right)=\left(\delta \mathscr{E}_{\rho}\right) U_{\rho} \quad \text { in } \quad \Omega_{T}
$$

$$
\left.\delta U_{\rho}\right|_{t=0}=0
$$

Since $\delta \rho$ belongs to $C_{\#}^{\left(r_{0}\right)}\left(\Gamma_{T}\right)$, the compatibility condition of order [ $\left.r_{0} / 2\right]$ of $\left(1_{\partial U}\right)-\left(4_{\partial U}\right)$ is satisfied. Then, by Lemmas 3.A.7, 3.B. 2 and 3.B.3 and (4.4), the function $\delta U_{\rho}$ is well-defined in $C_{\dot{\#}}^{\left(r_{0}\right)}\left(\bar{\Omega}_{T}\right)$ and satisfies

$$
\begin{equation*}
\left|\delta U_{\rho}\right|_{\left(r_{0}\right)} \leqq C|\delta \rho|_{\left(r_{0}\right)} . \tag{7.2}
\end{equation*}
$$

Here $C$ is a constant bounded with $\left|a_{0}\right|_{r_{0}}$ and $\left|b_{0}\right|_{\left(r_{0}\right)}$.
Further we observe that if $\rho+\delta \rho$ belongs to $V_{T}$, then

$$
\begin{equation*}
\left|U_{\rho+\delta \rho}-U_{\rho}-\delta U_{\rho}\right|_{\left(r_{0}\right)} \leqq C|\delta \rho|_{\left(r_{0}\right)}^{2} . \tag{7.3}
\end{equation*}
$$

Here $C$ is a constant bounded with $\left|a_{0}\right|_{r_{0}}$ and $\left|b_{0}\right|_{\left(r_{0}\right)}$. In fact we have:

$$
\begin{align*}
& \mathscr{L}_{\rho+\delta \rho}\left(U_{\rho+\delta \rho}-U_{\rho}-\delta U_{\rho}\right)  \tag{1}\\
& \quad=\left(\mathscr{E}_{\rho+\delta \rho}-\mathscr{E}_{\rho}-\delta \mathscr{E}_{\rho}\right) U_{\rho}+\left(\mathscr{E}_{\rho+\delta \rho}-\mathscr{E}_{\rho}\right) \delta U_{\rho} \quad \text { on } \quad \Omega_{T} . \tag{2}
\end{align*}
$$

By Lemmas 3.A.7, 3.B.2 and 3.B. 3 and by (4.4), (7.1) and (7.2), we obtain (7.3).
(B) For $\rho \in V_{T}$, define a linear operator $D \mathscr{F}(\rho): C_{\neq}^{\left(r_{0}\right)}\left(\Gamma_{T}\right) \rightarrow C_{\#}^{\left(r_{0}-2\right)}\left(\Gamma_{T}\right)$ by

$$
\begin{equation*}
D \mathscr{F}(\rho) \delta \rho=\partial_{t} \delta \rho+c_{0}\left[\left.\left(\partial_{\lambda} U_{\rho}\right)\right|_{\lambda=0}\right] \delta S_{\rho}+c_{0}\left[\left.\left[\partial_{\lambda}\left(\delta U_{\rho}\right)\right]\right|_{\lambda=0}\right\} S_{\rho} \tag{7.4}
\end{equation*}
$$

for $\delta \rho \in C_{\sharp}^{\left(r_{0}\right)}\left(\Gamma_{T}\right)$ (see (4.3)). Here

$$
\begin{equation*}
\delta S_{\rho}=[(\partial / \partial \rho) S] \delta \rho+\sum_{i=1}^{n-1}\left\{\left[\partial / \partial\left(\partial_{\omega_{i}} \rho\right) S\right\} \partial_{\omega_{i}} \delta \rho .\right. \tag{7.5}
\end{equation*}
$$

(See (4.2).) By (4.4), (7.2), (7.3) and (7.4), we have
(i) $|D \mathscr{F}(\rho) \delta \rho|_{\left(r_{0}-2\right)} \leqq C|\delta \rho|_{\left(r_{0}\right)}$ for $\delta \rho \in C_{\#}^{\left(r_{0}\right)}\left(\Gamma_{T}\right)$.
(ii) $|\mathscr{F}(\rho+\delta \rho)-\mathscr{F}(\rho)-D \mathscr{F}(\rho) \delta \rho|_{\left(r_{0}-1\right)} \leqq C|\delta \rho|_{\left(r_{0}\right)}^{2}$ for $\delta \rho \in C_{\#}^{\left(r_{0}\right)}\left(\Gamma_{T}\right)$ with $\rho+\delta \rho \in V_{T}$.
Here $C$ is a constant bounded with $\left|a_{0}\right|_{r_{0}}$ and $\left|b_{0}\right|_{\left(r_{0}\right)}$. This proves (II).
8. Technical lemmas II. We prove some lemmas which are used to verify (III). As in Section 3, by "a manifold", we mean both one with boundary and one without boundary.
(A) We begin with the following fact.

Lemma 8.A.1. Let $M$ be a compact $C^{\infty}$ manifold. Let $r \geqq 0$. Let $f$ belong to $C_{\sharp}^{(r)}(M \times[0, T])$. Then the function $(x, t) \mapsto \int_{0}^{t} f(x, \tau) d \tau$ belongs to $C_{\ddagger}^{(r)}(M \times[0, T])$ and satisfies

$$
\left|\int_{0}^{t} f(x, \tau) d \tau\right|_{(r), t} \leqq \int_{0}^{t}|f|_{(r), \tau} d \tau \quad \text { for } \quad t \in[0, T] .
$$

Proof. Extend $f$ to $M \times(-\infty, T$ ] by setting $f(x, t)=0$ for $(x, t) \in$ $M \times(-\infty, 0)$. Since $f \in C_{\neq}^{(r)}$, we have:
(i) The $C^{(r)}$ norm of $f$ is preseved in the above extension.
(ii) $\partial_{x}^{\alpha} \partial_{t}^{a}\left(\int_{0}^{t} f(x, \tau) d \tau\right)=\int_{0}^{t} \partial_{x}^{\alpha} \partial_{t}^{a} f(x, \tau) d \tau$ if $|\alpha|+2 \alpha \leqq[r]$.
(iii) $\int_{0}^{s} \partial_{x}^{\alpha} \partial_{t}^{a} f(x, \tau) d \tau=\int_{0}^{t} \partial_{x}^{\alpha} \partial_{t}^{a} f(x, \tau-t+s) d \tau \quad$ if $\quad|\alpha|+2 \alpha \leqq[r] \quad$ and $0 \leqq s \leqq t \leqq T$.
By (ii) and (iii)

$$
\begin{aligned}
& \left|\int_{0}^{t} \partial_{x}^{\alpha} \partial_{t}^{a} f(x, \tau) d \tau-\int_{0}^{s} \partial_{x}^{\alpha} \partial_{t}^{a} f(x, \tau) d \tau\right| /|t-s|^{\mu / 2} \\
& \quad \leqq \int_{0}^{t}\left(\left|\partial_{x}^{\alpha} \partial_{t}^{a} f(x, \tau)-\partial_{x}^{\alpha} \partial_{t}^{a} f(x, \tau-t+s)\right| /|t-s|^{\mu / 2}\right) d \tau
\end{aligned}
$$

if $|\alpha|+2 a \leqq[r], 0 \leqq s \leqq t \leqq T$ and $0<\mu<2$. Then the lemma follows from (i).

We introduce a special class of operators.
Definition 8.A.2. Let $M$ and $M^{\prime}$ be compact $C^{\infty}$ manifolds. Let $0 \leqq k \leqq r$. Let $C$ and $N$ be positive constants. We say that a linear operator $\mathscr{K}: C_{\sharp}^{(r)}(M \times[0, T]) \rightarrow C_{\sharp}^{(r)}\left(M^{\prime} \times[0, T]\right)$ is ( $k, r, C, N$ )-balanced (resp. ( $k, r, C, N$ )-integral-balanced) if:
(i) $|\mathscr{K} f|_{(k), t} \leqq C|f|_{(k), t}\left(\right.$ resp. $\left.|\mathscr{K} f|_{(k), t} \leqq C \int_{0}^{t}|f|_{(k), \tau} d \tau\right)$ for $t \in[0, T]$.
(ii) $|\mathscr{K} f|_{(r), t} \leqq C|f|_{(r), t}+C N|f|_{(k), t}\left(\right.$ resp. $|\mathscr{K} f|_{(r), t} \leqq C \int_{0}^{t}|f|_{(r), \tau} d \tau+$ $\left.C N \int_{0}^{t}|f|_{(k), \tau} d \tau\right)$ for $t \in[0, T]$.

Lemma 8.A.3. Let $M$ and $M^{\prime}$ be compact $C^{\infty}$ manifolds. Let $0 \leqq$ $k \leqq r$. Let $C$ and $N$ be positive constants. Let

$$
\mathscr{K}^{\prime}: C_{\sharp}^{(r)}(M \times[0, T]) \rightarrow C_{\sharp}^{(r)}\left(M^{\prime} \times[0, T]\right)
$$

be a ( $k, r, C, N$ )-balanced linear operator. Define a linear operator

$$
\mathscr{J}: C_{\#}^{(r)}(M \times[0, T]) \rightarrow C_{\#}^{(r)}\left(M^{\prime} \times[0, T]\right)
$$

$b y$

$$
(\mathscr{J} f)(x, t)=\int_{0}^{t}(\mathscr{K} f)(x, \tau) d \tau \quad \text { for } \quad(x, t) \in M^{\prime} \times[0, T]
$$

Then $\mathscr{J}$ is $(k, r, C, N)$-integral-balanced.
Proof. The lemma follows immediately from Lemma 8.A.1.
Lemma 8.A.4. Let $M$ and $M^{\prime}$ be compact $C^{\infty}$ manifolds. Let $0 \leqq$ $k \leqq r$. Let $C$ and $N$ be positive constants. Let $\mathscr{K}: C_{\#}^{(r)}(M \times[0, T]) \rightarrow$ $C_{\ddagger}^{(r)}(M \times[0, T])$ and $\mathscr{K}^{\prime}: C_{\ddagger}^{(r)}\left(M^{\prime} \times[0, T]\right) \rightarrow C_{\ddagger}^{(r)}(M \times[0, T])$ be $(k, r, C, N)-$ integral-balanced linear operators. Let $g$ belong to $C_{\neq}^{(r)}\left(M^{\prime} \times[0, T]\right)$. Then the equation $u-\mathscr{K}^{\prime} u=\mathscr{K}^{\prime} g$ has a unique solution $u \in C_{\#}^{(r)}(M \times[0, T])$. Further the linear operator $(I-\mathscr{K})^{-1} \mathscr{K}^{\prime}: C_{\#}^{(r)}\left(M^{\prime} \times[0, T]\right) \rightarrow C_{\#}^{(r)}(M \times[0, T])$
is ( $k, r, C^{\prime}, N$ )-integral-balanced, where $I$ denotes the identity operator and $C^{\prime}$ is a constant bounded with $C$.

Proof. By induction on $i$, we have:
(i) $\left|\mathscr{K}^{i} \mathscr{K}^{\prime} g\right|_{(k), t} \leqq C^{i+1}\left(t^{i} / i!\right) \int_{0}^{t}|g|_{(k), \tau} d \tau$ for $i=0,1, \cdots$ and $t \in$ $[0, T]$.
(ii) $\left|\mathscr{K}^{i} \mathscr{K}^{\prime} g\right|_{(r), t} \leqq C^{i+1}\left(t^{i} / i!\right) \int_{0}^{t}|g|_{(r), \tau} d \tau+(i+1) C^{i+1}\left(t^{i} / i!\right) N \int_{0}^{t}|g|_{(k), \tau} d \tau$ for $i=0,1, \cdots$ and $t \in[0, T]$.
Then, from the estimates of the Neumann series $u=\mathscr{K}^{\prime} g+\mathscr{K}_{\mathscr{K}} \mathscr{F}^{\prime} g+\cdots$, the lemma follows.
(B) Throughout (B), we use the following notations.
(a) $\left\{U_{1}, \cdots, U_{k}\right\}$ is a finite covering of $\Gamma_{0}$ by coordinate neighborhoods and let ( $\omega_{1}, \cdots, \omega_{n-1}$ ) be local coordinates.
(b) $\left\{\eta_{1}, \cdots, \eta_{k}\right\}$ is a $C^{\infty}$ partition of unity subordinate to $\left\{U_{1}, \cdots, U_{\kappa}\right\}$.
(c) $\left(N_{0}^{-}\right)_{\nu}=\left\{x \in N_{0}^{-} ; \omega(x) \in U_{\nu}\right\}$ for $\nu=1, \cdots, \kappa$.
(d) Let $\mathscr{H}_{1}$ be a differential operator on $\Gamma_{T}$ of the form $\mathscr{H}_{1}=\partial_{t}+$ $\sum_{i=1}^{n-1} H_{\nu, i}(\omega, t) \partial_{\omega_{i}}$ for $\nu=1, \cdots, \kappa$ and $(\omega, t) \in U_{\nu} \times[0, T]$.
(e) Let $\mathscr{P}_{1}$ be a differential operator on $\bar{\Omega}_{T}$ of the form:

$$
\begin{equation*}
\mathscr{P}_{1}=\partial_{t}+\chi_{0}(\lambda(x)) \sum_{i=1}^{n-1} H_{\nu, i}(\omega(x), t) \partial_{\omega_{i}}+\chi_{0}(\lambda(x)) H_{\lambda}(\omega(x), t) \partial_{\lambda} \tag{e.i}
\end{equation*}
$$

for $\nu=1, \cdots, \kappa$ and $(x, t) \in\left(N_{0}^{-}\right)_{\nu} \times[0, T]$. Here $\chi_{0}$ is the function introduced in Section 1 and $\partial_{\omega_{i}}$ is the partial differentiation with respect to $\omega_{i}$ in the ( $\omega_{1}, \cdots, \omega_{n-1}, \lambda$ ) coordinates.

$$
\begin{equation*}
\mathscr{P}_{1}=\partial_{t} \text { for } \quad(x, t) \in\left(\bar{\Omega}_{0}-N_{0}^{-}\right) \times[0, T] \tag{e.ii}
\end{equation*}
$$

Definition 8.B.1. Let $r \geqq 0$. Then $|H|_{(r)}=\sum_{v=1}^{k} \sum_{i=1}^{n-1}\left|\eta_{\nu} H_{\nu, i}\right|_{(r)}$, where $\eta_{\nu} H_{\nu, i}$ is regarded as a function on $R^{n-1} \times[0, T]$ in the canonical manner.

Definition 8.B.2. Let $0 \leqq k \leqq r$. Let $B$ be a positive constant. We say that $\mathscr{E}_{1}$ is a $(k, r, B)$-fine operator if:
(i) $\eta_{\nu} H_{\nu, i} \in C_{\sharp}^{(r)}\left(R^{n-1} \times[0, T]\right)$ for $\nu=1, \cdots, \kappa$ and $i=1, \cdots, n-1$.
(ii) $|H|_{(k)} \leqq B$.

We say that $\mathscr{P}_{1}$ is a $(k, r, B)-\lambda$-fine operator if:
(i) $\eta_{\nu} H_{\nu, i} \in C_{\sharp}^{(r)}\left(R^{n-1} \times[0, T]\right)$ for $\nu=1, \cdots, \kappa$ and $i=1, \cdots, n-1$.
(ii) $H_{\lambda} \in C_{\#}^{(r)}\left(\Gamma_{T}\right)$.
(iii) $|H|_{(k)}+\left|H_{\lambda}\right|_{(k)} \leqq B$.
(iv) $\quad H_{\lambda}(\omega, t) \geqq 0$ for $(\omega, t) \in \Gamma_{T}$.

Here $\eta_{\nu} H_{\nu, i}$ is regarded as a function on $\boldsymbol{R}^{n-1} \times[0, T]$ in the canonical manner.

Definition 8.B.3. Let $f$ be a mapping of $\Gamma_{T}$ into $\Gamma_{0}$. Then, for $i=1, \cdots, n$, a function $f_{i}$ on $\Gamma_{T}$ is the $i$-th component of $f$ when it is regarded as a mapping of $\Gamma_{T}$ into $\boldsymbol{R}^{n}$. Further, let $r \geqq 0$ and let $f_{1}, f_{2}, \cdots, f_{n}$ belong to $C^{(r)}\left(\Gamma_{T}\right)$. Then $|f|_{(r)}=\sum_{i=1}^{n}\left|f_{i}\right|_{(r)}$.

Definition 8.B.4. Let $I$ be an interval in $\boldsymbol{R}$. Let $D$ be a domain with piecewise- $C^{1}$ boundary in $\boldsymbol{R}^{n} \times I$. Let $r \geqq 0$. Let $f$ be a mapping of $\bar{D}$ into $\boldsymbol{R}^{n}$ whose components $f_{1}, f_{2}, \cdots, f_{n}$ belong to $C^{(r)}(\bar{D})$. Then $|f|_{(r)}=\sum_{i=1}^{n}\left|f_{i}\right|_{(r)}$.

Lemma 8.B.5. Let $2 \leqq k \leqq r$. Let B be a positive constant. Suppose that $\mathscr{H}_{1}$ is a $(k, r, B)$-fine operator. Define a mapping $\phi: \Gamma_{T} \rightarrow \Gamma_{0}$ so that the mapping $t \mapsto(\phi(\omega, t), t)$ is the characteristic curve of $\mathscr{\mathscr { H }}_{1}$ starting from $\omega \in \Gamma_{0}$ at $t=0$. Define a mapping $\psi: \Gamma_{T} \rightarrow \Gamma_{0}$ by $\phi(\psi(\omega, t), t)=\omega$ for $(\omega, t) \in \Gamma_{T}$. (Since $\Gamma_{0}$ is a compact $C^{\infty}$ manifold without boundary, $\phi$ and $\psi$ are well-defined.) Then we have:
(i) $\phi_{1}, \cdots, \phi_{n} \in C^{(r)}\left(\Gamma_{r}\right)$.
(ii) $\partial_{t} \phi_{1}, \cdots, \partial_{t} \phi_{n} \in C_{\sharp}^{(r-2)}\left(\Gamma_{T}\right)$.
(iii) $|\phi|_{(r)} \leqq C\left(1+|H|_{(r)}\right)$.

Here $C$ is a constant bounded with $r$ and $B$. Further, suppose that $0<T \leqq C_{B}^{-1}$ where $C_{B}$ is a constant determined by $B$. Then we have:
( $\mathrm{i}^{\prime}$ ) $\quad \psi_{1}, \cdots, \psi_{n} \in C^{(r)}\left(\Gamma_{T}\right)$.
(ii') $\partial_{t} \psi_{1}, \cdots, \partial_{t} \psi_{n} \in C_{\#}^{(r-2)}\left(\Gamma_{T}\right)$.
(iii') $|\psi|_{(r)} \leqq C^{\prime}\left(1+|H|_{(r)}\right)$.
Here $C^{\prime}$ is a constant bounded with $r$ and $B$.
Proof. Let $D$ be a bounded domain in $\boldsymbol{R}^{n}$ with $\Gamma_{0} \subset D$. Extend $\mathscr{H}_{1}$ to $\bar{D} \times[0, T]$ so that:
(i) The extended operator $\tilde{\mathscr{H}}_{1}$ is of the form $\tilde{\mathscr{H}}_{1}=\partial_{t}+\sum_{i=1}^{n} \tilde{H}_{i}(x, t) \partial_{x_{i}}$ with $\widetilde{H}_{i} \in C_{\sharp}^{(r)}(\bar{D} \times[0, T])$.
(ii) $\operatorname{supp} \widetilde{H}_{i} \subset D$ for $i=1, \cdots, n$.
(iii) $|\widetilde{H}|_{(q)} \leqq C_{1}|H|_{(q)}$ for $q \in[0, r]$.

Here $C_{1}$ is a constant bounded with $r$ and $|\widetilde{H}|_{(q)}=\sum_{i=1}^{n}\left|\widetilde{H}_{i}\right|_{(q)}$. Define a mapping $\tilde{\phi}: \bar{D} \times[0, T] \rightarrow \boldsymbol{R}^{n}$ by:

$$
\begin{equation*}
d_{t} \tilde{\phi}_{i}(x, t)=\widetilde{H}_{i}(\tilde{\phi}(x, t), t) \tag{*.1}
\end{equation*}
$$

for $i=1, \cdots, n$ and $(x, t) \in \bar{D} \times[0, T]$, where $\tilde{\phi}_{i}$ is the $i$-th component of $\tilde{\phi}$.

$$
\begin{equation*}
\tilde{\phi}(x, 0)=x \quad \text { for } \quad x \in \bar{D} \tag{*.2}
\end{equation*}
$$

Clearly $\tilde{\phi}$ is an extension of $\phi$ with $\tilde{\phi}(\bar{D} \times[0, T])=\bar{D}$. Since $H_{i} \in C^{1}$, the mapping $\tilde{\phi}$ is $C^{1}$. By (*.1), we easily observe that: If $2 \leqq q \leqq r$
and $\tilde{\phi}_{1}, \cdots, \tilde{\phi}_{n} \in C^{(q)}(\bar{D} \times[0, T])$, then $\partial_{t} \tilde{\phi}_{1}, \cdots, \partial_{t} \tilde{\phi}_{n} \in C_{\#}^{(q-2)}(\bar{D} \times[0, T])$.
Differentiating (*.1) and (*.2) with respect to $x_{j}$, we have:

$$
\begin{align*}
d_{t}\left[\partial_{x_{j}} \tilde{\phi}_{i}(x, t)\right]= & \sum_{n=1}^{n}\left[\partial_{x_{h}} \widetilde{H}_{i}(\tilde{\phi}(x, t), t)\right]\left[\partial_{x_{j}} \tilde{\phi}_{h}(x, t)\right]  \tag{1}\\
& \partial_{x_{j}} \tilde{\phi}_{i}(x, 0)=\delta_{i j} \tag{2}
\end{align*}
$$

Hence we have

$$
\begin{equation*}
\partial_{x_{j}} \tilde{\phi}_{i}(x, t)=\delta_{i j}+\sum_{h=1}^{n} \int_{0}^{t}\left[\partial_{x_{h}} \widetilde{H}_{i}(\dot{\phi}(x, \tau), \tau)\right]\left[\partial_{x_{j}} \tilde{\phi}_{h}(x, \tau)\right] d \tau \tag{**}
\end{equation*}
$$

for $i, j=1, \cdots, n$ and $(x, t) \in \bar{D} \times[0, T]$.
Let $W$ be the vector space of continuous mappings of $\bar{D} \times[0, T]$ into $\boldsymbol{R}^{n}$. Define a linear operator $\mathscr{G}: W \rightarrow W$ by

$$
(\mathscr{G} f)_{i}(x, t)=\sum_{n=1}^{n} \int_{0}^{t}\left[\partial_{x_{h}} \widetilde{H}_{i}(\tilde{\phi}(x, \tau), \tau)\right]\left[f_{h}(x, \tau)\right] d \tau
$$

for $i=1, \cdots, n$ and $(x, t) \in \bar{D} \times[0, T]$. For $i=1, \cdots, n$, define a mapping $\delta_{i}$ in $W$ by $\delta_{i}(x, t)=\left(\delta_{i 1}, \cdots, \delta_{i n}\right)$ for $(x, t) \in \bar{D} \times[0, T]$. Then, by ( $* *$ ), we have $\partial_{x_{i}} \tilde{\phi}=\delta_{i}+\mathscr{G} \delta_{i}+\mathscr{G}^{2} \delta_{i}+\cdots$. This and (*.1) give $|\tilde{\phi}|_{(1)} \leqq C_{2}$. Here $C_{2}$ is a constant bounded with $B$. Further, by Lemma 3.A. 9 and an argument similar to that in the proof of Lemma 8.A.4, we observe that if $1 \leqq q \leqq r-1$ and $\tilde{\phi}_{1}, \cdots, \tilde{\phi}_{n} \in C^{(q)}(\bar{D} \times[0, T])$, then $\partial_{x_{i}} \tilde{\phi}_{1}, \cdots, \partial_{x_{i}} \tilde{\phi}_{n} \in$ $C^{(q)}(\bar{D} \times[0, T])$ and $\left|\partial_{x_{i}} \tilde{\phi}\right|_{(q)} \leqq C_{3}\left(1+|H|_{(q+1)}+|\tilde{\phi}|_{(q)}\right)$. Here $C_{3}$ is a constant bounded with $r$ and $B$.

Note that if $f, \partial_{x_{1}} f, \cdots, \partial_{x_{n}} f, \partial_{t} f \in C^{(q)}(\bar{D} \times[0, T])$, then $f \in C^{(q+1)}(\bar{D} \times$ $[0, T])$ and $|f|_{(q+1)} \leqq C_{4}\left(|f|_{(q)}+\sum_{i=1}^{n}\left|\partial_{x_{i}} f\right|_{(q)}+\left|\partial_{t} f\right|_{(q)}\right)$. Consequently, with the aid of (*.1), we have: If $1 \leqq q \leqq r-1$ and $\tilde{\phi}_{1}, \cdots, \tilde{\phi}_{n} \in C^{(q)}(\bar{D} \times[0, T])$, then $\tilde{\phi}_{1}, \cdots, \tilde{\phi}_{n} \in C^{(q+1)}(\bar{D} \times[0, T])$ and $|\tilde{\phi}|_{(q+1)} \leqq C_{5}\left(1+|H|_{(q+1)}+|\tilde{\phi}|_{(q)}\right)$. Here $C_{5}$ is a constant bounded with $r$ and $B$. From this, we obtain (i)-(iii) in the lemma at once.

Define a mapping $\tilde{\psi}: \bar{D} \times[0, T] \rightarrow \bar{D}$ by

$$
\begin{equation*}
\tilde{\phi}(\tilde{\psi}(x, t), t)=x \quad \text { for } \quad(x, t) \in \boldsymbol{R}^{n} \times[0, T] \tag{***}
\end{equation*}
$$

Clearly $\tilde{\psi}$ is well-defined and is an extension of $\psi$. Since $H_{i}$ is $C^{1}$, the mapping $\tilde{\psi}$ is $C^{1}$.

By (*.1) and (*.2), we have $\tilde{\phi}_{i}(x, t)=x_{i}+\int_{0}^{t} \tilde{H}_{i}(\tilde{\phi}(x, \tau), \tau) d \tau$ for $i=$ $1, \cdots, n$ and $(x, t) \in \boldsymbol{R}^{n} \times[0, T]$. Hence there is a positive constant $C_{B}$ determined by $B$ such that if $0<T \leqq C_{B}^{-1}$, then $\left|\operatorname{det}\left[\partial_{x_{i}} \tilde{\phi}_{j}(x, t)\right]\right| \geqq 1 / 2$ for $(x, t) \in \bar{D} \times[0, T]$. On the other hand, by (***), we have:
(i) $\sum_{h=1}^{n}\left[\partial_{x_{h}} \tilde{\phi}_{i}(\tilde{\psi}(x, t), t)\right]\left[\partial_{x_{j}} \tilde{\psi}_{h}(x, t)\right]=\delta_{i j}$ for $i, j=1, \cdots, n$ and $(x, t) \in$ $\bar{D} \times[0, T]$.
(ii) $\quad \sum_{h=1}^{n}\left[\partial_{x_{n}} \tilde{\phi}_{i}(\tilde{\psi}(x, t), t)\right]\left[\partial_{t} \tilde{\psi}_{h}(x, t)\right]+\partial_{t} \tilde{\phi}_{i}(\tilde{\psi}(x, t), t)=0$ for $i=1, \cdots, n$ and $(x, t) \in \bar{D} \times[0, T]$.
Further, note that $\partial_{t} \tilde{\phi}_{i}(\tilde{\psi}(x, t), t)=\widetilde{H}_{i}(\tilde{\phi}(\tilde{\psi}(x, t), t), t)$. Then, with the aid of Lemmas 3.A. 9 and 3.A.10, by an argument similar to that in the proof of (i)-(iii) in the lemma, we obtain ( $\mathrm{i}^{\prime}$ )-(iii') in the lemma. This completes the proof of Lemma 8.B.5.

Lemma 8.B.6. Let $2 \leqq k \leqq r$. Let $B$ be a positive constant. Suppose that $0<T \leqq C_{B}^{-1}$ where $C_{B}$ is a constant determined by $B$. Suppose that $\mathscr{P}_{1}$ is a ( $k, r, B$ )- $\lambda$-fine operator. Define a mapping $\phi: \widetilde{\Omega}_{T} \rightarrow \bar{\Omega}_{0}$ so that:
(i) The mapping $t \mapsto(\phi(x, t), t)$ is the characteristic curve of $\mathscr{P}_{1}$ starting from $x \in \bar{\Omega}_{0}$ at $t=0$.
(ii) $\widetilde{\Omega}_{T}=\left\{(x, t) ; x \in \bar{\Omega}_{0}\right.$ and $\left.t \in\left[0, T_{x}\right]\right\}$.

Here $T_{x}$ is the largest number in $[0, T]$ such that $\phi(x, t)$ can be defined on $\left[0, T_{x}\right]$. (We may assume that the boundary of $\widetilde{\Omega}_{T}$ in $\boldsymbol{R}^{n} \times[0, T]$ is $C^{1}$.) Define a mapping $\psi: \bar{\Omega}_{T} \rightarrow \bar{\Omega}_{0}$ by $\phi(\psi(x, t), t)=x$ for $(x, t) \in \bar{\Omega}_{T}$. (Since $H_{\lambda}(\omega, t) \geqq 0$ for $(\omega, t) \in \Gamma_{T}$, the mapping $\psi$ is well-defined.) Then we have:
(i) $\phi_{1}, \cdots, \dot{\phi}_{n} \in C^{(r)}\left(\widetilde{\Omega}_{T}\right)$.
(ii) $\partial_{t} \phi_{1}, \cdots, \partial_{t} \phi_{n} \in C_{\#}^{(r-2)}\left(\widetilde{\Omega}_{T}\right)$.
(iii) $|\phi|_{(r)} \leqq C\left(1+|H|_{(r)}+\left|H_{2}\right|_{(r)}\right)$.
(iv) $\psi_{1}, \cdots, \psi_{n} \in C^{(r)}\left(\bar{\Omega}_{T}\right)$.
(v) $\partial_{t} \psi_{1}, \cdots, \partial_{t} \psi_{n} \in C_{\#}^{(r-2)}\left(\bar{\Omega}_{T}\right)$.
(vi) $|\psi|_{(r)} \leqq C\left(1+|H|_{(r)}+\left|H_{\lambda}\right|_{(r)}\right)$.

Here $C$ is a constant bounded with $r$ and $B$.
Proof. By an argument similar to the proof of Lemma 8.B.5, the lemma is proved.

Lemma 8.B.7. Let $2 \leqq k \leqq r$. Let $B, C$, and $N$ be positive constants. Suppose that:
(i) The operator $\mathscr{H}_{1}$ is $(k, r, B)$-fine.
(ii) A linear operator $\mathscr{K}: C_{\sharp}^{(r)}\left(\Gamma_{T}\right) \rightarrow C_{\#}^{(r)}\left(\Gamma_{T}\right)$ is $(k, r, C, N)$-balanced.
(iii) $g \in C_{\#}^{(r)}\left(\Gamma_{T}\right)$.
(iv) $0<T \leqq C_{B}^{-1}$.

Here $C_{B}$ is the constant introduced in Lemma 8.B.5. Consider the Cauchy problem:

$$
\begin{gather*}
\mathscr{H}_{1} u=\mathscr{K} u+g \text { on } \Gamma_{T} .  \tag{1}\\
\left.u\right|_{t=0}=0 . \tag{2}
\end{gather*}
$$

Then the problem (1)-(2) has a unique solution $u \in C_{\sharp}^{(r)}\left(\Gamma_{T}\right)$. Further the linear operator $\left(\mathscr{H}_{1}-\mathscr{K}\right)^{-1}$ : $C_{\#}^{(r)}\left(\Gamma_{T}\right) \rightarrow C_{\#}^{(r)}\left(\Gamma_{T}\right)$ is $\left(k, r, C^{\prime}, 1+N+|H|_{(r)}\right)$ -integral-balanced. Here $C^{\prime}$ is a constant bounded with $r$, $B$, and $C$.

Proof. We use the notations in Lemma 8.B.5. For $(\omega, t) \in \phi\left(U_{\nu} \times\right.$ $[0, T])$, we call the $n$-tuple $\left(\psi_{1}(\omega, t), \cdots, \psi_{n-1}(\omega, t), t\right)$ the characteristic coordinates of ( $\omega, t$ ). Here $\psi_{i}$ is the $i$-th component of the mapping $\psi: \phi\left(U_{\nu} \times[0, T]\right) \rightarrow U_{\nu}$ with respect to the local coordinates in $U_{\nu}$. As is well known, in the characteristic coordinates, (1) in the lemma is an ordinary differential equation. Then, from Lemmas 3.A.9, 8.A.3, 8.A.4 and 8.B.5, the lemma follows immediately.

Lemma 8.B.8. Let $2 \leqq k \leqq r$. Let $B, C$, and $N$ be positive constants. Suppose that:
(i) The operator $\mathscr{P}_{1}$ is $(k, r, B)-\lambda-f i n e$.
(ii) A linear operator $\mathscr{K}^{\prime}: C_{\ddagger}^{(r)}\left(\bar{\Omega}_{T}\right) \rightarrow C_{\sharp}^{(r)}\left(\bar{\Omega}_{T}\right)$ is ( $\left.k, r, C, N\right)$-balanced.
(iii) $g \in C_{\#}^{(r)}\left(\bar{\Omega}_{T}\right)$.
(iv) $0<T \leqq C_{B}^{-1}$.

Here $C_{B}$ is the constant introduced in Lemma 8.B.6. Consider the Cauchy problem:

$$
\begin{gather*}
\mathscr{P}_{1} u=\mathscr{K} u+g \quad \text { in } \quad \Omega_{T}  \tag{1}\\
\left.u\right|_{t=0}=0 \tag{2}
\end{gather*}
$$

Then the problem (1)-(2) has a unique solution $u \in C_{\#^{(r)}}\left(\bar{\Omega}_{T}\right)$. Further, the linear operator $\left(\mathscr{P}_{1}-\mathscr{K}^{\prime}\right)^{-1}: C_{\sharp}^{(r)}\left(\bar{\Omega}_{T}\right) \rightarrow C_{\sharp}^{(r)}\left(\bar{\Omega}_{T}\right)$ is $\left(k, r, C^{\prime}, 1+N+|H|_{(r)}+\right.$ $\left.\left|H_{\lambda}\right|_{(r)}\right)$-integral-balanced. Here $C^{\prime}$ is a constant bounded with $r, B$, and $C$.

Proof. We can easily obtain modifications of Lemmas 8.A. 3 and 8.A. 4 for the $C_{\sharp}^{(r)}\left(\bar{\Omega}_{T}\right)$ case. Then, by using Lemma 8.B. 6 in place of Lemma 8.B.5, the lemma is proved as in Lemma 8.B.7.

Lemma 8.B.9. Let $\mathscr{L}$ be a differential operator on $\bar{\Omega}_{T}$ which is expressed as in (B) in Section 3. Let $k$ and $i$ be integers with $2 \leqq k \leqq i$. Let $\sigma, B, C$, and $N$ be positive constants. Suppose that:
(i) The operator $\mathscr{L}$ is $\left(\varepsilon_{0}+k, \varepsilon_{0}+i, \sigma, B\right)$-\#-parabolic.
(ii) The operator $\mathscr{P}_{1}$ is $\left(\varepsilon_{0}+k, \varepsilon_{0}+i, B\right)-\lambda$-fine.
(iii) A linear operator $\mathscr{K}: C_{\#}^{\left(\varepsilon_{0}+i\right)}\left(\bar{\Omega}_{T}\right) \rightarrow C_{\#}^{\left(\varepsilon_{0}+i\right)}\left(\bar{\Omega}_{T}\right)$ is $\left(\varepsilon_{0}+k, \varepsilon_{0}+i\right.$, $C, N)$-balanced.
(iv) $g \in C_{\#}^{\left(\delta_{0}+i\right)}\left(\Gamma_{T}\right)$.
(v) $0<T \leqq C_{B}^{-1}$.

Here $C_{B}$ is a constant introduced in Lemma 8.B.6. Consider the problem:

$$
\begin{equation*}
\mathscr{L} u=0 \quad \text { in } \quad \Omega_{T} \tag{1}
\end{equation*}
$$

$$
\begin{equation*}
\left.u\right|_{t=0}=0 \tag{2}
\end{equation*}
$$

$$
\begin{equation*}
u=0 \quad \text { on } \quad J_{T} . \tag{3}
\end{equation*}
$$

$$
\begin{equation*}
\mathscr{P}_{1} u=\mathscr{K} u+g \quad \text { on } \quad \Gamma_{T} . \tag{4}
\end{equation*}
$$

Then the problem (1)-(4) has a solution $u \in C_{\xi^{\left(\varepsilon_{0}+i\right)}\left(\bar{\Omega}_{T}\right) \text {. Further, the }}$ mapping $g \mapsto u$, where $u$ is the obtained solution of (1)-(4), is an $\left(\varepsilon_{0}+k\right.$, $\left.\varepsilon_{0}+i, C^{\prime}, 1+N+|A|_{\left(\varepsilon_{0}+i\right)}+|H|_{\left(\varepsilon_{0}+i\right)}+\left|H_{2}\right|_{\left(\varepsilon_{0}+i\right)}\right)$-integral-balanced linear operator on $C_{\#}^{\left(\varepsilon_{0}+i\right)}\left(\Gamma_{T}\right)$ into $C_{\#}^{\left(\varepsilon_{0}+i\right)}\left(\bar{\Omega}_{T}\right)$. Here $C^{\prime}$ is a constant bounded with $i, \sigma, B$, and $C$ and $|A|_{\left(\varepsilon_{0}+i\right)}=\sum_{g, h=1}^{n}\left|A_{g h}\right|_{\left(\varepsilon_{0}+i\right)}+\sum_{h=1}^{n}\left|A_{g}\right|_{\left(\varepsilon_{0}+i\right)}+\left|A_{0}\right|_{\left(\varepsilon_{0}+i\right)}$.

Proof. If $u$ belongs to $C_{\#}^{\left(\varepsilon_{0}+i\right)}\left(\bar{\Omega}_{T}\right)$, then $u$ satisfies (1)-(4) if and only if it satisfies:

$$
\begin{gather*}
\mathscr{L} \mathscr{P}_{1} u=\left[\mathscr{L}, \mathscr{P}_{1}\right] u \text { in } \Omega_{T} . \\
\left.\left(\mathscr{P}_{1} u\right)\right|_{t=0}=0 . \\
\mathscr{P}_{1} u=0 \quad \text { on } J_{T} . \\
\mathscr{P}_{1} u=\mathscr{K} u+g \text { on } \Gamma_{T} .
\end{gather*}
$$

Here $\left[\mathscr{L}, \mathscr{P}_{1}\right]$ denotes the commutator, i.e., $\left[\mathscr{L}, \mathscr{P}_{1}\right]=\mathscr{L} \mathscr{P}_{1}-\mathscr{P}_{1} \mathscr{L}$. By Lemma 8.B.3, we observe that $\mathscr{P}_{1} u$ is expressed as $\mathscr{P}_{1} u=\mathscr{N} u+\mathscr{N} g$. Here $\mathscr{M}: C_{\#}^{\left(\varepsilon_{0}+i\right)}\left(\bar{\Omega}_{T}\right) \rightarrow C_{\#}^{\left(\varepsilon_{0}+i\right)}\left(\bar{\Omega}_{T}\right)$ (resp. $\mathscr{N}: C_{\#}^{\left(\varepsilon_{0}+i\right)}\left(\Gamma_{T}\right) \rightarrow C_{\sharp}^{\left(\varepsilon_{0}+i\right)}\left(\bar{\Omega}_{T}\right)$ ) is an $\left(\varepsilon_{0}+k, \varepsilon_{0}+i, C_{1}, N+|A|_{\left(\varepsilon_{0}+i\right)}+|H|_{\left(\varepsilon_{0}+i\right)}+\left|H_{\lambda}\right|_{\left(\varepsilon_{0}+i\right)}\right)$-balanced (resp. $\left(\varepsilon_{0}+k\right.$, $\left.\varepsilon_{0}+i, C_{2},|A|_{\left(\varepsilon_{0}+i\right)}\right)$-balanced) linear operator where $C_{1}$ (resp. $C_{2}$ ) is a constant bounded with $i, \sigma, B$ and $C$ (resp. $i, \sigma$ and $B$ ). Then, the lemma is proved by Lemma 8.B.8.
9. Inversion of $D \mathscr{F}(\rho)$. Finally we verify the condition (III) in Nash's implicit function theorem.

We solve the equation $D \mathscr{F}(\rho) \delta \rho=\delta G$ where $\rho$ and $\delta G$ are given and $\delta \rho$ is unknown. By the definition of $D \mathscr{F}(\rho)$, the above equation means:

$$
\mathscr{L}_{\rho}\left(\delta U_{\rho}\right)=\left(\delta \mathscr{E}_{\rho}\right) U_{\rho} \quad \text { in } \quad \Omega_{T}
$$

(2 ${ }_{\text {oU }}$ )

$$
\left.\delta U_{\rho}\right|_{t=0}=0 .
$$

$$
\begin{align*}
\delta U_{\rho} & =0 \quad \text { on } J_{T} . \\
\delta U_{\rho} & =0 \quad \text { on } \quad \Gamma_{T} . \\
\partial_{t} \delta \rho+c_{0}\left\{\left(\partial_{\lambda} U_{\rho}\right) \delta S_{\rho}\right. & \left.+\left[\partial_{\lambda}\left(\delta U_{\rho}\right)\right] S_{\rho}\right\}=\delta G \quad \text { on } \quad \Gamma_{T} .
\end{align*}
$$

Throughout this section, suppose that:
(i) $i=1, \cdots, 11$.
(ii) $\rho \in V_{T} \cap C_{F^{\left(r_{0}-4+4 i\right)}}\left(\Gamma_{T}\right)$.
(iii) $\delta G \in C_{\#}^{\left(r_{0}-7+4 i\right)}\left(\Gamma_{T}\right)$.

Under these hypotheses, we seek $\delta \rho \in C_{\sharp}^{\left(r_{0}-8+4 i\right)}\left(\Gamma_{T}\right)$ which satisfies $\left(1_{\delta U}\right)$ $\left(5_{\partial U}\right)$ when $i \geqq 2$ and set $\mathscr{J}(\rho) \delta G=\delta \rho$.
(A) We eliminate $\delta \rho$ from the problem.

For $\delta \rho \in C_{\sharp}^{\left(r_{0}\right)}\left(\Gamma_{T}\right)$, define a mapping $\delta e_{\rho}: \bar{\Omega}_{T} \rightarrow \boldsymbol{R}^{n} \times[0, T]$ by:

$$
\begin{align*}
& \delta e_{\rho}(x(\omega, \lambda), t)=\left(\partial_{\lambda} x\left(\omega, \lambda+\chi_{0}(\lambda) \rho(\omega, t)\right) \delta \rho(\omega, t), 0\right)  \tag{9.1.i}\\
& \text { for } \quad(x, t)=(x(\omega, \lambda), t) \in N_{0}^{-} \times[0, T] . \\
& \delta e_{\rho}(x, t)=(0,0) \text { for } \quad(x, t) \in\left(\bar{\Omega}_{0}-N_{0}^{-}\right) \times[0, T] . \tag{9.1.ii}
\end{align*}
$$

(See (1.2).) Let

$$
\begin{equation*}
u_{\rho}=U_{\rho} \circ e_{\rho}^{-1} \tag{9.2}
\end{equation*}
$$

Let

$$
\begin{equation*}
\delta V=\delta U_{\rho}-\left\langle\left(\operatorname{grad} u_{\rho}\right) \circ e_{\rho}, \delta e_{\rho}\right\rangle \tag{9.3}
\end{equation*}
$$

First we show that

$$
\begin{equation*}
\mathscr{L}_{\rho} \delta V=0 \quad \text { in } \quad \Omega_{T} \tag{9.4}
\end{equation*}
$$

Fix $\rho$ and $\delta \rho$. Let $D$ be a relatively compact subdomain with $C^{\infty}$ boundary in $\Omega_{T}$. Let $\varepsilon$ be a sufficiently small positive number. By (7.3), (9.2), and (9.3), $u_{\rho+\varepsilon \delta \rho} \circ e_{\rho+\varepsilon \delta \rho}-u_{\rho} \circ e_{\rho}-\varepsilon \delta V-\varepsilon\left\langle\left(\operatorname{grad} u_{\rho}\right) \circ e_{\rho}, \delta e_{\rho}\right\rangle=\varepsilon^{2} R_{\varepsilon}$, where $R_{\varepsilon}$ is bounded in $C^{\left(r_{0}\right)}\left(\bar{\Omega}_{T}\right)$. Hence, by (9.1),

$$
\begin{align*}
& u_{\rho+\varepsilon \delta \rho \rho} \circ e_{\rho+\varepsilon \delta \rho}-u_{\rho} \circ e_{\rho+\varepsilon \delta \rho}-\varepsilon \delta V  \tag{9.5}\\
& \quad=-\left[u_{\rho} \circ e_{\rho+\varepsilon \delta \rho}-u_{\rho} \circ e_{\rho}-\varepsilon\left\langle\left(\operatorname{grad} u_{\rho}\right) \circ e_{\rho}, \delta e_{\rho}\right\rangle\right]+\varepsilon^{2} R_{\varepsilon} \\
& \quad=\varepsilon^{2} R_{\varepsilon}^{\prime} \text { on } \bar{D},
\end{align*}
$$

where $R_{\varepsilon}^{\prime}$ is bounded in $C^{\left(r_{0}-2\right)}(\bar{D})$. Note that $\mathscr{L}_{\rho+\varepsilon \delta \rho}\left(u_{\rho+\varepsilon \delta \rho} \circ e_{\rho+\varepsilon \delta \rho}\right)=$ $\mathscr{L}_{\rho+\varepsilon \delta \rho}\left(u_{\rho} \circ e_{\rho+\varepsilon \delta \rho}\right)=0$ on $\bar{D}$. Hence, by (9.5), $\mathscr{L}_{\rho+\varepsilon \delta \rho} \delta V=\varepsilon R_{\varepsilon}^{\prime \prime}$ on $\bar{D}$, where $R_{\varepsilon}^{\prime \prime}$ is bounded in $C^{\left(r_{0}-4\right)}(\bar{D})$. Taking $\varepsilon \rightarrow 0$, we obtain (9.4).

Secondly we eliminate $\delta \rho$ from the boundary condition on $\Gamma_{T}$. In a neighborhood of $\Gamma_{T}$, we have $U_{\rho}(x(\omega, \lambda), t)=u_{\rho}\left(e_{\rho}(x(\omega, \lambda), t), t\right)=$ $u_{\rho}(x(\omega, \lambda+\rho(\omega, t)), t)$, so that $\left\langle\left(\operatorname{grad} u_{\rho}\right) \circ e_{\rho}, \delta e_{\rho}\right\rangle=\left[\left(\partial_{\lambda} u_{\rho}\right) \circ e_{\rho}\right] \delta \rho=\left(\partial_{\lambda} U_{\rho}\right) \delta \rho$. Hence ( $4_{\partial U}$ ) and ( $5_{\partial U}$ ) can be rewritten in the form:
( $4_{\text {oU }}$ )

$$
\delta V+\left(\partial_{\lambda} U_{\rho}\right) \delta \rho=0 \quad \text { on } \quad \Gamma_{T} .
$$

( $5_{\text {бU }}$ )

$$
\mathscr{H}_{\rho} \delta \rho+c_{0} S_{\rho} \partial_{\lambda} \delta V=\delta G \quad \text { on } \quad \Gamma_{T} .
$$

Here

$$
\begin{align*}
\mathscr{\mathscr { C }}_{\rho}= & \partial_{t}+c_{0}\left[\left.\left(\partial_{\lambda} U_{\rho}\right)\right|_{\lambda=0}\right] \sum_{j=1}^{n-1}\left\{\left[\partial / \partial\left(\partial_{\omega_{j}} \rho\right)\right] S\right\} \partial_{\omega_{j}}  \tag{9.6}\\
& +c_{0}\left\{\left[\left.\left(\partial_{\lambda} U_{\rho}\right)\right|_{\omega=0}\right][(\partial / \partial \rho) S]+\left[\left.\left(\partial_{\lambda}^{2} U_{\rho}\right)\right|_{\lambda=0}\right] S\right\} .
\end{align*}
$$

(See (7.5).) For $g \in C^{0}\left(\Gamma_{T}\right)$, define $\mathscr{H}_{\rho}^{-1} g$ by:

$$
\begin{gather*}
\mathscr{H}_{p}\left(\mathscr{H}_{\rho}^{-1} g\right)=g \quad \text { on } \quad \Gamma_{T} .  \tag{1}\\
\left.\left(\mathscr{H}_{\rho}^{-1} g\right)\right|_{t=0}=0 . \tag{2}
\end{gather*}
$$

Then, by ( $5_{\partial U}$ ), we have

$$
\begin{equation*}
\delta \rho=-c_{0} \mathscr{K}_{\rho}^{-1}\left(S_{\rho} \partial_{\lambda} \delta V\right)+\mathscr{H}_{\rho}^{-1} \delta G . \tag{9.7}
\end{equation*}
$$

Inserting this in ( $4_{\partial U}$ ), we eliminate $\delta \rho$ from the boundary condition on $\Gamma_{r}$. Consequently we have:

| $\left(1_{\delta V}\right)$ | $\mathscr{L}_{\rho} \delta V=0$ | in $\Omega_{T}$. |  |  |
| :--- | :---: | :--- | :--- | :--- |
| $\left(2_{\partial V}\right)$ | $\left.\delta V\right\|_{t=0}=0$. |  |  |  |
| $\left(3_{\partial V}\right)$ | $\delta V=0$ | on $\quad J_{T}$. |  |  |
| $\left(4_{\delta V}\right)$ | $\delta V-c_{0}\left(\partial_{\lambda} U_{\rho}\right) \mathscr{H}_{\rho}^{-1}\left(S_{\rho} \partial_{\lambda} \delta V\right)=-\left(\partial_{\lambda} U_{\rho}\right) \mathscr{H}_{\rho}^{-1} \delta G$ | on | $\Gamma_{T}$. |  |

This completes the elimination of $\delta \rho$.
(B) Extend $S_{\rho}$ to $\bar{\Omega}_{T}$ by:

$$
\begin{align*}
S_{\rho}(x(\omega, \lambda), t)= & \chi_{0}(\lambda) S_{\rho}(\omega, t)+1-\chi_{0}(\lambda)  \tag{9.8.i}\\
& \text { for } \quad(x, t)=(x(\omega, \lambda), t) \in N_{0}^{-} \times[0, T] . \\
S_{\rho}(x, t)=1 & \text { for } \quad(x, t) \in\left(\bar{\Omega}_{0}-N_{0}^{-}\right) \times[0, T] .
\end{align*}
$$

Note that the extended $S_{\rho}$ also has a positive constant lower bound. Extend $\mathscr{H}_{\rho}$ to $\bar{\Omega}_{T}$ by:

$$
\begin{align*}
\mathscr{H}_{\rho}= & \partial_{t}+c_{0} \chi_{0}(\lambda)\left[\left.\left(\partial_{\lambda} U_{\rho}\right)\right|_{\lambda=0}\right] \sum_{j=1}^{n-1}\left\{\left[\partial / \partial\left(\partial_{\omega_{j}} \rho\right)\right] S\right\} \partial_{\omega_{j}}  \tag{9.9.i}\\
& +c_{0} \chi_{0}(\lambda)\left\{\left[\left.\left(\partial_{\lambda} U_{\rho}\right)\right|_{\lambda=0}\right][(\partial / \partial \rho) S]+\left[\left.\left(\partial_{\lambda}^{2} U_{\rho}\right)\right|_{\lambda=0}\right] S\right\} \\
& \text { on } N_{0}^{-} \times[0, T] .
\end{align*}
$$

$$
\begin{equation*}
\mathscr{H}_{\rho}=\partial_{t} \quad \text { on } \quad\left(\bar{\Omega}_{0}-N_{0}^{-}\right) \times[0, T] . \tag{9.9.ii}
\end{equation*}
$$

(See (9.6).) For $g \in C^{0}\left(\bar{\Omega}_{T}\right)$, define $\mathscr{H}_{\rho}^{-1} g$ by:

$$
\begin{gather*}
\mathscr{H}_{\rho}\left(\mathscr{H}_{\rho}^{-1} g\right)=g \text { in } \Omega_{T} .  \tag{1}\\
\left.\left(\mathscr{H}_{\rho}^{-1} g\right)\right|_{t=0}=0 . \tag{2}
\end{gather*}
$$

Set

$$
\begin{equation*}
\delta X=\mathscr{H}_{\rho}^{-1}\left(S_{\rho} \delta V\right) . \tag{9.10}
\end{equation*}
$$

We transform ( $\left.1_{\partial V}\right)-\left(4_{\partial V}\right)$ into a problem for $\delta X$.
First we transform ( $1_{\partial V}$ ) into an equation for $\delta X$. By ( $2_{\partial V}$ ), we have $\left.\left(\mathscr{L}_{\rho} \delta X\right)\right|_{t=0}=0$. Hence, by $\left(1_{\text {万V }}\right)$, we have

$$
\begin{align*}
\mathscr{L}_{\rho} \delta X= & \mathscr{H}_{\rho}^{-1} \mathscr{H}_{\rho} \mathscr{L}_{\rho} \mathscr{H}_{\rho}^{-1}\left(S_{\rho} \delta V\right)  \tag{9.11}\\
= & \mathscr{H}_{\rho}^{-1}\left[\mathscr{H}_{\rho}, \mathscr{L}_{\rho}\right] \delta X+\mathscr{H}_{\rho}^{-1} \mathscr{L}_{\rho}\left(S_{\rho} \delta V\right) \\
= & \mathscr{H}_{\rho}^{-1}\left[\mathscr{H}_{\rho}, \mathscr{L}_{\rho}\right] \delta X+\mathscr{H}_{\rho}^{-1}\left[\left(\mathscr{L}_{\rho} S_{\rho}\right) S_{\rho}^{-1}\left(\mathscr{H}_{\rho} \delta X\right)\right] \\
& -2 \mathscr{H}_{\rho}^{-1}\left\{\sum_{g, h=1}^{n} A_{\rho, g h}\left[\partial_{x_{g}}\left(S_{\rho}\right)\right]\left[\partial_{x_{h}}\left(S_{\rho}^{-1}\left(\mathscr{H}_{\rho} \delta X\right)\right)\right]\right\} .
\end{align*}
$$

## Clearly

$$
\begin{align*}
\partial_{x_{h}}\left[S_{\rho}^{-1}\left(\mathscr{H}_{\rho} \delta X\right)\right]= & -S_{\rho}^{-2}\left[\partial_{x_{h}}\left(S_{\rho}\right)\right]\left(\mathscr{H}_{\rho} \delta X\right)  \tag{9.12}\\
& +S_{\rho}^{-1}\left(\left[\partial_{x_{h}}, \mathscr{H}_{\rho}\right] \delta X\right)+S_{\rho}^{-1}\left(\mathscr{H}_{\rho} \partial_{x_{h}} \delta X\right) .
\end{align*}
$$

Let $\mathscr{\mathscr { C }}_{\rho, 1}$ be the homogeneous part of the first order of $\mathscr{H}_{\rho}$. By the obvious formula $\mathscr{H}_{\rho}(f g)=\left(\mathscr{H}_{\rho, 1} f\right) g+\left(\mathscr{H}_{\rho} g\right)$, we obtain

$$
\begin{equation*}
\mathscr{H}_{\rho}^{-1}(F G)=F\left(\mathscr{H}_{\rho}^{-1} G\right)-\mathscr{H}_{\rho}^{-1}\left[\left(\mathscr{H}_{\rho, 1} F\right)\left(\mathscr{H}_{\rho}^{-1} G\right)\right] . \tag{9.13}
\end{equation*}
$$

By (9.11), (9.12) and (9.13), we have

$$
\begin{align*}
& \mathscr{L}_{\rho} \delta X-\left(\mathscr{L}_{\rho} S_{\rho}\right) S_{\rho}^{-1} \delta X-2 \sum_{g, h=1}^{n} A_{\rho, g h}\left[\partial_{x_{g}}\left(S_{\rho}\right)\right]\left[\partial_{x_{h}}\left(S_{\rho}\right)\right] S_{\rho}^{-2} \delta X  \tag{9.14}\\
&+2 \sum_{g, h=1}^{n} A_{\rho, g h}\left[\partial_{x_{g}}\left(S_{\rho}\right)\right] S_{\rho}^{-1}\left(\partial_{x_{h}} \delta X\right) \\
&= \mathscr{H}_{\rho}^{-1}\left\{\left[\mathscr{H}_{\rho}, \mathscr{L}_{\rho}\right] \delta X-\left[\mathscr{H}_{\rho, 1}\left(\left(\mathscr{L}_{\rho} S_{\rho}\right) S_{\rho}^{-1}\right)\right] \delta X\right. \\
&-2 \sum_{g, h=1}^{n}\left[\mathscr{H}_{\rho, 1}\left(A_{\rho, g_{h}}\left(\partial_{x_{g}}\left(S_{\rho}\right)\right)\left(\partial_{x_{h}}\left(S_{\rho}\right)\right) S_{\rho}^{-2}\right)\right] \delta X \\
& \quad-2 \sum_{g, h=1}^{n} A_{\rho, g h}\left[\partial_{x_{g}}\left(S_{\rho}\right)\right] S_{\rho}^{-1}\left(\left[\partial_{x_{h}}, \mathscr{H}_{\rho}\right] \delta X\right) \\
&\left.+2 \sum_{g, h=1}^{n}\left[\mathscr{H}_{\rho, 1}\left(A_{\rho, g_{h}}\left(\partial_{x_{g}}\left(S_{\rho}\right)\right) S_{\rho}^{-1}\right)\right]\left(\partial_{x_{h}} \delta X\right)\right\} .
\end{align*}
$$

This is the equation for $\delta X$.
Denote the left-hand side (resp. right-hand side) of (9.14) by $\tilde{\mathscr{L}}_{\rho} \delta X$ (resp. $\left.\mathscr{R}_{\rho} \delta X\right)$. Define a differential operator $\mathscr{P}_{p}$ on $\bar{\Omega}_{T}$ by
(9.15.i) $\quad \mathscr{P}_{\rho}=\mathscr{H}_{\rho}-c_{0} \chi_{0}(\lambda)\left[\left.\left(\partial_{\lambda} U_{\rho}\right)\right|_{\lambda=0}\right] S_{\rho} \partial_{\lambda} \quad$ on $\quad N_{0}^{-} \times[0, T]$.

$$
\begin{equation*}
\mathscr{P}_{\rho}=\partial_{t} \quad \text { on } \quad\left(\bar{\Omega}_{0}-N_{0}^{-}\right) \times[0, T] . \tag{9.15.ii}
\end{equation*}
$$

From (9.8) and (9.9), we observe that in a neighborhood of $\Gamma_{T}$ :
(i) $\partial_{\lambda}\left(S_{\rho} f\right)=S_{\rho}\left(\partial_{\lambda} f\right)$.
(ii) $\partial_{\lambda} \mathscr{H}_{\rho}^{-1}=\mathscr{H}_{\rho}^{-1} \partial_{\lambda}$.

Then the problem $\left(1_{\partial V}\right)-\left(4_{\partial V}\right)$ is transformed into the problem:
$\left(1_{\delta X}\right)$
( $2_{\text {}}^{\text {o }}$ )
( $3_{\delta X}$ )
( $4_{\delta X}$ )
$\tilde{\mathscr{L}}_{\rho} \delta X=\mathscr{R}_{\rho} \delta X$ in $\Omega_{T}$.

$$
\left.\delta X\right|_{t=0}=0
$$

$$
\delta X=0 \quad \text { on } \quad J_{T}
$$

$$
\mathscr{P}_{\rho} \delta X=-\left(\partial_{\lambda} U_{\rho}\right) S_{\rho} \mathscr{H}_{\rho}^{-1} \delta G \quad \text { on } \quad \Gamma_{T}
$$

This is the problem for $\delta X$.
(C) By the assumption (A.3) in Theorem and the maximum principle of the heat equation,

$$
\begin{equation*}
-c_{0}\left(\partial_{\lambda} U_{\rho}\right) S_{\rho}=-\hat{s}_{\theta}\left[\left(\partial_{\lambda} u_{\rho}\right) \circ e_{\rho}\right] S_{\rho} \geqq 0 \quad \text { on } \quad \Gamma_{T} . \tag{9.16}
\end{equation*}
$$

Now we can easily observe that:
(i) On $\Gamma_{T}$, if $i \neq 11$ (resp. $i=11$ ), then $\mathscr{H}_{\rho, 1}$ is an $\left(r_{0}-1, r_{0}-5+\right.$ $4 i, B$ )-fine (resp. ( $r_{0}-1, r_{0}+38, B$ )-fine) operator. See (9.6) and the assumption (T).
(ii) On $\bar{\Omega}_{T}$, if $i \neq 11$ (resp. $i=11$ ), then $\mathscr{H}_{\rho, 1}$ is an $\left(r_{0}-1, r_{0}-5+\right.$ $4 i, B)$ - $\lambda$-fine (resp. ( $r_{0}-1, r_{0}+38, B$ )- $\lambda$-fine) operator. See (9.9).
(iii) The operator $\tilde{\mathscr{L}}_{\rho}$ is $\left(r_{0}-3, r_{0}-7+4 i, \sigma_{0}, B^{\prime}\right)$-\#-parabolic. See (9.14).
(iv) The composition $\mathscr{H}_{\rho} \mathscr{R}_{\rho}$ is a spatial differential operator of the second order whose coefficients belong to $C_{\sharp}^{\left(r_{0}-9+4 i\right)}\left(\bar{\Omega}_{T}\right)$. See (9.14).
(v) If $i \neq 11$ (resp. $i=11$ ), then the homogeneous part of the first order of $\mathscr{P}_{\rho}$ is an ( $r_{0}-1, r_{0}-5+4 i, B$ )- $\lambda$-fine (resp. $\left(r_{0}-1, r_{0}+38, B\right)$ -$\lambda$-fine) operator (see (9.15) and (9.16)).
Here $B^{\prime}$ is a constant and $B$ is a constant bounded with $\left|a_{0}\right|_{r_{0}}$ and $\left|b_{0}\right|_{\left(r_{0}\right)}$.
(D) Note that we can apply Lemma 8.B. 9 to $\left(1_{\delta X}\right)-\left(4_{j X}\right)$ if the righthand side of $\left(1_{\partial X}\right)$ is replaced by 0 . To realize this idea, we introduce three linear operators as follows. The compatibility of these definitions is verified in (E).

Define a linear operator $\mathscr{V}_{\rho}: C_{\sharp}^{\left(r_{0}-7+4 i\right)}\left(\bar{\Omega}_{T}\right) \rightarrow C_{\sharp}^{\left(r_{0}-7+4 i\right)}\left(\bar{\Omega}_{T}\right)$ by:

$$
\begin{gather*}
\tilde{\mathscr{L}}_{\rho}\left(\mathscr{Y}_{\rho} \delta X\right)=\mathscr{R}_{\rho} \delta X \quad \text { in } \quad \Omega_{T}  \tag{2}\\
\left.\left(\mathscr{Y}_{\rho} \delta X\right)\right|_{t=0}=0 .  \tag{3}\\
\mathscr{Y}_{\rho} \delta X=0 \quad \text { on } \quad J_{T} .  \tag{z}\\
\partial_{\lambda}\left(\mathscr{U}_{\rho} \delta X\right)=0 \quad \text { on } \quad \Gamma_{T} . \tag{3}
\end{gather*}
$$

Here $\delta X$ belongs to $C_{\neq}^{\left(r_{0}-7+4 i\right)}\left(\bar{\Omega}_{T}\right)$. Define a linear operator $\mathscr{R}_{\rho}: C_{\#}^{\left(r_{0}-7+4 i\right)}\left(\bar{\Omega}_{T}\right) \rightarrow$ $C_{\ddagger}^{\left(r_{0}-7+4 i\right)}\left(\bar{\Omega}_{T}\right)$ by:

$$
\begin{equation*}
\tilde{\mathscr{L}}_{\rho}\left(\mathscr{F}_{\rho} \delta X\right)=0 \quad \text { in } \quad \Omega_{T} . \tag{x}
\end{equation*}
$$

$$
\begin{equation*}
\left.\left(\mathscr{Z}_{\rho} \delta X\right)\right|_{t=0}=0 \tag{X}
\end{equation*}
$$

Here $\delta X$ belongs to $C_{\ddagger}^{\left(r_{0}-7+4 i\right)}\left(\bar{\Omega}_{T}\right)$. Define a linear operator $\tilde{\mathscr{\Sigma}}_{\rho}: C_{\sharp}^{\left(r_{0}-7+4 i\right)}\left(\Gamma_{T}\right) \rightarrow$ $C_{\#}^{\left(r_{0}-7+4 i\right)}\left(\bar{\Omega}_{T}\right)$ by:

$$
\begin{equation*}
\tilde{\mathscr{L}}_{\rho}\left(\tilde{\mathscr{F}}_{\rho} \delta G\right)=0 \quad \text { in } \quad \Omega_{T} . \tag{x}
\end{equation*}
$$

$$
\begin{equation*}
\tilde{\tilde{z}}_{\rho} \delta G=0 \quad \text { on } \quad J_{T} . \tag{x}
\end{equation*}
$$

(4 ${ }^{\text {x }}$ )

$$
\mathscr{P}_{\rho}\left(\tilde{\mathscr{L}}_{\rho} \delta G\right)=-\left(\partial_{\lambda} U_{\rho}\right) S_{\rho} \mathscr{H}_{\rho}^{-1} \delta G \quad \text { on } \quad \Gamma_{T} .
$$

Here $\delta G$ belongs to $C_{\#}^{\left(r_{0}-7+4 i\right)}\left(\Gamma_{T}\right)$.
Clearly, for $\delta X$ in $C_{\neq}^{\left(r_{0}-7+4 i\right)}\left(\bar{\Omega}_{T}\right)$, the problem $\left(1_{i X}\right)-\left(4_{i X}\right)$ is equivalent to the equation

$$
\begin{equation*}
\delta X=\mathscr{V}_{\rho} \delta X+\mathscr{K}_{\rho} \delta X+\tilde{\mathscr{K}}_{\rho} \delta G . \tag{9.17}
\end{equation*}
$$

(E) Let $T$ be so small that:
(i) For $\mathscr{H}_{\rho, 1}$, we can use Lemmas 8.B.7 and 8.B.8.
(ii) For $\tilde{\mathscr{L}}_{\rho}$ and the homogeneous part of the first order of $\mathscr{P}_{\rho}$, we can use Lemma 8.B.9.

First we consider (1\%)-(4 ). By Lemma 8.B.8,

$$
\begin{equation*}
\left|\mathscr{R}_{\rho} \delta X\right|_{\left(r_{0}-9+4 i\right), t} \leqq C\left[|\delta X|_{\left(r_{0}-7+4 i\right), t}+\left(1+|\rho|_{\left(r_{0}-4+4 i\right), T}\right)|\delta X|_{\left(r_{0}-3\right), t}\right] \tag{9.18}
\end{equation*}
$$

for $t \in[0, T]$. Here $C$ is a constant bounded with $\left|a_{0}\right|_{r_{0}+39}$ and $\left|b_{0}\right|_{\left(r_{0}+39\right)}$. By Lemma 3.B. 3 and (9.18), the operator $\mathscr{Y}_{\rho}$ is well-defined and satisfies

$$
\begin{equation*}
\left|\mathscr{Y}_{\rho} \delta X\right|_{\left(r_{0}-7+4 i\right), t} \leqq C\left[|\delta X|_{\left(r_{0}-7+4 i\right), t}+\left(1+|\rho|_{\left(r_{0}-4+4 i\right), T}\right)|\delta X|_{\left(r_{0}-3\right), t}\right] \tag{9.19}
\end{equation*}
$$

for $t \in[0, T]$. The operator $\mathscr{Y}_{\rho}$ is $\left(r_{0}-3, r_{0}-7+4 i, C, 1+|\rho|_{\left(r_{0}-4+4 i\right), T}\right)-$ balanced. Here $C$ is a constant bounded with $\left|a_{0}\right|_{r_{0}+39}$ and $\left|b_{0}\right|_{\left(r_{0}+38\right)}$. From $\left(1_{3}\right)-(4 \%)$ we obtain:

$$
\begin{gather*}
\tilde{\mathscr{L}}_{\rho}\left(\mathscr{H}_{\rho} \mathscr{Y}_{\rho} \delta X\right)=\left[\tilde{\mathscr{L}}_{\rho}, \mathscr{\mathscr { L }}_{\rho}\right] \mathscr{Y}_{\rho} \delta X \text { in } \Omega_{T} .  \tag{1}\\
\left.\left(\mathscr{H}_{\rho} \mathscr{U}_{\rho} \delta X\right)\right|_{t=0}=0 . \tag{2}
\end{gather*}
$$

$$
\begin{equation*}
\mathscr{H}_{\rho} \mathscr{V}_{\rho} \delta X=0 \quad \text { on } \quad J_{T} . \tag{3}
\end{equation*}
$$

Hence, by Lemma 3.B. 3 and (9.19),

$$
\begin{equation*}
\left|\mathscr{K}_{\rho} \mathscr{Y}_{\rho} \delta X\right|_{\left(r_{0}-7+4 i\right), t} \leqq C\left[|\delta X|_{\left(r_{0}-7+4 i\right), t}+\left(1+|\rho|_{\left(r_{0}-4+4 i\right), T}\right)|\delta X|_{\left(r_{0}-3\right), t}\right] \tag{9.20}
\end{equation*}
$$

for $t \in[0, T]$. Here $C$ is a constant bounded with $\left|a_{0}\right|_{r_{0}+39}$ and $\left|b_{0}\right|_{\left(r_{0}+39\right)}$. By Lemma 8.B. 8 and (9.20),

$$
\begin{align*}
\left|\mathscr{Y}_{\rho} \delta X\right|_{\left(r_{0}-7+4 i\right), t} \leqq & C\left[\int_{0}^{t}|\delta X|_{\left(r_{0}-7+4 i\right), \tau} d \tau\right.  \tag{9.21}\\
& \left.+\left(1+|\rho|_{\left(r_{0}-4+4 i\right), T}\right) \int_{0}^{t}|\delta X|_{\left(r_{0}-3\right), \tau} d \tau\right]
\end{align*}
$$

for $t \in[0, T]$. The operator $\mathscr{Y}_{\rho}$ is $\left(r_{0}-3, r_{0}-7+4 i, C, 1+|\rho|_{\left(r_{0}-4+4 i\right), T}\right)-$ integral-balanced. Here $C$ is a constant bounded with $\left|a_{0}\right|_{r_{0}+38}$ and $\left|b_{0}\right|_{\left(r_{0}+38\right)}$.

Secondly we consider $\left(1_{\mathscr{x}}\right)-\left(4_{\check{x}}\right)$ and $\left(1_{\tilde{x}}\right)-(4 \tilde{x})$. By Lemma 8.B. 9 and (9.20), the operator $\mathscr{\mathscr { F }}_{\rho}$ is well-defined and satisfies

$$
\begin{align*}
\left|\mathscr{Z}_{\rho} \delta X\right|_{\left(r_{0}-7+4 i\right), t} \leqq & C\left[\int_{0}^{t}|\delta X|_{\left(r_{0}-7+4 i\right), \tau} d \tau\right.  \tag{9.22}\\
& \left.+\left(1+|\rho|_{\left(r_{0}-4+4 i\right), T}\right) \int_{0}^{t}|\delta X|_{\left(r_{0}-3\right), \tau} d \tau\right]
\end{align*}
$$

for $t \in[0, T]$. Here $C$ is a constant bounded with $\left|a_{0}\right|_{r_{0}+38}$ and $\left|b_{0}\right|_{\left(r_{0}+38\right)}$. By Lemmas 8.B.7 and 8.B.9, the operator $\tilde{\mathscr{F}}_{\rho}$ is well-defined and satisfies

$$
\begin{align*}
\left|\tilde{\mathscr{\Sigma}}_{\rho} \delta G\right|_{\left(r_{0}-7+4 i\right), t} \leqq & C\left[\int_{0}^{t}|\delta G|_{\left(r_{0}-7+4 i\right), \tau} d \tau\right.  \tag{9.23}\\
& \left.+\left(1+|\rho|_{\left(r_{0}-4+4 i\right), T}\right) \int_{0}^{t}|\delta G|_{\left(r_{0}-3\right), \tau} d \tau\right]
\end{align*}
$$

for $t \in[0, T]$. Here $C$ is a constant bounded with $\left|a_{0}\right|_{r_{0}+49}$ and $\left|b_{0}\right|_{\left(r_{0}+39\right)}$.
Hence, by Lemma 8.A.4, (9.21), (9.22) and (9.23), we obtain the solution $\delta X$ of (9.17) which satisfies

$$
\begin{equation*}
|\delta X|_{\left(r_{0}-7+4 i\right), T} \leqq C\left[|\delta G|_{\left(r_{0}-7+4 i\right), T}+\left(1+|\rho|_{\left(r_{0}-4+4 i\right), T}\right)|\delta G|_{\left(r_{0}-3\right), T}\right] \tag{9.24}
\end{equation*}
$$

Here $C$ is a constant bounded with $\left|a_{0}\right|_{r_{0}+39}$ and $\left|b_{0}\right|_{\left(r_{0}+39\right)}$.
(F) Define a linear operator $\mathscr{J}(\rho): C_{\#}^{\left(r_{0}-7+4 i\right)}\left(\Gamma_{T}\right) \rightarrow C_{\#}^{\left(r_{0}-8+4 i\right)}\left(\Gamma_{T}\right)$ by

$$
\begin{equation*}
\mathscr{J}(\rho) \delta G=-c_{0} \partial_{\lambda} \delta X+\mathscr{H}_{\rho}^{-1} \delta G \tag{9.25}
\end{equation*}
$$

where $\delta X$ is the solution of (9.17) in view of (9.7) and (9.10). By (9.24), (9.25) and (4.5.ii), we easily observe that:
(i) $D \mathscr{F}(\rho) \mathscr{J}(\rho) \delta G=\delta G$ when $i \geqq 2$.
(ii) $|\mathcal{J}(\rho) \delta G|_{\left(r_{0}-8+4 i\right)} \leqq C\left[|\delta G|_{\left(r_{0}-7+4 i\right)}+\left(1+|\rho|_{\left(r_{0}-4+4 i\right)}\right)|\delta G|_{\left(r_{0}-3\right)}\right]$.
(iii) $|\mathscr{J}(\rho) \mathscr{F}(\rho)|_{\left(r_{0}-8+4 i\right)} \leqq C\left(1+|\rho|_{\left(r_{0}-4+4 i\right)}\right)$.

Here $C$ is a constant bounded with $\left|a_{0}\right|_{r_{0}+38}$ and $\left|b_{0}\right|_{\left(r_{0}+39\right)}$. This proves (III). On the condition $i \geqq 2$ in (i), remember that we assume $\delta \rho$ belongs to $C_{\neq}^{\left(r_{0}\right)}\left(\Gamma_{T}\right)$ in (A).

Remark. The "right inverse" $\mathscr{J}(\rho)$ thus constructed is also the "left inverse" of $D \mathscr{F}(\rho)$.
(G) We have proved the conditions (I)-(III) in Nash's implicit function theorem for our setting. On the other hand, it is easy to see that

$$
|\mathscr{F}(\rho)|_{\left(r_{0}-2-\left(\varepsilon_{0} / 2\right)\right), T} \leqq C|\mathscr{F}(\rho)|_{\left(r_{0}-2\right), T} T^{\varepsilon_{0} / 4},
$$

where $C$ is a constant. Hence, by Nash's implicit function theorem, Theorem' is proved.
10. Elimination of the technical assumption. Finally, in this section, we show how our proof is modified when the assumption ( $T$ ) in Section 2 is replaced by the assumption (A.4) in Theorem in Section 1. Let $a_{0}$ be a Hestenes-Whitney extension of $a_{0}$ to $\boldsymbol{R}^{n}$. Define a function
$u_{0}$ on $\boldsymbol{R}^{n} \times\left[0, T_{0}\right)$ by solving

$$
\left(\partial_{t}-\Delta\right) u_{0}=0
$$

and $\left.u_{0}\right|_{t=0}=a_{0}$. Observe that the definitions of $e_{\rho}, S_{\rho}$ and $\mathscr{L}_{\rho}$ in Sections 2 and 4 can be extended for $\rho \in C^{\left(r_{0}\right)}\left(\Gamma_{T}\right)$ with small $|\rho|_{0}$. Recall that $a_{0} \in C^{r_{0}+43}\left(\bar{\Omega}_{0}\right)$. By the theory of nonlinear flrst-order equations (see e.g., Courant and Hilbert [8, Chapter 2]) and the weighted Hölder estimate for characteristic coordinate transforms (refer to Lemma 8.B.5), taking $T$ sufficiently small, we get $\rho_{0} \in C^{\left(r_{0}+40\right)}\left(\Gamma_{T}\right)$ with small $\left|\rho_{0}\right|_{0}$ such that $\left.\rho_{0}\right|_{t=0}=0$ and

$$
\partial_{t} \rho_{0}+c_{0}\left[\partial_{\lambda}\left(u_{0} \circ e_{\rho_{0}}\right)\right] S_{\rho_{0}}=0 \quad \text { on } \quad \Gamma_{T},
$$

because the right-hand side of the characteristic system (see [8, p. 97]) consists of $C^{\left(r_{0}+40\right)}$ functions. Then our problem is to seek $\rho \in V_{T}$ and $U \in C^{\left(r_{0}\right)}\left(\bar{\Omega}_{T}\right)$ which satisfy:

$$
\begin{equation*}
U=b_{0} \quad \text { on } \quad J_{T} \tag{3}
\end{equation*}
$$

$$
\begin{gather*}
\mathscr{L}_{\rho_{0}+\rho} U=0 \quad \text { in } \quad \Omega_{T} .  \tag{1}\\
\left.U\right|_{t=0}=a_{0} . \tag{2}
\end{gather*}
$$

$$
\begin{equation*}
U=0 \quad \text { on } \quad \Gamma_{T} . \tag{4}
\end{equation*}
$$

Here $\rho_{0}+\rho$ corresponds to $\rho$ in Theorem in Section 1. Taking $T$ and $\delta_{0}$ sufficiently small, we can suppose that $\mathscr{L}_{\rho_{0}+\rho}$ is parabolic and $S_{\rho_{0}+\rho}$ has a positive lower bound for $\rho \in V_{T}$. The assumption (A.4) implies that there exists a solution $U \in C^{\left(r_{0}+39\right)}\left(\bar{\Omega}_{T}\right)$ of the above parabolic system (1)-(4) for $\rho \in V_{T} \cap C_{\sharp}^{\left(r_{0}+39\right)}\left(\Gamma_{T}\right)$. Now we can define a mapping $\mathscr{F}: V_{T} \rightarrow C_{\sharp}^{\left(r_{0}-2\right)}\left(\Gamma_{T}\right)$ by

$$
\mathscr{F}(\rho)=\partial_{t}\left(\rho_{0}+\rho\right)+c_{0}\left(\partial_{\lambda} U\right) S_{\rho_{0}+\rho},
$$

where $U$ is the solution of the system (1)-(4). In this setting, replacing the heat operator by $\mathscr{L}_{\rho_{0}}$ and modifying the technical definitions and lemmas slightly, the rest of the proof can be developed in the same manner as in the case of the technical assumption ( T ).

Note added in proof. The author recently learned from the referee and a few other persons that there exist two announcements "On the classical solutions of the Stefan multidimensional problem, by B. M. Budak and M. Z. Moskal, Dokl. Akad. Nauk SSSR Tom, 184 (1969), 1263-1266= Soviet Math. Dokl., Vol. 10 (1969), 219-223" and "On classical solvability of the multidimensional Stefan problem, by A. M. Meirmanov, Dokl. Akad. Nauk SSSR Tom, 249 (1979)=Soviet Math. Dokl., Vol. 20 (1979), 1426-1429."

The former claims the local (-in-time) existence of the classical solutions for the initial value problem of the two-phase multidimensional Stefan problem, provided that the free boundary can be parametrized by flat space variables. However, no proofs are given. Furter it seems that the detailed proof has not been published yet. The latter claims that, for the initial value problem of the two-phase multidimensional Stefan problem, the local (-in-time) existence of the classical solutions can be proved by parabolic regularization method, provided that the initial normal gradient of the thermal distribution on the interface between ice and water has a positive lower bound. Also no detailed proofs are given.

## References

[1] L. A. Caffarelli, The regularity of free boundaries in higher dimensions, Acta Math. 139 (1977), 155-184.
[2] L. A. Caffarelli, Some aspects of the one-phase Stefan problem, Indiana Univ. Math. J. 27 (1978), 73-77.
[3] L. A. Caffarelli, The smoothness of the free boundary in a filtration problem, Arch. Rational Mech. Anal., to appear.
[4] L. A. Caffarelli and A. Friedman, Continuity of the temperature in the Stefan problem, Indiana Univ. Math. J. 28 (1979), 53-70.
[5] J. R. Cannon and C. D. Hill, Existence, uniqueness, stability, and monotone dependence in a Stefan problem for the heat equation, J. Math. Mech. 17 (1967), 1-19.
[6] J. R. Cannon, C. D. Hill and M. Primicerio, The one-phase Stefan problem for the heat equation with boundary temperature specification, Arch. Rational Mech. Anal. 39 (1970), 270-274.
[7] J. R. Cannon and M. Primicerio, Remarks on the one-phase Stefan problem for the heat equation with the flux prescribed on the fixed boundary, J. Math. Anal. Appl. 35 (1971), 361-373.
[8] R. Courant and D. Hilbert, Methods of mathematical physics, Volume II, Interscience Pub., New York, 1962.
[9] G. Duvaut, Résolution d'un problème de Stefan (Fusion d'un bloc de glace à zéro degré), C. R. Acad. Sci. Paris, Ser. A, 276 (1973), 1461-1463.
[10] A. Friedman, Free boundary problems for parabolic equations. I, Melting of solids, J. Math. Mech. 8 (1959), 499-517.
[11] A. Friedman, The Stefan problem in several space variables, Trans. Amer. Math. Soc. 133 (1968), 51-87.
[12] A. Friedman and D. Kinderlehrer, A one phase Stefan problem, Indiana Univ. Math. J. 24 (1975), 1005-1035.
[13] V. Guillemin, The Radon transform on Zoll surfaces, Adv. in Math. 22 (1976), 85-119.
[14] M.L. Gromov, Smoothing and inversion of differential operators, Math. USSR-Sb. 17 (1972), 381-435.
[15] R. S. Hamilton, Deformation of complex structures on manifolds with boundary. I: The stable case, J. Differential Geometry 12 (1977), 1-45.
[16] L. Hörmander, The boundary problem of physical geodesy, Arch. Rational Mech. Anal. 62 (1976), 1-52.
[17] S. L. Kamenomostskaja, On Stefan's problem (in Russian), Naučn Dokl. Vysš. Školy 1
(1958), No. 1, 60-62. See also Mat. Sb. 53 (1961), 489-514.
[18] D. Kinderlehrer and L. Nirenberg, Regularity in free boundary problems, Ann. Scuola Norm. Sup. Pisa Cl. Sci., Ser. 4, 4 (1977), 373-391.
[19] D. Kinderlefrer and L. Nirenberg, The smoothness of the free boundary in the one phase Stefan problem, Comm. Pure Appl. Math. 31 (1978), 257-282.
[20] S. Klainerman, Global existence for nonlinear wave equations, Comm. Pure Appl. Math. 33 (1980), 43-101.
[21] O. A. Ladyzenskaja, V.A. Solonnikov and N. N. Uralceva, Linear and quasilinear equations of parabolic type, Translations of mathematical monographs, Vol. 23, Amer. Math. Soc., Providence, R. I., 1968.
[22] J. Moser, A new technique for the construction of solutions of nonlinear differential equations, Proc. Nat. Acad. Sci. U.S.A. 47 (1961), 1824-1831.
[23] J. NASH, The imbedding problem for Riemannian manifolds, Ann. of Math. 63 (1956), 20-63.
[24] J. NASH, Analyticity of the solutions of implicit function problem with analytic data, Ann. of Math. 84 (1966), 345-355.
[25] T. Nogi, The Stefan problem (in Japanese), Sugaku 30 (1978), 1-11.
[26] O. A. Oleinik, A method of solution of the general Stefan problem, Dokl. Akad. Nauk SSSR 135 (1960), 1054-1057. =Soviet Math. Dokl. 1 (1960), 1350-1354.
[27] L. I. Rubinstein, On the solution of Stefan's problem (in Russian), Izv. Akad. Nauk SSSR Ser. Geograf. Geofiz. 11 (1947), 37-54.
[28] L.I. Rubinstein, On the determination of the position of the boundary which separates two phases in the one-dimensional problem of Stefan (in Russian), Dokl. Akad. Nauk SSSR 58 (1947), 217-220.
[29] L.I. Rubinstein, On some nonlinear problems arising from the Fourier equation (in Russian), Dissertation, Moscow State University, Moscow, 1957.
[30] L. I. Rubinstein, The Stefan problem, Translations of mathematical monographs, Vol. 27, Amer. Math. Soc., Providence, R. I., 1972.
[31] D. G. Schaeffer, A stability theorem for the obstacle problem, Adv. in Math. 16 (1975), 34-47.
[32] D. G. Schaeffer, The capacitor problem, Indiana Univ. Math. J. 24 (1975), 1143-1167.
[33] D. G. Schaeffer, Supersonic flow past a nearly straight wedge, Duke Math. J. 43 (1976), 637-670.
[34] J. T. Schwartz, Nonlinear functional analysis, Gordon and Breach, New York-LondonParis, 1969.
[35] F. Sergeraert, Une généralisation du théorème des fonctions implicites de Nash, C.R. Acad. Sci. Paris, Ser. A, 270 (1970), 861-863.
[36] J. Stefan, Uber einige Probleme der Theorie der Wärmeleitung, S.-B. Wien. Akad. Mat. Natur. 98 (1889), 473-484.
[37] J. Stefan, Uber die Diffusion von Sauren und Basen gegen einander, S.-B. Wien. Akad. Mat. Natur. 98 (1889), 616-634.
[38] J. Stefan, Úber die Theorie der Eisbildung insbesondere über die Eisbildung in Polarmeere, S.-B. Wien. Akad. Mat. Natur. 98 (1889), 965-983.
[39] J. Stefan, Uber die Verdampfung und die Auflösung als Vorgänge der Diffusion, S.-B. Wien. Akad. Mat. Natur. 98 (1889), 1418-1442.
[40] M. Yamaguti and T. Nogi, The Stefan problem (in Japanese), Sangyo Tosho, Tokyo, 1977.
[41] E. Zehnder, Generalized implicit function theorem with applications to some small divisor problems, I, Comm. Pure Appl. Math. 28 (1975), 91-140.

Mathematical Institute<br>Hokkaido University<br>Sapporo, 060 Japan

