# ESTIMATES OF FUNCTIONS WITH VANISHING PERIODIZATIONS 

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#### Abstract

We prove that if a function $f \in L^{p}\left(\mathbb{R}^{d}\right)$ has vanishing periodizations then $\|f\|_{p^{\prime}} \lesssim\|f\|_{p}$, provided $1 \leq p<2 d /(d+2)$ and $d \geq 3$.


## 1. Introduction

Let $f \in L^{1}\left(\mathbb{R}^{d}\right)$. Define a family of its periodizations with respect to a rotated integer lattice by

$$
\begin{equation*}
g_{\rho}(x)=\sum_{\nu \in \mathbb{Z}^{d}} f(\rho(x-\nu)) \tag{1}
\end{equation*}
$$

for all rotations $\rho \in \mathrm{SO}(d)$. We have the trivial estimate $\left\|g_{\rho}\right\|_{1} \leq\|f\|_{1}$, and $\widehat{g_{\rho}}(m)=\hat{f}(\rho m)$, where $m=\left(m_{1}, \ldots, m_{d}\right) \in \mathbb{Z}^{d}$. The author has shown recently that $g_{\rho}$ is in $L^{2}\left([0,1]^{d} \times \mathrm{SO}(d)\right)$ if and only if $f \in L^{2}\left(\mathbb{R}^{d}\right)$, provided $d \geq 5$. The requirement $f \in L^{1}\left(\mathbb{R}^{d}\right)$ can be replaced by $f \in L^{p}\left(\mathbb{R}^{d}\right)$ for a certain range of $p$; for details see [6] and [7].

The main object of our research are functions $f$ whose periodizations $g_{\rho}$ vanish identically for a.e. rotations $\rho \in \mathrm{SO}(d)$. This property is equivalent to the statement that $\hat{f}$ vanishes on all spheres of radius $|m|=\left(m_{1}^{2}+\cdots+m_{d}^{2}\right)^{1 / 2}$, where $m \in \mathbb{Z}^{d}$. Such functions are closely related to the Steinhaus tiling set problem (see [4] and [5]): Does there exists a (measurable) set $E \subset \mathbb{R}^{d}$ such that every rotation and translation of $E$ contains exactly one integer lattice point? M. Kolountzakis [4] showed that if $f \in L^{1}$ and $|x|^{\alpha} f(x) \in L^{1}$ for a certain $\alpha>0$ and $f$ has constant periodizations, then $\hat{f} \in L^{1}$ in the case of dimension $d=2$. Kolountzakis and Wolff [5, Theorem 1] proved that if the periodizations of a function from $L^{1}\left(\mathbb{R}^{d}\right)$ are constant, then the function is continuous and, in fact, bounded, provided that the dimension $d$ is at least three. Here we generalize the latter result for functions $f$ in $L^{1}\left(\mathbb{R}^{d}\right) \cap L^{p}\left(\mathbb{R}^{d}\right)$ :

[^0]Theorem 1. Let $d \geq 3$ and let $f \in L^{1}\left(\mathbb{R}^{d}\right) \cap L^{p}\left(\mathbb{R}^{d}\right), 1 \leq p<2 d /(d+2)$, have identically vanishing periodizations. Then $f \in L^{p^{\prime}}\left(\mathbb{R}^{d}\right)$, and

$$
\|f\|_{p^{\prime}} \leq C\|f\|_{p}
$$

where $C$ depends only on $d$ and $p$.
The main reason for the condition $d \geq 3$ is due to the famous result of Lagrange stating that every positive integer can be represented as a sum of four squares, and that every integer of the form $8 k+1$ can be written as a sum of three squares. Since relatively few integers can be represented as sums of two squares, we will show in Section 3 that the result of Kolountzakis and Wolff does not hold if $d=2$. This is why there is no analogous theorem for $d=2$. Another reason why the dimension $d$ has to be at least 3 is because we consider the family of periodizations with respect to the group of rotations $\mathrm{SO}(d)$. This leads to estimates involving the decay of spherical harmonics. For $d=2$ the rate of decay is not fast enough, although it is almost fast enough. In the case $d=2$ the range for $p$ in the theorem becomes $1 \leq p<1$, and hence is empty.

REmARK 1. There is no essential difference between the case of identically vanishing periodizations and the case where the functions $g_{\rho}$ are trigonometric polynomials of uniformly bounded degrees for all $\rho \in \mathrm{SO}(d)$.

Corollary 1. If $p \leq r \leq p^{\prime}$, then under the conditions of Theorem 1 we have

$$
\|f\|_{r} \leq C\|f\|_{p}
$$

where $C$ depends only on $d$ and $p$.
We will show in Section 3 that the range of $r$ in this result is sharp.
We will use the notation $x \lesssim y$ if $x \leq C y$ for some constant $C>0$ independent from $x$ and $y$, and we write $x \sim y$ if $x \lesssim y$ and $y \lesssim x$ both hold.

## 2. Proof of the theorem

We define functions $h, h_{1}, h_{2}: \mathbb{R}^{d} \times \mathbb{R}^{+} \rightarrow \mathbb{C}$ by

$$
\begin{align*}
h(y, t) & =\int \hat{f}(\xi) e^{i 2 \pi y \cdot \xi} d \sigma_{t}(\xi)  \tag{2}\\
& =\int_{\mathbb{R}^{d}} f(x) \widehat{d \sigma_{t}}(y-x) d x \\
& =\int_{\mathbb{R}^{d}} f(y-x) \widehat{d \sigma_{t}}(x) d x, \\
h_{1}(y, t) & =\int_{|x| \leq 1} f(y-x) \widehat{d \sigma}_{t}(x) d x, \tag{3}
\end{align*}
$$

$$
\begin{equation*}
h_{2}(y, t)=\int_{|x|>1} f(y-x) \widehat{d \sigma_{t}}(x) d x \tag{4}
\end{equation*}
$$

where $d \sigma_{t}$ is the Lebesgue surface measure on a sphere of radius $t$. Clearly, $h=h_{1}+h_{2}$. To proceed further we will need certain technical estimates involving the functions $h_{1}$ and $h_{2}$; these are given in two lemmas below. The proof of the theorem itself begins after Remark 2 following Lemma 2. The Fourier transforms in the two lemmas below are taken with respect to variable $t$, except in the second part of the proof of Lemma 2. The $L^{p^{\prime}}$ norms are taken with respect to the variable $y$. We will use some techniques of Kolountzakis and Wolff [5] and Kovrijkine [6], [7].

Lemma 1. Let $q: \mathbb{R} \rightarrow \mathbb{R}$ be a Schwartz function supported in $[1 / 2,2]$, let $f \in L^{p}\left(\mathbb{R}^{d}\right)$, where $1 \leq p \leq 2$, and let $b \in[0,1)$. Define $H_{1, N}: \mathbb{R}^{d} \times \mathbb{R} \rightarrow \mathbb{C}$ by

$$
H_{1, N}(y, t)=\frac{1}{\sqrt{t+b}} h_{1}(y, \sqrt{t+b}) q\left(\frac{\sqrt{t+b}}{N}\right)
$$

Then

$$
\begin{equation*}
\sum_{l \geq 0} \sum_{\nu \neq 0}\left\|\hat{H}_{1,2^{l}}(y, \nu)\right\|_{p^{\prime}} \leq C\|f\|_{p} \tag{5}
\end{equation*}
$$

where $C$ depends only on $q$ and $d$.
Proof. It will be enough to show that

$$
\begin{equation*}
\sum_{\nu \neq 0}\left\|\hat{H}_{1, N}(y, \nu)\right\|_{p^{\prime}} \leq \frac{C\|f\|_{p}}{N} \tag{6}
\end{equation*}
$$

We have

$$
\begin{equation*}
\left|\hat{H}_{1, N}(y, \nu)\right| \leq \frac{C}{|\nu|^{k}} \int\left|\frac{\partial^{k}}{\partial t^{k}} H_{1, N}(y, t)\right| d t \tag{7}
\end{equation*}
$$

for $\nu \neq 0$. Applying Minkowski's inequality to (7) we obtain

$$
\begin{equation*}
\left\|\hat{H}_{1, N}(y, \nu)\right\|_{p^{\prime}} \leq \frac{C}{|\nu|^{k}} \int\left\|\frac{\partial^{k}}{\partial t^{k}} H_{1, N}(y, t)\right\|_{L^{p^{\prime}}(d y)} d t \tag{8}
\end{equation*}
$$

We need to estimate the integrand on the right side of (8). To do so we will first estimate the $L^{p^{\prime}}$ norm of derivatives of $h_{1}(y, t)$ when $t \geq 1$. We have

$$
\begin{equation*}
\left\|\frac{\partial^{k}}{\partial t^{k}} h_{1}(y, t)\right\|_{p^{\prime}} \lesssim t^{d-1}\|f\|_{p} \tag{9}
\end{equation*}
$$

with an implicit constant depending only on $k$ and $d$. In order to obtain (9), we rewrite the definition (3) of $h_{1}$ as

$$
\begin{aligned}
h_{1}(y, t) & =\int_{|x| \leq 1} f(y-x) \widehat{d \sigma_{t}}(x) d x \\
& =t^{d-1} \int_{\mathbb{R}^{d}} f(y-x) \cdot \chi_{\{|x| \leq 1\}} \int_{|\xi|=1} e^{-i 2 \pi t x \cdot \xi} d \sigma(\xi) d x
\end{aligned}
$$

differentiate the last expression $k$ times, and apply Young's inequality.
We can easily prove by induction that

$$
\begin{equation*}
\frac{d^{k}}{d t^{k}}\left(\frac{h_{1}(\sqrt{t+b})}{\sqrt{t+b}}\right)=\sum_{i=0}^{k} C_{i, k} \frac{h_{1}^{(i)}(\sqrt{t+b})}{(\sqrt{t+b})^{2 k+1-i}} \tag{10}
\end{equation*}
$$

Combining (10) and (9) we obtain for $t \sim N^{2}$

$$
\begin{equation*}
\left\|\frac{\partial^{k}}{\partial t^{k}}\left(\frac{h_{1}(y, \sqrt{t+b})}{\sqrt{t+b}}\right)\right\|_{p^{\prime}} \leq C N^{d-k-2}\|f\|_{p} \tag{11}
\end{equation*}
$$

with $C$ depending only on $k$ and $d$.
Since $q((\sqrt{t+b}) / N)=q\left(\sqrt{t^{\prime}+b^{\prime}}\right)=\tilde{q}\left(t^{\prime}\right)$ with $t^{\prime}=t / N^{2}$ and $b^{\prime}=b / N^{2}$ and $\tilde{q}\left(t^{\prime}\right)$ is a Schwartz function supported in $t^{\prime} \sim 1$, we have

$$
\begin{equation*}
\left|\frac{d^{k}}{d t^{k}} q\left(\frac{(\sqrt{t+b})}{N}\right)\right|=N^{-2 k}\left|\frac{d^{k}}{d t^{\prime k}} \tilde{q}\left(t^{\prime}\right)\right| \leq C N^{-2 k} \tag{12}
\end{equation*}
$$

with $C$ depending only on $k$ and $q$.
Now $q((\sqrt{t+b}) / N)$ is supported in $t \sim N^{2}$. Hence we obtain from (11) and (12)

$$
\begin{align*}
\left\|\frac{\partial^{k}}{\partial t^{k}} H_{1, N}(y, t)\right\|_{p^{\prime}} & =\left\|\frac{d^{k}}{d t^{k}}\left(\frac{h_{1}(y, \sqrt{t+b})}{\sqrt{t+b}} q\left(\frac{\sqrt{t+b}}{N}\right)\right)\right\|_{p^{\prime}}  \tag{13}\\
& \leq C N^{d-2-k}\|f\|_{p}
\end{align*}
$$

with $C$ depending only on $k, d$ and $q$. Since $H_{1, N}(y, t)$ is also supported in $t \sim N^{2}$, we have

$$
\int\left\|\frac{\partial^{k}}{\partial t^{k}} H_{1, N}(y, t)\right\|_{L^{p^{\prime}}(d y)} d t \leq C N^{d-k}\|f\|_{p}
$$

Substituting this estimate into (8) we obtain

$$
\begin{equation*}
\left\|\hat{H}_{1, N}(y, \nu)\right\|_{p^{\prime}} \leq \frac{C N^{d-k}\|f\|_{p}}{|\nu|^{k}} \tag{14}
\end{equation*}
$$

for every $\nu \neq 0$.
Summing (14) over all $\nu \neq 0$ and putting $k=d+1$ we obtain our desired result

$$
\sum_{\nu \neq 0}\left\|\hat{H}_{1, N}(y, \nu)\right\|_{p^{\prime}} \leq \frac{C\|f\|_{p}}{N}
$$

where $C$ depends only on $q$ and $d$. The assertion of the lemma follows by summing over dyadic values $N$.

The next lemma will be proven using the methods of the Stein-Tomas restriction theorem (see [1, p. 104]).

Lemma 2. Let $q: \mathbb{R} \rightarrow \mathbb{R}$ be a Schwartz function supported in $[1 / 2,2]$, let $f \in L^{p}\left(\mathbb{R}^{d}\right)$, where $1 \leq p<2 d /(d+2)$ and let $b \in[0,1)$. Define $H_{2, N}$ : $\mathbb{R}^{d} \times \mathbb{R} \rightarrow \mathbb{C}$ by

$$
H_{2, N}(y, t)=\frac{1}{\sqrt{t+b}} h_{2}(y, \sqrt{t+b}) q\left(\frac{\sqrt{t+b}}{N}\right)
$$

Then we have

$$
\begin{equation*}
\sum_{\nu \neq 0}\left\|\sum_{l \geq 0} \hat{H}_{2,2^{l}}(y, \nu)\right\|_{p^{\prime}} \leq C\|f\|_{p} \tag{15}
\end{equation*}
$$

with $C$ depending only on $p, q$ and $d$.
Proof. We have

$$
\begin{align*}
& \hat{H}_{2, N}(y, \nu)=\int H_{2, N}(y, t) e^{-i 2 \pi \nu t} d t  \tag{16}\\
& =2 e^{i 2 \pi \nu b} \int N q(t) h_{2}(y, t N) e^{-i 2 \pi \nu(N t)^{2}} d t \\
& =2 e^{i 2 \pi \nu b} \int N q(t) e^{-i 2 \pi \nu(N t)^{2}} \int_{|x|>1} f(y-x) \widehat{d \sigma_{N t}}(x) d x d t \\
& =2 e^{i 2 \pi \nu b} \int_{|x|>1} f(y-x) \int N q(t) e^{-i 2 \pi \nu(N t)^{2}}(N t)^{d-1} \widehat{d \sigma}(N t x) d t d x \\
& =\left(D_{N, \nu} * f\right)(y),
\end{align*}
$$

where

$$
\begin{equation*}
D_{N, \nu}(x)=2 e^{i 2 \pi \nu b} \chi_{\{|x|>1\}} \int N q(t) e^{-i 2 \pi \nu(N t)^{2}}(N t)^{d-1} \widehat{d \sigma}(N t x) d t \tag{17}
\end{equation*}
$$

Set

$$
\begin{equation*}
K_{\nu}(x)=\sum_{l \geq 0} D_{2^{l}, \nu}(x) \tag{18}
\end{equation*}
$$

We need to estimate

$$
\left\|\sum_{l \geq 0} \hat{H}_{2,2^{l}}(y, \nu)\right\|_{p^{\prime}}=\left\|K_{\nu} * f\right\|_{p^{\prime}}
$$

If $p^{\prime}=\infty$ or $p^{\prime}=2$, then

$$
\begin{aligned}
\left\|K_{\nu} * f\right\|_{\infty} & \leq\left\|K_{\nu}\right\|_{\infty}\|f\|_{1} \\
\left\|K_{\nu} * f\right\|_{2} & \leq\left\|\hat{K}_{\nu}\right\|_{\infty}\|f\|_{2}
\end{aligned}
$$

We first show that

$$
\begin{align*}
\left\|K_{\nu}\right\|_{\infty} & \leq\left\|\sum_{l \geq 0}\left|D_{2^{l}, \nu}\right|(x)\right\|_{\infty}  \tag{19}\\
& \leq C|\nu|^{-d / 2}
\end{align*}
$$

To this end we need to estimate $D_{N, \nu}$.
We will use the well-known fact that $\widehat{d \sigma}(x)=\operatorname{Re}(B(|x|))$ with $B(r)=$ $a(r) e^{i 2 \pi r}$ and $a(r)$ satisfying

$$
\begin{equation*}
\left|a^{k}(r)\right| \leq \frac{C}{r^{(d-1) / 2+k}} \tag{20}
\end{equation*}
$$

with $C$ depending only on $k$ and $d$. We now estimate the integral in (17) with $B(|x|)$ instead of $\widehat{d \sigma}(x)$ :

$$
\text { (21) } \begin{array}{rl}
\int N & q(t) e^{-i 2 \pi \nu(N t)^{2}}(N t)^{d-1} a(N|x| t) e^{i 2 \pi N|x| t} d t \\
& =\frac{N^{\frac{d+1}{2}}}{|x|^{\frac{d-1}{2}}} \int q(t) e^{-i 2 \pi \nu(N t)^{2}} t^{d-1} a(N|x| t)(N|x|)^{\frac{d-1}{2}} e^{i 2 \pi N|x| t} d t \\
& =\frac{N^{\frac{d+1}{2}}}{|x|^{\frac{d-1}{2}}} e^{i 2 \pi \frac{|x|^{2}}{4 \nu}} \int q(t) a(N|x| t)(N|x|)^{\frac{d-1}{2}} t^{d-1} e^{-i 2 \pi \nu N^{2}\left(t-\frac{|x|}{2 \nu N}\right)^{2}} d t \\
& =\frac{N^{\frac{d+1}{2}}}{|x|^{\frac{d-1}{2}}} e^{i 2 \pi \frac{|x|^{2}}{4 \nu}} \int \phi(t,|x|) e^{-i 2 \pi \nu N^{2}\left(t-\frac{|x|}{2 \nu N}\right)^{2}} d t
\end{array}
$$

where $\phi(t,|x|)=q(t) a(N|x| t)(N|x|)^{(d-1) / 2} t^{d-1}$ is a Schwartz function with respect to the variable $t$ supported in $[1 / 2,2]$, which, by $(20)$, is bounded, together with each derivative, uniformly in $t,|x| \geq 1$, and $N$. Note that we used here the fact that $N|x| \geq 1$. We can say even more. Let $|x|=c \cdot r$, where $c \geq 2$ and $r \geq 1 / 2$. Then all partial derivatives of $\phi(t, c \cdot r)$ with respect to $t$ and $r$ are also bounded uniformly in $t, r, c$ and $N$. Hence $\phi(t, c \cdot t)$ is a Schwartz function supported in $[1 / 2,2]$ which is bounded, together with each derivative, uniformly in $t, c$ and $N$. We will use this fact later to estimate $\hat{K}_{\nu}$ 。

Fix some $x$ with $|x| \geq 1$. In the calculations below we will write $\phi(t)$ instead of $\phi(t,|x|)$ for simplicity. From the method of stationary phase (see
[3, Theorem 7.7.3]) it follows that if $k \geq 1$ then

$$
\begin{gather*}
\left|\int \phi(t) e^{-i 2 \pi \nu N^{2}\left(t-\frac{|x|}{2 \nu N}\right)^{2}} d t-\sum_{j=0}^{k-1} c_{j}\left(\nu N^{2}\right)^{-j-1 / 2} \phi^{(2 j)}\left(\frac{|x|}{2 \nu N}\right)\right|  \tag{22}\\
\leq c_{k}\left(|\nu| N^{2}\right)^{-k-1 / 2}
\end{gather*}
$$

with some constants $c_{j}$.
Since $\phi$ is supported in $[1 / 2,2]$, we conclude from (22) that

$$
\left|\int \phi(t) e^{-i 2 \pi \nu N^{2}\left(t-\frac{|x|}{2 \nu N}\right)^{2}} d t\right| \leq \begin{cases}C\left(|\nu| N^{2}\right)^{-1 / 2} & \text { if } N \in\left[\frac{|x|}{4 \nu}, \frac{|x|}{\nu}\right]  \tag{23}\\ C_{k}\left(|\nu| N^{2}\right)^{-k-1 / 2} & \text { if } N \notin\left[\frac{|x|}{4 \nu}, \frac{|x|}{\nu}\right] .\end{cases}
$$

Replacing in (17) $\widehat{d \sigma}(x)$ by $(B(|x|)+\bar{B}(|x|)) / 2$, it follows from (23) that

$$
\left|D_{N, \nu}(x)\right| \leq \frac{N^{\frac{d+1}{2}}}{|x|^{\frac{d-1}{2}}} \begin{cases}C\left(|\nu| N^{2}\right)^{-1 / 2} & \text { if } N \in\left[\frac{|x|}{4|\nu|}, \frac{|x|}{|\nu|}\right]  \tag{24}\\ C_{k}\left(|\nu| N^{2}\right)^{-k-1 / 2} & \text { if } N \notin\left[\frac{|x|}{4|\nu|}, \frac{|x|}{|\nu|}\right]\end{cases}
$$

The number of dyadic $N \in\left[\frac{|x|}{4 \nu}, \frac{|x|}{\nu}\right]$ is at most 3. Therefore choosing $k \geq$ $(d-1) / 2$ and summing (24) over all dyadic $N$ we have

$$
\left|K_{\nu}(x)\right| \leq \sum_{l \geq 0}\left|D_{2^{l}, \nu}(x)\right| \leq C|\nu|^{-d / 2}
$$

with $C$ depending only on $d$ and $q$. Thus we have proved (19).
We now show that

$$
\begin{equation*}
\left\|\hat{K}_{\nu}\right\|_{\infty} \leq\left\|\sum_{l \geq 0}\left|\hat{D}_{2^{l}, \nu}\right|(y)\right\|_{\infty} \leq C \tag{25}
\end{equation*}
$$

Since supp $\phi \in[1 / 2,2]$, we can rewrite (22) using a stronger version of the method of stationary phase (see [3, Theorems 7.6.4, 7.6.5, 7.7.3]).

$$
\begin{aligned}
& \left|\int \phi(t) e^{-i 2 \pi \nu N^{2}\left(t-\frac{|x|}{2 \nu N}\right)^{2}} d t-\sum_{j=0}^{k-1} c_{j}\left(\nu N^{2}\right)^{-j-1 / 2} \phi^{(2 j)}\left(\frac{|x|}{2 \nu N}\right)\right| \\
& \quad \leq \frac{c_{k}\left(|\nu| N^{2}\right)^{-k-1 / 2}}{\max \left(1, \frac{|x|}{8 N|\nu|}\right)^{k}}
\end{aligned}
$$

where the numbers $c_{j}$ are suitable constants. Therefore, for $\nu>0$,

$$
\begin{equation*}
D_{N, \nu}(x)=\chi_{\{|x|>1\}} \frac{N^{\frac{d+1}{2}}}{|x|^{\frac{d-1}{2}}} e^{i 2 \pi \frac{|x|^{2}}{4 \nu}} \sum_{j=0}^{k-1} c_{j}\left(\nu N^{2}\right)^{-j-1 / 2} \phi^{(2 j)}\left(\frac{|x|}{2 \nu N}\right)+\phi_{k}(x) \tag{26}
\end{equation*}
$$

where

$$
\left|\phi_{k}(x)\right| \leq \chi_{\{|x|>1\}} \frac{N^{\frac{d+1}{2}}}{|x|^{\frac{d-1}{2}}} \frac{c_{k}\left(|\nu| N^{2}\right)^{-k-1 / 2}}{\max \left(1, \frac{|x|}{8 N|\nu|}\right)^{k}} .
$$

If $\nu<0$ then we simply replace $\phi^{(2 j)}(|x| /(2 \nu N))$ by $\bar{\phi}^{(2 j)}(-|x| /(2 \nu N))$. We further assume that $\nu>0$. Choosing $k \geq(d+2) / 2$ we have

$$
\begin{equation*}
\left\|\hat{\phi}_{k}\right\|_{\infty} \leq\left\|\phi_{k}\right\|_{1}=\int_{|x| \leq 8 \nu N}\left|\phi_{k}\right| d x+\int_{|x|>8 \nu N}\left|\phi_{k}\right| d x \leq \frac{C}{N} \tag{27}
\end{equation*}
$$

where $C$ depends only on $d$ and $q$. We can ignore the factor $\chi_{\{|x|>1\}}$ in front of the sum in (26) because if $|x| /(2 \nu N) \in[1 / 2,2]$, then $|x| \geq \nu N \geq 1$. We will only consider the term $j=0$ in the sum; the other terms can be treated similarly. The Fourier transform of

$$
\frac{N^{\frac{d+1}{2}}}{|x|^{\frac{d-1}{2}}} e^{i 2 \pi \frac{|x|^{2}}{4 \nu}}\left(\nu N^{2}\right)^{-1 / 2} \phi\left(\frac{|x|}{2 \nu N}\right)
$$

at a point $y$ is equal to

$$
\begin{align*}
& N^{\frac{d+1}{2}}(2 \nu N)^{\frac{d+1}{2}}\left(\nu N^{2}\right)^{-1 / 2} \int_{\mathbb{R}^{d}} \psi(|x|) e^{i 2 \pi \nu N^{2}|x|^{2}} e^{-i 2 \pi 2 \nu N x \cdot y} d x  \tag{28}\\
&=C\left(\nu N^{2}\right)^{d / 2} e^{-i 2 \pi \nu|y|^{2}} \int_{\mathbb{R}^{d}} \psi(|x|) e^{i 2 \pi \nu N^{2}\left|x-\frac{y}{N}\right|^{2}} d x
\end{align*}
$$

where $\psi(t)=\phi(t, 2 \nu N t) t^{-(d-1) / 2}$ is a Schwartz function supported in [1/2,2] whose derivatives and the function itself are bounded uniformly in $t, \nu$ and $N$ (see the remark after (21)). The same holds for the partial derivatives of $\psi(|x|)$. Applying the stationary phase method for $\mathbb{R}^{d}$ (see [3, Theorem 7.7.3]), we get

$$
\left|\int_{\mathbb{R}^{d}} \psi(|x|) e^{i 2 \pi \nu N^{2}\left|x-\frac{y}{N}\right|^{2}} d x\right| \leq \begin{cases}C\left(\nu N^{2}\right)^{-d / 2} & \text { if } N \in\left[\frac{|y|}{2}, 2|y|\right]  \tag{29}\\ C_{k}\left(\nu N^{2}\right)^{-k-d / 2} & \text { if } N \notin\left[\frac{|y|}{2}, 2|y|\right]\end{cases}
$$

Therefore the absolute value of (28) can be bounded from above by

$$
\leq \begin{cases}C & \text { if } N \in\left[\frac{|y|}{2}, 2|y|\right]  \tag{30}\\ C_{k}\left(\nu N^{2}\right)^{-k} & \text { if } N \notin\left[\frac{|y|}{2}, 2|y|\right]\end{cases}
$$

Similar inequalities hold for the Fourier transforms of the other terms in the sum in (26). The number of dyadic values $N \in[|y| / 2,2|y|]$ is bounded by 3 . Using (27), choosing $k \geq 1$ in (30), and summing over all dyadic $N$, we obtain

$$
\begin{equation*}
\sum_{l \geq 0}\left|\hat{D}_{2^{l}, \nu}(y)\right| \leq C \tag{31}
\end{equation*}
$$

with $C$ depending only on $d$ and $q$, provided $\nu \neq 0$. Thus we have proved (25).

Using (19) and (25) and interpolating between $p=1$ and $p=2$, we obtain

$$
\begin{equation*}
\left\|K_{\nu} * f\right\|_{p^{\prime}} \leq C|\nu|^{-\alpha_{p}}\|f\|_{p} \tag{32}
\end{equation*}
$$

where $\alpha_{p}=(d / 2)(2-p) / p$. Note that $\alpha_{p}>1$ if $p<2 d /(d+2)$. Summing (32) over all $\nu \neq 0$ yields the desired inequality

$$
\sum_{\nu \neq 0}\left\|\sum_{l \geq 0} \hat{H}_{2,2^{l}}(y, \nu)\right\|_{p^{\prime}} \leq C\|f\|_{p}
$$

Remark 2. It is clear from the proof that we have the same inequality if the summation over $l \geq 0$ is replaced by a summation over any subset of the nonnegative integers.

We are now in a position to proceed with the proof of the theorem. Let $q: \mathbb{R} \rightarrow \mathbb{R}$ be a fixed nonnegative Schwartz function supported in $[1 / 2,2]$ such that

$$
q(t)+q(t / 2)=1
$$

when $t \in[1,2]$. It follows that

$$
\begin{equation*}
\sum_{l \geq 0} q\left(\frac{t}{2^{l}}\right)=1 \tag{33}
\end{equation*}
$$

when $t \geq 1$. Define

$$
q_{0}(t)=1-\sum_{l \geq 0} q\left(\frac{t}{2^{l}}\right)
$$

for $t \geq 0$. It is clear that $q_{0}(|x|)$ is a Schwartz function supported in $|x| \leq 1$. Let $\psi(t)=q_{0}(t)+q(t)$. Then

$$
\psi_{k}(t)=\psi\left(\frac{t}{2^{k}}\right)=q_{0}(t)+\sum_{l \geq 0}^{k} q\left(\frac{t}{2^{l}}\right)
$$

and $\psi(|x|)$ is a Schwartz function supported in $|x| \leq 2$ such that $\psi(|x|)=1$ if $|x| \leq 1$. Therefore

$$
\int \hat{f}(x) e^{2 \pi x \cdot y} \psi\left(\frac{|x|}{2^{k}}\right) d x=\left(f * \widehat{\psi_{k}}\right)(y)
$$

converges to $f$ in $L^{p}$ as $k \rightarrow \infty$. To prove that $f \in L^{p^{\prime}}$ and $\|f\|_{p^{\prime}} \lesssim\|f\|_{p}$ it will be enough to show that

$$
\left\|f * \widehat{\psi_{k}}\right\|_{p^{\prime}} \leq C\|f\|_{p}
$$

an application of Fatou's lemma to a subsequence of $f * \widehat{\psi_{k}}$ converging a.e. to $f$ will then yield the assertion.

We have

$$
\begin{align*}
\left(f * \widehat{\psi_{k}}\right)(y) & =\left(f * \widehat{q_{0}}\right)(y)+\sum_{l \geq 0}^{k} \int \hat{f}(x) e^{2 \pi x \cdot y} q\left(\frac{|x|}{2^{l}}\right) d x  \tag{34}\\
& =\left(f * \widehat{q_{0}}\right)(y)+\sum_{l \geq 0}^{k} \int_{0}^{\infty} q\left(\frac{t}{2^{l}}\right) \int \hat{f}(\xi) e^{i 2 \pi y \cdot \xi} d \sigma_{t}(\xi) d t \\
& =\left(f * \widehat{q_{0}}\right)(y)+\sum_{l \geq 0}^{k} \int_{0}^{\infty} q\left(\frac{t}{2^{l}}\right) h(y, t) d t .
\end{align*}
$$

By Young's inequality we have

$$
\begin{equation*}
\left\|f * \widehat{q_{0}}\right\|_{p^{\prime}} \lesssim\|f\|_{p} \tag{35}
\end{equation*}
$$

for $1 \leq p \leq 2$. It thus remains to estimate the sum over $l$.
A well-known result in number theory due to Lagrange states that every positive integer can be represented as a sum of four squares (see [2, p. 25]). Moreover, there exists an infinite arithmetic progression of positive integers (e.g., integers of the form $8 n+1$ ) which can be represented as sums of three squares (see [2, p. 38]). We will only use the latter result. By rescaling we can assume that $\hat{f}$ vanishes on all spheres of radius $\sqrt{n+b}$, where $n$ is a nonnegative integer and $0<b<1$ is a fixed number. Therefore $h(y, \sqrt{n+b})=$ 0 for all $y \in \mathbb{R}^{d}$. Making a change of variables and keeping in mind that $q$ is supported in $[1 / 2,2]$, we rewrite the terms in the sum as follows:

$$
\int_{0}^{\infty} q\left(\frac{t}{N}\right) h(y, t) d t=\int \frac{1}{2 \sqrt{t+b}} q\left(\frac{\sqrt{t+b}}{N}\right) h(y, \sqrt{t+b}) d t
$$

An application of Poisson's summation formula gives

$$
\begin{aligned}
0 & =\sum_{n} \frac{1}{\sqrt{n+b}} q\left(\frac{\sqrt{n+b}}{N}\right) h(y, \sqrt{n+b}) \\
& =\sum_{\nu}\left(\frac{1}{\sqrt{t+b}} q\left(\frac{\sqrt{t+b}}{N}\right) h(y, \sqrt{t+b})\right)^{\wedge}(\nu) \\
& =\int \frac{1}{\sqrt{t+b}} q\left(\frac{\sqrt{t+b}}{N}\right) h(y, \sqrt{t+b}) d t+\sum_{\nu \neq 0} \hat{H}_{1, N}(y, \nu)+\sum_{\nu \neq 0} \hat{H}_{2, N}(y, \nu)
\end{aligned}
$$

where

$$
H_{i, N}(y, t)=\frac{1}{\sqrt{t+b}} q\left(\frac{\sqrt{t+b}}{N}\right) h_{i}(y, \sqrt{t+b}), \quad i=1,2
$$

Applying Lemmas 1 and 2, along with Remark 2, we can bound the sum by

$$
\begin{aligned}
&\left\|\sum_{l \geq 0}^{k} \int_{0}^{\infty} q\left(\frac{t}{2^{l}}\right) h(y, t) d t\right\|_{p^{\prime}} \leq \sum_{l \geq 0} \sum_{\nu \neq 0}\left\|\hat{H}_{1,2^{l}}(y, \nu)\right\|_{p^{\prime}} \\
&+\sum_{\nu \neq 0}\left\|\sum_{l \geq 0}^{k} \hat{H}_{2,2^{l}}(y, \nu)\right\|_{p^{\prime}} \\
& \leq C\|f\|_{p}
\end{aligned}
$$

Combining (34), (35), and the last inequality, we obtain the desired inequality

$$
\left\|f * \widehat{\psi_{k}}\right\|_{p^{\prime}} \leq C\|f\|_{p}
$$

from which the statement of the theorem follows.
Remark 3. We say that a function $f \in L^{p}$ has vanishing periodizations if there exists a sequence of Schwartz functions $f_{k}$ with vanishing periodizations converging to $f$ in $L^{p}$. It follows from Theorem 1 that $f \in L^{p^{\prime}}$ and the functions $f_{k}$ converge to $f$ in $L^{p^{\prime}}$ if $d \geq 3$ and $1 \leq p<2 d /(d+2)$.

## 3. Counterexamples and open questions

When $d=1$ or $d=2$, Theorem 1 does not apply. The case $d=1$ is not interesting. We can easily construct examples of functions $f$ with vanishing periodizations such that their $L^{p}$ norms are not bounded by their $L^{q}$ norms, for any given pair $p \neq q$.

We now show that, when $d=2$, the assertion of Theorem 1 does not hold. More precisely, Lemma 3 below shows that if $1 \leq p<2$, then the inequality

$$
\|f\|_{p^{\prime}} \lesssim\|f\|_{p}
$$

does not hold for functions with vanishing periodizations. This lemma deals with a sequence of functions $f_{n}$ such that $\hat{f}_{n}$ vanishes on all circles of radius $\sqrt{l^{2}+k^{2}}$. Denote by $X_{2}$ the Banach space of functions from $L^{1}\left(\mathbb{R}^{2}\right)$ whose Fourier transforms vanish on all circles of radius $\sqrt{l^{2}+k^{2}}$, i.e.,

$$
X_{2}=\left\{f \in L^{1}\left(\mathbb{R}^{2}\right): \hat{f}(\mathbf{r})=0 \text { if }|\mathbf{r}|=\sqrt{l^{2}+k^{2}},(k, l) \in \mathbb{Z}^{2}\right\}
$$

The lemma depends in crucial way on the following fact from number theory (see [2, p. 22]):

The number of integers in $[n, 2 n]$ which can be represented as sums of two squares is $n \epsilon_{n}$, where $\epsilon_{n} \lesssim 1 / \ln ^{1 / 2} n \rightarrow 0$ as $n \rightarrow \infty$.

We only need that $\lim _{n \rightarrow \infty} \epsilon_{n}=0$.

Lemma 3. Let $1 \leq p<2$ and $d=2$. Then there exists a sequence of Schwartz functions $f_{n} \in X_{2}$ such that

$$
\lim _{n \rightarrow \infty} \frac{\left\|f_{n}\right\|_{p^{\prime}}}{\left\|f_{n}\right\|_{p}}=\infty
$$

Proof. Let $a_{1}<a_{2}<a_{3}<\cdots$ be the enumeration of the numbers $a_{m}=$ $\sqrt{l^{2}+k^{2}}$ in ascending order, and set $\delta_{m}=a_{m+1}-a_{m}$. As mentioned above, the number of $a_{m}$ in the interval $[\sqrt{n}, 2 \sqrt{n}]$ is $n \epsilon_{n}$. Let $a_{m_{0}}$ and $a_{m_{1}}$ denote, respectively, the smallest and largest elements $a_{m}$ in this interval. Then

$$
\sum_{m=m_{0}}^{m_{1}-1} \delta_{m}=a_{m_{1}}-a_{m_{0}} \sim \sqrt{n}
$$

Let

$$
\begin{equation*}
\delta=\frac{C}{\sqrt{n} \epsilon_{n}} \tag{36}
\end{equation*}
$$

with a small enough constant $C>0$ so that if

$$
M=\left\{m_{0} \leq m<m_{1}: \delta_{m} \geq \delta\right\}
$$

then

$$
\sqrt{n} \lesssim \sum_{m \in M} \delta_{m}
$$

This is possible since $m_{1}-m_{0} \sim n \epsilon_{n}$. Choose coordinate axes $x$ and $y$. We will construct functions $\hat{f}_{n}$ supported in $\bigcup_{m \in M} R_{m}$, where $R_{m}$ is a largest possible rectangle inscribed between circles of radius $a_{m}$ and $a_{m+1}$ with sides parallel to the coordinate axes. Then $R_{m}$ is of size $\sim \delta_{m} \times \sqrt{\delta_{m} a_{m}} \gtrsim \delta_{m} \times \sqrt{\delta \sqrt{n}} \gtrsim$ $\delta_{m} \times 1$. We split each rectangle $R_{m}$ further into $\left[\delta_{m} / \delta\right.$ ] smaller rectangles $r$ of the same size $\sim \delta \times 1$. The number of these rectangles $r$ is

$$
\begin{equation*}
N=\sum_{m \in M}\left[\frac{\delta_{m}}{\delta}\right] \sim \sum_{m \in M} \frac{\delta_{m}}{\delta} \sim \frac{\sqrt{n}}{1 / \sqrt{n} \epsilon_{n}}=n \epsilon_{n} \tag{37}
\end{equation*}
$$

since $\delta_{m} \geq \delta$ for $m \in M$. Enumerate these rectangles by $r_{k}, k=1, \ldots, N$. Let $r_{k}$ be centered at $\left(\lambda_{k}, 0\right)$. It is clear that $\left|\lambda_{k}-\lambda_{l}\right| \geq \delta$ for $k \neq l$. Let $\phi$ be a nonnegative Schwartz function on $\mathbb{R}$ supported in $[-1 / 2,1 / 2]$. Then $\check{\phi}(x) \geq C>0$ if $x$ is small enough. Define $\hat{f}_{n}$ by

$$
\begin{equation*}
\hat{f}_{n}(x, y)=\sum_{k=1}^{N} \phi\left(\frac{x-\lambda_{k}}{\delta}\right) \phi(y) \tag{38}
\end{equation*}
$$

The $k$ th term in (38) is supported in $r_{k}$. Therefore, $\hat{f}_{n}$ is a Schwartz function supported in $\bigcup_{m \in M} R_{m}$. Hence $\hat{f}_{n}$ vanishes on all circles of radius $a_{l}$. Taking
the inverse Fourier transform of (38), we get

$$
\begin{equation*}
f_{n}(\xi, \eta)=\delta \check{\phi}(\xi \delta) \check{\phi}(\eta) \sum_{k=1}^{N} e^{i \lambda_{k} \xi} \tag{39}
\end{equation*}
$$

Assume first that $p^{\prime}<\infty$. Then

$$
\begin{aligned}
\int\left|f_{n}(\xi, \eta)\right|^{p^{\prime}} d \xi d \eta & \geq\|\check{\phi}\|_{p^{\prime}}^{p^{\prime}} \delta^{p^{\prime}} \int_{|\xi| \leq\left(100^{-1}\right) / \sqrt{n}}|\check{\phi}(\xi \delta)|^{p^{\prime}}\left|\sum_{k=1}^{N} e^{i \lambda_{k} \xi}\right|^{p^{p^{\prime}}} d \xi \\
& \gtrsim \delta^{p^{\prime}} N^{p^{p^{\prime}}} \frac{1}{\sqrt{n}} \sim(\sqrt{n})^{p^{\prime}-1}
\end{aligned}
$$

where the second step follows from the bound

$$
\left|\sum_{k=1}^{N} e^{i \lambda_{k} \xi}\right| \geq\left|\sum_{k=1}^{N} \cos \left(\lambda_{k} \xi\right)\right| \gtrsim N
$$

since $\left|\lambda_{k} \xi\right| \leq 1 / 50$, and the third step follows from (36) and (37). Therefore

$$
\begin{equation*}
\left\|f_{n}\right\|_{p^{\prime}} \gtrsim(\sqrt{n})^{1 / p} \tag{40}
\end{equation*}
$$

If $p^{\prime}=\infty$ we obtain in a similar way that

$$
\begin{equation*}
\left\|f_{n}\right\|_{\infty} \geq\left|f_{n}(0)\right| \gtrsim \sqrt{n} \tag{41}
\end{equation*}
$$

We now estimate the $L^{p}$ norm from above. Set

$$
g(x)=\sum_{k=1}^{N} e^{i\left(\lambda_{k} / \delta\right) \xi}
$$

Since $\left|\left(\lambda_{k}-\lambda_{l}\right) / \delta\right| \geq \delta / \delta=1$ for $k \neq l$, we have

$$
\int_{I}|g|^{2} \sim N
$$

for any interval $I$ of length $4 \pi$ (see [8, Theorem 9.1]). Therefore,

$$
\begin{equation*}
\int_{I}|g|^{p} \leq|I|^{1-2 / p}\left(\int_{I}|g|^{2}\right)^{p / 2} \lesssim N^{p / 2} \tag{42}
\end{equation*}
$$

for any interval $I$ of length $4 \pi$. Since $\check{\phi}$ is a Schwartz function, we have

$$
|\check{\phi}(x)| \lesssim \frac{1}{1+x^{2}}
$$

Therefore

$$
\begin{aligned}
\int\left|f_{n}(\xi, \eta)\right|^{p} d \xi d \eta & =\|\check{\phi}\|_{p}^{p} \delta^{p-1} \int|\check{\phi}(\xi)|^{p} \cdot\left|\sum_{k=1}^{N} e^{i\left(\lambda_{k} / \delta\right) \xi}\right|^{p} d \xi \\
& =C \delta^{p-1} \sum_{l=-\infty}^{\infty} \int_{l 4 \pi}^{(l+1) 4 \pi}|\check{\phi}(\xi)|^{p} \cdot|g(\xi)|^{p} d \xi \\
& \lesssim \delta^{p-1} \sum_{l=-\infty}^{\infty} \frac{1}{\left(1+l^{2}\right)^{p}} N^{p / 2} \\
& \lesssim \sqrt{n} \epsilon_{n}^{1-p / 2}
\end{aligned}
$$

where the last step follows from (36) and (37). Hence

$$
\begin{equation*}
\left\|f_{n}\right\|_{p} \lesssim(\sqrt{n})^{1 / p} \epsilon_{n}^{(2-p) / 2 p} \tag{43}
\end{equation*}
$$

Dividing (40) by (43) we obtain the desired result

$$
\frac{\left\|f_{n}\right\|_{p^{\prime}}}{\left\|f_{n}\right\|_{p}} \geq \frac{(\sqrt{n})^{1 / p}}{(\sqrt{n})^{1 / p} \epsilon_{n}^{(2-p) /(2 p)}}=\frac{1}{\epsilon_{n}^{(2-p) /(2 p)}} \rightarrow \infty
$$

as $n \rightarrow \infty$ since $p<2$.
Corollary 2. There exists a function $f \in X_{2}$ such that

$$
\|f\|_{L^{\infty}(D(0,1))}=\infty
$$

Proof. It follows immediately from the lemma and (41) that if $p=1$ then

$$
\sup _{f \in X_{2}} \frac{\|f\|_{L^{\infty}(D(0,1))}}{\|f\|_{1}}=\infty
$$

We claim that there exists a function $f \in X_{2}$ such that $\|f\|_{L^{\infty}(D(0,1))}=\infty$. Suppose, to get a contradiction, that this is not true. Then the restriction operator

$$
T:\left.f \rightarrow f\right|_{D(0,1)}
$$

maps $X_{2}$ to $L^{\infty}(D(0,1))$. Note that if $f_{n} \rightarrow f$ in $L^{1}$ and $f_{n} \rightarrow g$ in $L^{\infty}(D(0,1))$, then $f=g$ a.e. on $D(0,1)$. An application of the Closed Graph Theorem shows that $T$ is a bounded operator acting from $X_{2}$ to $L^{\infty}(D(0,1))$. This contradicts Corollary 2, and thus proves our claim.

Obviously, this function $f$ is not continuous. Therefore the theorem of Kolountzakis and Wolff mentioned in the Introduction does not hold for dimension 2.

REMARK 4. It is an open problem whether, for $f \in X_{2}$, the inequality

$$
\|f\|_{r} \lesssim\|f\|_{p}
$$

holds when $1 \leq p<2$ and $p<r<p^{\prime}$.
We now show that the range of $r$ in Corollary 1 is sharp. We need to consider two cases, $r>p^{\prime}$ and $r<p$. In the first case the argument is similar to the one given in the previous lemma, and we therefore give only a sketch. We will deal with a sequence of functions $f_{n}$ such that the functions $\hat{f}_{n}$ vanish on all circles of radius $\sqrt{m_{1}^{2}+\cdots+m_{d}^{2}}$. Denote by $X_{d}$ the Banach space of functions from $L^{1}\left(\mathbb{R}^{d}\right)$ whose Fourier transforms vanish on all circles of radius $\sqrt{m_{1}^{2}+\cdots+m_{d}^{2}}$, i.e.,

$$
X_{d}=\left\{f \in L^{1}\left(\mathbb{R}^{d}\right): \hat{f}(\mathbf{r})=0 \text { if }|\mathbf{r}|=\sqrt{m_{1}^{2}+\cdots+m_{d}^{2}},\left(m_{1}, \ldots, m_{d}\right) \in \mathbb{Z}^{d}\right\}
$$

We will construct a sequence of Schwartz functions $f_{n}$ with Fourier transforms supported outside of spheres of radius $\sqrt{m}$. Therefore these functions automatically belong to $X_{d}$.

Lemma 4. Let $1<p \leq 2$ and $r>p^{\prime}$. Then there exists a sequence of Schwartz functions $f_{n} \in X$ such that

$$
\lim _{n \rightarrow \infty} \frac{\left\|f_{n}\right\|_{r}}{\left\|f_{n}\right\|_{p}}=\infty
$$

Proof. A maximal rectangle inscribed between spheres of radius $\sqrt{n}$ and $\sqrt{n+1}$ has dimensions $\sim(1 / \sqrt{n}) \times 1 \times 1 \times \cdots \times 1$. Let $r_{k}$ denote parallel identical rectangles inscribed between spheres of radius $\sqrt{n+k}$ and $\sqrt{n+k+1}$, for $k=0,1, \ldots, n-1$, with dimensions $\sim(1 / \sqrt{n}) \times 1 \times 1 \times \cdots \times 1$, and centered at $\left(\lambda_{k}, 0,0, \ldots, 0\right)$. It is clear that $\lambda_{k+1}-\lambda_{k} \sim 1 / \sqrt{n}$. Let $\phi$ be a nonnegative Schwartz function on $\mathbb{R}$ supported in $[-1 / 100,1 / 100]$. We have $\check{\phi}(x) \geq C>0$ when $x$ is small enough. Define $\hat{f}_{n}$ by

$$
\begin{equation*}
\hat{f}_{n}\left(x_{1}, x_{2}, \ldots, x_{d}\right)=\sum_{k=0}^{n-1} \phi\left(\left(x_{1}-\lambda_{k}\right) \sqrt{n}\right) \prod_{l=2}^{d} \phi\left(x_{l}\right) \tag{44}
\end{equation*}
$$

The $k$ th term in (44) is supported in $r_{k}$. Therefore, $\hat{f}_{n}$ is a Schwartz function vanishing on all spheres of radius $\sqrt{m}$. Taking the inverse Fourier transform of (44), we get

$$
\begin{equation*}
f_{n}\left(y_{1}, y_{2}, \ldots, y_{d}\right)=\prod_{l=2}^{d} \check{\phi}\left(y_{l}\right) \frac{1}{\sqrt{n}} \check{\phi}\left(\frac{y_{1}}{\sqrt{n}}\right) \sum_{k=0}^{n-1} e^{i \lambda_{k} y_{1}} \tag{45}
\end{equation*}
$$

Arguing as in the proof of Lemma 3, we obtain

$$
\left\|f_{n}\right\|_{r} \gtrsim(\sqrt{n})^{1 / r^{\prime}}
$$

and

$$
\left\|f_{n}\right\|_{p} \lesssim(\sqrt{n})^{1 / p}
$$

Therefore,

$$
\frac{\left\|f_{n}\right\|_{r}}{\left\|f_{n}\right\|_{p}} \gtrsim(\sqrt{n})^{\left(1 / p^{\prime}\right)-(1 / r)} \rightarrow \infty
$$

as $n \rightarrow \infty$, since $r>p^{\prime}$.
The case when $r<p$ is very simple. Let

$$
\hat{f}(x)=\phi\left(\frac{x-x_{0}}{\epsilon}\right)
$$

where $\phi$ is a Schwartz function supported in $B^{d}(0,1)$ so that $\hat{f}$ is supported in a small ball $B^{d}\left(x_{0}, \epsilon\right)$ placed between two fixed spheres of radius $\sqrt{n}$ and $\sqrt{n+1}$. Then $f(y)=\epsilon^{d} \check{\phi}(\epsilon y)$ and

$$
\frac{\|f\|_{r}}{\|f\|_{p}} \sim \frac{\epsilon^{d / r^{\prime}}}{\epsilon^{d / p^{\prime}}} \rightarrow \infty
$$

as $\epsilon \rightarrow 0$, since $r<p$. Note that we did not impose any restriction on $p$ here.
We now show that Theorem 1 does not hold if $p>2$. More precisely, let $p>2$ and $r \neq p$. Then the following inequality is not true for functions with vanishing periodizations:

$$
\|f\|_{r} \lesssim\|f\|_{p}
$$

Since we have already dealt with the case when $r<p$, we only need to consider the case $r>p$. The argument is almost the same as in the proof of Lemma 4. We construct a sequence of Schwartz functions $f_{n}$ with Fourier transforms vanishing on all spheres of radius $\sqrt{m}$ and such that $\left\|f_{n}\right\|_{r} \gtrsim(\sqrt{n})^{1 / r^{\prime}}$ and $\left\|f_{n}\right\|_{p} \leq\left\|\hat{f}_{n}\right\|_{p^{\prime}} \lesssim(\sqrt{n})^{1 / p^{\prime}}$. Therefore

$$
\frac{\left\|f_{n}\right\|_{r}}{\left\|f_{n}\right\|_{p}} \gtrsim(\sqrt{n})^{(1 / p)-(1 / r)} \rightarrow \infty
$$

Remark 5. Since Theorem 1 trivially holds for $p=2$, it is natural to expect that it also holds for $1 \leq p \leq 2$. However, the question whether the theorem holds for $2 d /(d+2) \leq p<2$ is still open.

Another interesting question is whether the inequality

$$
\begin{equation*}
\|\hat{f}\|_{p} \lesssim\|f\|_{p} \tag{46}
\end{equation*}
$$

holds for some range of $p<2$ if $f$ has vanishing periodizations. It would then follow that

$$
\begin{equation*}
\|\hat{f}\|_{r} \lesssim\|f\|_{p} \tag{47}
\end{equation*}
$$

for $p \leq r \leq p^{\prime}$. From Theorem 1 we see that (47) holds when $2 \leq r \leq p^{\prime}$, $1 \leq p<2 d /(d+2)$ and $d \geq 3$, since $\|f\|_{2} \lesssim\|f\|_{p}$.

Our final open question is whether the following inequalities are true for functions with not necessarily vanishing periodizations $g_{\rho}$ :

$$
\|f\|_{p^{\prime}} \lesssim\|f\|_{p}+\|g\|_{p^{\prime}}
$$

and

$$
\|g\|_{p^{\prime}} \lesssim\|f\|_{p}+\|f\|_{p^{\prime}}
$$

for some range of $p \leq 2 d /(d+1)$, where

$$
\|g\|_{p^{\prime}}=\left(\int_{\rho \in \operatorname{SO}(d)}\left\|g_{\rho}\right\|_{p^{\prime}}^{p} d \rho\right)^{1 / p}
$$

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