BOUNDEDNESS OF THE LITTLEWOOD-PALEY g-FUNCTION ON $Lip_{\alpha}(R^n)$ (0 < α < 1)

BY

SHILIN WANG

I. Introduction

Let \mathbb{R}^n be *n*-dimensional Euclidean space. We say that a scalar valued function ψ on \mathbb{R}^n is a Littlewood-Paley function if the following three conditions are satisfied:

$$\psi \in L(\mathbb{R}^n), \quad \int_{\mathbb{R}^n} \psi(x) \, dx = 0, \tag{1.1}$$

$$|\psi(x)| \le c(1+|x|)^{-(n+1)},$$
 (1.2)

$$|\psi(x+y) - \psi(x)| \le \frac{c|y|^{\epsilon}}{(1+|x|)^{n+1+\epsilon}}, \quad |y| \le \frac{|x|}{2}, \text{ some } \epsilon > 0.$$
 (1.3)

For instance, these conditions are satisfied by the well-known functions

$$\psi(x) = \frac{\partial}{\partial t} \frac{c_n t}{\left(t^2 + |x|^2\right)^{(n+1)/2}} \bigg|_{t=1}$$

and

$$\psi_j(x) = \frac{\partial}{\partial x_j} \frac{1}{(1+|x|^2)^{(n+1)/2}} \quad (j=1,2,\ldots,n)$$

with $\varepsilon = 1$, and higher order derivatives of these functions of order k give examples with $\varepsilon = k$.

For a fixed Littlewood-Paley function ψ , let

$$g_{\psi}(f)(x) = g(f)(x) = \left\{ \int_0^\infty |f * \psi_t(x)|^2 \frac{dt}{t} \right\}^{1/2}$$
(1.4)

denote the Littlewood-Paley g-function of f. The Littlewood-Paley g-function

Received July 14, 1987.

^{© 1989} by the Board of Trustees of the University of Illinois Manufactured in the United States of America

of f is well defined under some assumptions on f, although it may be infinite on a set of positive measure.

As in [1], page 312, we have

$$\|g(f)\|_{p} \le C \|f\|_{p} \quad (1$$

where C is independent of f.

Recently, Wang Silei [2] discussed the end-point cases of this result when $f \in L^{\infty}(\mathbb{R}^n)$ or $f \in BMO(\mathbb{R}^n)$.

The aim of this paper is to study the behaviour of g(f) when $f \in Lip_{\alpha}(\mathbb{R}^n)$ (0 < α < 1).

THEOREM 1. If $f \in Lip_{\alpha}(\mathbb{R}^n)$, $0 < \alpha < \min\{\varepsilon, 1\}$, and $g(f)(x_0) < \infty$ for a single point x_0 , then $g(f) \in Lip_{\alpha}(\mathbb{R}^n)$ and

$$\|g(f)\|_{\Lambda_{\alpha}} \leq C \|f\|_{\Lambda_{\alpha}},$$

where $||f||_{\Lambda_{\alpha}}$ denotes the Lip_{α} norm of f, and C is a constant depending only on n and α .

The assumption concerning the finiteness of $g(f)(x_0)$ is essential. In fact, consider the classical Littlewood-Paley g-function

$$g(f)(x) = \left\{\int_0^\infty t |\nabla u(x,t)|^2 dt\right\}^{1/2},$$

where

$$p_t(y) = \frac{c_n t}{\left(t^2 + |y|^2\right)^{(n+1)/2}}, \quad c_n = \frac{\Gamma\left(\frac{n+1}{2}\right)}{\pi^{(n+1)/2}},$$
$$u(x,t) = \int_{R^u} p_t(y) f(x-y) \, dy \quad (t > 0),$$
$$|\nabla u(x,t)|^2 = \left|\frac{\partial u}{\partial t}\right|^2 + \sum_{j=1}^n \left|\frac{\partial u}{\partial x_j}\right|^2.$$

THEOREM 2. There exists a function $f \in Lip_{\alpha}(\mathbb{R}^n)$ for all α ($0 < \alpha < 1$), such that $g(f)(x) = \infty$ everywhere where g is the classical Littlewood-Paley g-function.

532

The author extends most hearty thanks to professor A. Torchinsky for his encouragement and frequent assistance during the writing this paper. The author also thanks the referee for some valuable suggestions that improved the presentation of the results.

We say that a locally integrable function f on \mathbb{R}^n is Lipschitz α ($0 < \alpha < 1$), and denote this by $f \in Lip_{\alpha}$; if there is a constant c so that $|f(x) - f(y)| \le c|x - y|^{\alpha}$ for every x, y in \mathbb{R}^n , the smallest such constant c is called the Lip_{α} norm of f and is denoted by $||f||_{\Lambda_{\alpha}}$.

LEMMA 1 [1, page 213]. Suppose that f is a locally integrable function on \mathbb{R}^n , $0 < \alpha < 1$. Then the following four statements are equivalent and the constants appearing on the right-hand side of each are also equivalent.

(i) $|f(x) - f(y)| \le c_1 |x - y|^{\alpha}$, all x, y in \mathbb{R}^n . (ii)

$$\sup \frac{1}{|Q|^{1+\alpha/n}} \int_Q |f(x) - f_Q| \, dx = c_2 < \infty,$$

where the supremum ranges over all finite cubes Q in \mathbb{R}^n whose sides are parallel to the axes, |Q| is the Lebesgue measure of Q, and f_Q denotes the mean value of f over Q, namely,

$$f_Q = \frac{1}{|Q|} \int_Q f(y) \, dy.$$

(We will consider only these cubes in what follows.)

(iii) $|f(x) - f_Q| \le c_3 |Q|^{\alpha/n}$, all $x \in \mathbb{R}^n$, $Q \subset \mathbb{R}^n$. (iv)

$$\sup\left\{\frac{1}{|Q|^{1+\alpha p/n}}\int_{Q}|f(x)-f_{Q}|^{p} dx\right\}^{1/p}=c_{4}<\infty, \quad 1\leq p<\infty,$$

where the supremum is taken over the same range as in (ii). The equivalence is to be understood to mean that each f which satisfies (ii), (iii) or (iv), can be modified in a set of measure zero so as to coincide with a continuous function which satisfies (i) as well.

We will use freely c_1, c_2, c_3, c_4 as the Lip_{α} norm $||f||_{\Lambda_{\alpha}}$ of f.

Remark. In Lemma 1, if (ii) holds for some constant in place of f_Q then it also holds for f_Q .

LEMMA 2. Suppose that $f \in Lip_{\alpha}(\mathbb{R}^n)$, $\beta > 0$, $0 < \alpha < \min\{\beta, 1\}$, Q is a cube with center x_0 , edge-length h. Then

$$\int_{R^{n} \smallsetminus 4\sqrt{n}\,Q} \frac{|f(x) - f_{4\sqrt{n}\,Q}|}{h^{n+\beta} + |x - x_{0}|^{n+\beta}} dx \le Ch^{\alpha-\beta} ||f||_{\Lambda_{\alpha}},\tag{1.6}$$

C a constant depending only on n, α and β .

Using Lemma 1, the proof of Lemma 2 is similar to a result of Fefferman-Stein's [3], and we omit it.

LEMMA 3. Suppose that $f \in Lip_{\alpha}(\mathbb{R}^n)$ $(0 < \alpha < 1)$, Q is a cube and x_0 is in its interior, then $g(f\chi_0)(x_0) < \infty$.

Proof. To see this we use different estimates on $\psi_t^*(f\chi_Q)(x_0)$ depending on the size of t. If $t \ge 1$ we observe that

$$|\psi_t * (f\chi_Q)(x_0)| \le C \int_{|y| \le C} \frac{t}{(t+|y|)^{n+1}} dy \le Ct^{-n}.$$

This uses only the fact that $f\chi_Q$ is bounded and has compact support. As for $t \leq 1$, we write

$$\psi_t * (f\chi_Q)(x_0) = \int_{|y| \le d} (f(x_0 - y) - f(x_0))\psi_t(y) \, dy$$
$$+ \int_{|y| \ge d} (f(x_0 - y)\chi_Q(x_0 - y) - f(x_0))\psi_t(y) \, dy$$

where we choose d > 0 but $d \le distance$ from x_0 to the complement of Q. Then

$$\left|\int_{|y| \leq d} (f(x_0 - y) - f(x_0))\psi_t(y) \, dy\right| \leq C \int |y|^{\alpha} |\psi_t(y)| \, dy \leq Ct^{\alpha}$$

since $f \in Lip_{\alpha}(\mathbb{R}^n)$ and

$$\left| \int_{|y| \ge d} \left(f(x_0 - y) \chi_Q(x_0 - y) - f(x_0) \right) \psi_t(y) \, dy \right|$$

$$\le C \int_{|y| \ge d} \frac{t}{\left(t + |y|\right)^{n+1}} \, dy$$

$$\le Ct,$$

since fx_Q is bounded.

These estimates easily imply $g(f\chi_Q)(x_0) < \infty$.

II. Proof of Theorem 2

First we consider the case n = 1, and define the function

$$f(x) = \begin{cases} 1, & x \ge 1, \\ x, & 0 < x < 1, \\ 0, & x \le 0. \end{cases}$$

To estimate $|f(x_1) - f(x_2)|$, without loss of generality, we assume that $x_1 < x_2$. If $x_1 < 0$ or $x_2 > 1$, observing that f is constant on $(-\infty, 0]$ and $[1, +\infty)$, by the triangle inequality, we may reduce these cases to $x_1 = 0$ or $x_2 = 1$. That is, we need only to deal with the special case $0 \le x_1, x_2 \le 1$. In this case, $|x_1 - x_2| \le 1$, so

$$|f(x_1) - f(x_2)| = |x_1 - x_2| \le |x_1 - x_2|^{\alpha}$$
 for $0 < \alpha < 1$.

We will prove that $g(f)(x) = \infty$ everywhere. A simple calculation gives

$$\begin{aligned} \frac{\partial u}{\partial x}(x,t) &= -2c_1 \int_{\mathbb{R}^1} f(x-y) \frac{ty}{\left(t^2+y^2\right)^2} dy \\ &= -2c_1 \int_{-\infty}^{x-1} \frac{ty}{\left(t^2+y^2\right)^2} dy - 2c_1 \int_{x-1}^x \frac{\left(x-y\right)ty}{\left(t^2+y^2\right)^2} dy \\ &= c_1 \left(\arctan \frac{x}{t} - \arctan \frac{x-1}{t} \right) \quad (t > 0). \end{aligned}$$

Now, by the mean value theorem, we have

$$\operatorname{arctg} \frac{x}{t} - \operatorname{arctg} \frac{x-1}{t} = \frac{1}{t} \cdot \frac{1}{1+\xi^2} \quad \left(t > 0, \frac{x}{t} > \xi > \frac{x-1}{t}\right),$$

and consequently, if we let $\tilde{x}^2 = \max\{x^2, (x-1)^2\}$, it follows that

$$\begin{aligned} \frac{\partial u}{\partial x} &= \frac{c_1}{t} \cdot \frac{1}{1+\xi^2} \ge c_1 \cdot \frac{1}{t} \cdot \frac{1}{1+\max\left\{\frac{x^2}{t^2}, \frac{(x-1)^2}{t^2}\right\}} \\ &= \frac{c_1 t}{t^2 + \tilde{x}^2} \quad (t > 0). \end{aligned}$$

Thus,

$$g(f)(x) = \left\{\int_0^\infty t |\nabla u(x,t)|^2 dt\right\}^{1/2} \ge c_1 \left\{\int_0^\infty \frac{t^3}{(t^2 + \tilde{x}^2)^2} dt\right\}^{1/2} = \infty.$$

Next, for the case $n \ge 2$, for $x = (x_1, x_2, ..., x_n) \in \mathbb{R}^n$, define the function

$$f(x) = \begin{cases} 1, & x_n \ge 1, \\ x_n, & 0 < x_n < 1, \\ 0, & x_n \le 0. \end{cases}$$

A similar proof can be used to show that $f \in Lip_{\alpha}(\mathbb{R}^n)$ and $g(f)(x) = \infty$ everywhere, and the Theorem 2 follows.

III. Proof of Theorem 1

We first prove that $g(f)(x) < \infty$ a.e. on \mathbb{R}^n . It suffices to prove that $g(f)(x) < \infty$ a.e. on each cube Q_0 containing x_0 . Now, suppose that a cube Q_0 with edge-length h_0 is chosen such that x_0 is in its interior, and write

$$f(x) = f_{4\sqrt{n}Q_0} + (f(x) - f_{4\sqrt{n}Q_0})\chi_{4\sqrt{n}Q_0}(x) + (f(x) - f_{4\sqrt{n}Q_0})\chi_{R^n \smallsetminus 4\sqrt{n}Q_0}(x) = f_1(x) + f_2(x) + f_3(x).$$

Note that the fact that $\int_{\mathbb{R}^n} \psi(y) \, dy = 0$ implies that

$$g(const.) = 0$$
, so $g(f_1) \equiv 0$. (3.2)

As for $f_2(x)$, using (1.5) and Lemma 2 (iv) with p = 2, we have

$$\begin{split} \int_{Q_0} |g(f_2)(x)|^2 \, dx &\leq \int_{\mathbb{R}^n} |g(f_2)(x)|^2 \, dx \\ &\leq C \int_{\mathbb{R}^n} |f_2(x)|^2 \, dx \\ &= C \int_{4\sqrt{n}Q_0} |f(x) - f_{4\sqrt{n}Q_0}|^2 \, dx \\ &\leq C |Q_0|^{1+(2/n)\alpha} ||f||_{\Lambda_\alpha}^2, \end{split}$$
(3.3)

C a constant, not necessarily the same at each occurrence, depending only on n and α . Thus by the Cauchy-Schwarz inequality,

$$\int_{Q_0} |g(f_2)(x)| \, dx \le |Q_0|^{1/2} \left(\int_{Q_0} |g(f_2)(x)|^2 \, dx \right)^{1/2} \le C |Q_0|^{1+\alpha/n} ||f||_{\Lambda_{\alpha}}.$$
(3.4)

536

It also follows from (3.4) that

$$g(f_2)(x) < \infty \quad \text{a.e. on } Q_0. \tag{3.5}$$

Now we consider $g(f_3)(x)$, $x \in Q_0$. We have $g(f_3)(x)$ bounded by the sum of

$$I_1(x) = \left\{ \int_0^{h_0} |\psi_t * f_3(x) - \psi_t * f_3(x_0)|^2 \frac{dt}{t} \right\}^{1/2}$$

and

$$I_2(x) = \left\{ \int_{h_0}^{\infty} |\psi_t * f_3(x) - \psi_t * f_3(x_0)|^2 \frac{dt}{t} \right\}^{1/2},$$

and

$$g(f_3)(x_0) = \left\{\int_0^\infty |\psi_t * f_3(x_0)|^2 \frac{dt}{t}\right\}^{1/2}.$$

Applying Lemma 3 to the function $f - f_{4\sqrt{n}Q_0}$ and the cube $4\sqrt{n}Q_0$, $g(f_2)(x_0) < \infty$, and so

$$g(f_3)(x_0) \le g(f_2)(x_0) + g(f)(x_0) < \infty.$$
 (3.6)

By (1.2),

$$I_{1}(x) = \left\{ \int_{0}^{h_{0}} \left[\int_{\mathbb{R}^{n} \setminus 4\sqrt{n} Q_{0}} [\psi_{t}(x-y) - \psi_{t}(x_{0}-y)] \times [f(y) - f_{4\sqrt{n} Q_{0}}] dy \right]^{2} \frac{dt}{t} \right\}^{1/2} \\ \leq \left\{ \int_{0}^{h_{0}} \left[\int_{\mathbb{R}^{n} \setminus 4\sqrt{n} Q_{0}} [|\psi_{t}(x-y)| + |\psi_{t}(x_{0}-y)|] \times |f(y) - f_{4\sqrt{n} Q_{0}}| dy \right]^{2} \right\}^{1/2} \\ \leq C \left\{ \int_{0}^{h_{0}} t \left[\int_{\mathbb{R}^{n} \setminus 4\sqrt{n} Q_{0}} \left[\frac{1}{(t+|x-y|)^{n+1}} + \frac{1}{(t+|x_{0}-y|)^{n+1}} \right] \times |f(y) - f_{4\sqrt{n} Q_{0}}| dy \right]^{2} dt \right\}^{1/2}.$$
(3.7)

The conditions $x_0, x \in Q_0$ and $y \notin 4\sqrt{n}Q_0, 0 < t < h_0$, imply that

$$|x - y| \ge C|x_0 - y| \ge C(|x_0 - y| + h_0)$$

and

$$\frac{1}{(|x_0 - y| + t)^{n+1}} \le \frac{C}{|x_0 - y|^{n+1} + h_0^{n+1}},$$
$$\frac{1}{(|x - y| + t)^{n+1}} \le \frac{C}{|x_0 - y|^{n+1} + h_0^{n+1}}.$$

Therefore, we obtain

$$I_1(x) \leq C \left\{ \int_0^{h_0} t \left[\int_{R^n \smallsetminus 4\sqrt{n} Q_0} \frac{|f(y) - f_{4\sqrt{n} Q_0}|}{|x_0 - y|^{n+1} + h_0^{n+1}} dy \right]^2 dt \right\}^{1/2}.$$

By Lemma 2 with $\varepsilon = 1$,

$$I_{1}(x) \leq C \left\{ \int_{0}^{h_{0}} t \cdot h_{0}^{2\alpha-2} \|f\|_{\Lambda_{\alpha}}^{2} dt \right\}^{1/2} \leq Ch_{0}^{\alpha} \|f\|_{\Lambda_{\alpha}} \leq C \|Q_{0}\|^{\alpha/n} \|f\|_{\Lambda_{\alpha}}.$$
(3.8)

Now we estimate $I_2(x)$. The fact that $x_0, x_1 \in Q_0$ and $y \notin 4\sqrt{n}Q_0$ implies that

$$|(x-y) - (x_0 - y)| \le \frac{1}{2}|x_0 - y|,$$

and by (1.3),

$$\begin{aligned} |\psi_t(x-y) - \psi_t(x_0 - y)| &\leq ct^{-n} \Big| \frac{x - x_0}{t} \Big|^{\epsilon} \cdot \frac{1}{\left(1 + \frac{|x_0 - y|}{t}\right)^{n+1+\epsilon}} \\ &\leq ct \cdot h_0^{\epsilon} \cdot \frac{1}{\left(t + |x_0 - y|\right)^{n+1+\epsilon}}. \end{aligned} (3.9)$$

538

Applying Minkowski's inequality we have

$$I_{2}(x) = \left\{ \int_{h_{0}}^{\infty} \left| \int_{R^{n} \smallsetminus 4\sqrt{n} Q_{0}} [\psi_{t}(x-y) - \psi_{t}(x_{0}-y)] \right| \times \left[f(y) - f_{4\sqrt{n} Q_{0}} \right] dy \right|^{2} \frac{dt}{t} \right\}^{1/2} \\ \leq \int_{R^{n} \smallsetminus 4\sqrt{n} Q_{0}} \left[\int_{h_{0}}^{\infty} |\psi_{t}(x-y) - \psi_{t}(x_{0}-y)|^{2} \\ \times |f(y) - f_{4\sqrt{n} Q_{0}}|^{2} \frac{dt}{t} \right]^{1/2} dy \\ \leq \int_{R^{n} \smallsetminus 4\sqrt{n} Q_{0}} h_{0}^{e} |f(y) - f_{4\sqrt{n} Q_{0}}| \\ \times \left[\int_{h_{0}}^{\infty} \frac{t^{2}}{(t+|x_{0}-y|)^{2(n+1+e)}} \cdot \frac{dt}{t} \right]^{1/2} dy, \quad (3.10)$$

and for $x_0 \in Q_0$, $y \in \mathbb{R}^n \setminus 4\sqrt{n} Q_0$,

$$\int_{h_0}^{\infty} \frac{t^2}{(t+|x_0-y|)^{2(n+1+\epsilon)}} \cdot \frac{dt}{t} \bigg]^{1/2}$$

$$\leq C \bigg[\int_{h_0}^{\infty} \frac{t}{(t^2+|x_0-y|^2)^{(n+1+\epsilon)}} dt \bigg]^{1/2}$$

$$\leq C \bigg[\frac{1}{(h_0^2+|x_0-y|^2)^{n+\epsilon}} \bigg]^{1/2}$$

$$\leq \frac{C}{|x_0-y|^{n+\epsilon}+h_0^{n+\epsilon}}.$$
(3.11)

Then we use Lemma 2 to obtain

$$I_{2}(x) \leq C \int_{\mathbb{R}^{n} \setminus 4\sqrt{n} Q_{0}} h_{0}^{e} \frac{|f(y) - f_{4\sqrt{n} Q_{0}}|}{|x_{0} - y|^{n+e} + h_{0}^{n+e}} dy$$

$$\leq C h_{0}^{e} \cdot h_{0}^{\alpha-e} ||f||_{\Lambda_{\alpha}} \leq C |Q_{0}|^{\alpha/n} ||f||_{\Lambda_{\alpha}}.$$
(3.12)

Combining (3.6), (3.8) and (3.12), we have $g(f_3)(x) < \infty$, and so

$$g(f)(x) \le g(f_2)(x) + g(f_3)(x) < \infty$$
 a.e. on Q_0

Next let Q be any cube with edge-length h. Write

$$f(x) = f_{4\sqrt{n}Q} + (f(x) - f_{4\sqrt{n}Q})\chi_{4\sqrt{n}Q}(x) + (f(x) - f_{4\sqrt{n}Q})\chi_{R^n \smallsetminus 4\sqrt{n}Q}(x)$$

= $f_1(x) + f_2(x) + f_3(x).$ (3.13)

Then $g(f_1)(x) \equiv 0$, and repeating the process to prove (3.4) with Q_0 replaced by Q, we have

$$\int_{Q} |g(f_2)(x)| \, dx \le C |Q|^{1+\alpha/n} ||f||_{\Lambda_{\alpha}}.$$
(3.14)

It has been shown that $g(f)(x) < \infty$ a.e., so that

$$g(f_3)(x) \leq g(f_2)(x) + g(f)(x) < \infty \qquad a.e..$$

Therefore, there must be a point $\bar{x} \in Q$ for which $g(f_3)(\bar{x}) < \infty$. Repeating the proof of (3.8) and (3.12) with Q_0 replaced by Q, we have

$$\left\{\int_0^\infty |\psi_t * f_3(x) - \psi_t * f_3(\bar{x})|^2 \frac{dt}{t}\right\}^{1/2} \le C |Q|^{\alpha/n} ||f||_{\Lambda_{\alpha}}, \quad x \in Q.$$

Thus for $x \in Q$,

$$|g(f_3)(x) - g(f_3)(\bar{x})| \le \left[\int_0^\infty |\psi_t * f_3(x) - \psi_t * f_3(\bar{x})|^2 \frac{dt}{t}\right]^{1/2} \le C |Q|^{\alpha/n} ||f||_{\Lambda_\alpha}$$

and consequently

$$\int_{Q} |g(f_3)(x) - g(f_3)(\bar{x})| \, dx \le C |Q|^{1+\alpha/n} ||f||_{\Lambda_{\alpha}}. \tag{3.15}$$

Therefore, we have

$$\begin{split} \int_{Q} |g(f)(x) - g(f_{3})(\bar{x})| \, dx \\ &\leq \int_{Q} |g(f)(x) - g(f_{3})(x)| \, dx + \int_{Q} |g(f_{3})(x) - g(f_{3})(\bar{x})| \, dx \\ &\leq \int_{Q} |g(f_{2})(x)| \, dx + \int_{Q} |g(f_{3})(x) - g(f_{3})(\bar{x})| \, dx \\ &\leq C |Q|^{1+\alpha/n} ||f||_{\Lambda_{\alpha}}. \end{split}$$
(3.16)

Finally, by the remark after Lemma 1, for each cube Q,

$$\frac{1}{|\mathcal{Q}|^{1+\alpha/n}}\int_{\mathcal{Q}}|g(f)(x)-(g(f))_{\mathcal{Q}}|\,dx\leq C||f||_{\Lambda_{\alpha}},$$

namely,

$$\|g(f)\|_{\Lambda_{\alpha}} \leq C \|f\|_{\Lambda_{\alpha}},$$

C a constant depending only on n and α . This completes the proof of Theorem 1.

References

- 1. C.L. FEFFERMANN and E. M. STEIN, H^p-spaces of several variables, Acta Math., vol. 129 (1972), pp. 137–193.
- 2. A. TORCHINSKY, Real-variable methods in harmonic analysis, Academic Press, San Diego, Calif., 1986.
- SILEI WANG, Some properties of Littlewood-Paley's g-function, Scientia Sinica, Series A, vol. 28 (1985), pp. 252–262.

Indiana University Bloomington, Indiana