# 152. Some Radii of a Solid Associated with Polyharmonic Equations 

By Ichizo Yotsuya<br>Osaka Technical College

(Comm. by Kinjirô Kunugi, m. J. a., Nov. 12, 1973)

Introduction. In the preceding paper [1], we treated some quantities of a bounded domain in $R^{2}$ which we called polyharmonic inner radii. In the present paper, we deal with the similar quantities of a bounded domain in $R^{3}$ which is bounded by finite number of regular surfaces. G. Pólya and G. Szegö [2] defined the inner radius of a bounded domain using the Green's function of the domain relative to the Laplace's equation $\Delta u=0$ and they calculated the inner radius of a nearly spherical domain. The aim of this paper is to extend the above results. In the first place, we obtain the Green's function of a sphere relative to the $n$-harmonic equation $\Delta^{n} u=0$ and define the $n$-harmonic inner radius of a bounded domain. In the next place, we compute the $n$-harmonic inner radius of a nearly spherical domain and it is noticeable that it is monotonously decreasing with respect to integer $n$.

1. Inner radii associated with polyharmonic equations.

We use the following notations in this section. Let $V$ be a bounded domain in $R^{3}, S$ the surface of $V, P_{0}$ an inner point of $V, P$ the variable point in $V$ and $r$ the distance from $P_{0}$ to $P$.

Definition 1. If a function $u(P)$ satisfies the following two conditions, $u(P)$ is called the Green's function of $V$ with the pole $P_{0}$ relative to the $n$-harmonic equation $\Delta^{n} u=0$.
(1) In a neighborhood of $P_{0}, u(P)$ has the form

$$
u(P)=r^{2 n-3}+h_{n}(P)
$$

where $h_{n}(P)$ satisfies the equation $\Delta^{n} h_{n}=0$ in $V$ and all its derivatives of order $\leqq 2 n-1$ are continuous in $V+S$.
(2) All the normal derivatives of order $\leqq n-1$ of $u(P)$ vanish on $S$.

We can find the Green's function relative to the equation $\Delta^{n} u=0$ for a sphere in the explicit form.

Theorem 1. Let $V$ be the sphere of radius $R$ with the center $O$. If $P_{0} \neq O$, denoting $\rho$ the distance from $O$ to $P_{0}, P_{0}^{\prime}$ the inversion of $P_{0}$ with respect to $S$ and $r^{\prime}$ the distance from $P_{0}^{\prime}$ to $P$, the Green's function $G_{n}\left(P, P_{0}\right)$ of $V$ with the pole $P_{0}$ relative to the equation $\Delta^{n} u=0$ is as
follows,

$$
\begin{aligned}
& G_{1}\left(P, P_{0}\right)=\frac{1}{r}-\frac{R}{\rho r^{\prime}}, \\
& \begin{aligned}
G_{n}\left(P, P_{0}\right)=-\frac{R}{2 \rho r^{\prime}} & \left\{r^{2 n-4}\left(r-\frac{\rho}{R} r^{\prime}\right)^{2}\right. \\
& \left.-\sum_{k=2}^{n-1} \frac{(2 k-2)!}{2^{2 k-2} k!(k-1)!} r^{2 n-2 k-2}\left(r^{2}-\frac{\rho^{2}}{R^{2}} r^{\prime 2}\right)^{k}\right\}
\end{aligned}
\end{aligned}
$$

( $n \geqq 2$ ).
And $P_{0}=O$, we put $\rho r^{\prime}=R^{2}$ in the above equalities.
Proof. The case of $n=1$ is well known. Obviously the function $G_{n}\left(P, P_{0}\right)(n \geqq 2)$ satisfies the condition (1). We write the function $G_{n}\left(P, P_{0}\right)$ in the form

$$
G_{n}\left(P, P_{0}\right)=-\frac{R r^{2 n-2}}{2 \rho r^{\prime}}\left[\left(1-\frac{\rho r^{\prime}}{R r}\right)^{2}-\sum_{k=2}^{n-1} \frac{(2 k-2)!}{2^{2 k-2} k!(k-1)!}\left\{1-\left(\frac{\rho r^{\prime}}{R r}\right)^{2}\right\}^{k}\right]
$$

and we put

$$
x=\left(\frac{\rho r^{\prime}}{R r}\right)^{2} .
$$

Then $x$ is equal to 1 on $S$, and we can rewrite

$$
G_{n}\left(P, P_{0}\right)=-\frac{R r^{2 n-2}}{2 \rho r^{\prime}}\left\{(1-\sqrt{x})^{2}-\sum_{k=2}^{n-1} \frac{(2 k-2)!}{2^{2 k-2} k!(k-1)!}(1-x)^{k}\right\} .
$$

If we put

$$
f_{n}(x)=(1-\sqrt{x})^{2}-\sum_{k=2}^{n-1} \frac{(2 k-2)!}{2^{2 k-2} k!(k-1)!}(1-x)^{k},
$$

then

$$
f_{n}^{(\alpha)}(1)=0, \quad 0 \leqq \alpha \leqq n-1 .
$$

From this the condition (2) follows. The theorem is thus proved.
Given a domain $V$ and an inner point $P_{0}$ of $V$, G. Pólya and G. Szegö [2] defined the inner radius $r_{P_{0}}$ of $V$ with respect to the point $P_{0}$ as follows; when the Green's function $G\left(P, P_{0}\right)$ of $V$ with the pole $P_{0}$ relative to the equation $\Delta u=0$ is

$$
G\left(P, P_{0}\right)=\frac{1}{r}+h(P)
$$

they put

$$
\frac{1}{r_{P_{0}}}=-h\left(P_{0}\right) .
$$

Now we define the $n$-harmonic inner radius of a domain $V$ associated with $n$-harmonic equation $\Delta^{n} u=0$.

Definition 2. If the Green's function of a domain $V$ with the pole $P_{0}$ relative to the equation $\Delta^{n} u=0$ is

$$
r^{2 n-3}+h_{n}(P)
$$

then we put

$$
\frac{1}{r_{P_{0}, 1}}=-h_{1}\left(P_{0}\right)
$$

$$
\frac{(2 n-4)!}{2^{2 n-3}(n-1)!(n-2)!} r_{P_{0}, n}^{2 n-3}=\left|h_{n}\left(P_{0}\right)\right| \quad(n \geqq 2)
$$

we call $r_{P_{0}, n}$ the $n$-harmonic inner radius of the domain $V$ with respect to the point $P_{0}$.

Remark. When the domain $V$ is the sphere of radius $R$, we compute the ordinary inner radius and the $n$-harmonic inner radius of the sphere with respect to the point $P_{0}$, which are the same value

$$
\frac{R^{2}-\rho^{2}}{R}
$$

for an arbitrary integer $n$.
2. Inner radii of a nearly spherical domain.

In this section, we consider the radii of a nearly spherical domain defined in the former section.

Definition 3. Let

$$
\begin{equation*}
r=1+\rho(\theta, \varphi) \tag{1}
\end{equation*}
$$

be the equation of the surface of a domain in spherical coordinates $r$, $\theta$ and $\varphi$, where $\rho(\theta, \varphi)$ represents the infinitesimal variation of the unit sphere. We call the domain bounded by (1) the nearly spherical domain.

We consider the series

$$
\begin{equation*}
\rho(\theta, \varphi)=\sum_{k=0}^{\infty} X_{k}(\theta, \varphi), \tag{2}
\end{equation*}
$$

where the term $X_{k}(\theta, \varphi)$ represents a surface harmonic of degree $k$ which has infinitesimal coefficients of the first order.

Terms of higher infinitesimal than the second order are neglected in all the discussions of this section.

Lemma. When the surface harmonic $X_{1}(\theta, \varphi)$ in the series (2) is $X_{1}(\theta, \varphi)=(a \cos \varphi+b \sin \varphi) \sin \theta+c \cos \theta$,
neglecting terms of higher than the first order, the position of the centroid $C$ of the nearly spherical domain $r<1+\rho(\theta, \varphi)$ is $(a, b, c)$.

This lemma was given by G. Pólya and G. Szegö [2] and they obtained the ordinary inner radius $r_{C}$ of the nearly spherical domain with respect to the centroid $C$ as follows,

$$
\begin{equation*}
r_{C}=1+X_{0}+\frac{1}{4 \pi}\left\{\int\left[X_{1}(\theta, \varphi)\right]^{2} d S-\sum_{k=2}^{\infty}(k+1) \int\left[X_{k}(\theta, \varphi)\right]^{2} d S\right\} \tag{3}
\end{equation*}
$$

where the integral is extended over the surface of the unit sphere of which $d S$ is an element. As an extention of (3), we prove the following theorem.

Theorem 2. For an arbitrary positive integer $n$, the $n$-harmonic inner radius $r_{C, n}$ of the nearly spherical domain $r<1+\rho(\theta, \varphi)$ with respect to the centroid $C$ is

$$
\begin{equation*}
r_{C, n}=1+X_{0}+\frac{1}{4 \pi}\left\{\int\left[X_{1}(\theta, \varphi)\right]^{2} d S-\sum_{k=2}^{\infty}(n k-n+2) \int\left[X_{k}(\theta, \varphi)\right]^{2} d S\right\} \tag{4}
\end{equation*}
$$

Consequently, $r_{C, n}$ decreases monotonously with respect to $n$.
Proof. We seek the Green's function $G_{n}(P, C)$ of the nearly spherical domain with the pole $C$ relative to the equation $\Delta^{n} u=0$ in the form

$$
\begin{aligned}
G_{n}(P, C)= & -\frac{1}{2}\left\{r^{\prime 2 n-4}\left(r^{\prime}-1\right)^{2}-\sum_{k=2}^{n-1} \frac{(2 k-2)!}{2^{2 k-2} k!(k-1)!} r^{\prime 2 n-2 k-2}\left(r^{\prime 2}-1\right)^{k}\right\} \\
& +p(r, \theta, \varphi)+q(r, \theta, \varphi), \\
p(r, \theta, \varphi)= & \sum_{k=0}^{\infty} \sum_{i=0}^{n-1} r^{(k+2 i} S_{k}^{(i)}(\theta, \varphi), \\
q(r, \theta, \varphi)= & \sum_{k=0}^{\infty} \sum_{i=0}^{n-1} r^{k+2 i} T_{k}^{(i)}(\theta, \varphi) .
\end{aligned}
$$

Here $r^{\prime}$ is the distance from $C(a, b, \dot{c})$ to the point $P, S_{k}^{(t)}(\theta, \varphi)$ and $T_{k}^{(i)}(\theta, \varphi)$ are surface harmonics of degree $k$ with first and second order coefficients respectively. The $n$-harmonic inner radius $r_{C, n}$ is determined by

$$
\begin{aligned}
& \frac{(2 n-4)!}{2^{2 n-3}(n-1)!(n-2)!} r_{c, n}^{2 n-3} \\
& \quad=\left|(-1)^{n-1} \frac{(2 n-4)!}{2^{2 n-3}(n-1)!(n-2)!}+p\left(r_{0}, \theta_{0}, \varphi_{0}\right)+q\left(r_{0}, \theta_{0}, \varphi_{0}\right)\right|
\end{aligned}
$$

where $r_{0}, \theta_{0}$ and $\varphi_{0}$ are spherical coordinates of the centroid $C$. And so we have

$$
\begin{equation*}
r_{C, n}^{2 n-3}=1+(-1)^{n-1} \frac{2^{2 n-3}(n-1)!(n-2)!}{(2 n-4)!}\left\{S_{0}^{(0)}+\left(r S_{1}^{(0)}\right)+T_{0}^{(0)}\right\} \tag{5}
\end{equation*}
$$

where the term $\left(r S_{1}^{(0)}\right)$ has to be taken at $C$. If $\nu$ denotes the normal to the boundary of the nearly spherical domain, the condition

$$
\frac{\partial^{m} G_{n}}{\partial \nu^{m}}=0
$$

on the boundary can be replaced by

$$
\frac{\partial^{m} G_{n}}{\partial r^{m}}=0
$$

Let $\gamma$ be the variable angle between the radii $r$ and $r_{0}$. In view of lemma, we have

$$
r_{0} \cos \gamma=X_{1}(\theta, \varphi)
$$

Now the boundary conditions are

$$
\begin{aligned}
& \frac{\partial^{\alpha}}{\partial r^{\alpha}} p(1, \theta, \varphi)+\rho(\theta, \varphi) \frac{\partial^{\alpha+1}}{\partial r^{\alpha+1}} p(1, \theta, \varphi)+\frac{\partial^{\alpha}}{\partial r^{\alpha}} q(1, \theta, \varphi)=0 \\
& \frac{\partial^{n-2}}{\partial r^{n-2}} p(1, \theta, \varphi)+\rho(\theta, \varphi) \frac{\partial^{n-1}}{\partial r^{n-1}} p(1, \theta, \varphi)+\frac{\partial^{n-2}}{\partial r^{n-2}} q(1, \theta, \varphi) \\
& \quad=\frac{(2 n-2)!}{2^{n}(n-1)!}\left\{\rho(\theta, \varphi)-r_{0} \cos \gamma\right\}^{2}, \\
& \frac{\partial^{n-1}}{\partial r^{n-1}} p(1, \theta, \varphi)+\rho(\theta, \varphi) \frac{\partial^{n}}{\partial r^{n}} p(1, \theta, \varphi)+\frac{\partial^{n-1}}{\partial r^{n-1}} q(1, \theta, \varphi)
\end{aligned}
$$

$$
\begin{aligned}
= & \frac{(2 n-2)!}{2^{n+1}(n-1)!}\left[\left(n^{2}+n-6\right)\left\{\rho(\theta, \varphi)-r_{0} \cos \gamma\right\}^{2}+2 r_{0}^{2} \sin \gamma\right] \\
& +\frac{(2 n-2)!}{2^{n-1}(n-1)!}\left\{\rho(\theta, \varphi)-r_{0} \cos \gamma\right\} .
\end{aligned}
$$

The first order terms yield

$$
\begin{align*}
\frac{\partial^{\alpha}}{\partial r^{\alpha}} p(1, \theta, \varphi) & =0 \quad 0 \leqq \alpha \leqq n-2  \tag{7}\\
\frac{\partial^{n-1}}{\partial r^{n-1}} p(1, \theta, \varphi) & =\frac{(2 n-2)!}{2^{n-1}(n-1)!}\left\{\rho(\theta, \varphi)-r_{0} \cos \gamma\right\}
\end{align*}
$$

So that

$$
\begin{equation*}
p(r, \theta, \varphi)=\frac{(2 n-2)!}{2^{2 n-2}\{(n-1)!\}^{2}}\left(r^{2}-1\right)^{n-1}\left\{X_{0}+\sum_{k=2}^{\infty} r^{k} X_{k}(\theta, \varphi)\right\} \tag{8}
\end{equation*}
$$

in particular

$$
\begin{equation*}
S_{0}^{(0)}=(-1)^{n-1} \frac{(2 n-2)!}{2^{2 n-2}\{(n-1)!\}^{2}} X_{0}, \quad\left(r S_{1}^{(0)}\right)=0 \tag{9}
\end{equation*}
$$

We consider the second order terms. The mean value of the function $q(r, \theta, \varphi)$ on the surface of the unit sphere is equal to $\sum_{i=0}^{n-1} r^{2 i} T_{0}^{(i)}$. By the first equations of (6) and (7) we have

$$
\frac{\partial^{\alpha}}{\partial r^{\alpha}} q(1, \theta, \varphi)=0 \quad 0 \leqq \alpha \leqq n-3
$$

so that it must be the form

$$
\begin{equation*}
\sum_{i=0}^{n-1} r^{2 i} T_{0}^{(i)}=\left(r^{2}-1\right)^{n-2}\left(A r^{2}+B\right) \tag{10}
\end{equation*}
$$

where $A$ and $B$ are constants. Comparing the constant coefficients of $q(r, \theta, \varphi)$ and (10), we obtain
(11) $\quad T_{0}^{(0)}=(-1)^{n-2} B$.

We take now the mean values of second order terms of (6) and find

$$
\begin{aligned}
& A+B=-\frac{(2 n-2)!}{2^{2 n-2}(n-1)!(n-2)!}\left\{X_{0}^{2}+\sum_{k=2}^{\infty} \frac{1}{4 \pi} \int\left[X_{k}(\theta, \varphi)\right]^{2} d S\right\} \\
&(n+2) A+(n-2) B \\
&=-\frac{(2 n-2)!}{2^{2 n-2}\{(n-1)!\}^{2}} {\left[\left(n^{2}-3 n+6\right)\left\{X_{0}^{2}+\sum_{k=2}^{\infty} \frac{1}{4 \pi} \int\left[X_{k}(\theta, \varphi)\right]^{2} d S\right\}\right.} \\
&\left.+4 n \sum_{k=2}^{\infty} \frac{k}{4 \pi} \int\left[X_{k}(\theta, \varphi)\right]^{2} d S-\frac{4}{4 \pi} \int\left[X_{1}(\theta, \varphi)\right]^{2} d S\right] .
\end{aligned}
$$

Consequently

$$
\begin{align*}
B=\frac{(2 n-2)!}{2^{2 n-2}\{(n-1)!\}^{2}} & {\left[(-n+2)\left\{X_{0}^{2}+\sum_{k=2}^{\infty} \frac{1}{4 \pi} \int\left[X_{k}(\theta, \varphi)\right]^{2} d S\right\}\right.}  \tag{12}\\
& \left.+n \sum_{k=2}^{\infty} \frac{k}{4 \pi} \int\left[X_{k}(\theta, \varphi)\right]^{2} d S-\frac{1}{4 \pi} \int\left[X_{1}(\theta, \varphi)\right]^{2} d S\right] .
\end{align*}
$$

By virtue of (5), (9), (11) and (12) we find

$$
r_{C, n}=1+X_{0}+\frac{1}{4 \pi}\left\{\int\left[X_{1}(\theta, \varphi)\right]^{2} d S-\sum_{k=2}^{\infty}(n k-n+2) \int\left[X_{k}(\theta, \varphi)\right]^{2} d S\right\}
$$

This is the desired equation.

## References

[1] S. Ogawa, T. Kayano, and I. Yotsuya: Some radii associated with polyharmonic equations. Proc. Japan Acad., 47 (1), 44-49 (1971).
[2] G. Pólya and G. Szegö: Isoperimetric Inequalities in Mathematical Phisics. Princeton Univ. (1951).
[3] I. N. Vekua: New Methods for Solving Elliptic Equations. North-Holland, Amsterdam (1967).

