# REDUCING THE ORDER OF THE LAGRANGEAN FOR A CLASSICAL FIELD IN CURVED SPACE-TIME

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We show how the Lagrangean L can be replaced by another,  $L^*$ , having the same extremals, but having only first order derivatives and being in fact a first degree polynomial in these derivatives.

1. Introduction. Whittaker [4] showed how to reduce the order and degree of a Lagrangean to 1 in the case of one-dimensional "space-time". As he points out, this leads instantly to Hamilton's canonical formalism. Rodrigues [3] showed how to do this without using coordinates in the configuration space. "Reducing the order" is not an adequate description of the construction since when the order is 1 (as it usually is) it still takes some work to make it of the first degree.

In [1] we treated the case of true (several dimensional) space-time. However, we took it to be  $\mathbb{R}^4$ , and we used coordinates in the field space.

Such a theorem is not usable when several coordinate systems must be used in space-time M. This is because the  $L^*$  doing the desired things is not unique.

Our construction of  $L^*$  here depends on the choice of an affine connection  $\Gamma$  and a volume element  $\omega$  in M. The result is not independent of the  $\Gamma$  and  $\omega$  chosen, but it is independent of coordinates.

2.  $\pi$ -manifolds. In this paper, a suitable order N of differentiability is assumed.

Let M be a manifold which we will call space-time. Let P be another manifold. We will call it a  $\pi$ -manifold if there is defined on it a regular map  $\pi: P \to M$ .

If P and Q are two  $\pi$ -manifolds, let F(P, Q) be the class of all maps f of open sets in P into Q for which

(2.1) 
$$\pi(f(p)) = \pi(p)$$

whenever f(p) is defined. Let p be a point of P and let  $f, g \in F(P, Q)$ . Say  $f \equiv g$  at p if f and g agree up to the Nth order at p. Let  $J^{N}(P, Q)$  be

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the set of equivalence classes. An element C of  $J^{N}(P, Q)$  may be represented by a pair (p, f) where C is an equivalence at p and  $f \in C$ . Define  $\pi(C) = \pi(p) = \pi(f(p))$ . Then  $J^{N}(P, Q)$  is also a  $\pi$ -space.

Let  $X^1, \ldots, x^m, y^1, \ldots, y^n$  be coordinates in P and Q respectively. Then  $(y^i)_{\lambda \cdots \nu}$  shall be the function defined in  $J^N(P, Q)$  be saying that

(2.2) 
$$(y^i)_{\lambda\cdots\nu}(C) = \frac{\partial^k (y^i \circ f)}{\partial x^{\lambda}\cdots \partial x^{\nu}}(p).$$

Here k must not exceed N.

A particularly useful kind of coordinates  $x^1, \ldots, x^m$  are coordinates exponential at a point p. For this one must select an affine connection  $\Gamma$ for P. Let  $T^1(P)_p$  be the tangent space to P at p. Select a *linear* coordinate system in  $T^1(P)p$  and transfer these to P using the exponential map defined by  $\Gamma$  (see [2].)

3. Lagrangeans. In classical mechanics, a Lagrangean is a function defined on  $\mathbf{R} \times T^1(Q)$ , the latter being the tangent bundle of configuration space. Now  $\mathbf{R} \times T^1(Q)$  is naturally isomorphic with  $J^1(M, M \times Q)$  where  $M = \mathbf{R}$ , and  $M \times Q$  has the projection on M. Since we want to use 2.1 we define Lagrangeans in this *milieu*.

Let *M* be a manifold of dimension *m*. Let *S* be a  $\pi$ -manifold. If coordinates  $t^1, \ldots, t^m$  are chosen for *M*, then  $t^1 \circ \pi, \ldots, t^m \circ \pi$  are (independent, by the regularity of  $\pi$ ) variables in *S*. We will abbreviate them to  $t^1, \ldots, t^m$ .

A Lagrangean (of order at most N) is an m-form  $\Lambda$  defined on  $J^{N}(M, S)$  such that in terms of coordinates  $t^{1}, \ldots, t^{m}$  in M,

(3.1)  $\Lambda = L dt^1 \wedge \cdots \wedge dt^m.$ 

Let the class of these Lagrangeans be called  $\mathscr{L}^{N}(S)$ .

3.2. THEOREM. Let M and S be as above. Let  $\Gamma$  be an affine connection for M. Let  $\omega$  be a volume element for M. Let K be the cotangent bundle of  $J^{N}(M, S)$ . Then  $\Gamma$ ,  $\omega$  define a mapping

(3.3) 
$$\mathscr{L}^{N}(S) \to \mathscr{L}^{1}(K), \quad \Lambda \to \Lambda^{*}.$$

This mapping is linear and 1:1

3.4.  $L^*$  is a polynomial of the first degree in the derivatives (necessarily only of the first order), and

3.5.  $\Lambda^*$  has the same extremals as  $\Lambda$ .

The precise (and natural) meaning of 3.5 is given in §5 below, together with the proof of 3.2.

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Property 3.4 is useful for the following construction.

Let  $x^1, \ldots, x^n$  (plus those  $t^1, \ldots, t^m$ ) be coordinates in S.  $\Lambda^* = L^* dt^1 \wedge \cdots \wedge dt^m$  and by 3.4,  $L^*$  is a sum of a term -H depending only on the  $x^i$  and  $t^{\lambda}$ , and a sum of terms like  $a(x^i)_{\lambda}$ . Thus  $\Lambda^*$  is a sum of  $-H dt^1 \wedge \cdots \wedge dt^m$  and of terms like  $a(x^i)_{\lambda} dt^1 \wedge \cdots \wedge dt^m$ . This latter term is congruent modulo  $dx^i - (x^i)_{\lambda} dt^{\lambda}$  to  $a dt^1 \wedge \cdots \wedge dx^i \wedge \cdots dt^m$ ,  $dx^i$  being in the  $\lambda$ th place. Thus  $\Lambda^*$  is congruent to an *m*-form involving only the  $x^i$  and  $t^{\lambda}$ , that is, an *m*-form on *S*.

# 4. A vector field $U_{\Gamma}$ on $J^{N}(M, S)$ .

4.1. THEOREM. Let  $\Gamma$  be an affine connection for M. Using  $\Gamma$  one can construct a vector field  $U_{\Gamma}$  on  $J^{N}(M, S)$ . Let  $A \in J^{N}(M, S)$  and let  $t^{1}, \ldots, t^{m}$  be exponential coordinates at  $\pi(A)$  in M. Use  $t^{1}, \ldots, t^{m}$  together with some further coordinates  $x^{1}, \ldots, x^{n}$  as coordinates in S. Then at A,

(4.2) 
$$U_{\Gamma} = (x^{i})_{\lambda} \frac{\partial}{\partial (x^{i})_{\lambda}} + 2(x^{i})_{\lambda \mu} \frac{\partial}{\partial (x^{i})_{\lambda \mu}} + \dots + N(x^{i})_{\lambda \dots \nu} \frac{\partial}{\partial (x^{i})_{\lambda \dots \nu}}.$$

A sum is intended in 4.2. For example, by

$$(x^i)_{\lambda\cdots\mu}\frac{\partial}{\partial(x^i)_{\lambda\cdots\mu}}$$

where there are k indices, we mean the sum over all sets of k indices such that  $1 \le \lambda \le \cdots \le \mu \le m$ .

We now prove 4.1. Let A = (a, f) be a point of  $J^{N}(M, S)$ . Express f in terms of  $t^{1}, \ldots, t^{m}$ :

$$f = \varphi(t^1, \ldots, t^m).$$

For sufficiently small real s and  $|t^{\lambda}|$ ,  $f_s$  can be defined by

$$f_s = \varphi(e^s t^1, \dots, e^s t^m).$$

Define the "moving" point  $A_s$  in  $J^N(M, S)$  as  $(a, f_s)$ . We define  $U_{\Gamma}$  at A to be the tangent to the curve  $s \to A_s$  for s = 0.

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We must show that its  $(x^i)_{\lambda\cdots\mu}$  component is correctly represented in 4.2. Now

$$(x^{i})_{\lambda\cdots\mu}(A_{s}) = \frac{\partial^{k}}{\partial t^{\lambda}\cdots\partial t^{\mu}} [x^{i}\circ f_{s}]\Big|_{t=0}$$
$$= e^{ks} \frac{\partial^{k}}{\partial t^{\lambda}\cdots\partial t^{\mu}} [x^{i}\circ f]\Big|_{t=0} = e^{ks} (x^{i})_{\lambda\cdots\mu}(A).$$

Then we take d/ds of this for s = 0, obtaining

 $k(x^i)_{\lambda\cdots\mu}$ 

as 4.2 asserts. The construction of  $A_s$  is clearly independent of the coordinates.

4.3. COROLLARY. Let K be the cotangent bundle  $T_1(J)$  of  $J = J^N(M, S)$ . Then  $\Gamma$  defines a real-valued function on K whose expression in terms of coordinates is

(4.4) 
$$p_i^{\lambda}\overline{(x^i)_{\lambda}} + 2p_i^{\lambda\mu}\overline{(x^i)_{\lambda\mu}} + \cdots + Np_i^{\lambda\cdots\nu}\overline{(x^i)_{\lambda\cdots\nu}}.$$

Here  $p_i^{\lambda \cdots \mu}$  is the "momentum" coordinate dual to the configuration coordinate  $(x^i)_{\lambda \cdots \mu}$  in J.

To prove 4.3, we need only show that 4.4 is independent of the coordinate representation. A point *B* is a pair (*A*, *g*) where *A* is a point of *J* as before, and *g* is a map from a neighborhood of *A* into **R**, where g(A) = 0. The value of  $p_i^{\lambda \cdots \mu}$  at *B* is

$$\frac{\partial g}{\partial (x^i)_{\lambda\cdots\mu}}\bigg|_A$$

The value of  $(x^i)_{\lambda \dots \mu}$  at B is just  $(x^i)_{\lambda \dots \mu}(A)$ . Here the latter is defined in J, and the former is the coordinate defined in  $T_1(J)$  as obtained from the latter and the projection  $\rho$ 

(4.5) 
$$T_1(J) = K$$
$$\downarrow$$
$$J = J^N(M, S)$$

Hence the value of 4.4 at B is  $U_{\Gamma}[g]$ , the result of applying the operator  $U_{\Gamma}$  to g, and is thus independent of coordinates.

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4.6. COROLLARY. Let J and K be as in 4.3. Let  $P = J^1(M, K)$ . Then  $\Gamma$  defines a real-valued function on P whose value at C = (c, h) is

(4.7) 
$$\overline{p_i^{\lambda}} Z_{\lambda}^i + \overline{p_i^{\lambda\mu}} Z_{\lambda\mu}^i + \cdots + \overline{p_i^{\lambda\cdots\nu}} Z_{\lambda\cdots\nu}^i$$

where  $Z^{i}_{\lambda \cdots \mu}$  is the sum of all

(4.8) 
$$(\overline{(x^i)_{\lambda\cdots\hat{\sigma}\cdots\mu}})_{\sigma}.$$

Several explanations are needed.

4.81. In forming the sum  $Z^i_{\lambda\cdots\mu}$  we select each index appearing in  $\lambda\cdots\mu$ , form 4.8, and add the results.  $\hat{\sigma}$  means delete ' $\sigma$ ' from the string  $\lambda\cdots\mu$ .

4.82. The  $p_i^{\overline{\lambda}\cdots\mu}$  is  $p_i^{\overline{\lambda}\cdots\mu}\circ\xi$  where  $\xi$  is the projection of  $P = J^1(M, K)$   $\rightarrow K$ . Hence, for an element C = (c, h) of P, wherein h maps a neighborhood of  $c \in M$  into K,

(4.83) 
$$p_i^{\lambda \cdots \mu}(C) = p_i^{\lambda \cdots \mu}(h(c))$$

which is

$$\left. \frac{\partial g}{\partial (x^i)_{\lambda\cdots\mu}} \right|_A$$

if  $h(c) = (A, g), A \in K$ .

4.84. The bar on the  $(x^i)_{\lambda \cdots \mu}$  means that (see 4.5)

$$(x^i)_{\lambda\cdots\mu}=(x^i)_{\lambda\cdots\mu}\circ\rho.$$

For any real variable z (and thus for  $(x^i)_{\lambda \dots}$ ) on K, 2.2 with N = 1 gives sense to  $z_{\sigma}$  for  $1 \le \sigma \le m$ .

To begin the proof of 4.6, we evaluate at C = (c, h) of P. By 2.2, the answer is

$$\frac{\partial}{\partial t^{\sigma}} \Big[ \overline{(x^i)_{\lambda \cdots \mu}} \circ h \Big] \Big|_c.$$

Let us write  $\partial_{\sigma}$  for  $\partial/\partial t^{\sigma}$ . So

$$Z^{i}_{\lambda\cdots\mu} = \sum \partial_{\sigma} \big[ (x^{i})_{\lambda\cdots\hat{\sigma}\cdots\mu} \circ \rho \circ h \big].$$

We need a somewhat more abstract version of 2.2. Let  $\pi_1(C) = p$ ,  $\pi_2(C) = h$  in 2.2. Replace  $x^{\lambda}$  by  $t^{\lambda}$  and  $y^i$  by  $x^i$ , since that is the notation for the present instance. Then 2.2 says

$$(x')_{\lambda\cdots\mu} = \left\{\partial_{\lambda}\cdots\partial_{\mu}(x'\circ\pi_2)\right\}\circ\pi_1.$$

Accordingly,

$$Z^{i}_{\lambda\cdots\mu} = \sum \partial_{\sigma} \Big[ \Big\{ \partial_{\lambda\cdots\hat{\sigma}\cdots\mu} \Big( x^{i} \circ \pi_{2} \Big) \Big\} \circ \pi_{1} \circ \rho \circ h \Big] \Big|_{c}.$$

We assert that  $\pi_1 \circ \rho \circ h$  is the identity map. First of all  $\pi_1$  is the map  $\pi$  for the  $\pi$ -space J. Then  $\rho$  (4.5) is a  $\pi$ -space morphism, for that is the way in which K is made into a  $\pi$ -space. So  $\pi_1 \circ \rho = \pi$ . But h has to satisfy 2.1, so  $\pi \circ h = \pi$ . But the  $\pi$  for M itself is the identity (on some neighborhood). Therefore

$$Z^{i}_{\lambda\cdots\mu} = \sum \partial_{\sigma} \{ \partial_{\lambda\cdots\hat{\sigma}\cdots\mu} (x^{i} \circ \pi_{2}) \} = k \partial_{\lambda\cdots\mu} (x^{i} \circ \pi_{2})$$

where k is the length of the string  $\lambda \cdots \mu$  with nothing deleted. So

$$Z^{i}_{\lambda\cdots\mu} = k\left\{\partial_{\lambda\cdots\mu}(x^{i}\circ\pi_{2})\right\}\circ\pi_{1}\circ\rho\circ h\Big|_{c} = k(x^{i})_{\lambda\cdots\mu}(A)$$

since h(c) = (A, g). Combine this with 4.83 and obtain that this Z term contributes

$$k(x^i)_{\lambda\cdots\mu}(A)rac{\partial g}{\partial (x^i)_{\lambda\cdots\mu}}$$

to the sum 4.7. This sum is evidently  $U_{\Gamma}[g]$  evaluated at A. Thus 4.7 is independent of the coordinates. Thus it defines the function for which 4.6 holds.

5. Proof of Theorem 3.2. Let  $\Lambda$  be given. Choose coordinates in M. Then 3.1 defines L, a function defined on  $J^{N}(M, S)$ . We have projections

$$J^{1}(M, K) = P$$

$$\xi \downarrow$$

$$K = T_{1}(J)$$

$$(5.1)$$

$$\rho \downarrow$$

$$J = J^{N}(M, S)$$

$$\downarrow$$

$$S$$

Hence L defines a function on  $J^1(M, K)$ , which we denote by L also, for simplicity. Using  $\xi$ , we can lift the function 4.4 up to a function  $\varphi$  on P. The function defined by 4.6 may be called  $\psi$ .

Let

(5.2) 
$$\omega = \lambda \, dt^1 \wedge \cdots \wedge dt^m,$$

 $\lambda \neq 0$ , be the volume element postulated in 3.2.

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We define

$$L^* = L + \lambda(\psi - \varphi) = L - \lambda \varphi + \lambda \psi,$$

and

$$\Lambda^* = L^* dt^1 \wedge \cdots \wedge dt^m.$$

We consider assertion 3.4. 4.8 is a derivative of first order. The coefficients in 4.7 are mere coordinates, and the Z's are linear in the derivatives 4.8. Hence 3.4 holds.

To show 3.5 we mention first that  $L^*$  has the form

$$L^{*} = L + P_{i}^{\lambda} \Big[ (\overline{(X^{i})}_{\lambda})_{\lambda} - \overline{(x^{i})_{\lambda}} \Big]$$
  
+  $P_{i}^{\lambda\mu} \Big[ (\overline{(x^{i})_{\lambda}})_{\mu} + (\overline{(x^{i})_{\mu}})_{\lambda} - 2\overline{(x^{i})_{\lambda\mu}} \Big]$   
+  $P_{i}^{\lambda\mu\nu} \Big[ (\overline{(x^{i})_{\mu\nu}})_{\lambda} + (\overline{(x^{i})_{\lambda\nu}})_{\mu} + (\overline{(x^{i})_{\lambda\mu}})_{\nu} - 3\overline{(x^{i})_{\lambda\mu\nu}} \Big] + \cdots,$ 

where  $P_i^{\lambda}$ ,  $P_i^{\lambda\mu}$ ,... are  $\lambda$  (see 5.2) times the  $p_i^{\lambda}$ ,  $p_i^{\lambda\mu}$ ,... lifted up to  $J^1(M, K)$  by the projections 5.1. It follows from [1, 3.4] that  $L^*$  has the "same" extremals as *l*. The meaning of *same* is as follows. An extremal for  $L^*$  gives us expressions

(5.3) 
$$x^i = f^i(t^1, \dots, t^m),$$

(5.4)  

$$(x^{i})_{\lambda} = u^{i}_{\lambda}(t^{1}, \dots, t^{m})$$

$$(x^{i})_{\lambda\mu} = u^{i}_{\lambda\mu}(t^{1}, \dots, t^{m})$$

$$\vdots$$

$$P^{i}_{\lambda} = g^{i}_{\lambda}(t^{1}, \dots, t^{m})$$

$$\vdots$$

If we abandon 5.4 and those following it, then 5.3 gives an extremal for L in the usual sense. It is shown in [1] that

$$u_{\lambda}^{i}=\frac{\partial f^{i}}{\partial t^{\lambda}},$$

and so forth.

6. Correction to [1].

(a) Delete 6.5. (Prop. 6.6. remains true, with the  $\varphi$  of 6.7.)

(b) Delete 6.9. (Better results are in the author's "The dynamic differential forms of the Klein-Gordon field and the conformal group", Jour. Geometry and Physics, Vol. 1, 1983.)

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