# A fixed point formula for compact almost complex manifolds 

By

Kenji Tsuboi


#### Abstract

In this paper, using the group structures of the spheres $S^{1}, S^{3}$ and the results of Atiyah-Patodi-Singer, Donnelly and Morita, we introduce a fixed point formula for periodic automorphisms of compact almost complex manifolds. Our main result is Theorem 1.3. The theorem is refined for a certain case if the almost complex manifold admits an EinsteinKähler metric.


## 1. Introduction and Main Theorem

Let $M$ be a compact $2 m$-dimensional almost complex manifold with the almost complex structure $J$ and $P \rightarrow M$ the associated principal $G L(m ; \mathbb{C})$ bundle of $M$. We call a diffeomorphism $\psi: M \longrightarrow M$ an automorphism of $M$ if $\psi$ commutes with $J$ and denote the topological group consisting of all automorphisms of $M$ by $A(M)$. The group $A(M)$ naturally acts on $P$ on the left.

Definition 1.1. Let $S(n)$ be the set of symmetric homogeneous polynomials in $x_{1}, x_{2}, \ldots, x_{m}$ of order $n$ with integral coefficients. Let

$$
\phi=\phi\left(\tau_{1}, \tau_{2}, \ldots, \tau_{m}\right)
$$

be any element of $S(n)$ where $\tau_{j}=\sigma_{j}\left(x_{1}, x_{2}, \ldots, x_{m}\right)$ is the $j$-th elmentary symmetric polynomial in $\left\{x_{i}\right\}_{i=1}^{m}$, whose degree is equal to $j$. Let $V_{\phi}$ be the element of the representation ring $R(G L(m ; \mathbb{C}))$ of $G L(m ; \mathbb{C})$ defined by

$$
\begin{aligned}
V_{\phi}=\phi\left(\tau_{1}, \tau_{2}, \ldots, \tau_{m}\right) \in & R(G L(m ; \mathbb{C})) \\
& \subset R\left(T^{m}\right)=\mathbb{Z}\left[t_{1}, t_{1}^{-1}, t_{2}, t_{2}^{-1}, \ldots, t_{m}, t_{m}^{-1}\right],
\end{aligned}
$$

where $T^{m}$ is the maximal torus of $G L(m ; \mathbb{C}), t_{i}: T^{m} \rightarrow S^{1}$ is the $i$-th factor projection and $\tau_{j}=\sigma_{j}\left(t_{1}-1, t_{2}-1, \ldots, t_{m}-1\right)$ is the $j$-th elementary symmetric polynomial in $t_{1}-1, t_{2}-1, \ldots, t_{m}-1$. Note that $\sigma_{j}=\sigma_{j}\left(t_{1}, t_{2}, \ldots, t_{m}\right)$ is isomorphic to the $G L(m ; \mathbb{C})$-representation $\wedge^{j} \mathbb{C}^{m}$. Hence, setting

$$
\hat{\phi}\left(\sigma_{1}, \sigma_{2}, \ldots, \sigma_{m}\right)=\phi\left(\tau_{1}, \tau_{2}, \ldots, \tau_{m}\right),
$$

we have

$$
V_{\phi}=\hat{\phi}\left(\wedge^{1} \mathbb{C}^{m}, \wedge^{2} \mathbb{C}^{m}, \ldots, \wedge^{m} \mathbb{C}^{m}\right) \in R(G L(m ; \mathbb{C}))
$$

Using this virtual $G L(m ; \mathbb{C})$-representation $V_{\phi}$, we can define a virtual complex vector bundle $E_{\phi}$ on $M$ by

$$
\begin{equation*}
E_{\phi}=P \times_{G L(m ; \mathbb{C})} V_{\phi}=\hat{\phi}\left(\wedge^{1} T M, \wedge^{2} T M, \ldots, \wedge^{m} T M\right) \in K(M) \tag{1.1}
\end{equation*}
$$

where $T M$ is the tangent bundle of $M$ and $K(M)$ is the $K$-group of $M$. Then the action of $A(M)$ on $P$ naturally defines the action of $A(M)$ on $E_{\phi}$ and $E_{\phi}$ is a virtual complex $A(M)$-vector bundle.

Definition 1.2. Let $a$ be any periodic element of $A(M), G$ the cyclic subgroup of $A(M)$ generated by $a$ and $\Omega$ the fixed point set of $a$ consisting of compact connected submanifolds $N$ of $M$. Then the restriction of $J$ defines an almost complex structure of $N$ and the Todd class $\operatorname{Td}(T N)$ of $T N$ is defined by

$$
\operatorname{Td}(T N)=\prod_{k=1}^{d} \frac{x_{k}}{1-e^{-x_{k}}} \in H^{*}(N ; \mathbb{C})
$$

where $2 d$ is the dimension of $N$ and $\prod_{k=1}^{d}\left(1+x_{k}\right)$ equals to the total Chern class of $T N$. Note that $\operatorname{Td}(T N)=1$ if $N$ is a point. On the other hand, a complex $G$-vector bundle $E$ over $N$ is decomposed into the direct sum of subbundles

$$
E=E_{1} \oplus E_{2} \oplus \cdots \oplus E_{s},
$$

where $a$ acts on the subbundle $E_{j}$ via multiplication by $e^{\sqrt{-1} \theta_{j}}$. Then we can define the characteristic class $\operatorname{Ch}(E, a)$ by

$$
\operatorname{Ch}(E, a)=\sum_{j=1}^{s} e^{\sqrt{-1} \theta_{j}} \operatorname{Ch}\left(E_{j}\right) \in H^{*}(N ; \mathbb{C})
$$

where $\operatorname{Ch}\left(E_{j}\right)$ is the Chern character of $E_{j}$. This definition is extended to the case of virtual vector bundles by

$$
\operatorname{Ch}(E-F, a)=\operatorname{Ch}(E, a)-\operatorname{Ch}(F, a) \in H^{*}(N ; \mathbb{C})
$$

and $\mathrm{Ch}(*, a)$ defines a ring homomorphism

$$
\mathrm{Ch}(*, a): K(N) \longrightarrow H^{*}(N ; \mathbb{C})
$$

namely, satisfies the following equalities:

$$
\begin{equation*}
\operatorname{Ch}(E \pm F, a)=\operatorname{Ch}(E, a) \pm \operatorname{Ch}(F, a), \quad \operatorname{Ch}(E \otimes F, a)=\operatorname{Ch}(E, a) \operatorname{Ch}(F, a) . \tag{1.2}
\end{equation*}
$$

We can also define the characteristic class $\mathfrak{U}(E, a)$ by

$$
\mathfrak{U}(E, a)=\prod_{j=1}^{s} \prod_{k=1}^{r_{j}} \frac{1}{1-e^{-x_{k}-\sqrt{-1} \theta_{j}}} \in H^{*}(N ; \mathbb{C})
$$

where $r_{j}=\operatorname{rank}\left(E_{j}\right)$ and $\prod_{k=1}^{r_{j}}\left(1+x_{k}\right)$ equals to the total Chern class of $E_{j}$.

Our main result is the following theorem.
Theorem 1.3. Let $\ell$ be 0,1 or 2 and $\phi$ any element of $S(n)$. Let $\psi$ be any periodic element of $A(M)$ and assume that the order of $\psi$ is $p$. Let $\gamma$ be any natural number which is prime to $p$. Let $\Omega(k)$ be the fixed point set of $\psi^{k}$ $(1 \leqq k \leqq p-1)$ consisting of compact connected almost complex manifolds $N$, $\nu(N, M)$ the normal bundle of $N$ in $M$ and $[N]$ the fundamental cycle of $N$. Then the equality

$$
\sum_{k=1}^{p-1} C_{\ell}(k, \gamma) \sum_{N \subset \Omega(k)} \operatorname{Ch}\left(\left.E_{\phi}\right|_{N}, \psi^{k}\right) \operatorname{Td}(T N) \mathfrak{U}\left(\nu(N, M), \psi^{k}\right)[N] \equiv 0 \quad(\bmod p)
$$

holds for any $n>m+\ell$, where
$C_{0}(k, \gamma)=1, \quad C_{1}(k, \gamma)=\frac{1}{1-e^{-2 \pi \sqrt{-1} \gamma k / p}}, \quad C_{2}(k, \gamma)=\frac{1}{\mid 1-e^{-2 \pi \sqrt{-1} \gamma k /\left.p\right|^{2}}}$.
Let $N$ be a connected component of the fixed point set of the action of a periodic automorphism $a$ of $M$. Assume that the restriction of the tangent bundle $T M$ to $N$ splits into the direct sum of complex line bundles

$$
\left.T M\right|_{N}=L_{1} \oplus \cdots \oplus L_{m}
$$

where $a$ acts on $L_{j}$ via multiplication by $e^{\sqrt{-1} \theta_{j}}$. Let $\sigma_{j}$ be the $j$-th elementary symmetric polynomial in $\left\{e^{\sqrt{-1}} \theta_{j} e^{c_{1}\left(L_{j}\right)}\right\}_{j=1}^{m}$ and $\tau_{j}$ the $j$-th elementary symmetric polynomial in $\left\{e^{\sqrt{-1} \theta_{j}} e^{c_{1}\left(L_{j}\right)}-1\right\}_{j=1}^{m}$. Then since

$$
\operatorname{Ch}\left(L_{j}, a\right)=e^{\sqrt{-1} \theta_{j}} e^{c_{1}\left(L_{j}\right)}
$$

it follows from (1.1) and (1.2) that

$$
\begin{equation*}
\operatorname{Ch}\left(\left.E_{\phi}\right|_{N}, a\right)=\hat{\phi}\left(\sigma_{1}, \sigma_{2}, \ldots, \sigma_{m}\right)=\phi\left(\tau_{1}, \tau_{2}, \ldots, \tau_{m}\right) . \tag{1.3}
\end{equation*}
$$

The next corollary is deduced from Theorem 1.3 and (1.3).
Corollary 1.4. Assume that $\Omega(k)$ in Theorem 1.3 consists of points $\left\{q_{s}\right\}_{s=1}^{N(k)}$ for any $k$. Then the automorphism $\psi^{k}$ acts on the tangent space $T_{q_{s}} M$ via multiplication by some periodic diagonal unitary matrix, which we assume is the diagonal matrix with diagonal entries $\left\{e^{2 \pi \sqrt{-1} h_{j s}^{k} / p}\right\}_{j=1}^{m}\left(h_{j s}^{k} \in \mathbb{Z}\right)$. Let $\tau_{j}$ be the $j$-th elementary symmetric polynomial in $\left\{e^{2 \pi \sqrt{-1} h_{j s}^{k} / p}-1\right\}_{j=1}^{m}$. Then under the notation in Theorem 1.3, the equality

$$
\sum_{k=1}^{p-1} C_{\ell}(k, \gamma) \sum_{s=1}^{N(k)} \phi\left(\tau_{1}, \tau_{2}, \ldots, \tau_{m}\right) \prod_{j=1}^{m} \frac{1}{1-e^{-2 \pi \sqrt{-1} h_{j s}^{k} / p}} \equiv 0 \quad(\bmod p)
$$

holds for any $n>m+\ell$.

Proof. For any $q_{s} \in \Omega(k)$, the tangent space $T_{q_{s}} M$ splits into the direct sum of $m$-copies of $\mathbb{C}^{1}$

$$
T_{q_{s}} M=\mathbb{C}_{1}^{1} \oplus \mathbb{C}_{2}^{1} \oplus \cdots \oplus \mathbb{C}_{m}^{1}
$$

where $\psi^{k}$ acts on $\mathbb{C}_{j}^{1}$ via multiplication by $e^{2 \pi \sqrt{-1}} h_{j s}^{k} / p$. Hence it follows from (1.3) that

$$
\operatorname{Ch}\left(\left.E_{\phi}\right|_{q_{i}}, \psi^{k}\right)=\phi\left(\tau_{1}, \tau_{2}, \ldots, \tau_{m}\right),
$$

where $\tau_{j}$ is the $j$-th elementary symmetric polynomial in $\left\{e^{2 \pi \sqrt{-1} h_{j s}^{k} / p}-1\right\}_{j=1}^{m}$. Moreover, since $\operatorname{Td}\left(T q_{s}\right)=1$ for any $s$, the equality in Corollary 1.4 immediately follows from the equality in Theorem 1.3.

Remark 1.5. As we will see in Remarks 3.3 and 4.2 , the equality in Theorem 1.3 does not hold in general if $n=m+\ell$.

Remark 1.6. The author does not know whether the equality in Theorem 1.3 holds for $\ell \geqq 3$ by introducing some appropriate $C_{\ell}(k, \gamma)$.

## 2. Proof of the Theorem

In this section we give the proof of Theorem 1.3. Let $G$ be the cyclic subgroup of $A(M)$ generated by $\psi$. We give a $G$-invariant Hermitian metric on $M$ and let $Q \longrightarrow M$ be the subbundle of $P$ consisting of unitary frames with respect to the metric. Let $\nabla$ be a $G$-invariant connection in $Q$. Then since $V_{\phi}$ is considered as a virtual representation of $U(m)$ and $E_{\phi}$ equals to $Q \times{ }_{U(m)} V_{\phi}$, the natural $U(m)$-invariant inner product in $V_{\phi}$ defines a $G$-invariant inner product in $E_{\phi}$ and $\nabla$ defines a unitary connection of $E_{\phi}$. The connection $\nabla$ also defines a $G$-invariant connection of the half spinor bundles $S^{ \pm}=Q \times_{U(m)} \Delta^{ \pm}$over $M$ where $\Delta^{ \pm}$are the half spin representations of $\operatorname{spin}^{c}(2 m)$. (For details of spinor bundles and $\operatorname{spin}^{c}$-Dirac operators, see [6].) Using the connections defined above, we can define the $G$-equivariant $\operatorname{spin}^{c}$-Dirac (Dolbeault) operator

$$
D: \Gamma\left(S^{+} \otimes E_{\phi}\right) \longrightarrow \Gamma\left(S^{-} \otimes E_{\phi}\right)
$$

and it follows from the Riemann-Roch theorem (see (4.3) in [2]) that
(2.1) $\operatorname{Index}(D):=\operatorname{dim} \operatorname{ker}(D)-\operatorname{dim} \operatorname{coker}(D)=\int_{M} \operatorname{Ch}\left(E_{\phi}, \nabla\right) \operatorname{Td}(T M, \nabla)$,
where $\operatorname{Ch}\left(E_{\phi}, \nabla\right)$ is the Chern character form of $E_{\phi}$ with respect to $\nabla, \operatorname{Td}(T M$, $\nabla)$ is the Todd form of $T M$ with respect to $\nabla$. Here for any $1 \leqq j \leqq m$, we can see that

$$
\operatorname{Ch}\left(\wedge^{j} T M, \nabla\right)=\sigma_{j}\left(e^{x_{1}}, e^{x_{2}}, \ldots, e^{x_{m}}\right)
$$

where by definition the $j$-th Chern form $c_{j}(T M, \nabla)$ is the $j$-th elementary symmetric polynomial in $x_{1}, x_{2}, \ldots, x_{m}$. Hence it follows from (1.1) and (1.2) that

$$
\operatorname{Ch}\left(E_{\phi}, \nabla\right)=\hat{\phi}\left(\sigma_{1}, \sigma_{2}, \ldots, \sigma_{m}\right)=\phi\left(\tau_{1}, \tau_{2}, \ldots, \tau_{m}\right)
$$

where $\tau_{j}$ is the $j$-th elementary symmetric polynomial in $e^{x_{1}}-1, e^{x_{2}}-1, \ldots$, $e^{x_{m}}-1$ for $1 \leqq j \leqq m$. Since

$$
\begin{align*}
\tau_{j} & =\sigma_{j}\left(e^{x_{1}}-1, e^{x_{2}}-1, \ldots, e^{x_{m}}-1\right)  \tag{2.2}\\
& =\sigma_{j}\left(x_{1}, x_{2}, \ldots, x_{m}\right)+\text { higher order terms } \\
& =c_{j}(T M, \nabla)+\text { higher order terms }
\end{align*}
$$

we have
$\operatorname{Ch}\left(E_{\phi}, \nabla\right)=\phi\left(c_{1}(T M, \nabla), c_{2}(T M, \nabla), \ldots, c_{m}(T M, \nabla)\right)+$ higher order terms and therefore it follows that

$$
\begin{equation*}
\int_{M} \operatorname{Ch}\left(E_{\phi}, \nabla\right) \operatorname{Td}(T M, \nabla)=0 \tag{2.3}
\end{equation*}
$$

because the order of $\phi$ is greater than $m$ and the dimension of $M$ is $2 m$. On the other hand, it follows from (4.6) in [2] that
(2.4) $\operatorname{Index}\left(D, \psi^{k}\right):=\operatorname{Tr}\left(\left.\psi^{k}\right|_{\operatorname{ker}(D)}\right)-\operatorname{Tr}\left(\left.\psi^{k}\right|_{\operatorname{coker}(D)}\right)$

$$
=\sum_{N \subset \Omega(k)} \operatorname{Ch}\left(\left.E_{\phi}\right|_{N}, \psi^{k}\right) \operatorname{Td}(T N) \mathfrak{U}\left(\nu(N, M), \psi^{k}\right)[N]
$$

for $1 \leqq k \leqq p-1$. Now let $V$ be any finite dimensional complex $G$-module and $\beta$ an eigenvalue of $\left.\psi\right|_{V}$. Then since $\beta^{p}=1$, it follows that

$$
\sum_{k=1}^{p} \beta^{k} \equiv 0 \quad(\bmod p)
$$

and hence it follows that

$$
\begin{equation*}
\sum_{k=1}^{p} \operatorname{Tr}\left(\left.\psi^{k}\right|_{V}\right) \equiv 0 \quad(\bmod p) \tag{2.5}
\end{equation*}
$$

Therefore we have

$$
\sum_{k=1}^{p} \operatorname{Index}\left(D, \psi^{k}\right) \equiv 0 \quad(\bmod p)
$$

and hence it follows from (2.1), (2.3) and (2.4) that

$$
\begin{aligned}
& \sum_{k=1}^{p-1} \sum_{N \subset \Omega(k)} \operatorname{Ch}\left(\left.E_{\phi}\right|_{N}, \psi^{k}\right) \operatorname{Td}(T N) \mathfrak{U}\left(\nu(N, M), \psi^{k}\right)[N] \\
& =\sum_{k=1}^{p-1} \operatorname{Index}\left(D, \psi^{k}\right)=\sum_{k=1}^{p} \operatorname{Index}\left(D, \psi^{k}\right) \equiv 0 \quad(\bmod p)
\end{aligned}
$$

because $\operatorname{Index}\left(D, \psi^{p}\right)=\operatorname{Index}(D)=0$. This completes the proof of the equality in Theorem 1.3 for $\ell=0$.

Now assume that $\ell=1$ or 2 and let $D^{2 \ell}$ and $\partial D^{2 \ell}=S^{2 \ell-1}$ be the unit disk and the unit sphere in $\mathbb{C}^{\ell}$ respectively. Let $\mathbb{H}$ be the set of quaternions, which is identified with $\mathbb{C}^{2}$ as follows:

$$
\mathbb{H} \ni a+b i+c j+d k=(a+b i)+(c+d i) j \longleftrightarrow(a+b i, c+d i) \in \mathbb{C}^{2} .
$$

Then $\mathbb{C}^{1}$ is contained in $\mathbb{H}$ by $a+b i=a+b i+0 j+0 k$. Let $\alpha:=e^{2 \pi \sqrt{-1} / p}$ be the primitive $p$-th root of 1 . Then $G$ acts on $\mathbb{H}$ by $\psi \cdot h=h \alpha^{\gamma}(h \in \mathbb{H})$, which corresponds to the $S U(2)$-transformation

$$
\mathbb{C}^{2} \ni\left(z_{1}, z_{2}\right) \longrightarrow\left(\alpha^{\gamma} z_{1}, \bar{\alpha}^{\gamma} z_{2}\right) \in \mathbb{C}^{2}
$$

under the identification above because $j \alpha^{\gamma}=\bar{\alpha}^{\gamma} j$. This $G$-action defines $G$ actions on $D^{2 \ell}, S^{2 \ell-1}$ for $\ell=1,2$. We give the standard metric on $S^{2 \ell-1}$, which is $G$-invariant, and give a $G$-invariant Hermitian metric on $D^{2 \ell}$ such that it is a product metric of $S^{2 \ell-1} \times[0, \delta)$ near $\partial D^{2 \ell}=S^{2 \ell-1}$. Here since $\ell$ equals to 1 or 2 , the sphere $S^{2 \ell-1}$ has a group structure. Actually the group structure of $S^{3}$ is induced from the multiplication in the quaternions $\mathbb{H}$ and $S^{1}$ is the subgroup of $S^{3}$ consisting of complex numbers. Using this group structure, we can construct a global orthonormal frame field $\left\{F_{A}^{1}, F_{A}^{2}, F_{A}^{3}\right\}_{A \in S^{3}}$ on $S^{3}$ as follows:

$$
F_{A}^{1}=i \cdot A, F_{A}^{2}=j \cdot A, F_{A}^{3}=k \cdot A \in \mathbb{H} .
$$

It is clear that $\left\{F_{A}^{1}\right\}_{A \in S^{1}}$ defines a global orthonormal frame field on $S^{1}$. Now considering the associativity of the multiplication in $\mathbb{H}$, we can see that the frame field above is invariant under the action of $G$. Hence the trivialization of the tangent bundle $T S^{3}$ :

$$
T S^{3} \ni\left(A, w=a F_{A}^{1}+b F_{A}^{2}+c F_{A}^{3}\right) \longrightarrow(A,(a, b, c)) \in S^{3} \times \mathbb{R}^{3}
$$

$\left(A \in S^{3}, w \in T_{A} S^{3}\right)$ is $G$-invariant and therefore the unique trivial spin ${ }^{c}$ structure of $S^{2 \ell-1}$ is $G$-invariant. Moreover $F_{A}^{0}:=A$ defines the outward unit normal vector field on $S^{2 \ell-1}$ and the trivialization of $\left.T D^{4}\right|_{S^{3}}$ :

$$
\begin{aligned}
\left.T D^{4}\right|_{S^{3}} & \ni\left(A, v=a F_{A}^{0}+b F_{A}^{1}+c F_{A}^{2}+d F_{A}^{3}\right) \\
& \longrightarrow(A,((a+b i),(c+d i))) \in S^{3} \times \mathbb{C}^{2}
\end{aligned}
$$

$\left(A \in S^{3}, v \in T_{A} D^{4}\right)$ is $G$-invariant. Therefore the quotient $\left(T S^{2 \ell-1}\right) / G$ is the trivial real vector bundle and the quotient $\left(\left.T D^{2 \ell}\right|_{S^{2 \ell-1}}\right) / G$ is the trivial complex vector bundle.

Set $X=M \times D^{2 \ell}$ and $Y=\partial X=M \times S^{2 \ell-1}$. Then the metric on $M$ and the metrics on $D^{2 \ell}, S^{2 \ell-1}$ define the $G$-invariant product metrics on $X, Y$ respectively and the $G$-actions on $D^{2 \ell}, S^{2 \ell-1}$ define the diagonal $G$-actions on $X, Y$ as follows:

$$
\begin{equation*}
\psi \cdot(q, h)=\left(\psi \cdot q, h \alpha^{\gamma}\right) \quad(q \in M, h \in \mathbb{H}) . \tag{2.6}
\end{equation*}
$$

Moreover the tangent bundle $T X, T Y$ splits as

$$
\begin{aligned}
& T X=q_{X}^{*} T M \oplus r_{X}^{*} T D^{2 \ell}=q_{X}^{*} T M \oplus \varepsilon_{\mathbb{C}}^{\ell} \\
& T Y=q_{Y}^{*} T M \oplus r_{Y}^{*} T S^{2 \ell-1}=q_{Y}^{*} T M \oplus \varepsilon^{2 \ell-1},
\end{aligned}
$$

where $q_{X}: X \longrightarrow M, q_{Y}: Y \longrightarrow M$ denote the first factor projections, $r_{X}: X \longrightarrow D^{2 \ell}, r_{Y}: Y \longrightarrow S^{2 \ell-1}$ denote the second factor projections and $\varepsilon_{\mathbb{C}}^{k}\left(\varepsilon^{k}\right)$ denotes the trivial complex (real) vector bundle of rank $k$ with a $G$-invariant trivialization. Therefore $\operatorname{spin}^{c}$-structures on $X, Y$ are defined by the $U(m)$-structures $q_{X}^{*} Q, q_{Y}^{*} Q$ respectively and connections $\nabla^{X}, \nabla^{Y}$ in $q_{X}^{*} Q, q_{Y}^{*} Q$ are induced from the connection $\nabla$ in $Q$. These connections $\nabla^{X}, \nabla^{Y}$ define $G$-invariant metric connections of $T X, T Y$, which are the direct sum of the connection $\nabla$ of $T M$ and the globally flat connections of the trivial bundles. These connections $\nabla^{X}, \nabla^{Y}$ also define $G$-invariant connections of the half spinor bundles $S_{X}^{ \pm}=q_{X}^{*} Q \times_{U(m)} \Delta^{ \pm}$over $X$ and a $G$-invariant connection of the spinor bundle $S_{Y}=\left.S_{X}^{+}\right|_{Y}=\left.S_{X}^{-}\right|_{Y}=q_{Y}^{*} Q \times_{U(m)} \triangle$ over $Y$ where $\triangle^{ \pm}$are the half spin representations of $\operatorname{spin}^{c}(2 m+2 \ell)$ and $\Delta$ is the spin representation of $\operatorname{spin}^{c}(2 m+2 \ell-1)$.

Set $E_{\phi, X}=q_{X}^{*} E_{\phi}=q_{X}^{*} Q \times_{U(m)} V_{\phi}$ and $E_{\phi, Y}=q_{Y}^{*} E_{\phi}=q_{Y}^{*} Q \times_{U(m)} V_{\phi}$. Then $E_{\phi, X}$ and $E_{\phi, Y}$ are virtual $G$-vector bundles with $G$-invariant unitary connections $\nabla^{X}, \nabla^{Y}$ and the restriction of $E_{\phi, X}$ to $Y$ coincides with $E_{\phi, Y}$. Using the spin ${ }^{c}$-structures and the connections defined above, we can define the $G$-equivariant $\operatorname{spin}^{c}$-Dirac operators

$$
\begin{aligned}
D_{X} & : \Gamma\left(S_{X}^{+} \otimes E_{\phi, X}\right) \longrightarrow \Gamma\left(S_{X}^{-} \otimes E_{\phi, X}\right), \\
D_{Y} & : \Gamma\left(S_{Y} \otimes E_{\phi, Y}\right) \longrightarrow \Gamma\left(S_{Y} \otimes E_{\phi, Y}\right) .
\end{aligned}
$$

Since the metric and the connection $\nabla^{X}$ is product near $\partial X=Y, D_{X}$ can be expressed as

$$
D_{X}=\sigma\left(\frac{\partial}{\partial u}+D_{Y}\right)
$$

on the collar $Y \times[0, \delta) \subset X$ where $u$ is the coordinate of $[0, \delta)$ and $\sigma$ is a bundle isomorphism defined by the Clifford multiplication (see [1]). Hence the following equality is deduced from (4.3) in [1] (see also (4.6) in [2] and Lemma 3.5 .4 in [6]):

$$
\begin{equation*}
\operatorname{Index}\left(D_{X}\right)=\int_{X} \operatorname{Ch}\left(E_{\phi, X}, \nabla^{X}\right) \operatorname{Td}\left(T X, \nabla^{X}\right)-\frac{1}{2}\left(\eta_{Y}+\operatorname{dim} \operatorname{ker} D_{Y}\right), \tag{2.7}
\end{equation*}
$$

where $\operatorname{Index}\left(D_{X}\right)$ is the index of $D_{X}$ with a certain global boundary condition, which is an integer, $\operatorname{Ch}\left(E_{\phi, X}, \nabla^{X}\right)$ is the Chern character form of $E_{\phi, X}$ with respect to $\nabla^{X}, \operatorname{Td}\left(T X, \nabla^{X}\right)$ is the Todd form of $T X$ with respect to $\nabla^{X}$ and $\eta_{Y}$ is the eta invariant of $D_{Y}$. (For details of eta invariants, see [1], [3].) Here the same argument as was used to prove (2.3) shows that

$$
\begin{equation*}
\int_{X} \operatorname{Ch}\left(E_{\phi, X}, \nabla^{X}\right) \operatorname{Td}\left(T X, \nabla^{X}\right)=0 \tag{2.8}
\end{equation*}
$$

because the order of $\phi$ is greater than $m+\ell$ and the dimension of $X$ is $2 m+2 \ell$. Therefore it follows from (2.7) that

$$
\begin{equation*}
\frac{1}{2} \eta_{Y}=-\operatorname{Index}\left(D_{X}\right)-\frac{1}{2} \operatorname{dim} \operatorname{ker} D_{Y} \tag{2.9}
\end{equation*}
$$

Let $O$ be the origin of $\mathbb{C}^{\ell}$. Then $M$ is regarded as an almost complex submanifold of $X$ by the identification of $M$ with $M \times\{O\}$ and hence $N$ is also regarded as an almost complex submanifold of $X$. Note that the fixed point set of the $G$-action on $X$ is contained in $M$ and coincides with the fixed point set of the $G$-action on $M$. Let $\nu(N, X)$ be the normal bundle of $N$ in $X$. Then $\nu(N, X)$ is decomposed into the direct sum of complex subbundles

$$
\nu(N, X)=\nu(N, M) \oplus \varepsilon_{\mathbb{C}}^{\ell}=\oplus_{j} \nu_{j}(N, M) \oplus \varepsilon_{\mathbb{C}}^{\ell}
$$

where $\psi^{k}$ acts on $\nu_{j}(N, M)$ via multiplication by $e^{\sqrt{-1} \theta_{j}}$ and acts on the trivial complex line bundle $\varepsilon_{\mathbb{C}}^{\ell}=N \times \mathbb{C}^{\ell}$ by

$$
\psi^{k} \cdot\left(q,\left(z_{1}, \ldots, z_{\ell}\right)\right)= \begin{cases}\left(q,\left(\alpha^{\gamma k} z_{1}\right)\right) & (\ell=1) \\ \left(q,\left(\alpha^{\gamma k} z_{1}, \bar{\alpha}^{\gamma k} z_{2}\right)\right) & (\ell=2)\end{cases}
$$

$\left(q \in N,\left(z_{1}, \ldots, z_{\ell}\right) \in \mathbb{C}^{\ell}\right)$. Hence the following equality is deduced from Theorem 1.2 in [3] (see also (4.6) in [2] and Lemma 3.5.4 in [6]):

$$
\begin{align*}
& \operatorname{Index}\left(D_{X}, \psi^{k}\right)  \tag{2.10}\\
& =\sum_{N \subset \Omega(k)} \operatorname{Ch}\left(\left.E_{\phi}\right|_{N}, \psi^{k}\right) \operatorname{Td}(T N) \mathfrak{U}\left(\nu(N, M), \psi^{k}\right) C_{\ell}(k, \gamma)[N] \\
& \quad-\frac{1}{2}\left\{\eta_{Y}\left(\psi^{k}\right)+\operatorname{Tr}\left(\left.\psi^{k}\right|_{\operatorname{ker} D_{Y}}\right)\right\}
\end{align*}
$$

for $1 \leqq k \leqq p-1$, where $\operatorname{Index}\left(D_{X}, \psi^{k}\right)$ is the index of $D_{X}$ with a certain global boundary condition evaluated at $\psi^{k}$, namely,

$$
\operatorname{Index}\left(D_{X}, \psi^{k}\right):=\operatorname{Tr}\left(\left.\psi^{k}\right|_{\text {ker } D_{X}}\right)-\operatorname{Tr}\left(\left.\psi^{k}\right|_{\text {coker } D_{X}}\right),
$$

$\eta_{Y}\left(\psi^{k}\right)$ is the eta invariant of $D_{Y}$ evaluated at $\psi^{k}$ and

$$
C_{1}(k, \gamma)=\frac{1}{1-\alpha^{-\gamma k}}, \quad C_{2}(k, \gamma)=\frac{1}{1-\alpha^{-\gamma k}} \frac{1}{1-\bar{\alpha}^{-\gamma k}}=\frac{1}{\left|1-\alpha^{-\gamma k}\right|^{2}}
$$

Note that $\operatorname{Index}\left(D_{X}, \psi^{p}\right), \eta_{Y}\left(\psi^{p}\right)$ coincide with $\operatorname{Index}\left(D_{X}\right), \eta_{Y}$ in (2.7) respectively.

Since the restriction of the $G$-action to $Y$ is free and preserves the metric and the $\operatorname{spin}^{c}$-structure of $Y$, the quotient space $M_{S}=Y / G$ is a smooth manifold with the metric and the $\operatorname{spin}^{c}$-structure inherited from those of $Y$. The quotient space $X / G$ also has the metric and the spin ${ }^{c}$-structure inherited from those of $X$ near $\partial(X / G)=M_{S}$, whose restriction to $M_{S}$ coincides with those of $M_{S}$. Moreover the $G$-invariant metric connections $\nabla^{Y}, \nabla^{X}$ of $T Y, T X$ define
a metric connection $\nabla^{S}$ of $T M_{S}$, a unitary connection $\nabla^{X / G}$ of $T(X / G)$ near $M_{S}$ respectively. We can show that $M_{S}$ is the boundary of an almost complex manifold $W$ as follows. Let $\varepsilon^{1}$ be the normal bundle of $S^{2 \ell-1}$ in $\mathbb{C}^{\ell}$, which has a $G$-invariant trivialization, and $\varepsilon_{S}^{1}$ the quotient bundle $\left(r_{Y}^{*} \varepsilon^{1}\right) / G$. Note that both of $\varepsilon^{1}$ and $\varepsilon_{S}^{1}$ are trivial real line bundles. Since $T S^{2 \ell-1} \oplus \varepsilon^{1}=\left.T D^{2 \ell}\right|_{S^{2 \ell-1}}$ has the standard complex structure, which is invariant under the action of $G$,

$$
\begin{aligned}
T M_{S} \oplus \varepsilon_{S}^{1} & \cong\left(q_{Y}^{*} T M \oplus r_{Y}^{*} T S^{2 \ell-1} \oplus r_{Y}^{*} \varepsilon^{1}\right) / G \\
& \cong\left(q_{Y}^{*} T M \oplus r_{Y}^{*}\left(\left.T D^{2 \ell}\right|_{S^{2 \ell-1}}\right)\right) / G
\end{aligned}
$$

has a complex structure. Hence the $(2 m+2 \ell-1)$-dimensional compact manifold $M_{S}$ is stably almost complex manifold and therefore it follows from the result of Morita [8] that there exists a compact $(2 m+2 \ell)$-dimensional almost complex manifold $W$ such that $\partial W=M_{S}$ and $W=X / G$ near $M_{S}$ as an almost complex manifold with Hermitian metric. The Hermitian metric of $X / G$ near $M_{S}$ is extended to a Hermitian metric on $W$. Let $Q^{W}$ be the principal $U(m+\ell)$-bundle of unitary frames on $W$. Then the connection $\nabla^{X / G}$ extends to a unitary connection $\nabla^{W}$ in $Q^{W}$. On the other hand, we can see that $\left.T W\right|_{M_{S}}=\left.(T X / G)\right|_{M_{S}}$ is orthogonally decomposed into

$$
\begin{equation*}
\left.T W\right|_{M_{S}} \cong\left(q_{Y}^{*} T M \oplus r_{Y}^{*}\left(\left.T D^{2 \ell}\right|_{S^{2 \ell-1}}\right)\right) / G \cong(T M)_{S} \oplus \varepsilon_{\mathbb{C}}^{\ell} \tag{2.11}
\end{equation*}
$$

where $(T M)_{S}$ is the vector bundle over $M_{S}$ defined by $(T M)_{S}=\left(q_{Y}^{*} T M\right) / G$ and $\varepsilon_{\mathbb{C}}^{\ell}$ is the trivial complex line bundle of rank $\ell$. Then the connection $\nabla^{W}$ splits according to (2.11) as

$$
\begin{equation*}
\left.\nabla^{W}\right|_{T M_{S}}=\left.\nabla^{X / G}\right|_{T M_{S}}=\nabla^{(T M)_{S}} \oplus \nabla^{0} \tag{2.12}
\end{equation*}
$$

where $\nabla^{(T M)_{S}}$ denotes the connection of $(T M)_{S}$ naturally defined by $\nabla$ and $\nabla^{0}$ denotes the globally flat connection of $\varepsilon_{\mathbb{C}}^{\ell}$. Now let $V_{\phi}^{W}$ be the element of the representation ring $R(U(m+\ell))$ defined by

$$
\begin{aligned}
V_{\phi}^{W}=\phi\left(\tau_{1}, \tau_{2}, \ldots, \tau_{m}\right) \in & R(U(m+\ell)) \\
& \subset \mathbb{Z}\left[t_{1}, t_{1}^{-1}, \ldots, t_{m}, t_{m}^{-1}, \ldots, t_{m+\ell}, t_{m+\ell}^{-1}\right]
\end{aligned}
$$

where $\tau_{j}=\sigma_{j}\left(t_{1}-1, \ldots, t_{m}-1, t_{m+1}-1, \ldots, t_{m+\ell}-1\right)$ and set

$$
E_{\phi}^{W}=Q^{W} \times_{U(m+\ell)} V_{\phi}^{W}
$$

Then the connection $\nabla^{W}$ naturally defines a unitary connection of $E_{\phi}^{W}$ and the $E_{\phi}^{W}$-valued $\operatorname{spin}^{c}$-Dirac operator $D_{W}$ is defined.

On the other hand, the quotient bundle $E_{\phi, S}=E_{\phi, Y} / G$ is a virtual complex vector bundle over $M_{S}$ with a unitary connection and the $G$-equivariant Dirac operator $D_{Y}$ naturally defines a differential operator $D_{S}$, which is the $E_{\phi, S}$-valued $\operatorname{spin}^{c}$-Dirac operator on $M_{S}$. Since $Q_{S}=\left(q_{Y}^{*} Q\right) / G$ is the unitary frame bundle associated to $(T M)_{S}$, it follows from (2.11) and (2.12) that
$\left.Q^{W}\right|_{M_{S}}$ is reducible to $Q_{S}$ with the connection. Since $V_{\phi}^{W}$ is isomorphic to $V_{\phi}$ as a virtual $U(m)$-representation, it follows that

$$
\begin{aligned}
\left.E_{\phi}^{W}\right|_{M_{S}} & \cong\left(\left.Q^{W}\right|_{M_{S}}\right) \times_{U(m+\ell)} V_{\phi}^{W} \cong Q_{S} \times_{U(m)} V_{\phi}^{W} \cong Q_{S} \times_{U(m)} V_{\phi} \\
& \cong q_{Y}^{*}\left(Q \times_{U(m)} V_{\phi}\right) / G=\left(q_{Y}^{*} E_{\phi}\right) / G=E_{\phi, Y} / G=E_{\phi, S},
\end{aligned}
$$

where $\cong$ denotes the isomorphism as a virtual vector bundle with an inner product and a unitary connection. Hence, on the collar $M_{S} \times[0, \delta) \subset W, D_{W}$ can be expressed as

$$
D_{W}=\sigma\left(\frac{\partial}{\partial u}+D_{S}\right),
$$

where $u$ is the coordinate of $[0, \delta)$ and $\sigma$ is a bundle isomorphism defined by the Clifford multiplication. Hence the following equality is deduced from (4.3) in [1] as well as in (2.7):
(2.13) $\operatorname{Index}\left(D_{W}\right)=\int_{W} \operatorname{Ch}\left(E_{\phi}^{W}, \nabla^{W}\right) \operatorname{Td}\left(T W, \nabla^{W}\right)-\frac{1}{2}\left(\eta_{S}+\operatorname{dim} \operatorname{ker} D_{S}\right)$,
where $\operatorname{Index}\left(D_{W}\right)$ is the index of $D_{W}$ with a certain global boundary condition, $\operatorname{Ch}\left(E_{\phi}^{W}, \nabla^{W}\right)$ is the Chern character form of $E_{\phi}^{W}, \operatorname{Td}\left(T W, \nabla^{W}\right)$ is the Todd form of $T W$ and $\eta_{S}$ is the eta invariant of $D_{S}$. Here since the spin ${ }^{c}$-structure of $M_{S}$ comes from the $U(m)$-structure of $Y$ which is naturally defined by that of $M$, the spinor bundle $S_{M_{S}}=S_{Y} / G$ on $M_{S}$ splits into $S_{M_{S}}=S_{M_{S}}^{+} \oplus S_{M_{S}}^{-}$ and $D_{S}$ splits into $D_{S}=D_{S}^{+} \oplus D_{S}^{-}$, where

$$
\begin{gather*}
D_{S}^{+}: \Gamma\left(S_{M_{S}}^{+} \otimes E_{\phi, S}\right) \longrightarrow \Gamma\left(S_{M_{S}}^{-} \otimes E_{\phi, S}\right),  \tag{2.14}\\
D_{S}^{-}=\left(D_{S}^{+}\right)^{*}: \Gamma\left(S_{M_{S}}^{-} \otimes E_{\phi, S}\right) \longrightarrow \Gamma\left(S_{M_{S}}^{+} \otimes E_{\phi, S}\right) .
\end{gather*}
$$

Hence we have

$$
\operatorname{dim} \operatorname{ker} D_{S}=\operatorname{dim} \operatorname{ker} D_{S}^{+}+\operatorname{dim} \operatorname{ker} D_{S}^{-}
$$

On the other hand, since the dimension of $Y$ is odd, it follows that

$$
\begin{equation*}
\operatorname{Index}\left(D_{S}^{+}\right)=\operatorname{dim} \operatorname{ker} D_{S}^{+}-\operatorname{dim} \operatorname{ker}\left(D_{S}^{+}\right)^{*}=0 \tag{2.15}
\end{equation*}
$$

(see Proposition 9.2 in [2]). Therefore we have

$$
\operatorname{dim} \operatorname{ker} D_{S}^{-}=\operatorname{dim} \operatorname{ker}\left(D_{S}^{+}\right)^{*}=\operatorname{dim} \operatorname{ker} D_{S}^{+}
$$

and hence it follows that

$$
\frac{1}{2} \operatorname{dim} \operatorname{ker} D_{S}=\operatorname{dim} \operatorname{ker} D_{S}^{+} \equiv 0 \quad(\bmod \mathbb{Z})
$$

Moreover it follows from (3.6) in [3] that

$$
\frac{1}{2} \eta_{S}=\frac{1}{p} \sum_{k=1}^{p} \frac{1}{2} \eta_{Y}\left(\psi^{k}\right)
$$

Hence it follows from (2.13) that

$$
\begin{equation*}
\frac{1}{p} \sum_{k=1}^{p-1} \frac{1}{2} \eta_{Y}\left(\psi^{k}\right)+\frac{1}{p} \frac{1}{2} \eta_{Y} \equiv \int_{W} \operatorname{Ch}\left(E_{\phi}^{W}, \nabla^{W}\right) \operatorname{Td}\left(T W, \nabla^{W}\right) \quad(\bmod \mathbb{Z}) \tag{2.16}
\end{equation*}
$$

Here it follows from (2.9) and (2.10) that
(2.17) $\frac{1}{p} \sum_{k=1}^{p-1} \frac{1}{2}\left(\eta_{Y}\left(\psi^{k}\right)\right)+\frac{1}{p} \frac{1}{2} \eta_{Y}$

$$
\begin{aligned}
= & \frac{1}{p} \sum_{k=1}^{p-1} C_{\ell}(k, \gamma) \sum_{N \subset \Omega(k)} \operatorname{Ch}\left(\left.E_{\phi}\right|_{N}, \psi^{k}\right) \operatorname{Td}(T N) \mathfrak{U}\left(\nu(N, M), \psi^{k}\right)[N] \\
& -\frac{1}{p} \sum_{k=1}^{p} \operatorname{Index}\left(D_{X}, \psi^{k}\right)-\frac{1}{p} \sum_{k=1}^{p} \frac{1}{2} \operatorname{Tr}\left(\left.\psi^{k}\right|_{\text {ker } D_{Y}}\right) .
\end{aligned}
$$

Here since the spin ${ }^{c}$-structure of $Y$ comes from the $U(m)$-structure of $M$, the spinor bundle $S_{Y}$ splits into $S_{Y}=S_{Y}^{+} \oplus S_{Y}^{-}$and $D_{Y}$ splits into $D_{Y}=D_{Y}^{+} \oplus D_{Y}^{-}$ where

$$
\begin{gathered}
D_{Y}^{+}: \Gamma\left(S_{Y}^{+} \otimes E_{\phi, Y}\right) \longrightarrow \Gamma\left(S_{Y}^{-} \otimes E_{\phi, Y}\right), \\
D_{Y}^{-}=\left(D_{Y}^{+}\right)^{*}: \Gamma\left(S_{Y}^{-} \otimes E_{\phi, Y}\right) \longrightarrow \Gamma\left(S_{Y}^{+} \otimes E_{\phi, Y}\right)
\end{gathered}
$$

as in (2.14). Here since $\psi^{k}(1 \leq k \leq p-1)$ acts freely on $Y$, it follows from the fixed point formula in [2] that

$$
\operatorname{Index}\left(D_{Y}^{+}, \psi^{k}\right):=\operatorname{Tr}\left(\left.\psi^{k}\right|_{\operatorname{ker} D_{Y}^{+}}\right)-\operatorname{Tr}\left(\left.\psi^{k}\right|_{\operatorname{ker}\left(D_{Y}^{+}\right)^{*}}\right)=0
$$

for any $1 \leqq k \leqq p-1$. Moreover, since the dimension of $Y$ is odd, it follows as in (2.15) that

$$
\operatorname{Index}\left(D_{Y}^{+}\right)=\operatorname{Tr}\left(\left.\psi^{p}\right|_{\operatorname{ker} D_{Y}^{+}}\right)-\operatorname{Tr}\left(\left.\psi^{p}\right|_{\operatorname{ker}\left(D_{Y}^{+}\right)^{*}}\right)=0
$$

and hence that

$$
\begin{aligned}
\sum_{k=1}^{p} & \frac{1}{2} \operatorname{Tr}\left(\left.\psi^{k}\right|_{\text {ker } D_{Y}}\right)=\sum_{k=1}^{p} \frac{1}{2}\left\{\operatorname{Tr}\left(\left.\psi^{k}\right|_{\text {ker } D_{Y}^{+}}\right)+\operatorname{Tr}\left(\left.\psi^{k}\right|_{\text {ker } D_{Y}^{-}}\right)\right\} \\
& =\sum_{k=1}^{p} \frac{1}{2}\left\{\operatorname{Tr}\left(\left.\psi^{k}\right|_{\text {ker } D_{Y}^{+}}\right)+\operatorname{Tr}\left(\left.\psi^{k}\right|_{\operatorname{ker}\left(D_{Y}^{+}\right)^{*}}\right)\right\}=\sum_{k=1}^{p} \operatorname{Tr}\left(\left.\psi^{k}\right|_{\text {ker } D_{Y}^{+}}\right) .
\end{aligned}
$$

Therefore it follows from (2.5) that

$$
\begin{align*}
\sum_{k=1}^{p} & \operatorname{Index}\left(D_{X}, \psi^{k}\right)+\sum_{k=1}^{p} \frac{1}{2} \operatorname{Tr}\left(\left.\psi^{k}\right|_{\operatorname{ker} D_{Y}}\right)  \tag{2.18}\\
& =\sum_{k=1}^{p} \operatorname{Index}\left(D_{X}, \psi^{k}\right)+\sum_{k=1}^{p} \operatorname{Tr}\left(\left.\psi^{k}\right|_{\text {ker } D_{Y}^{+}}\right) \equiv 0 \quad(\bmod p) .
\end{align*}
$$

Hence it follows from (2.16), (2.17) and (2.18) that

$$
\begin{align*}
& \frac{1}{p} \sum_{k=1}^{p-1} C_{\ell}(k, \gamma) \sum_{N \subset \Omega(k)} \operatorname{Ch}\left(\left.E_{\phi}\right|_{N}, \psi^{k}\right) \operatorname{Td}(T N) \mathfrak{U}\left(\nu(N, M), \psi^{k}\right)[N]  \tag{2.19}\\
& \equiv \equiv \int_{W} \operatorname{Ch}\left(E_{\phi}^{W}, \nabla^{W}\right) \operatorname{Td}\left(T W, \nabla^{W}\right) \quad(\bmod \mathbb{Z})
\end{align*}
$$

Here the same argument as was used to prove (2.3) shows that

$$
\begin{equation*}
\int_{W} \operatorname{Ch}\left(E_{\phi}^{W}, \nabla^{W}\right) \operatorname{Td}\left(T W, \nabla^{W}\right)=0 \tag{2.20}
\end{equation*}
$$

because the order of $\phi$ is greater than $m+\ell$ and the dimension of $W$ is $2 m+2 \ell$. Now the equality in Theorem 1.3 is deduced from (2.19) and (2.20). This completes the proof of Theorem 1.3.

## 3. Examples

In this section, applying Theorem 1.3, we give certain fixed point formulae for the standard torus $T^{2}$, the sphere $S^{6}$ and the complex projective space $\mathbf{C P} \mathbb{P}^{m}$, which can be verified by direct computation.

Example 3.1. Let $T^{2}$ be the standard torus defined by $T^{2}=\mathbb{C} /(\mathbb{Z}+$ $\sqrt{-1} \mathbb{Z})$. Let $\psi$ be the automorphism of $T^{2}$ defined by the $\pi / 2$-rotation with center at $(1+i) / 2$. Then the order of $\psi$ is 4 and the fixed point set $\Omega(k)$ of $\psi^{k}$ is as follows:

$$
\begin{aligned}
& \Omega(1)=\Omega(3)=\left\{A=\frac{1+i}{2}, B=1+i\right\} \\
& \Omega(2)=\left\{A=\frac{1+i}{2}, B=1+i, C=\frac{1}{2}+i, D=1+\frac{i}{2}\right\} .
\end{aligned}
$$

Set $\ell=2, \gamma=3$ and $\phi=x_{1}^{n}=\tau_{1}^{n} \in S(n)$. Since $\psi^{k}$ acts on $T_{A} T^{2}, T_{B} T^{2}$ via multiplication by $i^{k}$ for $1 \leqq k \leqq 3$ and $\psi^{2}$ acts on $T_{C} T^{2}, T_{D} T^{2}$ via multiplication by -1 , it follows from Corollary 1.4 that the equality

$$
\begin{align*}
& \frac{1}{\left|1-i^{-3}\right|^{2}}\left(2(i-1)^{n} \frac{1}{1-i^{-1}}\right)  \tag{3.1}\\
& \quad+\frac{1}{\left|1-i^{-6}\right|^{2}}\left(2\left(i^{2}-1\right)^{n} \frac{1}{1-i^{-2}}+2(-1-1)^{n} \frac{1}{1-(-1)^{-1}}\right) \\
& \quad+\frac{1}{\left|1-i^{-9}\right|^{2}}\left(2\left(i^{3}-1\right)^{n} \frac{1}{1-i^{-3}}\right) \equiv 0 \quad(\bmod 4)
\end{align*}
$$

holds for any $n>m+\ell=3$. The equality above can be easily verified as follows:

$$
\text { the left-hand side of } \begin{aligned}
(3.1) & =i(i-1)^{n-1}+(-2)^{n}+\overline{\left(i(i-1)^{n-1}\right)} \\
& =2 \operatorname{Re}\left(i(i-1)^{n-1}\right)+(-2)^{n} \equiv 0 \quad(\bmod 4),
\end{aligned}
$$

where Re denotes the real part because we can show that both of the real part and the imaginary part of $i(i-1)^{n-1}$ are even for $n \geqq 3$ by induction.

Example 3.2. Let $\mathbb{O}=\mathbb{H} \oplus \mathbb{H}$ be the set of octonions with multiplication defined by the rule

$$
x \cdot x^{\prime}=\left(q_{1}, q_{2}\right) \cdot\left(q_{1}^{\prime}, q_{2}^{\prime}\right) \equiv\left(q_{1} q_{1}^{\prime}-\overline{q_{2}^{\prime}} q_{2}, q_{2}^{\prime} q_{1}+q_{2} \overline{q_{1}^{\prime}}\right)
$$

for any $x, x^{\prime} \in \mathbb{O}($ see $[7])$. The conjugation $\bar{x}$ and the real part $\operatorname{Re}(x)$ of $x=\left(q_{1}, q_{2}\right) \in \mathbb{O}$ are defined by $\bar{x}=\left(\overline{q_{1}},-q_{2}\right)$ and $\operatorname{Re}(x)=\operatorname{Re}\left(q_{1}\right)$ respectively. Moreover the standard Euclidean inner product $\left\langle x, x^{\prime}\right\rangle$ and its norm $|x|$ are defined for $x, x^{\prime} \in \mathbb{O}$ by

$$
\left\langle x, x^{\prime}\right\rangle=\operatorname{Re}\left(x \cdot \overline{x^{\prime}}\right)=\operatorname{Re}\left(\bar{x} \cdot x^{\prime}\right), \quad|x|=\sqrt{\langle x, x\rangle}=\sqrt{x \cdot \bar{x}}=\sqrt{\bar{x} \cdot x}
$$

respectively. The map

$$
\mathbb{O} \ni\left(\left(z_{1}, z_{2}\right),\left(z_{3}, z_{4}\right)\right) \longrightarrow\left(z_{1}, z_{2}, z_{3}, z_{4}\right) \in \mathbb{C}^{4}
$$

gives an isomorphism as a complex vector space. We denote $\left(\left(z_{1}, z_{2}\right),\left(z_{3}, z_{4}\right)\right)$ by $\left(z_{1}, z_{2}, z_{3}, z_{4}\right)$ hereafter. Let $\operatorname{Im}(\mathbb{O})$ be the set of pure imaginary octonions, namely,

$$
\operatorname{Im}(\mathbb{O})=\{x \in \mathbb{O} \mid \bar{x}=-x\},
$$

which is isomorphic to $\mathbb{R}^{7}$ as a real vector space and $S^{6}$ the standard 6dimensional sphere defined by

$$
S^{6}=\left\{A=\left(z_{1}, z_{2}, z_{3}, z_{4}\right) \in \operatorname{Im}(\mathbb{O})| | A\left|=\left|z_{1}\right|^{2}+\left|z_{2}\right|^{2}+\left|z_{3}\right|^{2}+\left|z_{4}\right|^{2}=1\right\}\right.
$$

Then, for any point $A \in S^{6}$, the tangent space $T_{A} S^{6}$ is given by

$$
T_{A} S^{6}=\{B \in \operatorname{Im}(\mathbb{O}) \mid\langle A, B\rangle=0\}
$$

For any $A \in S^{6}, B \in T_{A} S^{6}$, set $J_{A}(B)=A \cdot B$. Then since the equality $\bar{x} \cdot(x \cdot y)=(\bar{x} \cdot x) \cdot y=|x|^{2} y$ holds for any $x, y \in \mathbb{O}$, we have

$$
J_{A}\left(J_{A}(B)\right)=A \cdot(A \cdot B)=-\bar{A} \cdot(A \cdot B)=-(\bar{A} \cdot A) \cdot B=-|A|^{2} B=-B,
$$

which implies that $J_{A}(B) \in T_{A} S^{6}$ and $J_{A}^{2}=-1$ because $\bar{A} \cdot(A \cdot B)=B$ implies that $\langle A, A \cdot B\rangle=\operatorname{Re}(\bar{A} \cdot(A \cdot B))=\operatorname{Re}(B)=0$. Hence this $J$ defines an almost complex structure of $S^{6}$. Let $p$ be any natural number, $\alpha=e^{2 \pi \sqrt{-1} / p}$ the primitive $p$-th root of 1 and $\psi$ the periodic $\mathbb{C}$-linear map of $\mathbb{O}$ of order $p$ defined by

$$
\left(\mathbb{O} \ni\left(q_{1}, q_{2}\right)=\left(z_{1}, z_{2}, z_{3}, z_{4}\right) \longrightarrow\left(q_{1}, \alpha q_{2}\right)=\left(z_{1}, z_{2}, \alpha z_{3}, \alpha z_{4}\right) .\right.
$$

Then $\psi$ maps $S^{6}$ to $S^{6}$ and we have

$$
\begin{aligned}
\psi(x) \cdot \psi(y) & =\left(q_{1}, \alpha q_{2}\right) \cdot\left(q_{1}^{\prime}, \alpha q_{2}^{\prime}\right)=\left(q_{1} q_{1}^{\prime}-\overline{\alpha q_{2}^{\prime}} \alpha q_{2}, \alpha q_{2}^{\prime} q_{1}+\alpha q_{2} \overline{q_{1}^{\prime}}\right) \\
& =\left(q_{1} q_{1}^{\prime}-\overline{q_{2}^{\prime}} \bar{\alpha} \alpha q_{2}, \alpha q_{2}^{\prime} q_{1}+\alpha q_{2} \overline{q_{1}^{\prime}}\right)=\left(q_{1} q_{1}^{\prime}-\overline{q_{2}^{\prime}} q_{2}, \alpha\left(q_{2}^{\prime} q_{1}+q_{2} \overline{q_{1}^{\prime}}\right)\right) \\
& =\psi(x \cdot y)
\end{aligned}
$$

for any $x=\left(q_{1}, q_{2}\right), y=\left(q_{1}^{\prime}, q_{2}^{\prime}\right) \in \mathbb{O}$. Hence it follows that

$$
J_{\psi(A)}\left(\psi_{*}(B)\right)=J_{\psi(A)}(\psi(B))=\psi(A) \cdot \psi(B)=\psi(A \cdot B)=\psi_{*}\left(J_{A}(B)\right)
$$

for any $A \in S^{6}, B \in T_{A} S^{6}$, which implies that $\psi$ commutes with $J$. Hence $\psi$ defines an automorphism of the almost complex manifold $S^{6}$. The fixed point set of $\psi^{k}$ is independent of $k$ and coincides with the standard 2-dimensional sphere

$$
S^{2}=\left\{\left(z_{1}, z_{2}, z_{3}, z_{4}\right) \in S^{6} \mid z_{3}=z_{4}=0\right\}
$$

for any $1 \leqq k \leqq p-1$. The normal bundle $\nu\left(S^{2}, S^{6}\right)$ is the trivial complex vector bundle of rank 2 and $\psi^{k}$ acts on $\nu\left(S^{2}, S^{6}\right)$ via multiplications by $\alpha^{k}$.

Set $\ell=0$ and $\phi=\left(x_{1} x_{2} x_{3}\right)^{n}=\tau_{3}^{n} \in S(3 n)$. Since $\left.T S^{6}\right|_{S^{2}}$ splits into the direct sum $T S^{2} \oplus \nu\left(S^{2}, S^{6}\right)$ and $\psi^{k}$ acts on $T S^{2}$ via multiplication by 1, it follows from (1.3) that

$$
\operatorname{Ch}\left(\left.E_{\phi}\right|_{S^{2}}, \psi^{k}\right)=\left\{\left(e^{c_{1}\left(T S^{2}\right)}-1\right)\left(\alpha^{k}-1\right)^{2}\right\}^{n}=\left\{\left(e^{2 x}-1\right)\left(\alpha^{k}-1\right)^{2}\right\}^{n}
$$

where $x$ denotes the positive generator of $H^{2}\left(S^{2}\right) \cong \mathbb{Z}$. Hence it follows from Theorem 1.3 that the equality

$$
\begin{align*}
\sum_{k=1}^{p-1} & \left\{\left(e^{2 x}-1\right)\left(\alpha^{k}-1\right)^{2}\right\}^{n} \frac{x}{1-e^{-x}}\left(\frac{1}{1-\alpha^{-k}}\right)^{2}\left[S^{2}\right]  \tag{3.2}\\
& =2^{n} \sum_{k=1}^{p-1}\left(\alpha^{k}-1\right)^{2 n}\left(\frac{1}{1-\alpha^{-k}}\right)^{2}\left(x^{n}+\text { higher order terms }\right)\left[S^{2}\right] \\
& \equiv 0 \quad(\bmod p)
\end{align*}
$$

holds for any $n$ such that $3 n>m+\ell=3$. The equality above can also be easily verified because $3 n>3$ implies that $n \geqq 2$ and hence that

$$
\left(x^{n}+\text { higher order terms }\right)\left[S^{n}\right]=0 .
$$

Remark 3.3. If $3 n=m+\ell=3 \Longleftrightarrow n=1$, it follows from (2.5) and (3.2) that

$$
\begin{aligned}
& \sum_{k=1}^{p-1}\left(e^{2 x}-1\right)\left(\alpha^{k}-1\right)^{2} \frac{x}{1-e^{-x}}\left(\frac{1}{1-\alpha^{-k}}\right)^{2}\left[S^{2}\right] \\
& \quad=2 \sum_{k=1}^{p-1} \alpha^{2 k}=2\left(\sum_{k=1}^{p} \alpha^{2 k}-1\right) \equiv-2 \neq 0 \quad(\bmod p)
\end{aligned}
$$

if $p \neq 2$.
Example 3.4. Let $M$ be the $m$-dimensional complex projective space $\mathbf{C P} \mathbb{P}^{m}, p$ any natural number, $\alpha=e^{2 \pi \sqrt{-1} / p}$ the primitive $p$-th root of 1 and $\psi$ the periodic automorphism of $\mathbf{C P}{ }^{m}$ of order $p$ defined by

$$
\mathbf{C P} \mathbb{P}^{m} \ni\left[z_{0}: z_{1}: \cdots: z_{m}\right] \longrightarrow\left[\alpha z_{0}: z_{1}: \cdots: z_{m}\right] .
$$

Then the fixed point set of $\psi^{k}$ is independent of $k$ and coincides with the disjoint union of the point $q=[1: 0: \cdots: 0]$ and the hyperplane $\mathbf{C} \mathbb{P}^{m-1}$ defined by $z_{0}=0$. Set $\ell=1, \gamma=1$ and $\phi=\left(x_{1}+x_{2}+\cdots+x_{m}\right)^{n}=\tau_{1}^{n} \in S(n)$. Then it follows that $\phi=\left(t_{1}+t_{2}+\cdots+t_{m}-m\right)^{n}$ and hence that

$$
E_{\phi}=\otimes^{n}\left(T \mathbf{C} \mathbb{P}^{m}-\varepsilon_{\mathbb{C}}^{m}\right),
$$

where $\psi^{k}$ acts on the trivial bundle $\varepsilon_{\mathbb{C}}^{m}$ via multiplication by 1 . Here $\psi^{k}$ acts on $\nu\left(q, \mathbf{C P}^{m}\right) \cong \mathbb{C}^{m}$ via multiplication by $\alpha^{-k}$ and hence we have

$$
\operatorname{Ch}\left(\left.T \mathbf{C P}^{m}\right|_{q}, \psi^{k}\right)=m \alpha^{-k}, \quad \operatorname{Td}(T q)=1, \quad \mathfrak{U}\left(\nu(q, M), \psi^{k}\right)=\left(\frac{1}{1-\alpha^{k}}\right)^{m}
$$

On the other hand, the normal bundle $\nu\left(\mathbf{C P}^{m-1}, \mathbf{C} \mathbb{P}^{m}\right)$ is isomorphic to the restriction of the hyperplane bundle $L$ to $\mathbf{C P} \mathbb{P}^{m-1}$ and $\psi^{k}$ acts on $\nu\left(\mathbf{C P}{ }^{m-1}, \mathbf{C P}{ }^{m}\right)$ $\left.\cong L\right|_{\mathbf{C P}^{m-1}}$ via multiplication by $\alpha^{k}$. Let $x$ be the positive generator of $H^{2}\left(\mathbf{C P}^{m-1}\right) \cong \mathbb{Z}$ which equals to the first Chern class $c_{1}\left(\left.L\right|_{\mathbf{C P}^{m-1}}\right)$. Then since

$$
\left.T \mathbf{C} \mathbb{P}^{m}\right|_{\mathbf{C P}^{m-1}}=T \mathbf{C} \mathbb{P}^{m-1} \oplus\left(\left.L\right|_{\mathbf{C} \mathbb{P}^{m-1}}\right)
$$

where $\psi^{k}$ acts on $T \mathbf{C P}{ }^{m-1}$ via multiplication by 1 , it follows that
$\operatorname{Ch}\left(\left.T \mathbf{C P} \mathbb{P}^{m}\right|_{\mathbf{C P}^{m-1}}, \psi^{k}\right)=m e^{x}-1+\alpha^{k} e^{x}, \quad \operatorname{Td}\left(T \mathbf{C P}^{m-1}\right)=\left(\frac{x}{1-e^{-x}}\right)^{m}$,
$\mathfrak{U}\left(\nu\left(\mathbf{C P}^{m-1}, \mathbf{C P}^{m}\right), \psi^{k}\right)=\frac{1}{1-\alpha^{-k} e^{-x}}$.

Hence it follows from Theorem 1.3 that the equality

$$
\begin{equation*}
\sum_{k=1}^{p-1} \frac{1}{1-\alpha^{-k}}\left(m \alpha^{-k}-m\right)^{n}\left(\frac{1}{1-\alpha^{k}}\right)^{m}+\varphi(x)\left[\mathbf{C P}^{m-1}\right] \equiv 0 \quad(\bmod p) \tag{3.3}
\end{equation*}
$$

holds for any $n>m+\ell=m+1$, where

$$
\varphi(x)=\sum_{k=1}^{p-1} \frac{1}{1-\alpha^{-k}}\left(m e^{x}-1+\alpha^{k} e^{x}-m\right)^{n}\left(\frac{x}{1-e^{-x}}\right)^{m} \frac{1}{1-\alpha^{-k} e^{-x}} .
$$

We can verify (3.3) as follows.
$\varphi(x)\left[\mathbf{C P}^{m-1}\right]=x^{m-1}$-coefficient of $\varphi(x)$
$=x^{-1}$-coefficient of

$$
\begin{aligned}
& \frac{\varphi(x)}{x^{m}}=\sum_{k=1}^{p-1} \frac{\left(m e^{x}-1+\alpha^{k} e^{x}-m\right)^{n}}{\left(1-\alpha^{-k}\right)\left(1-e^{-x}\right)^{m}\left(1-\alpha^{-k} e^{-x}\right)} \\
= & \sum_{k=1}^{p-1} \frac{\left(\alpha^{k}\right)^{2}}{\alpha^{k}-1}\left(\left(m+\alpha^{k}\right) e^{x}-1-m\right)^{n} \frac{\left(e^{x}\right)^{m}}{\left(e^{x}-1\right)^{m}} \frac{1}{\alpha^{k} e^{x}-1} e^{x} \\
= & \frac{1}{2 \pi i} \oint_{C(x)} \sum_{k=1}^{p-1} \frac{\left(\alpha^{k}\right)^{2}}{\alpha^{k}-1}\left(\left(m+\alpha^{k}\right) e^{x}-1-m\right)^{n} \frac{\left(e^{x}\right)^{m}}{\left(e^{x}-1\right)^{m}} \frac{1}{\alpha^{k} e^{x}-1} e^{x} d x
\end{aligned}
$$

$(C(x)$ is a sufficiently small counterclockwise simple loop around $0 \in \mathbb{C})$
$=\frac{1}{2 \pi i} \oint_{C(y)} \sum_{k=1}^{p-1} \frac{\left(\alpha^{k}\right)^{2}}{\alpha^{k}-1}\left(\left(m+\alpha^{k}\right)(y+1)-1-m\right)^{n} \frac{(y+1)^{m}}{y^{m}} \frac{1}{\alpha^{k}(y+1)-1} d y$
( $y=e^{x}-1, C(y)$ is a counterclockwise simple loop around $0 \in \mathbb{C}$ )
$=y^{-1}$-coefficient of

$$
\begin{aligned}
& \sum_{k=1}^{p-1} \frac{\left(\alpha^{k}\right)^{2}}{\alpha^{k}-1}\left(\left(m+\alpha^{k}\right)(y+1)-1-m\right)^{n} \frac{(y+1)^{m}}{y^{m}} \frac{1}{\alpha^{k}(y+1)-1} \\
= & y^{m-1} \text {-coefficient of } \\
& \sum_{k=1}^{p-1} \frac{\left(\alpha^{k}\right)^{2}}{\alpha^{k}-1}\left(\left(m+\alpha^{k}\right) y+\alpha^{k}-1\right)^{n}(y+1)^{m} \frac{1}{\alpha^{k}-1+\alpha^{k} y}
\end{aligned}
$$

$=y^{m-1}$-coefficient of

$$
\sum_{k=1}^{p-1} \frac{\left(\alpha^{k}\right)^{2}}{\left(\alpha^{k}-1\right)^{2}} \sum_{i=0}^{n}\binom{n}{i}\left(m+\alpha^{k}\right)^{i} y^{i}\left(\alpha^{k}-1\right)^{n-i} \sum_{j=0}^{m}\binom{m}{j} y^{j} \sum_{s=0}^{\infty}\left(\frac{-\alpha^{k} y}{\alpha^{k}-1}\right)^{s}
$$

$=y^{m-1}$-coefficient of

$$
\begin{aligned}
& \sum_{k=1}^{p-1} \sum_{i=0}^{n} \sum_{j=0}^{m} \sum_{s=0}^{\infty}\binom{n}{i}\binom{m}{j}(-1)^{s}\left(\alpha^{k}\right)^{s+2}\left(m+\alpha^{k}\right)^{i}\left(\alpha^{k}-1\right)^{n-i-s-2} y^{i+j+s} \\
= & \sum_{k=1}^{p-1} \sum_{i=0}^{m-1} \sum_{s=0}^{m-1-i}\binom{n}{i}\binom{m}{m-1-i-s}(-1)^{s}\left(\alpha^{k}\right)^{s+2}\left(m+\alpha^{k}\right)^{i}\left(\alpha^{k}-1\right)^{n-i-s-2} \\
= & \sum_{k=1}^{p-1}\left(\alpha^{k}-1\right) R\left(\alpha^{k}\right)
\end{aligned}
$$

where $R(z)$ is an integral polynomial defined by

$$
R(z)=\sum_{i=0}^{m-1} \sum_{s=0}^{m-1-i}\binom{n}{i}\binom{m}{m-1-i-s}(-1)^{s} z^{s+2}(m+z)^{i}(z-1)^{n-i-s-3} .
$$

(Note that $n>m+1$ and $i+s \leqq m-1$ imply that $n-i-s-3 \geqq 0$.)
Now since $\left(\alpha^{ \pm \nu}\right)^{p}=1$ for any nonnegative integer $\nu$, it follows that

$$
\sum_{k=1}^{p-1}\left(\alpha^{ \pm k}\right)^{\nu}=\sum_{k=1}^{p-1}\left(\alpha^{ \pm \nu}\right)^{k}=\sum_{k=1}^{p}\left(\alpha^{ \pm \nu}\right)^{k}-1 \equiv-1 \quad(\bmod p)
$$

(see (2.5)). Hence, for any integral polynomial $Q(z)$, we can see that

$$
\begin{equation*}
\sum_{k=1}^{p-1} Q\left(\alpha^{k}\right) \equiv \sum_{k=1}^{p-1} Q\left(\alpha^{-k}\right) \equiv-Q(1) \quad(\bmod p) \tag{3.5}
\end{equation*}
$$

and therefore it follows from (3.4) that

$$
\varphi(x)\left[\mathbf{C P}^{m-1}\right]=\sum_{k=1}^{p-1}\left(\alpha^{k}-1\right) R\left(\alpha^{k}\right) \equiv-(1-1) R(1)=0 \quad(\bmod p)
$$

On the other hand, it follows from (3.5) that

$$
\begin{aligned}
& \sum_{k=1}^{p-1} \frac{1}{1-\alpha^{-k}}\left(m \alpha^{-k}-m\right)^{n}\left(\frac{1}{1-\alpha^{k}}\right)^{m} \\
& =\sum_{k=1}^{p-1}\left(-m^{n}\right)\left(\alpha^{-k}\right)^{m}\left(\alpha^{-k}-1\right)^{n-m-1} \equiv m^{n} \cdot 1^{m} \cdot(1-1)^{n-m-1}=0 \quad(\bmod p)
\end{aligned}
$$

because $n-m-1>0$. Hence the equality (3.7) is verified.

## 4. Relation to the Einstein-Kähler metrics

If the Ricci form $\rho(\omega)$ of the Kähler form $\omega$ on a Kähler manifold $M$ is a constant multiple of $\omega, M$ is called an Einstein-Kähler manifold and the metric corresponding to $\omega$ is called an Einstein-Kähler metric. In this section, we refine the result of Theorem 1.3 for $\ell=1, \gamma=1$ and $\phi=\tau_{1}^{m+1}$ in the case that $M$ is an Einstein-Kähler manifold.

Let $M$ be an $m$-dimensional complex manifold and $A(M)$ the complex Lie group consisting of all biholomorphic automorphisms of $M$. Assume that the periodic element $\psi \in A(M)$ of order $p$ is contained in the identity component of $A(M)$ and hence is expressed as $\psi=\exp v$ by a holomorphic vector field $v$ on $M$. Then using the result of Futaki in [4] and the result in [9], we can prove the next theorem.

Theorem 4.1. If $M$ admits an Einstein-Kähler metric, then under the notation in Theorem 1.3 the equality
$\sum_{k=1}^{p-1} C_{1}(k, 1) \sum_{N \subset \Omega(k)} \operatorname{Ch}\left(\left.E_{\phi}\right|_{N}, \psi^{k}\right) \operatorname{Td}(T N) \mathfrak{U}\left(\nu(N, M), \psi^{k}\right)[N] \equiv 0 \quad(\bmod p)$
holds for $\phi=\left(x_{1}+x_{2}+\cdots+x_{m}\right)^{m+1}=\tau_{1}^{m+1} \in S(m+1)$.

Proof. Let $G$ be the cyclic subgroup of $A(M)$ generated by $\psi$. Set $\ell=1$, $\gamma=1$ and $\phi=\left(x_{1}+x_{2}+\cdots+x_{m}\right)^{m+1}=\tau_{1}^{m+1} \in S(m+1)$. Then $G$ acts freely on $Y=M \times S^{1}$ by

$$
\psi \cdot(q, z)=(\psi \cdot q, z \alpha) \quad(q \in M, z \in \mathbb{C})
$$

Let $M_{S}$ be the quotient space $Y / G$ and $W$ the $(2 m+2)$-dimensional almost complex manifold whose boundary is $M_{S}$ as in Section 2. Then it follows from Theorem 1.6 and Lemma 2.1 in [9] that the equality

$$
\begin{equation*}
f(v) \equiv \int_{W} c_{1}\left(T W, \nabla^{W}\right)^{m+1} \quad(\bmod \mathbb{Z}) \tag{4.2}
\end{equation*}
$$

holds, where $f(v)$ is the Futaki invariant of $v$ (see [4]). Since

$$
\operatorname{Ch}\left(E_{\phi}^{W}, \nabla^{W}\right)=c_{1}\left(T W, \nabla^{W}\right)^{m+1}+\text { higher order terms }
$$

(see (2.2)) it follows from (4.2) that

$$
f(v) \equiv \int_{W} \operatorname{Ch}\left(E_{\phi}^{W}, \nabla^{W}\right) \operatorname{Td}\left(E_{\phi}^{W}, \nabla^{W}\right) \quad(\bmod \mathbb{Z})
$$

Therefore it follows from (2.19) and the equality above that

$$
\begin{equation*}
f(v) \equiv \frac{1}{p} \sum_{k=1}^{p-1} C_{1}(k, 1) \sum_{N \subset \Omega(k)} \operatorname{Ch}\left(\left.E_{\phi}\right|_{N}, \psi^{k}\right) \operatorname{Td}(T N) \mathfrak{U}\left(\nu(N, M), \psi^{k}\right)[N] \tag{4.3}
\end{equation*}
$$

$(\bmod \mathbb{Z})$.
On the other hand, Futaki proved in [4] (see also [5]) that $f(v)=0$ for any holomorphic vector field $v$ if $M$ admits an Einstein-Kähler metric. Hence it follows that
$\frac{1}{p} \sum_{k=1}^{p-1} C_{1}(k, 1) \sum_{N \subset \Omega(k)} \operatorname{Ch}\left(\left.E_{\phi}\right|_{N}, \psi^{k}\right) \operatorname{Td}(T N) \mathfrak{U}\left(\nu(N, M), \psi^{k}\right)[N] \equiv 0 \quad(\bmod \mathbb{Z})$
if $M$ admits an Einstein-Kähler metric. This completes the proof of Theorem 4.1.

Remark 4.2. Note that it follows from Theorem 1.3 that the equality (4.1) holds for any almost complex manifold $M$ if $\phi \in S(n)$ and $n>m+\ell=$ $m+1$. On the other hand, the equality in Theorem 1.3 does not hold in general if $n=m+\ell$. For example, let $M$ be the blowing-up of $\mathbf{C P}^{2}$ at one point. Then as was seen in [9] (see Theorem 1.6 and p. 215 in [9]) there exists a periodic biholomorphic automorphism $\psi=\exp v \in A(M)$ such that $f(v)$ is not an integer. Hence it follows from (4.3) that
$\sum_{k=1}^{p-1} C_{1}(k, 1) \sum_{N \subset \Omega(k)} \operatorname{Ch}\left(\left.E_{\phi}\right|_{N}, \psi^{k}\right) \operatorname{Td}(T N) \mathfrak{U}\left(\nu(N, M), \psi^{k}\right)[N] \not \equiv 0 \quad(\bmod p)$, where $\phi=\left(x_{1}+x_{2}+\cdots+x_{m}\right)^{m+1}=\tau_{1}^{m+1} \in S(m+1)=S(m+\ell)$.

Example 4.3. Let $M=\mathbf{C P}{ }^{m}$ and $\psi$ the periodic automorphism defined in Example 3.4. Then the equality (4.1) holds for any periodic $\psi \in$ $A\left(\mathbf{C P}^{m}\right)$ because $A\left(\mathbf{C P}^{m}\right)$ is connected and $\mathbf{C P}{ }^{m}$ admits an Einstein-Kähler metric. In fact it follows as in (3.3) that the equality
the left-hand side of (4.1)

$$
=\sum_{k=1}^{p-1} \frac{1}{1-\alpha^{-k}}\left(m \alpha^{-k}-m\right)^{m+1}\left(\frac{1}{1-\alpha^{k}}\right)^{m}+\varphi(x)\left[\mathbf{C P}^{m-1}\right] \equiv 0 \quad(\bmod p)
$$

holds, where

$$
\varphi(x)=\sum_{k=1}^{p-1} \frac{1}{1-\alpha^{-k}}\left(m e^{x}-1+\alpha^{k} e^{x}-m\right)^{m+1}\left(\frac{x}{1-e^{-x}}\right)^{m} \frac{1}{1-\alpha^{-k} e^{-x}} .
$$

Hence it follows from the same argument as in Example 3.4 and (3.5) that
the left-hand side of (4.1)

$$
\begin{aligned}
\equiv & m^{m+1} \cdot 1^{m} \cdot(1-1)^{m+1-m-1} \\
& -\sum_{i=0}^{m-1} \sum_{s=0}^{m-1-i}\binom{m+1}{i}\binom{m}{m-1-i-s} \\
& \times(-1)^{s} 1^{s+2}(m+1)^{i}(1-1)^{m+1-i-s-2} \quad(\bmod p) \\
= & m^{m+1}-\sum_{i=0}^{m-1}\binom{m+1}{i}\binom{m}{0}(-1)^{m-1-i}(m+1)^{i}(1-1)^{0} \\
= & m^{m+1} \\
& -\sum_{i=0}^{m+1}\binom{m+1}{i}(-1)^{m+1-i}(m+1)^{i}+(m+1)^{m+1}-\binom{m+1}{m}(m+1)^{m} \\
= & m^{m+1}-\sum_{i=0}^{m+1}\binom{m+1}{i}(-1)^{m+1-i}(m+1)^{i}=m^{m+1}-\{(m+1)-1\}^{m+1} \\
= & m^{m+1}-m^{m+1}=0 .
\end{aligned}
$$

Thus the equality (4.1) holds for $M=\mathbf{C P}{ }^{m}$ and

$$
\psi: \mathbf{C P}^{m} \ni\left[z_{0}: z_{1}: \cdots: z_{m}\right] \longrightarrow\left[\alpha z_{0}: z_{1}: \cdots: z_{m}\right]
$$

Acknowledgement. The author is grateful to Professor Akito Futaki for valuable information.

## References

[1] M. F. Atiyah, V. K. Patodi and I. M. Singer, Spectral asymmetry and Riemannian geometry I, Math. Proc. Cambridge Philos. Soc., 77 (1975), 43-69.
[2] M. F. Atiyah and I. M. Singer, The index of elliptic operators III, Ann. of Math., 87 (1968), 546-604.
[3] H. Donnelly, Eta invariants for G-spaces, Indianna Math. J., 27 (1978), 889-918.
[4] A. Futaki, An obstruction to the existence of Einstein-Kähler metrics, Invent. Math., 73 (1983), 437-443.
[5] A. Futaki and S. Morita, Invariant polynomials of the automorphism group of a compact complex manifold, J. Differential Geom., 21 (1985), 135-142.
[6] P. B. Gilkey, Invariance Theory, the Heat Equation, and the Atiyah-Singer Index Theorem, Math. Lecture Series 11, Publish or Perish, 1984.
[7] R. Harvey and H. B. Lawson, Jr., A constellation of minimal varieties defined over the group $G_{2}$, Lecture Notes in Pure and App. Math. 48 (1978), 43-59.
[8] S. Morita, Almost complex manifolds and Hirzebruch invariant for isolated singularities in complex spaces, Math. Ann., 211 (1974), 245-260.
[9] K. Tsuboi, The lifted Futaki invariants and the spin $^{c}$-Dirac operators, Osaka J. Math., 32 (1995), 207-225.

