

## 2. An Asymptotic Formula for the Eigenvalues of the Laplacian in a Domain with a Small Hole

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**§ 1. Introduction.** This note is a continuation of our previous paper [1]. Let  $\Omega$  be a bounded domain in  $\mathbb{R}^3$  with  $C^\infty$  boundary  $\gamma$  and  $w$  be a fixed point in  $\Omega$ . For any sufficiently small  $\varepsilon > 0$ , let  $B_\varepsilon$  be the ball defined by  $B_\varepsilon = \{z \in \Omega; |z - w| < \varepsilon\}$ . Let  $\Omega_\varepsilon$  be the bounded domain defined by  $\Omega_\varepsilon = \Omega \setminus \bar{B}_\varepsilon$ . Then  $\partial\Omega_\varepsilon = \gamma \cup \partial B_\varepsilon$ .

Let  $0 > \mu_1(\varepsilon) \geq \mu_2(\varepsilon) \geq \dots$  be the eigenvalues of the Laplacian in  $\Omega_\varepsilon$  under the Dirichlet condition on  $\partial\Omega_\varepsilon$ . Let  $0 > \mu_1 \geq \mu_2 \geq \dots$  be the eigenvalues of the Laplacian in  $\Omega$  under the Dirichlet condition on  $\gamma$ . We arrange them repeatedly according to their multiplicities.

The aim of this note is to give an asymptotic expression of  $\mu_j(\varepsilon)$  as  $\varepsilon$  tends to zero. We need some notations to state the main result.

Let  $G(x, y)$  be the Green function of the Laplacian in  $\Omega$  satisfying

$$\begin{aligned} \Delta_x G(x, y) &= -\delta(x - y) & x, y \in \Omega, \\ G(x, y)|_{x \in \gamma} &= 0 & y \in \Omega. \end{aligned}$$

Then the Robin constant  $\tau$  ( $=\tau(w)$ ) at  $w$  is defined by

$$\tau = \lim_{x \rightarrow w} (G(x, w) - (4\pi)^{-1} |x - w|^{-1}).$$

Let  $G$  be the Green operator defined by

$$(1.1) \quad (Gf)(x) = \int_{\Omega} G(x, y) f(y) dy$$

for  $x \in \Omega$ .

We have the following

**Theorem 1.** Fix  $j$ . Assume that the multiplicity of  $\mu_j$  is one, then

$$(1.2) \quad \begin{aligned} \mu_j(\varepsilon) - \mu_j &= -(\tau + (4\pi\varepsilon)^{-1})^{-1} \varphi_j(w)^2 \\ &\quad - (\tau + (4\pi\varepsilon)^{-1})^{-2} e_j(w) \varphi_j(w) + O(\varepsilon^{5/2}) \end{aligned}$$

as  $\varepsilon$  tends to zero. Here  $\varphi_j(x)$  denotes the eigenfunction of the Laplacian under the Dirichlet condition on  $\gamma$  satisfying

$$\int_{\Omega} \varphi_j(x)^2 dx = 1.$$

And here

$$(1.3) \quad e_j(w) = \lim_{x \rightarrow w} (G(x, w) \varphi_j(w) + \psi(x)),$$

where  $\psi \in L^2(\Omega)$  is the unique solution of

$$(1.4) \quad ((G + (1/\mu_j))\psi)(x) = -(1/\mu_j)G(x, w)\varphi_j(w) - (1/\mu_j^2)\varphi_j(w)^2\varphi_j(x)$$

and

$$(1.5) \quad \int_{\Omega} \psi(x)\varphi_j(x)dx=0.$$

**Remarks.** The remainder term in (1.2) is not uniform with respect to  $j$ . The above formula (1.2) is a refinement of the formula (1.2) in [1]. See also [3], [2]. From (1.4) it is easily seen that  $G(x, w)\varphi_j(w) + \psi(x)$  is continuous with respect to  $x$ .

In § 2 we give a rough sketch of our proof of Theorem 1. Details of this paper will appear elsewhere.

### § 2. Sketch of proof of Theorem 1.

**Step 1.** Let  $\tau$  be as before. Put  $\varepsilon^* = (\tau + (4\pi\varepsilon)^{-1})^{-1}$ . Put  $\beta_\varepsilon = \{x \in \Omega; G(x, w) > \varepsilon^{*-1}\}$  and  $w_\varepsilon = \Omega \setminus \bar{\beta}_\varepsilon$ . Then it is easy to see that there exists a constant  $C (> 0)$  independent of  $\varepsilon$  satisfying

$$(2.1) \quad \Omega_{\varepsilon + C\varepsilon^3} \subset \omega_\varepsilon \subset \Omega_{\varepsilon - C\varepsilon^3}.$$

Let  $0 > \tilde{\mu}_1(\varepsilon) \geq \dots \geq \tilde{\mu}_j(\varepsilon) \dots$  be the eigenvalues of the Laplacian in  $\omega_\varepsilon$  under the Dirichlet condition on  $\partial\omega_\varepsilon$ . If we prove

$$(2.2) \quad \tilde{\mu}_j(\varepsilon) - \mu_j = -\varepsilon^* \varphi_j(w)^2 - (\varepsilon^*)^2 e_j(w) \varphi_j(w) + O(\varepsilon^{5/2})$$

when the multiplicity of  $\mu_j$  is one, then we get (1.2) because of (2.1). Thus we have only to prove (2.2) to obtain Theorem 1.

**Step 2.** Put  $G^{(1)}(x, y) = G(x, y)$  and

$$G^{(k)}(x, y) = \int_{\Omega} G^{(k-1)}(x, z)G(z, y)dz$$

inductively. We define the symbol  $\langle \nabla_w, \nabla_w \rangle$  by the following:

$$\langle \nabla_w a(x, w), \nabla_w b(y, w) \rangle = \sum_{i=1}^3 \frac{\partial a}{\partial w_i}(x, w) \frac{\partial b}{\partial w_i}(y, w).$$

We now introduce the integral kernel  $h_\varepsilon(x, y)$  by the following:

$$(2.3) \quad \begin{aligned} h_\varepsilon(x, y) = & G^{(3)}(x, y) - \varepsilon^* \sum_{k=1}^3 G^{(k)}(x, w)G^{(4-k)}(y, w) \\ & - 4\pi\varepsilon^3 \sum_{k=1}^3 \langle \nabla_w G^{(k)}(x, w), \nabla_w G^{(4-k)}(y, w) \rangle \xi_\varepsilon(x)\xi_\varepsilon(y) \\ & + (\varepsilon^*)^2 G^{(2)}(w, w) \sum_{k=1}^2 G^{(k)}(x, w)G^{(3-k)}(y, w) \\ & + (\varepsilon^*)^2 G^{(3)}(w, w)G(x, w)G(y, w), \end{aligned}$$

where  $\xi_\varepsilon(x) \in C^\infty(\Omega)$  is a function satisfying  $|\xi_\varepsilon(x)| \leq 1$ ,  $\xi_\varepsilon(x) = 0$  for  $x \in \beta_{\varepsilon/2}$ ,  $\xi_\varepsilon(x) = 1$  for  $x \in \omega_{(3/4)\varepsilon}$ .

Let  $H_\varepsilon$  be the operator given by

$$(2.4) \quad (H_\varepsilon f)(x) = \int_{\omega_\varepsilon} h_\varepsilon(x, y)f(y)dy$$

for  $x \in \omega_\varepsilon$ . And let  $G_\varepsilon$  be the operator given by

$$(2.5) \quad (G_\varepsilon g)(x) = \int_{\omega_\varepsilon} G_\varepsilon(x, y)g(y)dy$$

for  $x \in \omega_\varepsilon$ , where  $G_\varepsilon(x, y)$  be the Green kernel of the Laplacian in  $\omega_\varepsilon$  under the Dirichlet condition on  $\partial\omega_\varepsilon$ .

We have the following

**Proposition 1.** *There exists a constant  $C > 0$  independent of  $\varepsilon$*

such that

$$(2.6) \quad \|\mathbf{G}_\varepsilon^3 - \mathbf{H}_\varepsilon\|_{L^2(\omega_\varepsilon)} \leq C\varepsilon^{5/2}$$

holds, where  $\|T\|_{L^2(\omega_\varepsilon)}$  denotes the operator norm of  $T$ .

We need hard and laborious calculations including  $L^p$ -spaces to get (2.6).

Step 3. Let  $\tilde{\mathbf{H}}_\varepsilon$  be the operator given by

$$(2.7) \quad (\tilde{\mathbf{H}}_\varepsilon g)(x) = \int_\Omega h_\varepsilon(x, y)g(y)dy$$

for  $x \in \Omega$ .

We here construct an approximate eigenvalue and an approximate eigenfunction of  $\tilde{\mathbf{H}}_\varepsilon$ . Put  $\lambda = -1/\mu_j$  and  $\lambda_1 = -3\lambda^4\varphi_j(w)^2$ . We consider the following equations :

$$(2.8) \quad (\mathbf{G} - \lambda)\Phi(x) = \lambda\varphi_j(w)(\mathbf{G}(x, w) - \lambda\varphi_j(w)\varphi_j(x)),$$

$$(2.9) \quad \int_\Omega \Phi(x)\varphi_j(x)dx = 0.$$

Since  $\mathbf{G} - \lambda$  is the operator of Fredholm type, the unique solution  $\Phi$  in  $L^2(\Omega)$  of (2.8), (2.9) exists. Put

$$\lambda_2 = \lambda^2\varphi_j(w)^2(2\lambda\mathbf{G}^{(2)}(w, w) + \mathbf{G}^{(3)}(w, w)) - \lambda\varphi_j(w) \sum_{k=0}^2 \lambda^k (\mathbf{G}^{3-k}\Phi)(w).$$

Then consider the equations :

$$(2.10) \quad \begin{aligned} (\mathbf{G}^3 - \lambda^3)\Psi(x) &= \lambda_2\varphi_j(x) + \lambda_1\Phi(x) + \sum_{k=0}^2 \mathbf{G}^{(k+1)}(x, w)(\mathbf{G}^{3-k}\Phi)(w) \\ &\quad - \left( \sum_{k=1}^2 \lambda^{3-k}\mathbf{G}^{(k)}(x, w)\mathbf{G}^{(2)}(w, w) \right. \\ &\quad \left. + \lambda\mathbf{G}(x, w)\mathbf{G}^{(3)}(w, w) \right)\varphi_j(w) \end{aligned}$$

and

$$(2.11) \quad \int_\Omega \Psi(x)\varphi_j(x)dx = 0.$$

We see that (2.10), (2.11) have the unique solution  $\Psi$  in  $L^2(\Omega)$ . Now put  $\tilde{\lambda}(\varepsilon) = \lambda^3 + \varepsilon^*\lambda_1 + (\varepsilon^*)^2\lambda_2$  and  $\tilde{\varphi}(\varepsilon) = \varphi_j + \varepsilon^*\Phi + (\varepsilon^*)^2\Psi$ . We have the following

**Proposition 2.** *There exists a constant  $C > 0$  independent of  $\varepsilon$  such that*

$$(2.12) \quad \|(\tilde{\mathbf{H}}_\varepsilon - \tilde{\lambda}(\varepsilon))\tilde{\varphi}(\varepsilon)\|_{L^2(\Omega)} \leq C\varepsilon^{5/2}$$

holds.

Let  $\chi(\varepsilon)$  be the characteristic function of  $\omega_\varepsilon$ . Then we can prove

$$(2.13) \quad \|\chi(\varepsilon)\tilde{\mathbf{H}}_\varepsilon\tilde{\varphi}(\varepsilon) - \mathbf{H}_\varepsilon(\chi(\varepsilon)\tilde{\varphi}(\varepsilon))\|_{L^2(\omega_\varepsilon)} \leq C\varepsilon^{5/2}$$

for a constant  $C$  independent of  $\varepsilon$ .

By (2.6), (2.12) and (2.13) we get the following

**Proposition 3.** *The inequality*

$$(2.14) \quad \|(\mathbf{G}_\varepsilon^3 - \tilde{\lambda}(\varepsilon))\chi(\varepsilon)\tilde{\varphi}(\varepsilon)\|_{L^2(\omega_\varepsilon)} \leq C\varepsilon^{5/2}$$

holds.

As a consequence of (2.14), we conclude the following

**Proposition 4.** *There exists at least one eigenvalue  $\lambda^*(\varepsilon)$  of  $G_\varepsilon$  satisfying*

$$(2.15) \quad \lambda^*(\varepsilon)^3 = \lambda^3 + \varepsilon^* \lambda_1 + (\varepsilon^*)^2 \lambda_2 + O(\varepsilon^{5/2})$$

*as  $\varepsilon$  tends to zero.*

By the result of Rauch-Taylor [4], we see that there exists exactly one eigenvalue  $\lambda^*(\varepsilon)$  of  $G_\varepsilon$  satisfying (2.15). Therefore, (2.2) is proved.

### References

- [1] Ozawa, S.: Singular Hadamard's variation of domains and eigenvalues of the Laplacian. Proc. Japan Acad., **56A**, 306–310 (1980).
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- [4] Rauch, J., and M. Taylor: Potential and scattering theory on wildly perturbed domains. J. Funct. Anal., **18**, 27–59 (1975).