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SHARP EXPONENTIAL INTEGRABILITY FOR TRACES OF MONOTONE SOBOLEV FUNCTIONS

PEKKA PANKKA, PIETRO POGGI-CORRADINI AND KAI RAJALA

Abstract. We answer a question posed in [12] on exponential integrability of functions of restricted *n*-energy. We use geometric methods to obtain a sharp exponential integrability result for boundary traces of monotone Sobolev functions defined on the unit ball.

§1. Introduction

The following result answered a problem of A. Beurling, mentioned by J. Moser in a famous paper [10].

THEOREM A. (Chang-Marshall (1985), [1]) There is a universal constant $C < \infty$ so that if f is analytic in the unit disc \mathbb{D} , f(0) = 0, and

(1.1)
$$\int_{\mathbb{D}} |f'(z)|^2 \frac{\mathrm{d}A(z)}{\pi} \le 1,$$

then

$$\int_{0}^{2\pi} \exp\left(|f^{\star}(e^{i\theta})|^2\right) \mathrm{d}\theta \le C,$$

where f^* is the trace of f on $\partial \mathbb{D}$, i.e., $f^*(\zeta) = \lim_{t \uparrow 1} f(t\zeta)$ for \mathcal{H}^1 -a.e. $\zeta \in \partial \mathbb{D}$.

This result is moreover "sharp" in the following sense: the Beurling functions,

$$B_a(z) := \left(\log \frac{1}{1 - az}\right) \left(\log \frac{1}{1 - a^2}\right)^{-1/2} \quad 0 < a < 1$$

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are analytic in \mathbb{D} , satisfy $B_a(0) = 0$ and (1.1), and have the property that for any given $\alpha > 1$, one can choose a so that the integral

$$\int_0^{2\pi} \exp\left(\alpha |B_a(e^{i\theta})|^2\right) \mathrm{d}\theta$$

is as large as desired.

The following is an easy corollary of the Chang-Marshall Theorem.

COROLLARY A. There is a universal constant $C < \infty$ so that if $u : \mathbb{D} \to \mathbb{R}$ is harmonic with u(0) = 0 and

$$\int_{\mathbb{D}} |\nabla u(z)|^2 \, \frac{\mathrm{d}A(z)}{\pi} \le 1,$$

then

$$\int_0^{2\pi} \exp\left(u^* (e^{i\theta})^2\right) \mathrm{d}\theta \le C,$$

where u^* is the trace of u on $\partial \mathbb{D}$, i.e., $u^*(\zeta) = \lim_{t \uparrow 1} u(t\zeta)$ for \mathcal{H}^1 -a.e. $\zeta \in \partial \mathbb{D}$.

This can also be shown to be sharp by considering the real parts of the Beurling functions.

In [12] the last two authors generalized the Chang-Marshall theorem to quasiregular mappings in all dimensions. They asked in [12] whether Corollary A also generalizes, perhaps substituting "harmonic" with "nharmonic". In this note we show that this is indeed possible. The key concept is that of a monotone Sobolev function, whose definition we recall below, and which is quite general, and includes for instance n-harmonic functions.

§2. Main results

Let Ω be an open and connected set. For a continuous function $u: \Omega \to \mathbb{R}$, we define the oscillation of u on a compact set $K \subset \Omega$ by

$$\underset{K}{\operatorname{osc}} u = \max_{x,y \in K} |u(x) - u(y)|.$$

We say that $u: \Omega \to \mathbb{R}$ is monotone if $\operatorname{osc}_{\partial B} u = \operatorname{osc}_{\overline{B}} u$ for all *n*-balls *B* compactly contained in Ω .

By integrating the gradient over radial segments and changing variables, we see that, for a continuous $u: B^n \to \mathbb{R}$ in the Sobolev space $W^{1,n}(B^n)$, the radial limit

$$\tilde{u}(y) = \lim_{r \to 1} u(ry)$$

exists at \mathcal{H}^{n-1} -a.e. point $y \in S^{n-1}$. We denote by \tilde{u} the almost everywhere defined trace of u. Moreover, we denote the L^p -norm of a p-integrable $g: \Omega \to \mathbb{R}^n$ by $\|g\|_p = \|g\|_{\Omega,p}$. The surface measure $\mathcal{H}^{n-1}(S^{n-1})$ of the unit sphere S^{n-1} is ω_{n-1} . The notations $B^n(r) = B^n(0,r)$, $B^n = B^n(1)$ for n-dimensional balls will be used.

THEOREM 1. There exists a constant C = C(n) > 0 so that if $u \in W^{1,n}(B^n)$ is a non-constant continuous monotone function such that u(0) = 0, then

(2.2)
$$\int_{S^{n-1}} \exp\left(\alpha(|\tilde{u}(y)|/\|\nabla u\|_n)^{n/(n-1)}\right) \mathrm{d}\mathcal{H}^{n-1}(y) \le C,$$

where

(2.3)
$$\alpha = (n-1) \left(\frac{\omega_{n-1}}{2}\right)^{1/(n-1)}$$

The continuity assumption in Theorem 1 is of technical nature. By a theorem of Manfredi [8], so-called weakly monotone functions in $W^{1,n}$ are always continuous and monotone in the above sense. In general, $W^{1,n}$ functions need not be continuous.

The monotonicity assumption in Theorem 1 cannot be dropped altogether, since the *n*-capacity of a point is zero. Indeed, if we define $u_i: B^n \to \mathbb{R}$,

$$u_i(x) = \begin{cases} \frac{\log(1/|x|)}{\log i}, & 1/i \le |x| < 1, \\ 1, & 0 \le |x| < 1/i, \end{cases}$$

and $v_i = 1 - u_i$, we see that $v_i(0) = 0$, $\tilde{v}_i = 1$ on the unit sphere, and $\|\nabla v_i\|_n \to 0$ as $i \to \infty$.

Our method of proof for Theorem 1 has a similar geometric flavor as in [9] and in [12], and the end-game is again to appeal to Moser's original one-dimensional proof. However, the so-called "egg-yolk" property, which was the hardest part to establish in the two papers cited above, can be quickly established in our present case. It might come as a surprise then that Theorem 1 is sharp, as we will see in Theorem 2 below, as opposed to the situation in [12].

A function $u \in W^{1,p}_{loc}(\Omega)$ is called *p*-harmonic, 1 , if

$$\int_{\Omega} |\nabla u|^{p-2} \nabla u \cdot \nabla \phi \, \mathrm{d}x = 0$$

for every C^{∞} -smooth test function ϕ with compact support in Ω , see [6]. Since *p*-harmonic functions are continuous and satisfy the maximum principle ([6, 6.5]), they are, in particular, monotone.

The next result shows that the constant α in Theorem 1 is sharp.

THEOREM 2. Let α be as in Theorem 1. There exists a sequence of *n*-harmonic functions $u_i \in W^{1,n}(B^n)$ satisfying $\|\nabla u_i\|_n \leq 1$ and $u_i(0) = 0$, so that

$$\int_{S^{n-1}} \exp\left(\beta |\tilde{u}_i(y)|^{n/(n-1)}\right) \mathrm{d}\mathcal{H}^{n-1}(y) \to \infty \quad as \ i \to \infty$$

whenever $\beta > \alpha$.

§3. Proof of Theorem 1

In this section we assume that u satisfies the assumptions of Theorem 1. Moreover, by considering balls $B^n(0, 1-1/j)$, for j large, and using Fatou's lemma, we may assume that the function u in Theorem 1 is defined in a neighborhood of the unit ball.

LEMMA 3. There exists a constant $r_0 = r_0(n) > 0$ so that if $M_0 := \max_{\bar{B}^n(r_0)} |u|$, then

$$\int_{\{|u| \le M_0\}} |\nabla u|^n \,\mathrm{d}x \ge M_0^n.$$

Proof. For 0 < r < 1 let $m := \max_{\bar{B}^n(r)} |u|$ and set $v := \min\{|u|, m\}$. By monotonicity, and since u(0) = 0, $\operatorname{osc}_{S^{n-1}(t)} v = m$ for every $t \ge r$. By the Sobolev embedding theorem on spheres, see e.g. [5, Lemma 1] or [11], there exists a constant a depending only on n such that

$$\int_{B^n \setminus \bar{B}^n(r)} |\nabla v|^n \, \mathrm{d}x = \int_r^1 \left(\int_{S^{n-1}(t)} |\nabla v|^n \, \mathrm{d}\mathcal{H}^{n-1} \right) \mathrm{d}t$$
$$\geq \int_r^1 \frac{\left(\operatorname{osc}_{S^{n-1}(t)} v \right)^n}{at} \, \mathrm{d}t = \frac{m^n}{a} \log \frac{1}{r}.$$

The claim follows by choosing $r_0 := \exp(-a)$.

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Let Γ be a family of paths in an open and connected set Ω . The *n*modulus $M_n(\Gamma)$ of Γ is defined as follows:

$$\mathsf{M}_n(\Gamma) = \inf_{\rho} \int_{\Omega} \rho^n \, \mathrm{d}x,$$

where $\rho: \Omega \to [0, \infty]$ is an admissible function for Γ , i.e. a Borel function satisfying

(3.4)
$$\int_{\gamma} \rho \, \mathrm{d}s \ge 1$$

for every locally rectifiable $\gamma \in \Gamma$. The family of all paths joining two sets $A, B \subset \overline{\Omega}$ in Ω is denoted by $\Delta(A, B; \Omega)$. We say that a given property holds for *n*-almost every path in a path family Γ if the property holds for all paths in $\Gamma \setminus \Gamma_0$, where Γ_0 is a subfamily of Γ having *n*-modulus zero.

LEMMA 4. For every $r \in (0,1)$, there exists a constant c = c(n,r), so that

(3.5)
$$\mathcal{H}^{n-1}\big(\{y \in S^{n-1} : |u(y)| \ge s\}\big) \le c \exp\big(-\alpha I_M^s(u)\big)$$

for $s \geq M$. Here α is as in (2.3), $M = M(r, u) = \max_{S^{n-1}(r)} |u|$, and

$$I_M^s(u) = \int_M^s \frac{\mathrm{d}t}{\left(\int_{\{|u|=t\}} |\nabla u|^{n-1} \,\mathrm{d}\mathcal{H}^{n-1}\right)^{1/(n-1)}}.$$

Proof. Fix $r \in (0,1)$ and s > M = M(r,u). Write

$$E = E_s := \{ y \in S^{n-1} : |u(y)| \ge s \}$$

and

$$U_M := \{ x \in B^n : M \le |u(x)| \le s \}.$$

Also, in what follows we write

(3.6)
$$A_t := \int_{\{|u|=t\}} |\nabla u|^{n-1} \, \mathrm{d}\mathcal{H}^{n-1}.$$

The fact that A_t is a Borel function of t is standard, see for instance [2] p. 117. By the coarea formula, cf. [7], and the *n*-integrability of $|\nabla u|$, A_t is an integrable function of t.

We construct an admissible function ρ for $\Delta(B^n(r), E; B^n)$ as follows: Let $I = I^s_M(u)$, and set

$$\rho(x) := \frac{|\nabla u(x)|}{IA_{|u(x)|}^{1/(n-1)}} \chi_{U_M}(x).$$

Recall that, by Fuglede's theorem [3, Theorem 3], u is absolutely continuous on *n*-almost every path. So, for *n*-almost every $\gamma \in \Delta(B^n(r), E; B^n)$ parameterized by arc length $\ell(\gamma)$, we have, by change of variables

$$\int_{\gamma} \rho \,\mathrm{d}s = \int_{0}^{\ell(\gamma)} \frac{|\nabla u(\gamma(t))|}{IA_{|u(\gamma(t))|}^{1/(n-1)}} \chi_{U_{M}}(\gamma(t)) \,\mathrm{d}t$$

$$\geq I^{-1} \int_{0}^{\ell(\gamma)} \frac{|(u \circ \gamma)'(t)|}{A_{|(u \circ \gamma)(t)|}^{1/(n-1)}} \chi_{U_{M}}(\gamma(t)) \,\mathrm{d}t \geq I^{-1} \int_{M}^{s} \frac{\mathrm{d}t}{A_{t}^{1/(n-1)}} = 1.$$

Thus ρ is an admissible function for $\Delta(B^n(r), E; B^n)$, by the definition of *n*-modulus. By the coarea formula, we have

$$\begin{split} \mathsf{M}_{n}(\Delta(B^{n}(r),E;B^{n})) &\leq \int_{U_{M}} \rho^{n} \, \mathrm{d}x = I^{-n} \int_{U_{M}} \frac{|\nabla u(x)|^{n}}{A_{|u(x)|}^{n/(n-1)}} \, \mathrm{d}x \\ &= I^{-n} \int_{M}^{s} \int_{\{|u|=t\}} \frac{|\nabla u(y)|^{n-1}}{A_{t}^{n/(n-1)}} \, \mathrm{d}\mathcal{H}^{n-1}(y) \, \mathrm{d}t \\ &= I^{-n} \int_{M}^{s} \frac{A_{t}}{A_{t}^{n/(n-1)}} \, \mathrm{d}t = I^{1-n}. \end{split}$$

By the conformal invariance of n-modulus, taking inversion with respect to the unit sphere yields

$$2\mathsf{M}_n(\Delta(B^n(r), E; B^n)) \ge \mathsf{M}_n(\Delta(S^{n-1}(r) \cup S^{n-1}(1/r), E; \mathbb{R}^n)).$$

By spherical symmetrization and [4, Theorem 4],

$$2I^{1-n} \ge \mathsf{M}_n(\Delta(S^{n-1}(r) \cup S^{n-1}(1/r), E; \mathbb{R}^n))$$
$$\ge \omega_{n-1} \left(\log \frac{c_2}{\mathcal{H}^{n-1}(E)^{1/(n-1)}}\right)^{1-n},$$

where c_2 depends only on n and r. See [12] for further details. This implies (3.5).

Proof of Theorem 1. We will use the following result of Moser [10, Equation (6)]: If $\omega: [0, \infty) \to [0, \infty)$ is absolutely continuous and satisfies $\omega(0) = 0, \omega' \ge 0$ almost everywhere, and

$$\int_0^\infty (\omega'(t))^n \, \mathrm{d}t \le 1,$$

then

(3.7)
$$\int_0^\infty \exp\left(\omega(t)^{n/(n-1)} - t\right) \mathrm{d}t \le C,$$

where C > 0 depends only on *n*. By scaling invariance of (2.2), we may assume that

(3.8)
$$\int_{B^n} |\nabla u|^n \,\mathrm{d}x = 1.$$

Moreover, we fix $r = r_0$ and $M = M_0$ as in Lemma 3. Then, in particular, $M \leq 1$.

By the Cavalieri principle,

$$\int_{S^{n-1}} \exp\left(\alpha |u(x)|^{n/(n-1)}\right) \mathrm{d}\mathcal{H}^{n-1}(x)$$
$$= \omega_{n-1} + \frac{\alpha n}{n-1} \int_0^\infty s^{1/(n-1)} \mathcal{H}^{n-1}(E_s) \exp\left(\alpha s^{n/(n-1)}\right) \mathrm{d}s,$$

where

$$E_s = \{ y \in S^{n-1} : |u(y)| \ge s \}.$$

Then, by Lemma 4, it suffices to bound

(3.9)
$$\int_0^{\|u\|_{\infty}} s^{1/(n-1)} \exp\left(\alpha (s^{n/(n-1)} - I_M^s(u))\right) \mathrm{d}s,$$

where $||u||_{\infty} = \max_{y \in S^{n-1}} |u(y)|$, and $I^s_M(u) = 0$ for 0 < s < M. We define a function $\psi : [0, \infty) \to [0, \infty)$,

$$\psi(s) = \begin{cases} \mu s, & 0 < s < M\\ \alpha I_M^s(u) + \mu M, & M \le s \le \|u\|_{\infty}\\ \alpha I_M^{\|u\|_{\infty}}(u) + \mu M, & s > \|u\|_{\infty}, \end{cases}$$

where

(3.10)
$$\mu = \alpha \left(\frac{M}{\int_{\{|u| \le M\}} |\nabla u|^n \, \mathrm{d}x}\right)^{1/(n-1)}.$$

Then, by Lemma 3, $\mu M \leq \alpha$, and thus we may consider

(3.11)
$$\int_0^{\|u\|_{\infty}} s^{1/(n-1)} \exp\left(\alpha s^{n/(n-1)} - \psi(s)\right) \mathrm{d}s$$

instead of (3.9). We define ϕ by $\phi(y) = \psi^{-1}(y)$ for $0 < y < \|\psi\|_{\infty}$, and $\phi(y) = \|u\|_{\infty}$ for $y \ge \|\psi\|_{\infty}$. Then, changing variables $y = \psi(s)$ in (3.11) yields

(3.12)
$$\int_0^\infty \exp(\alpha \phi(y)^{n/(n-1)} - y) \phi'(y) \phi(y)^{1/(n-1)} \, \mathrm{d}y.$$

Integrating by parts, we then have that (3.12) equals $C_1(n) + C_2(n)T$,

$$T = \int_0^\infty \exp((\alpha^{(n-1)/n}\phi(y))^{n/(n-1)} - y) \, \mathrm{d}y.$$

Now, since ϕ is absolutely continuous and increasing, and $\phi(0) = 0$, Theorem 1 follows from Moser's result (3.7) if we can show that

(3.13)
$$\int_0^\infty \left(\alpha^{(n-1)/n} \phi'(y) \right)^n \, \mathrm{d}y \le 1.$$

We have

$$\alpha^{(n-1)/n}\phi'(y) = \begin{cases} \alpha^{(n-1)/n}\mu^{-1}, & 0 < y < \mu M \\ \alpha^{-1/n}A^{1/(n-1)}_{\phi(y)}, & \mu M < y < \|\psi\|_{\infty} \\ 0, & y > \|\psi\|_{\infty}, \end{cases}$$

where $A_{\phi(y)}$ as in (3.6). Hence,

(3.14)
$$\alpha^{n-1} \int_0^\infty \phi'(y)^n \, \mathrm{d}y = \alpha^{n-1} \mu^{1-n} M + \alpha^{-1} \int_{\mu M}^{\|\psi\|_\infty} A_{\phi(y)}^{n/(n-1)} \, \mathrm{d}y.$$

By our choice of μ , the first term equals $\int_{\{|u| \le M\}} |\nabla u|^n \, dx$. Also, by changing variables $\phi(y) = s$ in the right hand integral, and using the coarea formula,

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we have

(3.15)
$$\alpha^{-1} \int_{\mu M}^{\|\psi\|_{\infty}} A_{\phi(y)}^{n/(n-1)} \, \mathrm{d}y = \int_{\mu M}^{\|\psi\|_{\infty}} A_{\phi(y)} \phi'(y) \, \mathrm{d}y$$
$$= \int_{M}^{\|u\|_{\infty}} A_s \, \mathrm{d}s = \int_{\{|u| \ge M\}} |\nabla u|^n \, \mathrm{d}x.$$

Combining (3.14), (3.15), (3.10) and (3.8) then yields (3.13). The proof is complete.

§4. Proof of Theorem 2

Fix $\beta > \alpha$. For notational convenience, we consider first functions in $B^n(e_n, 1)$ instead of B^n . Fix $2 \leq i \in \mathbb{N}$, and denote $\varepsilon = \varepsilon_i = i^{-1}$. Define $v = v_i \colon B^n(-\varepsilon e_n, 2+\varepsilon) \to \overline{\mathbb{R}}$,

$$v(x) = -\log|x + \varepsilon e_n|.$$

Then v is n-harmonic in $B^n(-\varepsilon e_n, 2+\varepsilon) \setminus \{-\varepsilon e_n\}$. We first show that

(4.16)
$$\int_{B^n(e_n,1)} |\nabla v|^n \, \mathrm{d}x \le \frac{\omega_{n-1}}{2} \log \frac{2+\varepsilon}{\varepsilon}$$

Clearly,

$$\int_{B^n(e_n,1)} |\nabla v|^n \, \mathrm{d}x \le \frac{1}{2} \int_A |\nabla v|^n \, \mathrm{d}x,$$

where

$$A = B^n(-\varepsilon e_n, 2+\varepsilon) \setminus \bar{B}^n(-\varepsilon e_n, \varepsilon).$$

Since

$$|\nabla v(x)|^n = |x + \varepsilon e_n|^{-n},$$

we have

$$\frac{1}{2} \int_{A} |\nabla v|^{n} \,\mathrm{d}x = \frac{1}{2} \int_{B^{n}(0,2+\varepsilon) \setminus \bar{B}^{n}(0,\varepsilon)} |x|^{-n} \,\mathrm{d}x = \frac{\omega_{n-1}}{2} \log \frac{2+\varepsilon}{\varepsilon}.$$

Hence (4.16) holds.

To study exponential integrability of v, set

$$\gamma = \beta \left(\frac{\omega_{n-1}}{2}\log\frac{2+\varepsilon}{\varepsilon}\right)^{1/(1-n)}$$

and $\tau = \gamma/(n-1)$. By the choice of γ , (4.16), and the Cavalieri principle,

(4.17)
$$\int_{S^{n-1}(e_n,1)} \exp\left(\beta \left(|v|/\|\nabla v\|_n\right)^{n/(n-1)}\right) \mathrm{d}\mathcal{H}^{n-1}$$
$$\geq \omega_{n-1} + \frac{n\gamma}{n-1} \int_0^\infty \mathcal{H}^{n-1}(E_s) s^{1/(n-1)} \exp\left(\gamma s^{n/(n-1)}\right) \mathrm{d}s,$$

where

$$E_s = \{x \in S^{n-1}(e_n, 1) : |v(x)| \ge s\}.$$

Since

$$E_s = S^{n-1}(e_n, 1) \cap \bar{B}^n(-\varepsilon e_n, \exp(-s))$$
$$\cup S^{n-1}(e_n, 1) \setminus B^n(-\varepsilon e_n, \exp(s)),$$

we have

(4.18)
$$\mathcal{H}^{n-1}(E_s) \ge C(n)(\exp(-s))^{n-1} = C(n)\exp((1-n)s)$$

for $0 \le s \le \log(1/(2\varepsilon))$.

Combining (4.17) and (4.18) yields

$$\begin{aligned} \frac{1}{C(n)} \int_{S^{n-1}(e_n,1)} \exp(\beta(|v|/\|\nabla v\|_n)^{n/(n-1)}) \, \mathrm{d}\mathcal{H}^{n-1} \\ &\geq \frac{n\gamma}{n-1} \int_0^{\log(1/(2\varepsilon))} s^{1/(n-1)} \exp(\gamma s^{n/(n-1)} + (1-n)s) \, \mathrm{d}s \\ &= n\tau \int_0^{\log(1/(2\varepsilon))} s^{1/(n-1)} \exp((n-1)(\tau s^{n/(n-1)} - s)) \, \mathrm{d}s \\ &= \int_0^{\log(1/(2\varepsilon))} (n\tau s^{1/(n-1)} - (n-1)) \exp((n-1)(\tau s^{n/(n-1)} - s)) \, \mathrm{d}s \\ &+ (n-1) \int_0^{\log(1/(2\varepsilon))} \exp((n-1)(\tau s^{n/(n-1)} - s)) \, \mathrm{d}s \\ &\geq \exp\left((n-1)\left(\tau\left(\log\frac{1}{2\varepsilon}\right)^{n/(n-1)} - \log\frac{1}{2\varepsilon}\right)\right) - 1. \end{aligned}$$

Since

$$\left(\log\frac{2+\varepsilon}{\varepsilon}\right)^{1/(1-n)} \left(\log\frac{1}{2\varepsilon}\right)^{1/(n-1)} \ge 1-\delta(\varepsilon),$$

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where $\delta(\varepsilon) \to 0$ as $\varepsilon \to 0$, we have

(4.19)
$$\int_{S^{n-1}(e_n,1)} \exp\left(\beta (|v|/\|\nabla v\|_n)^{n/(n-1)}\right) \mathrm{d}\mathcal{H}^{n-1} \ge C(n)\varepsilon^{-T} - C(n),$$

where

$$T = (\beta - \alpha)(2/\omega_{n-1})^{1/(n-1)} - \delta'(\varepsilon),$$

and $\delta'(\varepsilon) \to 0$ when $\varepsilon \to 0$.

To prove Theorem 2, we consider the sequence $u_i \colon \overline{B}^n \to \mathbb{R}$,

$$u_i(x) = v_i(x+e_n) - v_i(e_n),$$

where $v_i(e_n) = -\log(1 + \varepsilon_i) \ge -\log 2$ for all *i*. We fix *M* such that

$$\beta' = \beta \left(\frac{M - \log 2}{M}\right)^{n/(n-1)} > \alpha.$$

Set also $E_i = \{y \in S^{n-1}(e_n, 1) : |v_i(y)| \ge M\}$. Then

$$\beta |v_i(y) - v_i(e_n)|^{n/(n-1)} \ge \beta' |v_i(y)|^{n/(n-1)}$$

on E_i for every *i*. Thus

$$\begin{split} &\int_{S^{n-1}} \exp(\beta(|u_i|/\|\nabla u_i\|_n)^{n/(n-1)}) \, \mathrm{d}\mathcal{H}^{n-1} \\ &= \int_{S^{n-1}(e_n,1)} \exp(\beta(|v_i(y) - v_i(e_n)|/\|\nabla v_i\|_n)^{n/(n-1)}) \, \mathrm{d}\mathcal{H}^{n-1}(y) \\ &\geq \int_{E_i} \exp(\beta'(|v_i(y)|/\|\nabla v_i\|_n)^{n/(n-1)}) \, \mathrm{d}\mathcal{H}^{n-1}(y) \\ &\geq \int_{S^{n-1}(e_n,1)} \exp(\beta'(|v_i(y)|/\|\nabla v_i\|_n)^{n/(n-1)}) \, \mathrm{d}\mathcal{H}^{n-1}(y) \\ &- \omega_{n-1} \exp(\beta'(M/\|\nabla v_i\|_n)^{n/(n-1)}). \end{split}$$

Since $\beta' > \alpha$ and $\varepsilon_i = i^{-1}$ in (4.19), the claim now follows from (4.19).

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Pekka Pankka Department of Mathematics University of Michigan Ann Arbor, MI 48109 USA

pankka@umich.edu

Pietro Poggi-Corradini Department of Mathematics Cardwell Hall Kansas State University Manhattan, KS 66506 USA pietro@math.ksu.edu Kai Rajala Department of Mathematics and Statistics P. O. Box 35 (MaD), FI-40014 Univ. of Jyväskylä Finland kirajala@maths.jyu.fi