Minkowskian product of Finsler spaces and Berwald connection

Dedicated to Professor Dr. Makoto Matsumoto on the occasion of his sixtieth birthday

By

Tsutomu Okada

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In the present paper, we shall first show Theorem 1: The Berwald connection of a Finsler space is the *h*-connection [1], [3] which is uniquely determined from the fundamental function by four axioms. A hint as to this theorem has been got from J. Grifone [2]. Next, using Theorem 1 we shall prove Theorem 2: The Berwald connection of any *Minkowskian product* of Finsler spaces coincides with that of the *Euclidean product* of these Finsler spaces. Finally, we shall obtain Theorem 3 on geodesics of the Minkowskian product of Finsler spaces.

The terminology and notations in the present paper are referred to M. Matsumoto's monograph [1].

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§1. Berwald connection

Let *M* be a C^{∞} manifold and *TM* the tangent bundle of *M*. And let (x^i) , i=1, 2, ..., n, be local coordinates of *M* and (x^i, y^i) the canonical local coordinates of *TM*. We denote by $s_0(TM)$ the zero section of *TM* and $(TM)_0 \equiv TM - s_0(TM)$. A positively homogeneous domain is defined as a subdomain *D* of *TM* such that if $y \in D$ then $ty \in D$ for any $t \in \mathbb{R}^+ - (0)$, where \mathbb{R}^+ is the non-negative real line.

Definition 1. A Finsler space (M, F) is a pair of a C^{∞} manifold M and a C^{1} function $F: TM \rightarrow \mathbb{R}^{+}$ such that

- (A) F(y) = 0 if and only if $y \in s_0(TM)$,
- (B) $F(ty) = t^2 F(y)$ for $t \in \mathbb{R}^+$
- (C) F is C^{∞} on a positively homogeneous domain D_1 ,

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(D) $g_{ij} \equiv \frac{\partial^2 F}{\partial y^i \partial y^j}$ is nondegenerate on a positively homogeneous domain $D_2(\subset D_1)$.

We call F the fundamental function of Finsler space (M, F).

Remark 1. It is obvious that if $D_1 = TM$ then (M, F) is a Riemann space. In §1 we may take $D_1 = D_2 = (TM)_0$, but in §2 and §3 those domains should be restricted as occasion calls. (See Remark 5 and 8)

Coefficients G_{kj}^i of the Berwald connection of a Finsler space (M, F) are given on D_2 by

$$G^{j} = \frac{1}{2} g^{ik} \left(\frac{\partial^{2} F}{\partial x^{r} \partial y^{k}} y^{r} - \frac{\partial F}{\partial x^{k}} \right),$$

$$G^{i}_{j} = \frac{\partial G^{i}}{\partial y^{j}}, \qquad G^{i}_{kj} = \frac{\partial G^{i}_{j}}{\partial y^{k}}.$$

Theorem 1. The Berwald connection of a Finsler space (M, F) is the hconnection which is uniquely determined on D_2 by the following four axioms:

- (a) The h-covariant derivatives of the fundamental function $F_{i}=0$.
- (b) The deflection tensor $D_i^i \equiv y_{i,i}^i = 0$.
- (c) The (v)hv-torsion tensor $P_{ik}^i = 0$.
- (d) The (h)h-torsion tensor $T^i_{jk} = 0$.

Proof. In terms of canonical local coordinates, we denote coefficients of the above *h*-connection by (F_{ik}^i, F_i^i) . Then these axioms are written as

(a) $F_{;i} = \frac{\partial F}{\partial x^i} - \frac{\partial F}{\partial y^i} F^i_i = 0,$

(b)
$$D_j^i = -F_j^i + y^i F_{lj}^i = 0,$$

- (c) $P_{jk}^i = \frac{\partial F_j^i}{\partial y^k} F_{kj}^i = 0,$
- (d) $T^{i}_{jk} = F^{i}_{jk} F^{i}_{kj} = 0.$

It is well-known that the Berwald connection (G_{jk}^i, G_j^i) satisfies these axioms. Now, from (c) and (d) we have

(1.1)
$$F_{jk}^{i} = F_{kj}^{i} = \frac{\partial F_{k}^{i}}{\partial y^{j}} = \frac{\partial F_{j}^{i}}{\partial y^{k}}.$$

The partial differentiation of (b) by y^k and (1.1) lead us to

(1.2)
$$\frac{\partial F_{l\,i}^i}{\partial y^k} y^l = 0.$$

Differentiating (a) by y^j and putting $g_{jl} = \frac{\partial^2 F}{\partial y^j \partial y^l}$, we get

(1.3)
$$\frac{\partial g_{jl}}{\partial x^i} y^l - g_{jl} F_i^l - g_{ls} y^s F_{jl}^l = 0.$$

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Moreover, the differentiation of (1.3) by y^k leads us to

(1.4)
$$\frac{\partial g_{jk}}{\partial x^i} - 2g_{jkl}F_l^l - g_{jl}F_{ki}^l - g_{kl}F_{ji}^l - g_{ls}y^s \frac{\partial F_{ji}^l}{\partial y^k} = 0$$

where we put $g_{jkl} = \frac{1}{2} \frac{\partial g_{jk}}{\partial y^{l}}$. Contracting (1.4) by y^{i} and using (1.1), (1.2) and (b), we obtain

(1.5)
$$\frac{\partial g_{jk}}{\partial x^i} y^i - 2g_{jkl}F^l_i y^i - g_{jl}F^l_k - g_{kl}F^l_j = 0.$$

Applying the Christoffel process [1] to (1.3) and (1.5), by (d) we see

 $\{i, k, j\}y^{i} - g_{jkl}F_{i}^{l}y^{i} - g_{kl}F_{j}^{l} = 0,$

where $\{i, k, j\} = \frac{1}{2} \left(\frac{\partial g_{jk}}{\partial x^i} + \frac{\partial g_{ki}}{\partial x^j} - \frac{\partial g_{ij}}{\partial x^k} \right)$. Therefore

(1.6)
$$F_{j}^{k} = \{_{ij}^{k}\} y^{i} - g_{jl}^{k} F_{i}^{l} y^{i},$$

where ${k \atop ij} = g^{kl}\{i, l, j\}$ and $g_{jl}^k = g^{ks}g_{jsl}$. In the following, the suffix 0 means contraction by y. Contracting (1.6) by y^j , we get

(1.7)
$$F_0^k = \{ {}_{00}^k \} = 2G^k.$$

Substituting (1.7) into (1.6), we have

(1.8)
$$F_{j}^{k} = \{ {}_{0j}^{k} \} - g_{lj}^{k} \{ {}_{00}^{l} \} = \frac{1}{2} \frac{\partial \{ {}_{00}^{k} \}}{\partial y^{j}} = G_{j}^{k}$$

Hence, from (c) we obtain

(1.9)
$$F_{ij}^{k} = \frac{1}{2} \frac{\partial^{2} \{ {}_{00}^{k} \}}{\partial y^{i} \partial y^{j}} = G_{ij}^{k}.$$

These three equations (1.7), (1.8) and (1.9) complete the proof.

Remark 2. Differentiating (1.5) by y^i and using (1.4), we see

(1.10)
$$2g_{ijk;0} + g_{i0} \frac{\partial F_{ji}^{l}}{\partial y^{k}} = 0.$$

Therefore, from (1.4) we obtain

$$(1.11) g_{jk;i} = -2g_{ijk;0}$$

This is a well-known equation, showing that the Berwald connection is not h-metrical in general although (a) is satisfied.

Remark 3. [M. Matsumoto] The equations (1.7) are also got directly as follows. Differentiating (a) by y^{j} , we see

(1.3)'
$$\frac{\partial^2 F}{\partial x^i \partial y^j} - \frac{\partial^2 F}{\partial y^k \partial y^j} F_i^k - \frac{\partial F}{\partial y^k} \frac{\partial F_i^k}{\partial y^j} = 0.$$

Contracting (1.3)' by y^i and using (b), (c) and (d), we have

(1.12)
$$\frac{\partial^2 F}{\partial x^0 \partial y^j} - \frac{\partial^2 F}{\partial y^k \partial y^j} F_0^k - \frac{\partial F}{\partial y^k} F_j^k = 0.$$

The last equations and (a) imply

(1.13)
$$g_{kj}F_0^k = \frac{\partial^2 F}{\partial x^0 \partial y^j} - \frac{\partial F}{\partial x^j}.$$

Thus we obtain

(1.7)
$$F_0^k = g^{kj} \left(\frac{\partial^2 F}{\partial x^0 \partial y^j} - \frac{\partial F}{\partial x^j} \right) = 2G^k.$$

§2. Minkowskian product of Finsler spaces

Let Ψ be a C^0 function: $\mathbf{R}^+ \times \mathbf{R}^+ \to \mathbf{R}^+$ such that

- (a) $\Psi(q^1, q^2) = 0$ if and only if $(q^1, q^2) = (0, 0)$,
- (β) $\Psi(tq^1, tq^2) = t\Psi(q^1, q^2)$ for $t \in \mathbb{R}^+$,
- (γ) Ψ is C^{∞} on $\mathbf{R}_0^+ \times \mathbf{R}_0^+$, where \mathbf{R}_0^+ is the positive real line,

(
$$\delta$$
) $\Psi_A \equiv \frac{\partial \Psi}{\partial q^A} \neq 0$ on $\mathbf{R}_0^+ \times \mathbf{R}_0^+$, $A = 1, 2,$

(
$$\epsilon$$
) $\Delta \equiv \Psi_1 \Psi_2 - 2\Psi \Psi_{12} \neq 0$ on $R_0^+ \times R_0^+$, where $\Psi_{12} = \frac{\partial^2 \Psi}{\partial q^1 \partial q^2}$

Remark 4. $\tilde{\Psi}(p^1, p^2) \equiv \Psi(\frac{1}{2}(p^1)^2, \frac{1}{2}(p^2)^2)$ is C^{∞} on $\mathbf{R}_0 \times \mathbf{R}_0$ if and only if $\Psi(q^1, q^2)$ is C^{∞} on $\mathbf{R}_0^+ \times \mathbf{R}_0^+$, where \mathbf{R}_0 is the real line except 0. Putting $G_{AB} = \frac{\partial^2 \tilde{\Psi}}{\partial p^A \partial p^B}$, we have

$$\Delta = G_{11}G_{22} - (G_{12})^2$$

on $\mathbf{R}_0 \times \mathbf{R}_0$. Moreover, putting $\Phi(p^1, p^2) = \sqrt{2\tilde{\psi}(p^1, p^2)}$, we get

$$\Delta = -\frac{\Phi^3}{p^1 p^2} \frac{\partial^2 \Phi}{\partial p^1 \partial p^2}$$

on $\mathbf{R}_0 \times \mathbf{R}_0$. Therefore, $\Delta = 0$ identically if and only if Φ is linear homogeneous with respect to p^1 and p^2 on $\mathbf{R}_0 \times \mathbf{R}_0$. For example, $\Psi = ((q^1)^r + (q^2)^r)^{\frac{1}{r}}$ satisfies all conditions (α) ~ (ε) provided $r \neq \frac{1}{2}$, but it does not satisfy (ε) if $r = \frac{1}{2}$.

Definition 2. A Finsler space (M, F) is called the *Minkowskian product* of Finsler spaces $(M_1, \stackrel{1}{F})$ and $(M_2, \stackrel{2}{F})$ with respect to a *product function* Ψ if $M = M_1 \times M_2$ and $F = \Psi(\stackrel{1}{F}, \stackrel{2}{F})$. Especially, in case of $\Psi = q^1 + q^2$ it is called the *Euclidean product* of these Finsler spaces.

Let (x^{a}) (resp. (x^{α})) be local coordinates of M_{1} (resp. M_{2}) and (x^{a}, y^{a}) (resp.

 (x^{α}, y^{α})) the canonical local coordinates of TM_1 (resp. TM_2). We shall refer to the local coordinates $(x^i) = (x^{\alpha}, x^{\alpha})$ adapted to the product manifold $M_1 \times M_2 = M$ and the canonical local coordinates $(x^i, y^i) = (x^{\alpha}, x^{\alpha}, y^{\alpha}, y^{\alpha})$ adapted to $T(M_1 \times M_2) = TM$. Throughout the remainder of the present paper, the ranges of indices are as follows: $a, b, c, ... = 1, 2, ..., r; \alpha, \beta, \gamma, ... = r+1, r+2, ..., r+s (=n); i, j, k, ... = 1, 2, ..., n; A, B = 1, 2.$

Assuming that $\stackrel{A}{F}$ is C^{∞} on $(TM_A)_0$ and putting $\stackrel{I}{g}_{ab} = \frac{\partial^2 \stackrel{I}{F}}{\partial y^a \partial y^b}$, $\stackrel{2}{g}_{\alpha\beta} = \frac{\partial^2 \stackrel{I}{F}}{\partial y^\alpha \partial y^\beta}$, the metric tensor g_{ij} of the Minkowskian product (M, F) of Finsler spaces $(M_1, \stackrel{I}{F})$ and $(M_2, \stackrel{2}{F})$ with respect to Ψ is written

(2.1)
$$\begin{cases} g_{ab} = \lambda_a^{\gamma} g_{cb}^{2}, \\ g_{a\beta} = \lambda_a^{\gamma} g_{\gamma\beta}^{2} = \lambda_{\beta}^{c} g_{ca}^{1}, \\ g_{\alpha\beta} = \lambda_{\alpha}^{\gamma} g_{\gamma\beta}^{2}, \end{cases}$$

on $(TM_1)_0 \times (TM_2)_0$, where

(2.2)
$$\begin{cases} \lambda_a^c = \Psi_1 \delta_a^c + \Psi_{11} y_a y^c, \quad \lambda_a^\gamma = \Psi_{12} y_a y^\gamma, \\ \lambda_\alpha^c = \Psi_{12} y_\alpha y^c, \quad \lambda_\alpha^\gamma = \Psi_2 \delta_\alpha^\gamma + \Psi_{22} y_\alpha y^\gamma, \\ y_a^l = \frac{\partial F}{\partial y^a}, \quad y_\alpha^l = \frac{\partial F}{\partial y^\alpha}, \quad \Psi_{AB} = \frac{\partial^2 \Psi}{\partial q^A \partial q^B}. \end{cases}$$

If conditions (δ) and (ε) are satisfied, then the reciprocal (ω_i^k) of (λ_i^j) is given by

(2.3)
$$\begin{cases} \omega_{b}^{c} = \frac{1}{\Psi_{1}} \delta_{b}^{c} + \Psi^{11} \overset{1}{\mathcal{Y}}_{b} \overset{r}{\mathcal{Y}}^{c}, \quad \omega_{b}^{\gamma} = \Psi^{12} \overset{1}{\mathcal{Y}}_{b} \overset{r}{\mathcal{Y}}^{\gamma}, \\ \omega_{\beta}^{c} = \Psi^{12} \overset{2}{\mathcal{Y}}_{\beta} \overset{r}{\mathcal{Y}}^{c}, \quad \omega_{\beta}^{\gamma} = \frac{1}{\Psi_{2}} \delta_{\beta}^{\gamma} + \Psi^{22} \overset{2}{\mathcal{Y}}_{\beta} \overset{r}{\mathcal{Y}}^{\gamma}, \end{cases}$$

on $(TM_1)_0 \times (TM_2)_0$, where

(2.4)
$$\Psi^{11} = \frac{-\Psi_{11}\Psi_2}{\Delta\Psi_1}, \quad \Psi^{12} = \frac{-\Psi_{12}}{\Delta}, \quad \Psi^{22} = \frac{-\Psi_{22}\Psi_1}{\Delta\Psi_2}.$$

Moreover, if $(\stackrel{1}{g_{ab}})$ (resp. $(\stackrel{2}{g_{a\beta}})$) is nondegenerate on $(TM_1)_0$ (resp. $(TM_2)_0$) and its reciprocal is denoted by $(\stackrel{1}{g^{bc}})$ (resp. $(\stackrel{2}{g^{\beta\gamma}})$), then the reciprocal (g^{jk}) of (g_{ij}) is written as

(2.5)
$$\begin{cases} g^{ac} = \int_{ab}^{ab} \omega_{b}^{c}, \\ g^{a\gamma} = \int_{ab}^{ab} \omega_{b}^{\gamma} = g^{2\gamma\beta} \omega_{\beta}^{a}, \\ g^{\alpha\gamma} = g^{2\alpha\beta} \omega_{\beta}^{\gamma} \end{cases}$$

on $(TM_1)_0 \times (TM_2)_0$.

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Remark 5. It is obvious [6] that if the fundamental function F of the Euclidean product (M, F) of Finsler spaces $(M_1, \stackrel{1}{F})$ and $(M_2, \stackrel{2}{F})$ is C^2 on $(TM)_0$, then all (M, F), $(M_1, \stackrel{1}{F})$ and $(M_2, \stackrel{2}{F})$ are Riemann spaces. In case of Minkowskian product, it is not true in general. For example, when the fundamental function $\stackrel{1}{F} = \frac{1}{2}\sqrt{u_{abcd}(x^a)y^ay^by^cy^d}$ (resp. $\stackrel{2}{F} = \frac{1}{2}\sqrt{v_{x\beta\gamma\delta}(x^a)y^ay^by^\gamma y^o}$) and the product function $\Psi = \sqrt{(q^1)^2 + (q^2)^2}$, the fundamental function $F = \Psi(\stackrel{1}{F}, \stackrel{2}{F})$ of Minkowskian product of $(M_1, \stackrel{1}{F})$ and $(M_2, \stackrel{2}{F})$ is C^2 on $(TM)_0$, but then (g_{ij}) is degenerate on $s_0(TM_1) \times (TM_2)_0 \cup (TM_1)_0 \times s_0(TM_2)$. (See Remark 8)

Now, we shall show that the Berwald connection of the Minkowskian product space is independent of Ψ .

Theorem 2. The Berwald connection of the Minkowskian product of Finsler spaces $(M_1, \stackrel{1}{F})$ and $(M_2, \stackrel{2}{F})$ with respect to any product function Ψ coincides with the Berwald connection of the Euclidean product of these Finsler spaces on $(TM_1)_0 \times (TM_2)_0$ provided that $\stackrel{1}{g}_{ab}$ (resp. $\stackrel{2}{g}_{\alpha\beta}$) exists and is nondegenerate on $(TM_1)_0$ (resp. $(TM_2)_0$).

Proof. In terms of canonical local coordinates, let $\overset{1}{G}_{bc}^{a}$ (resp. $\overset{2}{G}_{\beta\gamma}^{a}$) be coefficients of the Berwald connection of $(M_1, \overset{1}{F})$ (resp. $(M_2, \overset{2}{F})$). Next, in terms of the adapted canonical local coordinates, let G_{jk}^{i} be coefficients of the Berwald connection $B\Gamma^{e}$ of the Euclidean product of $(M_1, \overset{1}{F})$ and $(M_2, \overset{2}{F})$. Then it is obvious that on $(TM_1)_0 \times (TM_2)_0$

(2.6)
$$\begin{cases} G_{bc}^{a} = G_{bc}^{l}, \quad G_{\beta\gamma}^{\alpha} = G_{\beta\gamma}^{\alpha}, \\ G_{\beta\gamma}^{a} = G_{bc}^{\alpha} = G_{\betac}^{\alpha} = G_{\betac}^{\alpha} = 0. \end{cases}$$

Now, $B\Gamma^{e}$ is an *h*-connection on *M* which satisfies three axioms (*b*), (*c*) and (*d*) in in Theorem 1. As for (a), the *h*-covariant derivatives with respect to $B\Gamma^{e}$ of the fundamental function $F = \Psi(\vec{F}, \vec{F})$ of the Minkowskian product space are written as

(2.7)
$$\begin{cases} F_{;a} = \frac{\partial \Psi}{\partial q^{1}} \left(\frac{\partial F}{\partial x^{a}} - \frac{\partial F}{\partial y^{b}} G_{ca}^{b} y^{c} \right) - \frac{\partial \Psi}{\partial q^{2}} \frac{\partial F}{\partial y^{\beta}} G_{ka}^{\beta} y^{k}, \\ F_{;a} = \frac{\partial \Psi}{\partial q^{2}} \left(\frac{\partial F}{\partial x^{a}} - \frac{\partial F}{\partial y^{p}} G_{ya}^{\beta} y^{\gamma} \right) - \frac{\partial \Psi}{\partial q^{1}} \frac{\partial F}{\partial y^{b}} G_{ka}^{b} y^{k}. \end{cases}$$

Therefore $F_{;i} = 0$ are satisfied, because the *h*-covariant derivatives of the fundamental function $\stackrel{1}{F}$ (resp. $\stackrel{2}{F}$) with respect to the Berwald connection $(\stackrel{1}{G}_{bc}^{a})$ (resp. $(\stackrel{2}{G}_{\beta\gamma}^{a})$) of $(M_1, \stackrel{1}{F})$ (resp. $(M_2, \stackrel{2}{F})$) vanish and $G_{ka}^{b} = G_{ka}^{\beta} = 0$. Hence, from uniqueness of the Berwald connection in Theorem 1 this connection $B\Gamma^{e} = (G_{jk}^{i})$ is nothing but the Berwald connection of Finsler space $(M, F), M = M_1 \times M_2, F = \Psi(\stackrel{1}{F}, \stackrel{2}{F})$.

A Finsler space is called a Berwald space iff its Berwald connection is linear,

or the coefficients of the Berwald connection do not depend on y^i . Concerning Berwald spaces, the following are immediately deduced from Theorem 2.

Corollary 1. [7] If both Finsler spaces (M_1, F) and (M_2, F) are Berwald spaces, then the Minkowskian product of these spaces with respect to any Ψ is a Berwald space.

In particular we have

Corollary 2. If both (M_1, F) and (M_2, F) are Riemann spaces, then the Minkowskian product of these spaces with respect to any Ψ is a Berwald space and its Berwald connection is Levi-Civita's connection on the product Riemann manifold $M_1 \times M_2$.

§3. Geodesics of Minkowskian product space

A geodesic $x^{i}(t)$ with an affine parameter t of a Finsler space (M, F), or an extremal of the energy integral $\int F\left(x^{i}(t), \frac{dx^{i}}{dt}\right) dt$, satisfies Euler-Lagrange differential equations $e_{i}=0$, where

$$e_i = \frac{d}{dt} \left(\frac{\partial F}{\partial y^i} \right) - \frac{\partial F}{\partial x_i}, \quad y^i = \frac{dx^i}{dt}.$$

That t is an affine parameter is implied from

Lemma 1.
$$\frac{dF}{dt} = e_{i,y}^{i}$$
.
Proof. $e_{i,y}^{i} = \frac{d}{dt} \left(\frac{\partial F}{\partial y^{i}} y^{i} \right) - \left(\frac{\partial F}{\partial y^{i}} \frac{dy^{i}}{dt} + \frac{\partial F}{\partial x^{i}} y^{i} \right) = \frac{dF}{dt}$.

Now, let (M, F) be the Minkowskian product of Finsler spaces (M_1, F) and (M_2, F) with respect to Ψ . Preparing for the following theorem, stronger conditions for Ψ than (γ) , (δ) and (ε) must be introduced as follows;

- (sy) Ψ is C^{∞} on $(\mathbf{R}^+ \times \mathbf{R}^+)_0 \equiv \mathbf{R}^+ \times \mathbf{R}^+ (0, 0)$,
- $(s\delta)$ $\Psi_A \neq 0$ on $(\mathbf{R}^+ \times \mathbf{R}^+)_0$, A = 1, 2,
- (se) $\Delta \equiv \Psi_1 \Psi_2 2\Psi \Psi_{12} \neq 0$ on $(\mathbf{R}^+ \times \mathbf{R}^+)_0$.

Putting

$$\stackrel{1}{e_a} = \frac{d}{dt} \left(\frac{\partial F}{\partial y^a} \right) - \frac{\partial F}{\partial x^a}, \qquad \stackrel{2}{e_x} = \frac{d}{dt} \left(\frac{\partial F}{\partial y^z} \right) - \frac{\partial F}{\partial x^z},$$

we have

Lemma 2. (1) Under the condition $(s\gamma)$, if $e_a^1 = 0$ and $e_x^2 = 0$, then $e_i = 0$. (2) Under the conditions $(s\gamma)$, $(s\delta)$ and $(s\varepsilon)$, if $e_i = 0$, then $e_a^1 = 0$ and $e_x^2 = 0$. *Proof.* From $F = \Psi(\vec{F}, \vec{F})$, we see

(3.1)
$$\begin{cases} e_{a} = \Psi_{1} \overset{i}{e}_{a} + \left(\Psi_{11} \frac{d^{1}}{dt} + \Psi_{12} \frac{d^{2}}{dt}\right)^{1}_{\mathcal{Y}a}, \\ e_{z} = \Psi_{2} \overset{2}{e}_{z} + \left(\Psi_{21} \frac{d^{1}}{dt} + \Psi_{22} \frac{d^{2}}{dt}\right)^{2}_{\mathcal{Y}x}. \end{cases}$$

By Lemma 1 $d\vec{F}/dt = \vec{e}_a y^a$ and $d\vec{F}/dt = \vec{e}_a y^a$; therefore we have

(3.2)
$$\begin{cases} e_a = \lambda_a^b e_b^\dagger + \lambda_a^\beta e_\beta^2, \\ e_x = \lambda_a^b e_b^\dagger + \lambda_a^\beta e_\beta^2 \end{cases}$$

on $(TM)_0$ provided (sy) is satisfied (See (2.2)). Moreover, the last equations lead us to

(3.3)
$$\begin{cases} e_b = \omega_b^c e_c + \omega_b^{\gamma} e_{\gamma}, \\ \frac{2}{e_{\beta}} = \omega_{\beta}^c e_c + \omega_{\beta}^{\gamma} e_{\gamma}, \end{cases}$$

on $(TM)_0$ provided $(s\gamma)$, $(s\delta)$ and $(s\varepsilon)$ are satisfied (See (2.3)). (3.2) and (3.3) complete the proofs of (1) and (2) respectively.

Remark 6. $\Psi = ((q^1)^r + (q^2)^r)^{\frac{1}{r}}$, r = 2, 3, ..., satisfies (sy), but not (s δ) and (s ε). Following Ψ 's are examples satisfying all (sy), (s δ) and (s ε).

(1) $\Psi = q^1 + q^2$ (Euclidean)

(2)
$$\Psi = \frac{1}{2} \{ q^1 + q^2 + ((q^1)^r + (q^2)^r)^{\frac{1}{r}} \}, r = 2, 3, ..., [7]$$

(3)
$$\Psi = ((q^1)^2 + 3q^1q^2 + (q^2)^2)^{\frac{1}{2}}$$

(4)
$$\Psi = (2q^1 + 3q^2)^2/(q^1 + q^2)$$

(5) $\Psi = ((q^1)^2 + 3q^1q^2 + (q^2)^2)/(q^1 + q^2)$

Concerning geodesics of the Minkowskian product of two Finsler spaces, we obtain from Lemma 2 the following theorem, which is a generalization of the result [8] showed in case of the Minkowskian product of two Riemann spaces.

Theorem 3. Let a Finsler space (M, F) be the Minkowskian product of Finsler spaces $(M_1, \stackrel{1}{F})$ and $(M_2, \stackrel{2}{F})$ with respect to a product function Ψ .

(a) Under the condition (sy), any geodesic of (M_1, \vec{F}) (resp. (M_2, \vec{F})) is a geodesic of (M, F). That is to say, the Finsler space (M_1, \vec{F}) (resp. (M_2, \vec{F})) is a totally geodesic subspace of the Finsler space (M, F).

(b) Under the conditions $(s\gamma)$, $(s\delta)$ and $(s\varepsilon)$, the projection of any geodesic of (M, F) into M_1 (resp. M_2) is a geodesic of $(M_1, \stackrel{1}{F})$ (resp. $(M_2, \stackrel{2}{F})$).

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Proof (a) As \vec{F} is C^1 on TM_2 and positively homogeneous of degree 2 with respect to y^{α} , we see $\frac{\partial \vec{F}}{\partial y^{\alpha}} = 0$ on $s_0(TM_2)$. Therefore, we have $\vec{e}_{\alpha} = 0$ if $y^{\alpha}(t) = 0$ identically.

A geodesic $(x^{a}(t), x^{z})$ $(x^{z}$ being constant) with an affine parameter t of the subspace (M_{1}, F) satisfies $e_{a}^{1} = 0$, and $\frac{dx^{z}}{dt} = 0$ imply $e_{a}^{2} = 0$ by the above notice. Hence, by Lemma 2 (1) we obtain $e_{i} = 0$. Thus the curve $(x^{a}(t), x^{z})$ is a geodesic with an affine parameter t of (M, F). As to geodesic of (M_{2}, F) , it is proved similarly.

(b) We consider a geodesic $(x^{a}(t), x^{z}(t))$ with an affine parameter of (M, F). By Lemma 2 (2) $e_i = 0$ imply $e_a^{1} = 0$. Hence, the curve $(x^{a}(t), x^{z}) (x^{z}$ being constant) projected into M_1 is a geodesic with an affine parameter t of (M_1, F) . The proof as to the curve projected into M_2 is done similarly.

Remark 7. If \vec{F} is C^2 on $(TM_A)_0$, A = 1, 2, then we have

(3.4)
$$\begin{cases} \frac{1}{e_a} = \int_{a}^{1} \frac{dy^b}{dt} + 2G_a^1 & \text{on} \quad (TM_1)_0, \\ \frac{2}{e_x} = \int_{a}^{2} \frac{dy^\beta}{dt} + 2G_x^2 & \text{on} \quad (TM_2)_0, \end{cases}$$

where

$$2\dot{G}_{a}^{1} = \frac{\partial^{2}\dot{F}}{\partial x^{b}\partial y^{a}} y^{b} - \frac{\partial\dot{F}}{\partial x^{a}}, \quad 2\dot{G}_{\alpha}^{2} = \frac{\partial^{2}\dot{F}}{\partial x^{b}\partial y^{\alpha}} y^{\beta} - \frac{\partial\dot{F}}{\partial x^{\alpha}}$$

The equations (3.4) together with (3.2) and (2.1) lead us to

(3.5)
$$\begin{cases} G_a = \lambda_a^b G_b^1 + \lambda_a^\beta G_\beta^2, \\ G_a = \lambda_a^b G_b^1 + \lambda_a^\beta G_\beta^2, \end{cases}$$

where

$$2G_i = \frac{\partial^2 F}{\partial x^j \partial y^i} y^j - \frac{\partial F}{\partial x^i}.$$

Moreover, if $\begin{pmatrix} 1 \\ g_{ab} \end{pmatrix}$ (resp. $\begin{pmatrix} 2 \\ g_{x\beta} \end{pmatrix}$) is nondegenerate on $(TM_1)_0$ (resp. $(TM_2)_0$), then from (3.5) and (2.5) we obtain

$$(3.6) G^a = \overset{1}{G}{}^a, \quad G^a = \overset{2}{G}{}^a$$

on $(TM_1)_0 \times (TM_2)_0$, where

$$G^{i} = g^{ij}G_{j}, \quad G^{a} = g^{ab}G^{b}_{b}, \quad G^{a} = g^{a}{}^{a}{}^{b}G^{c}_{b},$$

This is another proof of Theorem 2.

Remark 8. If Ψ satisfies $(s\gamma)$, $(s\delta)$ and $F = \Psi(F^{1}, F^{2})$ is C^{2} on $(TM)_{0}$, then both

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 $(M_1, \stackrel{1}{F})$ and $(M_2, \stackrel{2}{F})$ are Riemann spaces. Hence, by Corollary 2 of Theorem 2 the Minkowskian product (M, F) of these spaces is Berwald space with Levi-Civita's connection.

Because
$$y'_{a} \equiv \frac{\partial F}{\partial y^{a}} = \frac{1}{\Psi_{1}} \frac{\partial F}{\partial y^{a}} \left(\text{resp. } y'_{x} \equiv \frac{\partial F}{\partial y^{x}} = \frac{1}{\Psi_{2}} \frac{\partial F}{\partial y^{x}} \right)$$

are C^1 on $(TM)_0$ under the above conditions. Therefore, y_a^1 (resp. y_α^2) are C^1 on TM_1 (resp. TM_2). Thus F^1 (resp. F^2) is C^2 on TM_1 (resp. TM_2).

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FACULTY OF ENGINEERING, DOSHISHA UNIVERSITY