ASYMPTOTIC BEHAVIOR OF STRONGLY DAMPED NONLINEAR BEAM EQUATIONS

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ABSTRACT. For $\xi \in [1/2, 3/4]$, the existence of the global attractor for the evolutionary equation corresponding to the following strongly damped nonlinear beam equation $(1 + \beta A^{1/2})u_{tt} + \delta A^{1/2}u_t + \alpha Au + g(||u||_{\xi-1/4}^2)A^{1/2}u = f, t > 0$, has been studied in $\mathcal{D}_H(A^{\xi}) \times \mathcal{D}_H(A^{\xi-1/4})$. Such an equation is related to a nonlinear beam equation as well as Timoshenko's equation.

The main difficulty of our work comes from the terms $\beta A^{1/2} u_{tt}$ and $g(||u||_{\xi-1/4}^2)A^{1/2}u$, representing the rotational inertia of the beam and the tension within the beam due to its extensibility, respectively. We overcome the difficulty of introducing the solution, bounded absorbing set, and κ -contracting property by carefully using the fractional power theory and suitable time-uniform a priori estimates.

1. Introduction. The asymptotic behavior and global dynamics for nonlinear beam equations of the following form

(1.1)
$$u_{tt} + \delta u_t + \alpha u_{xxxx} + \left(a_0 + a_1 \int_0^1 |u_x|^2 dx\right) u_{xx} = f, \quad t > 0,$$

where $a_0 \in \mathbb{R}$ and $\delta, \alpha, a_1 > 0$, have been carefully studied by many researchers, cf., [4, 8, 13, 14]. The asymptotic behavior for different types of nonlinear beam equations has been extensively studied in the literature [2, 3, 5, 6, 7, 9, 12].

In [1, 11, 15] and the references therein, Timoshenko's equation

(1.2)
$$u_{tt} + \beta u_{xxtt} + \alpha u_{xxxx} + \left(\int_0^1 |u_x|^2 dx\right) u_{xx} = f, \quad t > 0,$$

has been studied, where $\alpha, \beta > 0$.

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Here, we introduce a nonlinear equation related to equations (1.1) and (1.2). Indeed, consider the Hilbert space $H = L^2(0, 1)$ whose inner product and norm are denoted by $\langle \cdot, \cdot \rangle_0$ and $\|\cdot\|_0$, respectively. For a given $\xi \in [1/2, 3/4]$, we shall study the existence of the global attractor for an abstract evolutionary equation corresponding to the following nonlinear strongly damped beam equation:

$$(1.3) \quad (1+\beta A^{1/2})u_{tt} + \delta A^{1/2}u_t + \alpha Au + g(\|u\|_{\xi-1/4}^2)A^{1/2}u = f, \quad t > 0,$$

with initial conditions

(1.4)
$$u(0) = u_0 \text{ and } u_t(0) = u_1,$$

where the closed linear operator A, the functions f and g, and the parameters α, β , and δ , are given by assumptions (A1)–(A4) below.

Taking into account assumption (A4) and the theory of fractional powers discussed below, it is clear that, for $\beta = 0$, $\xi = 1/2$, damping δu_t and $g(u) = a_0 + a_1 ||A^{1/4}u||_0^2$, equation (1.3) becomes identical to equation (1.1), and for $\xi = 1/2$, $\delta = 0$ and $g(u) = ||A^{1/4}u||_0^2$, equation (1.3) becomes identical to equation (1.2). From the definition of the closed linear operator A given by assumption (A4) below, it is clear that the hinged boundary conditions are considered in this work.

The nonlinear beam equation (1.3) can be considered as a mathematical model for small transversal vibrations of an extensible beam, where u(x,t) is the transverse deflection of the beam, δ and β are the parameters related to the structural strong damping and the rotational inertia of the beam, respectively. The nonlinear term in the equation is the tension within the beam due to its extensibility, and the function fstands for an external input.

Within this work, we assume the following:

(A1) $\alpha, \delta, \beta > 0$ and $\xi \in [1/2, 3/4]$.

(A2) $f \in H = L^2(0, 1)$, where the inner product and norm in H are denoted by $\langle \cdot, \cdot \rangle_0$ and $\|\cdot\|_0$, respectively.

(A3) $g: [0, \infty) \to \mathbb{R}$ is locally Lipschitz continuous and there exist $r, K_0, K_1 \ge 0$, with $K_1 < \pi^2 \alpha$, such that

(1.5)
$$\int_{r}^{s} g(y) \, dy \ge -(K_0 + K_1 s), \quad g(s_1) \ge g(s_2),$$
for all $s \ge r, \quad s_1 \ge s_2 \ge r.$

Choosing

$$K_2 = K_2(r) = \frac{3}{2}r \max_{y \in [0,r]} |g(y)|$$

it is clear that, for $0 \le s < r$,

$$\int_{0}^{s} g(y) \, dy \ge -\int_{0}^{s} |g(y)| \, dy \ge -r \max_{y \in [0,r]} |g(y)| \ge -K_2,$$

$$\frac{1}{2} \int_{0}^{s} g(y) \, dy - sg(s) \le \frac{1}{2} \int_{0}^{s} |g(y)| \, dy + s|g(s)| \le \frac{3}{2}r \max_{y \in [0,r]} |g(y)| = K_2,$$

and, for $s \ge r$,

$$\int_0^s g(y) \, dy = \int_0^r g(y) \, dy + \int_r^s g(y) \, dy \ge -r \max_{y \in [0,r]} |g(y)| + \int_r^s g(y) \, dy \ge -(K_2 + K_0 + K_1 s),$$

$$\begin{split} \frac{1}{2} \int_0^s &g(y) \, dy - sg(s) = \frac{1}{2} \int_0^r &g(y) \, dy + \frac{1}{2} \int_r^s &g(y) \, dy - (s-r)g(s) - rg(s) \\ &= \frac{1}{2} \int_0^r &g(y) \, dy + \int_r^s \left(\frac{1}{2}g(y) - g(s)\right) dy - rg(s) \\ &\leq \frac{1}{2} r \max_{y \in [0,r]} |g(y)| - \frac{1}{2} \int_r^s g(y) \, dy - rg(r) \\ &\leq \frac{1}{2} (K_0 + K_1 s) + K_2. \end{split}$$

Hence, for $s \ge 0$, we have

(1.6)
$$\frac{1}{2} \int_0^s g(y) \, dy - sg(s) \le \frac{1}{2} (K_0 + K_1 s) + K_2,$$
$$\int_0^s g(y) \, dy \ge -(K_2 + K_0 + K_1 s).$$

(A4) $A : \mathcal{D}_H(A) \subset H \to H$ is the closed linear operator defined as follows

 $A\varphi = \varphi_{xxxx}, \quad \text{for all } \varphi \in \mathcal{D}_H(A),$

where the domain of A in H is given by

$$\mathcal{D}_H(A) = \{ \varphi \in H^4(0,1) : \varphi(0) = \varphi_{xx}(0) = \varphi(1) = \varphi_{xx}(1) = 0 \}.$$

Following [13], it can be shown that A is a densely defined, selfadjoint, positive definite operator with compact resolvent A^{-1} . The spectrum of A consists of the eigenvalues $\lambda_n = (n\pi)^4$, $n \in \mathbb{N}$, with eigenvectors $e_n = 2^{1/2} \sin(n\pi x)$, $n \in \mathbb{N}$.

Along the lines of the fractional power theory [10, 13], for $\xi \geq 0$, $V^{\xi} = \mathcal{D}_H(A^{\xi})$ is a Hilbert space where $V^0 = H$, whose inner product and norm are given by

$$\langle \varphi, \psi \rangle_{\xi} = \langle A^{\xi} \varphi, A^{\xi} \psi \rangle_{0}, \quad \|\varphi\|_{\xi} = \langle A^{\xi} \varphi, A^{\xi} \varphi \rangle_{0}^{1/2} = \|A^{\xi} \varphi\|_{0}$$

for all φ , $\psi \in V^{\xi}$. If $\xi > \eta \ge 0$, then V^{ξ} is compactly embedded in V^{η} and

(1.7)
$$\|\varphi\|_{\eta} \le \pi^{4(\eta-\xi)} \|\varphi\|_{\xi} \le \|\varphi\|_{\xi}, \text{ for all } \varphi \in V^{\xi}.$$

For $\xi, \eta \ge 0$, we have

(1.8)
$$A^{\xi}A^{\eta} = A^{\eta}A^{\xi} = A^{\xi+\eta} \quad \text{on } V^{\xi+\eta},$$
$$\langle A^{\xi}\varphi, \psi \rangle_0 = \langle \varphi, A^{\xi}\psi \rangle_0, \quad \text{for all } \varphi, \psi \in V^{\xi}.$$

Considering the linear operator $A : \mathcal{D}_{V^{\xi}}(A) \subset V^{\xi} \to V^{\xi}, \xi \geq 0$, where the domain of A in V^{ξ} is given by $\mathcal{D}_{V^{\xi}}(A) = V^{\xi+1}$, one can show that A is a sectorial operator and -A generates an analytic semigroup on V^{ξ} . Moreover, we have

$$A^{1/2}\varphi = -\varphi_{xx}, \text{ for all } \varphi \in V^{1/2},$$
$$\|\varphi\|_{\xi} = \left\|\frac{\partial^{4\xi}\varphi}{\partial x^{4\xi}}\right\|_{0}, \text{ for all } \varphi \in V^{\xi}, \ 4\xi = 0, 1, 2, \dots.$$

For $\xi, \eta \geq 0$, we shall consider the Hilbert space $E_{\xi,\eta} = V^{\xi} \times V^{\eta}$, whose inner product and norm are given by:

$$\begin{split} \langle \Psi_0, \Psi_1 \rangle_{E_{\xi,\eta}} &= \langle \varphi_0, \varphi_1 \rangle_{\xi} + \langle \psi_0, \psi_1 \rangle_{\eta}, \\ \| \Psi_0 \|_{E_{\xi,\eta}} &= \langle \Psi_0, \Psi_0 \rangle_{E_{\xi,\eta}}^{1/2} \end{split}$$

for all $\Psi_i = (\varphi_i, \psi_i)^T \in E_{\xi,\eta}, i = 0, 1.$

The objective of this research is to prove the existence of the global attractor for an abstract dynamical system corresponding to the initial value problem (1.3)–(1.4) in $E_{\xi,\xi-1/4}$. The existence of a global attractor for the solution semiflow of a given dynamical system is very

important because it is the depository of all the longtime dynamics of such a system. The background material in regard to the theory of global attractors can be found in [8, 13].

Fixing $\xi \in [1/2, 3/4]$, this work is organized as follows. In Section 2, the initial value problem (1.3)–(1.4) will be represented as an abstract evolutionary equation in $E_{\xi,\xi-1/4}$ and the standard semigroup theory will be used to prove the existence of a local solution for this abstract equation. In Section 3, we shall prove that the solution exists globally and the semiflow generated by these solutions is bounded dissipative. In Section 4, we shall show that the solution semiflow is κ -contracting. Using these results, along the lines of the classical theory of global attractors, it follows that the solution semiflow possesses a global attractor.

The difficulty of studying the existence of the global attractor for the evolutionary equation related to the initial value problem (1.3)-(1.4) comes from the terms:

$$\beta A^{1/2} u_{tt}$$
 and $g(\|u\|_{\ell=1/4}^2) A^{1/2} u_{t}$

representing the rotational inertia of the beam and the tension within the beam due to its extensibility, respectively. Carefully using the fractional power theory and suitable time-uniform a priori estimates, we overcome the difficulty of presenting the solution semiflow, bounded absorbing set, and κ -contracting property for the solution semiflow.

2. Formulation and well posedness. Letting $v = u_t$, we can reformulate (1.3)–(1.4) as the following first-order evolution equation

$$B\Phi_t = C\Phi + F_{\xi}(\Phi), \quad t > 0$$

$$\Phi(0) = \Phi_0 = (u_0, u_1)^T,$$

where $\Phi = (u, v)^T$,

$$B = \begin{pmatrix} I & 0 \\ 0 & I + \beta A^{1/2} \end{pmatrix}, \qquad C = \begin{pmatrix} 0 & I \\ -\alpha A & -\delta A^{1/2} \end{pmatrix},$$
$$F_{\xi}(\Phi) = (0, J_{\xi}(u))^{T} \quad \text{and} \quad J_{\xi}(u) = f - g(\|u\|_{\xi - 1/4}^{2})A^{1/2}u$$

Using (1.7) and the assumption that g is locally Lipschitz continuous, it follows that, for K > 0, $\varphi, \psi \in V^{\xi}$ with $\xi \ge 1/4$,

 $\|\varphi\|_{\xi-1/4} \le K, \quad \|\psi\|_{\xi-1/4} \le K,$

there exist constants $k_i = k_i(K) > 0$, i = 0, 1, such that

(2.1)
$$|g(\|\varphi\|_{\xi-1/4}^2) - g(\|\psi\|_{\xi-1/4}^2)| \le k_0 |\|\varphi\|_{\xi-1/4}^2 - \|\psi\|_{\xi-1/4}^2 | \le k_1 \|\varphi - \psi\|_{\xi}.$$

Lemma 2.1. Given $\xi \in [1/2, 3/4]$, the nonlinear operator

$$B^{-1}F_{\xi}(\Phi) = (0, (I + \beta A^{1/2})^{-1}J_{\xi}(u))^{T}, \text{ for all } \Phi = (u, v)^{T}$$

is locally Lipschitz continuous in $E_{\xi,\xi-1/4}$ and $E_{3/4,1/2}$.

Proof. Given $\xi \in [1/2, 3/4]$, let $k_2 > 0$ and $\Phi_i = (u_i, v_i)^T \in E_{\xi, \xi - 1/4}$ be such that

$$\|\Phi_i\|_{E_{\xi,\xi-1/4}} \le k_2, \quad i = 0, 1.$$

Then, from equation (1.7) and the fact that $(I + \beta A^{1/2})^{-1}$ is a bounded operator from H into $V^{1/2}$, we find that there exist a constant $k_3 > 0$ such that

$$\begin{split} \|B^{-1}F_{\xi}(\Phi_{0}) - B^{-1}F_{\xi}(\Phi_{1})\|_{E_{\xi,\xi-1/4}} \\ &= \|(I + \beta A^{1/2})^{-1}[g(\|u_{0}\|_{\xi-1/4}^{2})A^{1/2}u_{0} - g(\|u_{1}\|_{\xi-1/4}^{2})A^{1/2}u_{1}]\|_{\xi-1/4} \\ &\leq \|(I + \beta A^{1/2})^{-1}[g(\|u_{0}\|_{\xi-1/4}^{2})A^{1/2}u_{0} - g(\|u_{1}\|_{\xi-1/4}^{2})A^{1/2}u_{1}]\|_{1/2} \\ &\leq k_{3}\|g(\|u_{0}\|_{\xi-1/4}^{2})A^{1/2}u_{0} - g(\|u_{1}\|_{\xi-1/4}^{2})A^{1/2}u_{1}\|_{0} \\ &\leq k_{3}(|g(\|u_{0}\|_{\xi-1/4}^{2})|\|u_{0} - u_{1}\|_{1/2} + |g(\|u_{0}\|_{\xi-1/4}^{2}) \\ &- g(\|u_{1}\|_{\xi-1/4}^{2})|\|u_{1}\|_{1/2}). \end{split}$$

From equations (1.7) and (2.1), along with the fact that g is locally Lipschitz continuous, it follows that there exists a constant $k_4 = k_4(k_2) > 0$ such that

$$\begin{aligned} \|B^{-1}F_{\xi}(\Phi_{0}) - B^{-1}F_{\xi}(\Phi_{1})\|_{E_{\xi,\xi-1/4}} \\ &\leq k_{4}\|u_{0} - u_{1}\|_{\xi} \leq k_{4}\|\Phi_{0} - \Phi_{1}\|_{E_{\xi,\xi-1/4}}. \end{aligned}$$

Thus, $B^{-1}F_{\xi}$ is locally Lipschitz continuous in $E_{\xi,\xi-1/4}$. Similarly, one can see that $B^{-1}F_{\xi}$ is locally Lipschitz continuous in $E_{3/4,1/2}$.

For $\xi \in [1/2, 3/4]$, consider the linear operator

$$B^{-1}C: \mathcal{D}_{E_{\xi,\xi-1/4}}(B^{-1}C) \subset E_{\xi,\xi-1/4} \to E_{\xi,\xi-1/4},$$

given by

$$B^{-1}C = \begin{pmatrix} 0 & I \\ -\alpha(I + \beta A^{1/2})^{-1}A & -\delta(I + \beta A^{1/2})^{-1}A^{1/2} \end{pmatrix},$$

whose domain in $E_{\xi,\xi-1/4}$ is given by $\mathcal{D}_{E_{\xi,\xi-1/4}}(B^{-1}C) = E_{1,\xi}$. It can be shown that $-B^{-1}C$ is a sectorial operator and $B^{-1}C$ is the infinitesimal generator of an analytic contraction semigroup in $E_{\xi,\xi-1/4}$. In this case, using Lemma 2.1, and the standard semigroup theory [13], we get the following result.

Lemma 2.2. For $\xi \in [1/2, 3/4]$, consider the initial value problem

(2.2)
$$\Phi_t = B^{-1}C\Phi + B^{-1}F_{\xi}(\Phi), \quad t > 0$$
$$\Phi(0) = \Phi_0 = (u_0, u_1)^T.$$

(a) Given $\Phi_0 = (u_0, u_1)^T \in E_{\xi, \xi - 1/4}$, equation (2.2) has a unique maximal solution

$$\Phi(t;\xi,\Phi_0) = (u(t), u_t(t))^T, t \in [0,\tau) \quad and \quad \tau = \tau(\Phi_0) > 0,$$

in $E_{\xi,\xi-1/4}$ satisfying (2.2) for $t \in [0,\tau)$, such that

(2.3)
$$(u(\cdot), u_t(\cdot))^T = C([0, \tau), E_{\xi, \xi - 1/4}) \cap C^1((0, \tau), E_{\xi, \xi - 1/4}) \cap C((0, \tau), E_{1,\xi}),$$

where $E_{1,\xi}$ is $\mathcal{D}_{E_{\xi,\xi-1/4}}(B^{-1}C)$. Moreover, if $\tau < \infty$, then

(2.4)
$$\lim_{t \to \tau^{-}} \| (u(t), u_t(t))^T \|_{E_{\xi, \xi^{-1/4}}} = \infty$$

(b) Given $\Phi_0 = (u_0, u_1)^T \in E_{3/4, 1/2}$, equation (2.2) has a unique maximal solution

$$\Phi(t;\xi,\Phi_0) = (u(t), u_t(t))^T, t \in [0,\tau) \quad and \quad \tau = \tau(\Phi_0) > 0,$$

in $E_{3/4,1/2}$ satisfying (2.2) for $t \in [0, \tau)$ such that

(2.5)
$$(u(\cdot), u_t(\cdot))^T = C([0, \tau), E_{3/4, 1/2}) \cap C^1((0, \tau), E_{3/4, 1/2}) \cap C((0, \tau), E_{1, 3/4}),$$

where $E_{1,3/4}$ is $\mathcal{D}_{E_{3/4,1/2}}(B^{-1}C)$. Moreover, if $\tau < \infty$, then

(2.6)
$$\lim_{t \to \tau^{-}} \| (u(t), u_t(t))^T \|_{E_{3/4, 1/2}} = \infty.$$

3. Global solutions and dissipativity. Before introducing Lemmas 3.2 and 4.1, we need the following remark to explain the reason for choosing the initial data Φ_0 in $E_{3/4,1/2}$ while studying in $E_{\xi,\xi-1/4}$ the global existence of solutions bounded dissipativity, and uniform κ -contracting property for the solution semiflow associated with the initial value problem (2.2).

Remark 3.1. In both lemmas, we need (1.3) to be satisfied as a first step to complete the proof. Reading (2.3) and (2.5) carefully, one can see that, for $\Phi_0 \in E_{\xi,\xi-1/4}$, the solution given by (2.3) satisfies (2.2), but it does not satisfy (1.3). For $\Phi_0 \in E_{3/4,1/2}$, the solution given by (2.5) satisfies (2.2), and it consequently satisfies (1.3).

The key idea for proving the following lemma is applying suitable time-uniform a priori estimates.

Lemma 3.2. In $E_{\xi,\xi-1/4}$, $\xi \in [1/2, 3/4]$, for each $\Phi_0 \in E_{\xi,\xi-1/4}$, the unique solution of the initial value problem (2.2) given by equation (2.3), exists globally and the associated semiflow $S_{\xi}(t)$ generated by these solutions is bounded dissipative.

Proof. Let

$$\xi \in [1/2, 3/4], \quad \Phi_0 \in E_{3/4, 1/2}$$

and

$$S_{\xi}(t)\Phi_0 = \Phi(t) = (u(t), \ u_t(t))^T$$

be the solution of equation (2.2) in $E_{3/4,1/2}$, given by equation (2.5). Recalling Remark 3.1, it is clear that equation (1.3) is satisfied for $t \in (0, \tau)$. For $\varepsilon > 0$, if we consider the inner product of equation (1.3) with $A^{2\xi-1}(u_t + \varepsilon u)$ in H, taking into account equation (1.8), we obtain

(3.1)
$$\frac{d}{dt}\Theta_{\xi}(t) + \Pi_{\xi}(t) = 0, \quad \text{for all } t \in (0,\tau),$$

where

(3.2)
$$\Theta_{\xi}(t) = \frac{1}{2} \|u_t\|_{\xi-1/2}^2 + \frac{\beta}{2} \|u_t\|_{\xi-1/4}^2 + \frac{\alpha}{2} \|u\|_{\xi}^2 + \frac{\varepsilon\delta}{2} \|u\|_{\xi-1/4}^2$$
$$- \langle f, A^{2\xi-1}u \rangle_0 + \frac{1}{2} \int_0^{\|u\|_{\xi-1/4}^2} g(y) \, dy$$
$$+ \varepsilon \langle u_t, u \rangle_{\xi-1/2} + \varepsilon\beta \langle u_t, u \rangle_{\xi-1/4}$$

and

(3.3)
$$\Pi_{\xi}(t) = (\delta - \varepsilon \beta) \|u_t\|_{\xi - 1/4}^2 - \varepsilon \|u_t\|_{\xi - 1/2}^2 + \varepsilon \alpha \|u\|_{\xi}^2 + \varepsilon g(\|u\|_{\xi - 1/4}^2) \|u\|_{\xi - 1/4}^2 - \varepsilon \langle f, A^{2\xi - 1}u \rangle_0.$$

From Young's inequality and (1.7), we find

$$\begin{aligned} (3.4) \\ |\langle u_t, u \rangle_{\xi-1/2}| &\leq \frac{1}{2} \|u_t\|_{\xi-1/2}^2 + \frac{1}{2} \|u\|_{\xi-1/2}^2 \leq \frac{1}{2} \|u_t\|_{\xi-1/4}^2 + \frac{1}{2} \|u\|_{\xi}^2 \\ |\langle u_t, u \rangle_{\xi-1/4}| &\leq \frac{1}{2} \|u_t\|_{\xi-1/4}^2 + \frac{1}{2} \|u\|_{\xi}^2, \\ |\langle f, A^{2\xi-1}u \rangle_0| &\leq \|f\|_0 \|u\|_{2\xi-1} \leq \frac{\|f\|_0^2}{2\varepsilon} + \frac{\varepsilon}{2} \|u\|_{\xi}^2. \end{aligned}$$

Using (3.2) and (3.4), we get

$$\begin{split} \Theta_{\xi}(0) &\leq \frac{1}{2} \|u_{t}(0)\|_{\xi-1/2}^{2} + \frac{\beta}{2} \|u_{t}(0)\|_{\xi-1/4}^{2} + \frac{\alpha}{2} \|u(0)\|_{\xi}^{2} \\ &+ \frac{\varepsilon}{2} \delta \|u(0)\|_{\xi-1/4}^{2} + \frac{\|f\|_{0}^{2}}{2\varepsilon} + \frac{\varepsilon}{2} \|u(0)\|_{\xi}^{2} + \frac{1}{2} \int_{0}^{\|u(0)\|_{\xi-1/4}^{2}} \\ &\cdot g(y) \, dy + \frac{\varepsilon}{2} (1+\beta) (\|u_{t}(0)\|_{\xi-1/4}^{2} + \|u(0)\|_{\xi}^{2}). \end{split}$$

From (1.7), it follows that

(3.5)
$$\Theta_{\xi}(0) \leq \frac{1}{2} (1 + \beta + \varepsilon(1 + \beta)) \|u_t(0)\|_{\xi - 1/4}^2 + \frac{\|f\|_0^2}{2\varepsilon} \\ + \frac{1}{2} (\alpha + \varepsilon(2 + \delta + \beta) + \max_{y \in [0, \|u(0)\|_{\xi}^2]} |g(y)|) \|u(0)\|_{\xi}^2.$$

Again, recalling (3.2) and (3.4), we have

$$\frac{1}{2} \|u_t(t)\|_{\xi-1/2}^2 + \frac{\beta}{2} \|u_t(t)\|_{\xi-1/4}^2 + \frac{\alpha}{2} \|u(t)\|_{\xi}^2 + \frac{\varepsilon\delta}{2} \|u(t)\|_{\xi-1/4}^2 \\
- \frac{\|f\|_0^2}{2\varepsilon} - \frac{\varepsilon}{2} \|u(t)\|_{\xi}^2 + \frac{1}{2} \int_0^{\|u(t)\|_{\xi-1/4}^2} g(y) \, dy \\
- \frac{1}{2} \varepsilon (1+\beta) (\|u_t(t)\|_{\xi-1/4}^2 + \|u(t)\|_{\xi}^2) \le \Theta_{\xi}(t),$$

$$\begin{split} \frac{1}{2} (\beta - \varepsilon (1 + \beta)) \| u_t(t) \|_{\xi - 1/4}^2 + \frac{1}{2} (\alpha - \varepsilon (2 + \beta)) \| u(t) \|_{\xi}^2 + \frac{1}{2} \int_0^{\| u(t) \|_{\xi - 1/4}^2} g(y) dy \\ & \leq \Theta_{\xi}(t) + \frac{\| f \|_0^2}{2\varepsilon}, \end{split}$$

From (1.6) and (1.7) we obtain

(3.6)

$$\frac{1}{2}(\beta - \varepsilon(1+\beta)) \|u_t(t)\|_{\xi-1/4}^2 + \frac{1}{2}(\alpha - \pi^{-2}K_1 - \varepsilon(2+\beta)) \|u(t)\|_{\xi}^2$$

$$\leq \Theta_{\xi}(t) + \frac{\|f\|_0^2}{2\varepsilon} + \frac{K_0}{2} + \frac{K_2}{2}.$$

Taking into account (3.2) and (3.3), we find

$$\begin{split} \varepsilon \Theta_{\xi}(t) - \Pi_{\xi}(t) &= \frac{3\varepsilon}{2} \|u_{t}\|_{\xi-1/2}^{2} + \left(\frac{3\varepsilon\beta}{2} - \delta\right) \|u_{t}\|_{\xi-1/4}^{2} - \frac{\varepsilon\alpha}{2} \|u\|_{\xi}^{2} \\ &+ \frac{\varepsilon^{2}\delta}{2} \|u\|_{\xi-1/4}^{2} + \frac{\varepsilon}{2} \int_{0}^{\|u\|_{\xi-1/4}^{2}} g(y) \, dy \\ &- \varepsilon g(\|u\|_{\xi-1/4}^{2}) \|u\|_{\xi-1/4}^{2} \\ &+ \varepsilon^{2} \langle u_{t}, u \rangle_{\xi-1/2} + \varepsilon^{2}\beta \langle u_{t}, u \rangle_{\xi-1/4}, \end{split}$$

using (3.4) it follows that

$$\begin{split} \varepsilon \Theta_{\xi}(t) - \Pi_{\xi}(t) &= \frac{3\varepsilon}{2} \|u_t\|_{\xi-1/2}^2 + \left(\frac{3\varepsilon\beta}{2} - \delta\right) \|u_t\|_{\xi-1/4}^2 - \frac{\varepsilon\alpha}{2} \|u\|_{\xi}^2 \\ &+ \frac{\varepsilon^2 \delta}{2} \|u\|_{\xi-1/4}^2 + \frac{\varepsilon}{2} \int_0^{\|u\|_{\xi-1/4}^2} g(y) \, dy \\ &- \varepsilon g(\|u\|_{\xi-1/4}^2) \|u\|_{\xi-1/4}^2 \\ &+ \frac{1}{2} \varepsilon^2 (1+\beta) (\|u_t\|_{\xi-1/4}^2 + \|u\|_{\xi}^2). \end{split}$$

From (1.6) and (1.7), we get

(3.7)

$$\varepsilon \Theta_{\xi}(t) - \Pi_{\xi}(t) \leq \left(\frac{\varepsilon}{2}(1+\beta)(\varepsilon+3) - \delta\right) \|u_t(t)\|_{\xi-1/4}^2 + \frac{\varepsilon}{2} [\varepsilon(\delta+1+\beta) + \pi^{-2}K_1 - \alpha] \|u(t)\|_{\xi}^2 + \varepsilon \left(\frac{K_0}{2} + K_2\right).$$

In (A3), we assumed that $K_1 < \pi^2 \alpha$. For sufficiently small $\varepsilon > 0$, inequality (3.7) becomes

$$\varepsilon \Theta_{\xi}(t) - \Pi_{\xi}(t) \le \varepsilon \left(\frac{K_0}{2} + K_2\right).$$

Using (3.1), we get

$$\frac{d}{dt}\Theta_{\xi}(t) + \varepsilon\Theta_{\xi}(t) \le \varepsilon \bigg(\frac{K_0}{2} + K_2\bigg),$$

that is,

$$\frac{d}{dt}(e^{\varepsilon t}\Theta_{\xi}(t)) \leq \varepsilon \left(\frac{K_0}{2} + K_2\right)e^{\varepsilon t}.$$

Integrating the above inequality from 0 into t, where $t \in (0, \tau)$, we find

(3.8)
$$\Theta_{\xi}(t) \le \Theta_{\xi}(0)e^{-\varepsilon t} + \frac{K_0}{2} + K_2, \quad \text{for all } t \in [0, \tau).$$

Again, since $K_1 < \pi^2 \alpha$, we can choose $\varepsilon > 0$ sufficiently small such that

$$\min\{\beta - \varepsilon(1+\beta), \alpha - \pi^{-2}K_1 - \varepsilon(2+\beta)\} > 0.$$

Choosing

$$k_{7} = \frac{1}{2} \min\{\beta - \varepsilon(1+\beta), \ \alpha - \pi^{-2}K_{1} - \varepsilon(2+\beta)\},\$$

$$k_{8} = \frac{1}{2} \max\left\{\begin{array}{c}1 + \beta + \varepsilon(1+\beta), \\ \alpha + \varepsilon(2+\delta+\beta) + \max_{y \in [0, \|u(0)\|_{\xi}^{2}]}|g(y)|\end{array}\right\},\$$

and using (3.5) and (3.6), it follows that

$$\Theta_{\xi}(0) \le k_8(\|u_t(0)\|_{\xi-1/4}^2 + \|u(0)\|_{\xi}^2) + \frac{\|f\|_0^2}{2\varepsilon}$$
$$= k_8 \|\Phi_0\|_{E_{\xi,\xi-1/4}}^2 + \frac{\|f\|_0^2}{2\varepsilon}$$

and

$$k_7 \|S_{\xi}(t)\Phi_0\|_{E_{\xi,\xi^{-1/4}}}^2 = k_7 (\|u_t(t)\|_{\xi^{-1/4}}^2 + \|u(t)\|_{\xi}^2)$$

$$\leq \Theta_{\xi}(t) + \frac{\|f\|_0^2}{2\varepsilon} + \frac{K_0}{2} + \frac{K_2}{2}.$$

In such a case, recalling (3.8) and the previous two inequalities, it follows that for $t \in [0, \tau)$,

(3.9)
$$||S_{\xi}(t)\Phi_{0}||^{2}_{E_{\xi,\xi-1/4}} \leq \frac{1}{k_{7}} \bigg[\bigg(k_{8} ||\Phi_{0}||^{2}_{E_{\xi,\xi-1/4}} + \frac{||f||^{2}_{0}}{2\varepsilon} \bigg) e^{-\varepsilon t} + k_{9} \bigg],$$

where

$$k_9 = \frac{\|f\|_0^2}{2\varepsilon} + K_0 + \frac{3K_2}{2}$$

Using (2.6) and (3.9) it is clear that for $\Phi_0 \in E_{3/4,1/2}$, the solution exists globally in $E_{3/4,1/2}$, and since $E_{3/4,1/2} \subset E_{\xi,\xi-1/4}$, such a solution exists globally in $E_{\xi,\xi-1/4}$. Using the fact that $E_{3/4,1/2}$ is dense in $E_{\xi,\xi-1/4}$, we find that for $\xi \in [1/2, 3/4]$ and $\Phi_0 \in E_{\xi,\xi-1/4}$, the unique solution (2.3) of (2.2) exists globally in $E_{\xi,\xi-1/4}$.

For R > 0 and $\Phi_0 \in E_{\xi,\xi-1/4}$ such that $\|\Phi_0\|_{E_{\xi,\xi-1/4}}^2 \leq R$, using (3.9), it follows that

(3.10)
$$||S_{\xi}(t)\Phi_0||^2_{E_{\xi,\xi-1/4}} \le 2k_9/k_7$$
, for all $t \ge t_1$,

where

$$t_1 = t_1(R) = \frac{1}{\varepsilon} \ln[(k_8 R + ||f||_0^2 / 2\varepsilon)/k_9].$$

Thus, the solution semiflow S(t) generated by the solutions of (2.2) is bounded dissipative in $E_{\xi,\xi-1/4}$.

4. Global attractor. Here we present the main result of this work, where we prove the existence of a global attractor for the solution semiflow generated by the solutions of (2.2) in $E_{\xi,\xi-1/4}$, for $\xi \in [1/2, 3/4]$. Indeed, we show that the solution semiflow is κ -contracting, then Lemma 3.2 and the classical theory of global attractors provide the result.

Lemma 4.1. In $E_{\xi,\xi-1/4}$, $\xi \in [1/2, 3/4]$, the semiflow $S_{\xi}(t)$ generated by the solutions of (2.2) is uniformly κ -contracting.

Proof. For $\xi \in [1/2, 3/4]$, let D be a bounded set in

$$E_{\xi,\xi-1/4}, \quad \Phi_{0i} \in E_{3/4,1/2} \cap D, \quad i = 1, 2,$$

and let

$$S_{\xi}(t)\Phi_{0i} = \Phi_i(t) = (u_i(t), u_{it}(t))^T, \quad i = 1, 2$$

be the solutions of (2.2) given by (2.5). Let

$$Y(t) = \Phi_1(t) - \Phi_2(t) = (y(t), y_t(t))^T,$$

where $y(t) = u_1(t) - u_2(t)$. Recalling Remark 3.1 and equation (1.3), we obtain

(4.1)
$$(1 + \beta A^{1/2})y_{tt} + \delta A^{1/2}y_t + \alpha Ay$$

= $g(||u_2||_{\xi - 1/4}^2)A^{1/2}u_2 - g(||u_1||_{\xi - 1/4}^2)A^{1/2}u_1.$

For t > 0, let us consider the inner product of the last equation with $A^{2\xi-1}(y_t + \varepsilon y)$ in $H, \varepsilon > 0$, we get

(4.2)
$$\frac{d}{dt}\Xi_{\xi}(t) + \Gamma_{\xi}(t) = G_{\xi}(t;u_1,u_2).$$

where

$$\Xi_{\xi}(t) = \frac{1}{2} \|y_t\|_{\xi-1/2}^2 + \frac{\beta}{2} \|y_t\|_{\xi-1/4}^2 + \frac{\alpha}{2} \|y\|_{\xi}^2 + \frac{\varepsilon\delta}{2} \|y\|_{\xi-1/4}^2 + \varepsilon\langle y_t, y\rangle_{\xi-1/2} + \varepsilon\beta\langle y_t, y\rangle_{\xi-1/4}$$
$$\Gamma_{\xi}(t) = (\delta - \varepsilon\beta) \|y_t\|_{\xi-1/4}^2 - \varepsilon \|y_t\|_{\xi-1/2}^2 + \varepsilon\alpha \|y\|_{\xi}^2,$$

and

$$G_{\xi}(t; u_1, u_2) = \langle g(\|u_2\|_{\xi^{-1/4}}^2) A^{1/2} u_2 - g(\|u_1\|_{\xi^{-1/4}}^2) A^{1/2} u_1, \ A^{2\xi^{-1}}(y_t + \varepsilon y) \rangle_0$$

Recalling (1.7) and (3.4) it follows that, for sufficiently small ε ,

(4.3)
$$\Xi_{\xi}(0) \leq \frac{1}{2}(1+\beta)(1+\varepsilon) \|y_t(0)\|_{\xi-1/4}^2 + \frac{1}{2}(\alpha+\varepsilon(1+\beta+\delta))\|y(0)\|_{\xi}^2$$

(4.4)
$$\Xi_{\xi}(t) \ge \frac{1}{2} (\beta - \varepsilon(1+\beta)) \|y_t(t)\|_{\xi-1/4}^2 + \frac{1}{2} (\alpha - \varepsilon(1+\beta)) \|y(t)\|_{\xi}^2,$$

(4.5)
$$\varepsilon \Xi_{\xi}(t) - \Gamma_{\xi}(t) \le 0.$$

In such a case, using (4.2) and (4.5), we obtain

$$\frac{d}{dt}\Xi_{\xi}(t) + \varepsilon \Xi_{\xi}(t) \le G_{\xi}(t; u_1, u_2), \quad \text{for all } t > 0.$$

That is,

(4.6)
$$\frac{d}{dt}(e^{\varepsilon t}\Xi_{\xi}(t)) \le e^{\varepsilon t}G_{\xi}(t;u_1,u_2), \quad \text{for all } t > 0.$$

From (3.9), it follows that there exists $k_{10} = k_{10}(D)$ such that

(4.7)
$$||S_{\xi}(t)\Phi_{0i}||_{E_{\xi,\xi-1/4}} = (||u_i(t)||_{\xi}^2 + ||u_{it}(t)||_{\xi-1/4}^2)^{1/2} \le k_{10},$$

for all $t \ge 0$, i = 1, 2. Taking into account (1.7), (2.1) and (4.7), there exists $k_{11} = k_{11}(D)$ such that for $t \ge 0$,

$$\begin{aligned} G_{\xi}(t;u_{1},u_{2}) &= \langle (g(\|u_{2}\|_{\xi-1/4}^{2}) - g(\|u_{1}\|_{\xi-1/4}^{2}))A^{1/2}u_{1} \\ &- g(\|u_{2}\|_{\xi-1/4}^{2})A^{1/2}y, \ A^{2\xi-1}(y_{t}+\varepsilon y)\rangle_{0}. \\ &= (g(\|u_{2}\|_{\xi-1/4}^{2}) - g(\|u_{1}\|_{\xi-1/4}^{2}))\langle u_{1},y_{t}+\varepsilon y\rangle_{\xi-1/4} \\ &- g(\|u_{2}\|_{\xi-1/4}^{2})\langle y,y_{t}+\varepsilon y\rangle_{\xi-1/4} \\ &\leq |g(\|u_{2}\|_{\xi-1/4}^{2}) - g(\|u_{1}\|_{\xi-1/4}^{2})|\|u_{1}\|_{\xi-1/4}\|y_{t}+\varepsilon y\|_{\xi-1/4} \\ &+ |g(\|u_{2}\|_{\xi-1/4}^{2})|\|y\|_{\xi-1/4}\|y_{t}+\varepsilon y\|_{\xi-1/4} \leq k_{11}\|y\|_{\xi-1/4}. \end{aligned}$$

Substituting the previous inequality into (4.6) and integrating the result from 0 into t > 0, we get

(4.8)
$$\Xi_{\xi}(t) \le e^{-\varepsilon t} \Xi_{\xi}(0) + \varepsilon^{-1} k_{11} \max_{l \in [0,t]} \|y(l)\|_{\xi - 1/4}.$$

We can choose $\varepsilon > 0$ sufficiently small such that

$$\min\{\beta - \varepsilon(1+\beta)\,\alpha - \varepsilon(1+\beta)\} > 0.$$

Let

$$k_{12} = \frac{1}{2} \max\{(1+\beta)(1+\varepsilon), \ \alpha + \varepsilon(1+\beta+\delta)\},\$$

$$k_{13} = \frac{1}{2} \min\{(\beta - \varepsilon(1+\beta)), \ (\alpha - \varepsilon(1+\beta))\} > 0.$$

In such a case, recalling (4.3) and (4.4), we find

(4.9)
$$\Xi_{\xi}(0) \le k_{12}(\|y_t(0)\|_{\xi-1/4}^2 + \|y(0)\|_{\xi}^2) = k_{12}\|\Phi_{01} - \Phi_{02}\|_{E_{\xi,\xi-1/4}}^2,$$

(4.10)

$$k_{13} \|S_{\xi}(t)\Phi_{01} - S_{\xi}(t)\Phi_{02}\|_{E_{\xi,\xi-1/4}}^2 = k_{13}(\|y_t(t)\|_{\xi-1/4}^2 + \|y(t)\|_{\xi}^2)^{1/2} \le \Xi_{\xi}(t).$$

Using (4.8)–(4.10), it follows that

$$\begin{split} \|S_{\xi}(t)\Phi_{01} - S_{\xi}(t)\Phi_{02}\|_{E_{\xi,\xi-1/4}}^{2} &\leq \frac{k_{12}}{k_{13}}e^{-\varepsilon t}\|\Phi_{01} - \Phi_{02}\|_{E_{\xi,\xi-1/4}}^{2} \\ &+ \frac{k_{11}}{\varepsilon k_{13}}\max_{l\in[0,t]}\|u_{1}(l) - u_{2}(l)\|_{\xi-1/4}, \quad \text{for all } t > 0, \end{split}$$

that is,

(4.11)

$$\begin{split} \|S_{\xi}(t)\Phi_{01} - S_{\xi}(t)\Phi_{02}\|_{E_{\xi,\xi-1/4}} &\leq \left(\frac{k_{12}}{k_{13}}\right)^{1/2} e^{-\varepsilon t/2} \|\Phi_{01} - \Phi_{02}\|_{E_{\xi,\xi-1/4}} \\ &+ \left(\frac{k_{11}}{\varepsilon k_{13}}\right)^{1/2} \max_{l \in [0,t]} \|u_1(l) - u_2(l)\|_{\xi-1/4}^{1/2}, \quad \text{for all } t > 0. \end{split}$$

Since $E_{3/4,1/2}$ is dense in $E_{3/4,1/2}$, one can see that the above inequality is satisfied for $\Phi_{0i} \in D$, i = 1, 2.

Now we shall apply the Ascoli-Arzela theorem in order to prove that for all t > 0,

$$\rho_t(\Phi_{01}, \Phi_{02}) = \left(\frac{k_{11}}{\varepsilon k_{13}}\right)^{1/2} \max_{l \in [0,t]} \|u_1(l) - u_2(l)\|_{\xi - 1/4}^{1/2}$$

is a precompact pseudometric in $E_{\xi,\xi-1/4}$. Indeed, given a bounded sequence $\{\Phi_0^n\}_{n=1}^{\infty}$ in $E_{\xi,\xi-1/4}$, let $S_{\xi}(t)\Phi_0^n = \Phi_n(t) = (u_n(t), u_{nt}(t))^T, n = 1, 2, \ldots$, be the solutions of (2.2). From (3.9), it is clear that there exists a constant $k_{14} = k_{14}(\{\Phi_0^n\}) > 0$ such that

(4.12)
$$\|S_{\xi}(t)\Phi_{0}^{n}\|_{E_{\xi,\xi-1/4}}^{2} = \|u_{n}(t)\|_{\xi}^{2} + \|u_{nt}(t)\|_{\xi-1/4}^{2} \le k_{14}^{2},$$

for all $t \ge 0$, $n = 1, 2, \ldots$. Using (4.12) and the fact that V^{ξ} is compactly embedded in $V^{\xi-1/4}$, we find that $\{u_n(t)\}_{n=1}^{\infty}$ is precompact in $V^{\xi-1/4}$, for all $t \ge 0$. From the Sobolev embedding theorem, we know that $V^{\xi-1/4}$ is compactly embedded in

 $C_B(0,1) = \{f : f \text{ is a continuous and bounded function on } (0,1)\},\$

that is, there exists a constant k_{15} such that

$$\sup_{\rho \in (0,1)} |\varphi(\rho)| \le k_{15} \|\varphi\|_{\xi-1/4}^2, \quad \text{for all } \varphi \in V^{\xi-1/4}.$$

From (4.12) and the last inequality, it follows that

$$\sup_{\rho \in (0,1)} |u_{nt}(\rho, t)| \le k_{16}, \quad \text{for all } t \ge 0, \ n = 1, 2, \dots,$$

where $k_{16} = k_{14}k_{15}$. In such a case, using the mean value theorem, we find

$$\|u_n(t_1) - u_n(t_2)\|_{\xi - 1/4} \le k_{16}|t_1 - t_2|,$$

for all $t_1, t_2 > 0, n = 1, 2, \dots,$

that is, the sequence $\{u_n(t)\}_{n=1}^{\infty}$, as $V^{\xi-1/4}$ -valued functions, is equicontinuous in t, for all $t \geq 0$. In such a case, using the Ascoli-Arzela theorem, there is a subsequence of $\{u_n(t)\}_{n=1}^{\infty}$ which is convergent in the Frechet space $L_{loc}^{\infty}[0,\infty;V^{\xi-1/4})$ and this is a Cauchy sequence with respect to $\rho_t(\Phi_{01}, \Phi_{02})$. Thus ρ_t is a precompact pseudometric on $E_{\xi,\xi-1/4}$.

Taking into account (4.11), since $\Psi(t) = (k_{12}/k_{13})^{1/2}e^{-\varepsilon t/2}$ is a nonincreasing function such that $\Psi(t) \to 0$ as $t \to \infty$ and we proved

that ρ_t is a precompact pseudometric on $E_{\xi,\xi-1/4}$. From [13, Lemma 22.5] we find that the semigroup $S_{\xi}(t)$ on $E_{\xi,\xi-1/4}$ is uniformly κ contracting.

Now we are ready to introduce the main result of this work.

Theorem 4.2. Given $\xi \in [1/2, 3/4]$, there exists a global attractor $\mathcal{A}_{\xi,\xi-1/4}$ for the solution semiflow $S_{\xi}(t)$ generated by the solutions of (2.2) in $E_{\xi,\xi-1/4}$ which attracts every bounded set in $E_{\xi,\xi-1/4}$.

Proof. From Lemmas 3.2 and 4.1, we know that for $\xi \in [1/2, 3/4]$, the semiflow $S_{\xi}(t)$ is bounded dissipative and κ -contracting in $E_{\xi,\xi-1/4}$. In such a case, following the classical theory of global attractors we get the result.

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