

SOLUTION OF THE SCHRÖDINGER EQUATION USING PADÉ APPROXIMANTS TO SUM ASYMPTOTIC SERIES

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In physics one is frequently interested in finding solutions to a Schrödinger Equation which contains a potential falling off as r^{-n} at large distances. An r^{-4} potential, for example, corresponds to the polarization of an atom by a distant charged particle. In the case of the interaction of an electron with an excited hydrogen atom there is an r^{-2} potential at large r . In these cases, the solutions for the long range behavior (Mathieu and Bessel functions, respectively) are known.

However, in many cases the solutions are not known, except in terms of a divergent series in $1/r$. A good example is the Coulomb three body problem where it has been only very recently that Stagat, Nuttall, and Hidalgo [1] have found such an asymptotic series for the three body wave function for the v th excited level of hydrogen in the presence of a distant electron for zero angular momentum.

For small distances the long ranged potential is usually modified by a different short ranged behavior. Our idea is to use Padé approximants to sum the asymptotic series in $1/r$ to a form which at large r might be used as the basis for a variational trial wave function. For small r we suggest using other wave functions and looking for a region of overlap where the two pieces of the wave function match.

Let us now consider radial Schrödinger equation:

$$(1) \quad \left(-\frac{d^2}{dr^2} + \frac{\ell(\ell+1)}{r^2} + V(r) - k^2 \right) \psi(r) = 0.$$

In most cases one knows (or pretends to know) the asymptotic form of the wave function, namely

$$(2) \quad \lim_{r \rightarrow \infty} \psi(r) = \frac{a_1}{r}.$$

This is the first term in a power series expansion in r^{-1} ,

$$(3) \quad \psi(r) = \sum_{n=1}^{\infty} \frac{a_n}{r^n}.$$

Substituting this expansion into the differential equation and equating like inverse powers of r , we may find relations among the a_n 's. For

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example $d^2/dr^2 - k^2$ will relate a_{n+2} to a_n . Using these recursion relations, and knowing the first coefficient a_1 , the entire series in r^{-n} may be generated.

Unfortunately this series is not useful. Generally one finds that the series is an asymptotic series with a zero radius of convergence. It is therefore a good patient for Padé approximants.

In order to test this idea of using Padé approximants to find solutions to the Schrödinger Equation, we tried the technique for the potential $V(r) = V_0/r^2$. The exact solution in this case is known, namely

$$(4) \quad \begin{aligned} \psi(r) &= r^{1/2} J_\nu(kr), \\ v^2 &= \ell(\ell + 1) + \frac{1}{4} + V_0. \end{aligned}$$

Since Gargantini and Henrici [2] have related $K_\nu(z)$ to a series of Stieltjes for $-\frac{1}{2} \leq v \leq \frac{1}{2}$, we choose to work with the $K_\nu(z)$ functions. For real v we have been able to show [3] that $K_\nu(z)$ is proportional to a series of Stieltjes plus a polynomial whose degree is equal to the integer part of v . Hence we were able to prove the convergence of the Padé approximants obtained from the asymptotic series. We also determined the rate of convergence.

If the potential is sufficiently attractive v can become imaginary. For complex v , $K_\nu(z)$ is not related to a series of Stieltjes nor any function for which a proof of convergence now exists. However, the Padé approximants converge at a rate similar to v real. For both real and imaginary v the number of Padé approximants required increases as $r \rightarrow 0$.

Our example for a $1/r^2$ potential suggests the value of using Padé approximants to sum an asymptotic series for $\psi(r)$ in general. However, the usefulness of this technique for computing physical observables such as cross sections involving integrals over $\psi(r)$ remains open to question.

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