# Minkowskian product of Finsler spaces and Berwald connection 

Dedicated to Professor Dr. Makoto Matsumoto on the occasion of his sixtieth birthday

## By

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In the present paper, we shall first show Theorem 1: The Berwald connection of a Finsler space is the $h$-connection [1], [3] which is uniquely determined from the fundamental function by four axioms. A hint as to this theorem has been got from J. Grifone [2]. Next, using Theorem 1 we shall prove Theorem 2: The Berwald connection of any Minkowskian product of Finsler spaces coincides with that of the Euclidean product of these Finsler spaces. Finally, we shall obtain Theorem 3 on geodesics of the Minkowskian product of Finsler spaces.

The terminology and notations in the present paper are referred to M. Matsumoto's monograph [1].

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## § 1. Berwald connection

Let $M$ be a $C^{\infty}$ manifold and $T M$ the tangent bundle of $M$. And let ( $x^{i}$ ), $i=1,2, \ldots, n$, be local coordinates of $M$ and $\left(x^{i}, y^{i}\right)$ the canonical local coordinates of $T M$. We denote by $s_{0}(T M)$ the zero section of $T M$ and $(T M)_{0} \equiv T M-s_{0}(T M)$. A positively homogeneous domain is defined as a subdomain $D$ of $T M$ such that if $y \in D$ then $t y \in D$ for any $t \in \boldsymbol{R}^{+}-(0)$, where $\boldsymbol{R}^{+}$is the non-negative real line.

Definition 1. A Finsler space $(M, F)$ is a pair of a $C^{\infty}$ manifold $M$ and a $C^{1}$ function $F: T M \rightarrow \boldsymbol{R}^{+}$such that
(A) $F(y)=0$ if and only if $y \in s_{0}(T M)$,
(B) $F(t y)=t^{2} F(y)$ for $t \in \boldsymbol{R}^{+}$
(C) $F$ is $C^{\infty}$ on a positively homogeneous domain $D_{1}$,
(D) $g_{i j} \equiv \frac{\partial^{2} F}{\partial y^{i} \partial y^{j}}$ is nondegenerate on a positively homogeneous domain $D_{2}\left(\subset D_{1}\right)$.
We call $F$ the fundamental function of Finsler space $(M, F)$.
Remark 1. It is obvious that if $D_{1}=T M$ then $(M, F)$ is a Riemann space. In $\S 1$ we may take $D_{1}=D_{2}=(T M)_{0}$, but in $\S 2$ and $\S 3$ those domains should be restricted as occasion calls. (See Remark 5 and 8)

Coefficients $G_{k j}^{i}$ of the Berwald connection of a Finsler space $(M, F)$ are given on $D_{2}$ by

$$
\begin{aligned}
& G^{j}=\frac{1}{2} g^{i k}\left(\frac{\partial^{2} F}{\partial x^{r} \partial y^{k}} y^{r}-\frac{\partial F}{\partial x^{k}}\right), \\
& G_{j}^{i}=\frac{\partial G^{i}}{\partial y^{j}}, \quad G_{k j}^{i}=\frac{\partial G_{i}^{i}}{\partial y^{k}} .
\end{aligned}
$$

Theorem 1. The Berwald connection of a Finsler space $(M, F)$ is the $h_{1}$ connection which is uniquely determined on $D_{2}$ by the following four axioms:
(a) The h-covariant derivatives of the fundamental function $F_{; i}=0$.
(b) The deflection tensor $D_{j}^{i} \equiv y_{; j}^{i}=0$.
(c) The (v)hv-torsion tensor $P_{j k}^{i}=0$.
(d) The (h)h-torsion tensor $T_{j k}^{i}=0$.

Proof. In terms of canonical local coordinates, we denote coefficients of the above $h$-connection by $\left(F_{j k}^{i}, F_{j}^{i}\right)$. Then these axioms are written as
(a) $\quad F_{; i}=\frac{\partial F}{\partial x^{i}}-\frac{\partial F}{\partial y} F_{i}^{\prime}=0$,
(b) $D_{j}^{i}=-F_{j}^{i}+y^{l} F_{i j}^{i}=0$,
(c) $\quad P_{j k}^{i}=\frac{\partial F_{j}^{i}}{\partial y^{k}}-F_{k j}^{i}=0$,
(d) $\quad T_{j k}^{i}=F_{j k}^{i}-F_{k j}^{i}=0$.

It is well-known that the Berwald connection ( $G_{j k}^{i}, G_{j}^{i}$ ) satisfies these axioms.
Now, from (c) and (d) we have

$$
\begin{equation*}
F_{j k}^{i}=F_{k j}^{i}=\frac{\partial F_{k}^{i}}{\partial y^{j}}=\frac{\partial F_{j}^{i}}{\partial y^{k}} \tag{1.1}
\end{equation*}
$$

The partial differentiation of (b) by $y^{\boldsymbol{k}}$ and (1.1) lead us to

$$
\begin{equation*}
\frac{\partial F_{l j}^{i}}{\partial y^{k}} y^{l}=0 \tag{1.2}
\end{equation*}
$$

Differentiating (a) by $y^{j}$ and putting $g_{j l}=\frac{\partial^{2} F}{\partial y^{j} \partial y^{l}}$, we get

$$
\begin{equation*}
\frac{\partial g_{j l}}{\partial x^{i}} y^{l}-g_{j l} F_{i}^{l}-g_{l s} y^{s} F_{j i}^{l}=0 \tag{1.3}
\end{equation*}
$$

Morec ver, the differentiation of (1.3) by $y^{k}$ leads us to

$$
\begin{equation*}
\frac{\partial g_{i k}}{\partial x^{i}}-2 g_{j k l} F_{i}^{l}-g_{j l} F_{k i}^{l}-g_{k l} F_{j i}^{l}-g_{l s} y^{s} \frac{\partial F_{j i}^{l}}{\partial y^{k}}=0, \tag{1.4}
\end{equation*}
$$

where we put $g_{j k l}=\frac{1}{2} \frac{\partial a_{j k}}{\partial y^{l}}$. Contracting (1.4) by $y^{i}$ and using (1.1), (1.2) and (b), we obtain

$$
\begin{equation*}
\frac{\partial g_{j k}}{\partial x^{i}} y^{i}-2 g_{j k l} F_{i}^{l} y^{i}-g_{j l} F_{k}^{l}-g_{k l} F_{j}^{l}=0 . \tag{1.5}
\end{equation*}
$$

Applying the Christoffel process [1] to (1.3) and (1.5), by (d) we see

$$
\{i, k, j\} y^{i}-g_{j k l} F_{i}^{\prime} y^{i}-g_{k l} F_{j}^{\prime}=0,
$$

where $\{i, k, j\}=\frac{1}{2}\left(\frac{\partial g_{j k}}{\partial x^{i}}+\frac{\partial g_{k i}}{\partial x^{j}}-\frac{\partial g_{i j}}{\partial x^{k}}\right)$. Therefore

$$
F_{j}^{k}=\left\{\begin{array}{l}
k  \tag{1.6}\\
i
\end{array}\right\} y^{i}-g_{j l}^{k} F_{i}^{l} y^{i},
$$

where $\left\{{ }_{i}^{k}\right\}=g^{k l}\{i, l, j\}$ and $g_{j l}^{k}=g^{k s} g_{j s l}$. In the following, the suffix 0 means contraction by $y$. Contracting (1.6) by $y^{j}$, we get

$$
F_{0}^{k}=\left\{\begin{array}{c}
k  \tag{1.7}\\
00
\end{array}\right\}=2 G^{k} .
$$

Substituting (1.7) into (1.6), we have

$$
F_{j}^{k}=\left\{\begin{array}{l}
k  \tag{1.8}\\
0_{j}
\end{array}\right\}-g_{l j}^{k}\left\{0_{0}^{l}\right\}=\frac{1}{2} \frac{\partial\left\{\left\{_{00}^{k}\right\}\right.}{\partial y^{j}}=G_{j}^{k} .
$$

Hence, from (c) we obtain

$$
F_{i j}^{k}=\frac{1}{2} \frac{\partial^{2}\left\{\begin{array}{l}
k  \tag{1.9}\\
\partial y^{i} \partial y^{j}
\end{array}\right.}{\partial y_{i j}} .
$$

These three equations (1.7), (1.8) and (1.9) complete the proof.
Remark 2. Differentiating (1.5) by $y^{i}$ and using (1.4), we see

$$
\begin{equation*}
2 g_{i j k ; 0}+g_{10} \frac{\partial F_{j i}^{l}}{\partial y^{k}}=0 . \tag{1.10}
\end{equation*}
$$

Therefore, from (1.4) we obtain

$$
\begin{equation*}
g_{j k ; i}=-2 g_{i j k ; 0} . \tag{1.11}
\end{equation*}
$$

This is a well-known equation, showing that the Berwald connection is not h-metrical in general although (a) is satisfied.

Remark 3. [M. Matsumoto] The equations (1.7) are also got directly as follows. Differentiating (a) by $y^{j}$, we see

$$
\begin{equation*}
\frac{\partial^{2} F}{\partial x^{i} \partial y^{j}}-\frac{\partial^{2} F}{\partial y^{k} \partial y^{j}} F_{i}^{k}-\frac{\partial F}{\partial y^{k}} \frac{\partial F_{i}^{k}}{\partial y^{j}}=0 . \tag{1.3}
\end{equation*}
$$

Contracting (1.3)' by $y^{i}$ and using (b), (c) and (d), we have

$$
\begin{equation*}
\frac{\partial^{2} F}{\partial x^{0} \partial y^{j}}-\frac{\partial^{2} F}{\partial y^{k} \partial y^{j}} F_{v}^{k}-\frac{\partial F}{\partial y^{k}} F_{j}^{k}=0 . \tag{1.12}
\end{equation*}
$$

The last equations and (a) imply

$$
\begin{equation*}
g_{k j} F_{0}^{k}=\frac{\partial^{2} F}{\partial x^{0} \partial y^{j}}-\frac{\partial F}{\partial x^{j}} . \tag{1.13}
\end{equation*}
$$

Thus we obtain

$$
\begin{equation*}
F_{0}^{k}=g^{k j}\left(\frac{\partial^{2} F}{\partial x^{0} \partial y^{j}}-\frac{\partial F}{\partial x^{j}}\right)=2 G^{k} . \tag{1.7}
\end{equation*}
$$

## § 2. Minkowskian product of Finsler spaces

Let $\Psi$ be a $C^{0}$ function: $\boldsymbol{R}^{+} \times \boldsymbol{R}^{+} \rightarrow \boldsymbol{R}^{+}$such that
( $\alpha$ ) $\Psi\left(q^{1}, q^{2}\right)=0$ if and only if $\left(q^{1}, q^{2}\right)=(0,0)$,
( $\beta$ ) $\Psi\left(t q^{1}, t q^{2}\right)=t \Psi\left(q^{1}, q^{2}\right)$ for $t \in \boldsymbol{R}^{+}$,
( $\gamma$ ) $\Psi$ is $C^{\infty}$ on $\boldsymbol{R}_{0}^{+} \times \boldsymbol{R}_{0}^{+}$, where $\boldsymbol{R}_{j}^{+}$is the positive real line,
( $\delta$ ) $\Psi_{A} \equiv \frac{\partial \Psi}{\partial q^{4}} \neq 0$ on $\boldsymbol{R}_{0}^{+} \times \boldsymbol{R}_{0}^{+}, A=1,2$,
(ع) $\Delta \equiv \Psi_{1} \Psi_{2}-2 \Psi \Psi_{12} \neq 0$ on $\boldsymbol{R}_{0}^{+} \times \boldsymbol{R}_{0}^{+}$, where $\Psi_{12}=\frac{\partial^{2} \Psi}{\partial q^{1} \partial q^{2}}$.
Remark 4. $\widetilde{\Psi}\left(p^{1}, p^{2}\right) \equiv \Psi\left(\frac{1}{2}\left(p^{1}\right)^{2}, \frac{1}{2}\left(p^{2}\right)^{2}\right)$ is $C^{\infty}$ on $\boldsymbol{R}_{0} \times \boldsymbol{R}_{0}$ if and only if $\Psi\left(q^{1}, q^{2}\right)$ is $C^{\infty}$ on $\boldsymbol{R}_{0}^{+} \times \boldsymbol{R}_{0}^{+}$, where $\boldsymbol{R}_{0}$ is the real line except 0 . Putting $G_{A B}=$ $\frac{\partial^{2} \tilde{\Psi}}{\partial p^{A} \partial p^{B}}$, we have

$$
\Delta=G_{11} G_{22}-\left(G_{12}\right)^{2}
$$

on $\boldsymbol{R}_{0} \times \boldsymbol{R}_{0}$. Moreover, putting $\Phi\left(p^{1}, p^{2}\right)=\sqrt{2}\left(p^{1}, p^{2}\right)$, we get

$$
\Delta=-\frac{\Phi^{3}}{p^{1} p^{2}} \frac{\partial^{2} \Phi}{\partial p^{1} \partial p^{2}}
$$

on $\boldsymbol{R}_{0} \times \boldsymbol{R}_{0}$. Therefore, $\Delta=0$ identically if and only if $\Phi$ is linear homogeneous with respect to $p^{1}$ and $p^{2}$ on $\boldsymbol{R}_{0} \times \boldsymbol{R}_{0}$. For example, $\Psi=\left(\left(q^{1}\right)^{r}+\left(q^{2}\right)^{r}\right)^{\frac{1}{r}}$ satisfies all conditions $(\alpha) \sim(\varepsilon)$ provided $r \neq \frac{1}{2}$, but it does not satisfy ( $\varepsilon$ ) if $r=\frac{1}{2}$.

Definition 2. A Finsler space ( $M, F$ ) is called the Minkowskian product of Finsler spaces $\left(M_{1}, \stackrel{1}{F}\right)$ and $\left(M_{2}, \stackrel{2}{F}\right)$ with respect to a product function $\Psi$ if $M=$ $M_{1} \times M_{2}$ and $F=\Psi(F, \stackrel{2}{F})$. Especially, in case of $\Psi=q^{1}+q^{2}$ it is called the Euclidean product of these Finsler spaces.

Let $\left(x^{a}\right)$ (resp. $\left(x^{\alpha}\right)$ ) be local coordinates of $M_{1}\left(\right.$ resp. $\left.M_{2}\right)$ and ( $x^{a}, y^{a}$ ) (resp.
$\left(x^{\alpha}, y^{\alpha}\right)$ ) the canonical local coordinates of $T M_{1}\left(\right.$ resp. $\left.T M_{2}\right)$. We shall refer to the local coordinates $\left(x^{i}\right)=\left(x^{a}, x^{\alpha}\right)$ adapted to the product manifold $M_{1} \times M_{2}=M$ and the canonical local coordinates $\left(x^{i}, y^{i}\right)=\left(x^{a}, x^{x}, y^{a}, y^{\alpha}\right)$ adapted to $T\left(M_{1} \times M_{2}\right)=$ $T M$. Throughout the remainder of the present paper, the ranges of indices are as follows : $a, b, c, \ldots=1,2, \ldots, r ; \alpha, \beta, \gamma, \ldots=r+1, r+2, \ldots, r+s(=n) ; i, j, k, \ldots=1$, $2, \ldots, n ; A, B=1,2$.

Assuming that $\stackrel{A}{F}$ is $C^{\infty}$ on $\left(T M_{A}\right)_{0}$ and putting $\stackrel{1}{g}_{a b}=\frac{\partial^{2} \stackrel{1}{F}}{\partial y^{a} \partial y^{b}}, \stackrel{2}{g}_{\alpha \beta}=\frac{\partial^{2} \stackrel{2}{F}}{\partial y^{\alpha} \partial y^{\beta}}$, the metric tensor $g_{i j}$ of the Minkowskian product $(M, F)$ of Finsler spaces $\left(M_{1}, \stackrel{1}{F}\right)$ and $\left(M_{2}, F\right)$ with respect to $\Psi$ is written

$$
\left\{\begin{array}{l}
g_{a b}=\lambda_{r}^{c} g_{c h}^{\prime},  \tag{2.1}\\
g_{a \beta}=\lambda_{a}^{2} g_{\gamma \beta}^{2}=i_{\beta}^{c} g_{c a}^{\prime}, \\
g_{x \beta}=i_{\alpha}^{*} g_{; \beta}^{2},
\end{array}\right.
$$

on $\left(T M_{1}\right)_{0} \times\left(T M_{2}\right)_{0}$, where

$$
\begin{aligned}
& \stackrel{1}{r_{a}}=\frac{\partial \stackrel{1}{F}}{\partial y^{a}}, \quad{ }_{r^{a}}^{r^{2}}=\frac{\partial \stackrel{2}{F}}{\partial y^{\alpha}}, \quad \Psi_{A B}=\frac{\partial^{2} \Psi}{\partial q^{A} \partial q^{B}} .
\end{aligned}
$$

If conditions $(\delta)$ and $(\varepsilon)$ are satisfied, then the reciprocal $\left(\omega_{j}^{k}\right)$ of $\left(\lambda_{i}^{j}\right)$ is given by
on $\left(T M_{1}\right)_{0} \times\left(T M_{2}\right)_{0}$, where

$$
\begin{equation*}
\Psi^{11}=\frac{-\Psi_{11} \Psi_{2}}{\Delta \Psi_{1}}, \quad \Psi^{12}=\frac{-\Psi_{12}}{\Delta}, \quad \Psi^{22}=\frac{-\Psi_{22} \Psi_{1}}{\Delta \Psi_{2}} . \tag{2.4}
\end{equation*}
$$

Moreover, if $\left(\stackrel{1}{g_{a b}}\right)$ (resp. $\left.\left(\underset{\alpha}{g_{\alpha \beta}}\right)\right)$ is nondegenerate on $\left(T M_{1}\right)_{0}$ (resp. $\left.\left(T M_{2}\right)_{0}\right)$ and its reciprocal is denoted by $\left(g^{b c}\right)\left(\right.$ resp. $\left(g^{2} g^{\beta j i}\right)$, then the reciprocal $\left(g^{j k}\right)$ of $\left(g_{i j}\right)$ is written as

$$
\left\{\begin{array}{l}
g^{a c}=g^{a b} \omega_{b}^{c},  \tag{2.5}\\
g^{a \gamma}=g^{\frac{1}{a b}} \omega_{b}^{\gamma}=g^{2}{ }^{z \beta} \omega_{\beta}^{a}, \\
g^{\alpha \gamma}=g^{2 \beta} \omega_{\beta}^{\gamma}
\end{array}\right.
$$

on $\left(T M_{1}\right)_{0} \times\left(T M_{2}\right)_{0}$.

Remark 5. It is obvious [6] that if the fundamental function $F$ of the Euclidean product $(M, F)$ of Finsler spaces $\left(M_{1}, \stackrel{\downarrow}{F}\right)$ and $\left(M_{2}, \stackrel{2}{F}\right)$ is $C^{2}$ on $(T M)_{0}$, then all $(M, F),\left(M_{1}, \stackrel{1}{F}\right)$ and $\left(M_{2}, \stackrel{2}{F}\right)$ are Riemann spaces. In case of Minkowskian product, it is not true in general. For example, when the fundamental function
 $\Psi=\sqrt{\left(q^{1}\right)^{2}+\left(q^{2}\right)^{2}}$, the fundamental function $F=\Psi\left(F,{ }_{F}^{2}\right)$ of Minkowskian product of $\left(M_{1}, F\right)$ and $\left(M_{2}, \stackrel{2}{F}\right)$ is $C^{2}$ on $(T M)_{0}$, but then $\left(g_{i j}\right)$ is degenerate on $s_{0}\left(T M_{1}\right) \times$ $\left(T M_{2}\right)_{0} \cup\left(T M_{1}\right)_{0} \times s_{0}\left(T M_{2}\right)$. (See Remark 8)

Now, we shall show that the Berwald connection of the Minkowskian product space is independent of $\psi$.

Theorem 2. The Berwald connection of the Minkowskian product of Finsler spaces $\left(M_{1}, \stackrel{1}{F}\right)$ and $\left(M_{2}, \stackrel{2}{F}\right)$ with respect to any product function $\Psi$ coincides with the Berwald connection of the Euclidean product of these Finsler spaces on $\left(T M_{1}\right)_{0} \times$ $\left(T M_{2}\right)_{0}$ provided that $\stackrel{1}{g}_{a b}\left(\right.$ resp. $\left.\stackrel{2}{g}_{\alpha \beta}\right)$ exists and is nondegenerate on $\left(T M_{1}\right)_{0}$ (resp. $\left.\left(T M_{2}\right)_{0}\right)$.

Proof. In terms of canonical local coordinates, let ${ }_{G}^{\dot{G}}{ }_{b c}^{a}$ (resp. ${\left.\underset{G}{\beta}{ }_{\gamma}^{\alpha}\right) \text { be coeffi- }}_{2}$ cients of the Berwald connection of $\left(M_{1}, \stackrel{1}{F}\right)\left(\right.$ resp. $\left.\left(M_{2}, \stackrel{2}{F}\right)\right)$. Next, in terms of the adapted canonical local coordinates, let $G_{j k}^{i}$ be coefficients of the Berwald connection $B \Gamma^{e}$ of the Euclidean product of $\left(M_{1}, \stackrel{\prime}{F}\right)$ and $\left(M_{2}, \stackrel{2}{F}\right)$. Then it is obvious that on $\left(T M_{1}\right)_{0} \times\left(T M_{2}\right)_{0}$

$$
\left\{\begin{array}{l}
G_{b c}^{a}=G_{b c}^{a}, \quad G_{\beta j}^{\alpha}=G_{\beta ;}^{\alpha},  \tag{2.6}\\
G_{\beta, j}^{a}=G_{b c}^{\alpha}=G_{\beta c}^{a}=G_{\beta c}^{\alpha}=0 .
\end{array}\right.
$$

Now, $B \Gamma^{e}$ is an $h$-connection on $M$ which satisfies three axioms (b), (c) and (d) in in Theorem 1. As for (a), the $h$-covariant derivatives with respect to $B \Gamma^{\circ}$ of the fundamental function $F=\Psi(\stackrel{1}{F}, \stackrel{2}{F})$ of the Minkowskian product space are written as

Therefore $F_{: i}=0$ are satisfied, because the $h$-covariant derivatives of the fundamental function $\stackrel{1}{F}$ (resp. $\stackrel{2}{F}$ ) with respect to the Berwald connection $\left(\dot{G}_{b c}^{a}\right)$ (resp. $\left(\stackrel{G}{G}_{\beta \gamma}^{\alpha}\right)$ ) of $\left(M_{1}, \stackrel{1}{F}\right)\left(\right.$ resp. $\left.\left(M_{2}, \stackrel{2}{F}\right)\right)$ vanish and $G_{k \alpha}^{b}=G_{k a}^{\beta}=0$. Hence, from uniqueness of the Berwald connection in Theorem 1 this connection $B \Gamma^{e}=\left(G_{j k}^{i}\right)$ is nothing but the Berwald connection of Finsler space $(M, F), M=M_{1} \times M_{2}, F=\Psi(\stackrel{1}{F}, \stackrel{2}{F})$.

A Finsler space is called a Berwald space iff its Berwald connection is linear,
or the coefficients of the Berwald connection do not depend on $y^{i}$. Concerning Berwald spaces, the following are immediately deduced from Theorem 2.

Corollary 1. [7] If both Finsler spaces $\left(M_{1}, \stackrel{1}{F}\right)$ and $\left(M_{2}, \stackrel{2}{F}\right)$ are Berwald spaces, then the Minkowskian product of these spaces with respect to any $\Psi$ is a Berwald space.

In particular we have
Corollary 2. If both $\left(M_{1}, \stackrel{\perp}{F}\right)$ and $\left(M_{2}, \stackrel{2}{F}\right)$ are Riemann spaces, then the Minkowskian product of these spaces with respect to any $\Psi$ is a Berwald space and its Berwald connection is Levi-Civita's connection on the product Riemann manifold $M_{1} \times M_{2}$.

## § 3. Geodesics of Minkowskian product space

A geodesic $x^{i}(t)$ with an affine parameter $t$ of a Finsler space $(M, F)$, or an extremal of the energy integral $\int F\left(x^{i}(t), \frac{d x^{i}}{d t}\right) d t$, satisfies Euler-Lagrange differential equations $e_{i}=0$, where

$$
e_{i}=\frac{d}{d t}\left(\frac{\partial F}{\partial y^{i}}\right)-\frac{\partial F}{\partial x_{i}}, \quad y^{i}=\frac{d x^{i}}{d t}
$$

That $t$ is an affine parameter is implied from
Lemma 1. $\frac{d F}{d t}=e_{i},{ }^{\prime}$.

$$
\text { Proof. } \quad e_{i} y^{i}=\frac{d}{d t}\left(\frac{\partial F}{\partial y^{i}} y^{i}\right)-\left(\frac{\partial F}{\partial y^{i}} \frac{d y^{i}}{d t}+\frac{\partial F}{\partial x^{i}} y^{i}\right)=\frac{d F}{d t} .
$$

Now, let $(M, F)$ be the Minkowskian product of Finsler spaces $\left(M_{1}, F\right)$ and $\left(M_{2}, F\right)$ with respect to $\Psi$. Preparing for the following theorem, stronger conditions for $\Psi$ than $(\gamma),(\delta)$ and $(\varepsilon)$ must be intrdoduced as follows;

$$
\begin{aligned}
& (s \gamma) \quad \Psi \text { is } C^{\infty} \quad \text { on } \quad\left(\boldsymbol{R}^{+} \times \boldsymbol{R}^{+}\right)_{0} \equiv \boldsymbol{R}^{+} \times \boldsymbol{R}^{+}-(0,0), \\
& (s \delta) \quad \Psi_{A} \neq 0 \quad \text { on }\left(\boldsymbol{R}^{+} \times \boldsymbol{R}^{+}\right)_{0}, \quad A=1,2, \\
& (s \varepsilon) \quad \Delta \equiv \Psi_{1} \Psi_{2}-2 \Psi \Psi_{12} \neq 0 \quad \text { on } \quad\left(\boldsymbol{R}^{+} \times \boldsymbol{R}^{+}\right)_{0} .
\end{aligned}
$$

Putting

$$
\stackrel{1}{e}_{a}=\frac{d}{d t}\left(\frac{\partial \dot{F}}{\partial y^{a}}\right)-\frac{\partial \dot{F}}{\partial x^{a}}, \quad \stackrel{2}{e}_{x}=\frac{d}{d t}\left(\frac{\partial F^{2}}{d y^{x}}\right)-\frac{\partial F^{2}}{\partial x^{\bar{z}}},
$$

we have
Lemma 2. (1) Under the condition ( $s \gamma$ ), if $e_{a}^{1}=0$ and ${\underset{e}{e}}_{2}^{2}=0$, then $e_{i}=0$. Under the conditions (s $\gamma$ ), ( $s \delta$ ) and $(s \varepsilon)$, if $e_{i}=0$, then $\stackrel{1}{e}_{a}=0$ and ${\underset{e}{e}}_{2}^{2}=0$.

Proof. From $F=\Psi(\stackrel{1}{F}, \stackrel{2}{F})$, we see

$$
\left\{\begin{array}{l}
e_{a}=\Psi_{1} e_{a}^{\prime}+\left(\Psi_{11} \frac{d \dot{F}^{\prime}}{d t}+\Psi_{12} \frac{d{ }^{2}}{d t}\right) y^{\prime}  \tag{3.1}\\
\left.e_{x}=\Psi_{2}{ }_{2}^{2} e_{x}+\left(\Psi_{21} \frac{d \dot{F}^{\prime}}{d t}+\Psi_{22} \frac{d^{2}}{d t}\right)\right)^{2}
\end{array}\right.
$$

By Lemma $1 d d^{\prime} / d t=\stackrel{1}{e_{a}}!^{\prime \prime}$ and $d{ }^{2} / d t=\stackrel{2}{e_{x}} y^{\prime x}$; therefore we have

$$
\left\{\begin{array}{l}
e_{a}=\lambda_{a}^{b} e_{b}^{1}+i_{a}^{B} e_{\beta}^{2},  \tag{3.2}\\
e_{x}=i_{\alpha}^{b} e_{b}+i_{\alpha}^{B} e_{\beta}^{2}
\end{array}\right.
$$

on $(T M)_{0}$ provided $(s \gamma)$ is satisfied (See (2.2)). Moreover, the last equations lead us to

$$
\left\{\begin{array}{l}
e_{b}^{\prime}=\omega_{b}^{c} e_{c}+\omega_{b}^{\ddot{0}} e_{\ddot{\prime}},  \tag{3.3}\\
{ }_{2}^{2}=\omega_{\beta}^{c} e_{c}+\omega_{\beta}^{\gamma} e_{\gamma}
\end{array}\right.
$$

on ( $T M)_{0}$ provided ( $s \gamma$ ), ( $s \delta$ ) and ( $s \varepsilon$ ) are satisfied (See (2.3)). (3.2) and (3.3) complete the proofs of (1) and (2) respectively.

Remark 6. $\psi=\left(\left(q^{1}\right)^{r}+\left(q^{2}\right)^{r}\right)^{\frac{1}{r}}, r=2,3, \ldots$, satisfies $(s \gamma)$, but not $(s \delta)$ and $(s \varepsilon)$. Following ' $\Psi$ 's are examples satisfying all $(s \gamma),(s \delta)$ and $(s \varepsilon)$.
(1) $\psi=q^{1}+q^{2}$ (Euclidean)
(2) $\Psi=\frac{1}{2}\left\{q^{1}+q^{2}+\left(\left(q^{1}\right)^{r}+\left(q^{2}\right)^{r}\right)^{\frac{1}{r}}\right\}, r=2,3, \ldots,[7]$
(3) $\Psi=\left(\left(q^{1}\right)^{2}+3 q^{1} q^{2}+\left(q^{2}\right)^{2}\right)^{\frac{1}{2}}$
(4) $\quad \Psi=\left(2 q^{1}+3 q^{2}\right)^{2} /\left(q^{1}+q^{2}\right)$
(5) $\quad \Psi=\left(\left(q^{1}\right)^{2}+3 q^{1} q^{2}+\left(q^{2}\right)^{2}\right) /\left(q^{1}+q^{2}\right)$

Concerning geodesics of the Minkowskian product of two Finsler spaces, we obtain from Lemma 2 the following theorem, which is a generalization of the result [8] showed in case of the Minkowskian product of two Riemann spaces.

Theorem 3. Let a Finsler space $(M, F)$ be the Minkowskian product of Finsler spaces $\left(M_{1}, \stackrel{1}{F}\right)$ and $\left(M_{2}, \stackrel{2}{F}\right)$ with respect to a product function $\Psi$.
(a) Under the condition ( $s \gamma$ ), any geodesic of $\left(M_{1}, \stackrel{1}{F}\right)\left(r e s p .\left(M_{2}, \stackrel{2}{F}\right)\right)$ is a geodesic of $(M, F)$. That is to say, the Finsler space $\left(M_{1}, \stackrel{1}{F}\right)\left(\operatorname{resp} .\left(M_{2}, \stackrel{2}{F}\right)\right)$ is a totally geodesic subspace of the Finsler space ( $M, F$ ).
(b) Under the conditions ( $s \gamma$ ), ( $s \delta$ ) and ( $(\varepsilon \varepsilon)$, the projection of any geodesic of $(M, F)$ into $M_{1}\left(\right.$ resp. $\left.M_{2}\right)$ is a geodesic of $\left(M_{1}, \stackrel{1}{F}\right)\left(\right.$ resp. $\left(M_{2}, \stackrel{2}{F}\right)$ ).

Proof (a) As $\stackrel{2}{F}$ is $C^{1}$ on $T M_{2}$ and positively homogeneous of degree 2 with respect to $y^{\alpha}$, we see $\frac{\partial F^{2}}{\partial y^{\alpha}}=0$ on $s_{0}\left(T M_{2}\right)$. Therefore, we have $e_{\alpha}^{2}=0$ if $y^{\alpha}(t)=0$ identically.

A geodesic $\left(x^{a}(t), x^{x}\right)\left(x^{x}\right.$ being constant) with an affine parameter $t$ of the subspace $\left(M_{1}, \stackrel{1}{F}\right)$ satisfies $\stackrel{1}{e_{a}}=0$, and $\frac{d x^{x}}{d t}=0$ imply $\stackrel{2}{e}_{\alpha}^{2}=0$ by the above notice. Hence, by Lemma $2(1)$ we obtain $e_{i}=0$. Thus the curve $\left(x^{a}(t), x^{x}\right)$ is a geodesic with an affine parameter $t$ of $(M, F)$. As to geodesic of $\left(M_{2}, \stackrel{2}{F}\right)$, it is proved similarly.
(b) We consider a geodesic $\left(x^{a}(t), x^{x}(t)\right)$ with an affine parameter of $(M, F)$. By Lemma 2 (2) $e_{i}=0$ imply $\mathfrak{c}_{a}^{\prime}=0$. Hence, the curve $\left(x^{a}(t), x^{\alpha}\right)\left(x^{\alpha}\right.$ being constant) projected into $M_{1}$ is a geodesic with an affine parameter $t$ of $\left(M_{1}, \stackrel{1}{F}\right)$. The proof as to the curve projected into $M_{2}$ is done similarly.

Remark 7. If $\hat{F}$ is $C^{2}$ on $\left(T M_{A}\right)_{0}, A=1,2$, then we have

$$
\left\{\begin{array}{lll}
e_{a}^{\prime}=\stackrel{1}{g}_{a b} \frac{d y^{b}}{d t}+2 G_{a}^{\prime} & \text { on } \quad\left(T M_{1}\right)_{0}  \tag{3.4}\\
\stackrel{2}{e}_{\alpha}=\stackrel{2}{g_{\alpha \beta}} \frac{d y^{\beta}}{d t}+2 G_{\alpha}^{2} & \text { on } \quad\left(T M_{2}\right)_{0}
\end{array}\right.
$$

where

$$
2 \stackrel{1}{G}_{a}=\frac{\partial^{2} \stackrel{1}{F}}{\partial x^{b} \partial y^{a}} y^{b}-\frac{\partial \stackrel{1}{F}}{\partial x^{a}}, \quad 2 \stackrel{2}{G}_{\alpha}^{2}=\frac{\partial^{2} \stackrel{2}{F}}{\partial x^{\beta} \partial y^{, x}} y^{\beta}-\frac{\partial F^{2}}{\partial x^{\alpha}} .
$$

The equations (3.4) together with (3.2) and (2.1) lead us to

$$
\left\{\begin{array}{l}
G_{a}=\lambda_{a}^{b} \dot{G}_{b}+\lambda_{a}^{\beta} \stackrel{2}{G}_{\beta},  \tag{3.5}\\
G_{x}=\lambda_{\alpha}^{b} G_{b}+\lambda_{\alpha}^{\beta} \mathbf{G}_{\beta},
\end{array}\right.
$$

where

$$
2 G_{i}=\frac{\partial^{2} F}{\partial x^{j} \partial y^{i}} y^{j}-\frac{\partial F}{\partial x^{i}} .
$$

Moreover, if $\left(\stackrel{1}{g}_{a b}\right)$ (resp. $\left.\left(\stackrel{2}{g}_{\alpha \beta}\right)\right)$ is nondegenerate on $\left(T M_{1}\right)_{0}$ (resp. $\left.\left(T M_{2}\right)_{0}\right)$, then from (3.5) and (2.5) we obtain

$$
\begin{equation*}
G^{a}=\dot{G}^{a}, \quad G^{\alpha}=\stackrel{2}{G^{\alpha}} \tag{3.6}
\end{equation*}
$$

on $\left(T M_{1}\right)_{0} \times\left(T M_{2}\right)_{0}$, where

$$
G^{i}=g^{i j} G_{j}, \quad \dot{G}^{a}=\stackrel{1}{g^{a} b} \dot{1}_{b}, \quad \stackrel{2}{G}^{\alpha}=\stackrel{2}{g^{\alpha \beta}} \stackrel{2}{G}_{\beta} .
$$

This is another proof of Theorem 2.
Remark 8. If $\psi$ satisfies $(s \gamma),(s \delta)$ and $F=\Psi(\stackrel{1}{F}, \stackrel{2}{F})$ is $C^{2}$ on $(T M)_{0}$, then both
$\left(M_{1}, \stackrel{1}{F}\right)$ and $\left(M_{2}, \stackrel{2}{F}\right)$ are Riemann spaces. Hence, by Corollary 2 of Theorem 2 the Minkowskian product $(M, F)$ of these spaces is Berwald space with Levi-Civita's connection.

$$
\text { Because }{\stackrel{1}{y^{\prime}}}_{a} \equiv \frac{\partial \stackrel{1}{F}}{\partial y^{a}}=\frac{1}{\Psi_{1}} \frac{\partial F}{\partial y^{a}}\left(\text { resp. } \stackrel{2}{1}_{x} \equiv \frac{\partial \stackrel{2}{F}}{\partial y^{x}}=\frac{1}{\Psi_{2}} \frac{\partial F}{\partial y^{x}}\right)
$$

are $C^{1}$ on $(T M)_{0}$ under the above conditions. Therefore, $\stackrel{1}{y}_{a}$ (resp. $\stackrel{2}{y}_{\alpha}$ ) are $C^{1}$ on $T M_{1}\left(\operatorname{resp} . T M_{2}\right)$. Thus $\stackrel{1}{F}(\operatorname{resp} . \stackrel{2}{F})$ is $C^{2}$ on $T M_{1}\left(\operatorname{resp} . T M_{2}\right)$.

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