# On the middle dimension cohomology of $A_{l}$ singularity 

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#### Abstract

. Let ( $V, o$ ) be a normal isolated singularity in a complex Euclidean space ( $C^{N}, o$ ). Let $M$ be the intersection of this singularity and the real hypersphere $S_{\epsilon}^{2 N-1}(o)$, centered at the origin $o$ with an $\epsilon$ radius. Then, naturally, this link $M$ admits a CR structure, induced from $V$, and the deformation theory of this CR structures has been studied in [1], [2],[3]. Especially in [3], a particular subspace of the infinitesimal deformation space is found, and we propose to study the relation between this subspace and simultaneous deformation. We note that: if the canonical line bundle of the CR structure is trivial, then the infinitesimal space of the deformation of CR structures is a part of the middle dimension cohomology. And in this line, we conjecture that $Z^{1}$, introduced in [3], might be related to the simultaneous deformation of isolated singularity ( $V, o)$ (see also [2]). We discuss this problem for $A_{l}$ singularities.


## §1. Motivation and $Z^{1}$ - space

Let $\left(V^{(n)}, o\right)$ be an isolated singularity in a complex eucliean space $\left(C^{N}, o\right)$. We consider the intersection

$$
M=S_{\epsilon}^{2 N-1}(o) \cap V
$$

Then $M$ is a compact non-singular real $2 n-1$ dimensional $C^{\infty}$ manifold, and a CR structure $\left(M,{ }^{0} T^{\prime \prime}\right)$ is induced from $V$, by ;

$$
{ }^{0} T^{\prime \prime}=C \otimes T M \cap T^{\prime \prime}(V-o)
$$

Here $T^{\prime \prime}(V-o)$ means the space consisting of type $(1,0)$ vectors on $V-o$. This pair $\left(M,{ }^{0} T^{\prime \prime}\right)$ is called a CR structure(or a CR manifold). For this CR structure, the deformation theory, related to the deformation theory
of isolated singularities $(V, o)$, is successfully developed by Kuranishi. After the great work of Kuranishi, we are interested in the mixed Hodge structure of CR manifolds. We take a supplement vector field $\zeta$ to ${ }^{0} T^{\prime \prime}+{ }^{0} T^{\prime}$, here ${ }^{0} T^{\prime}=\overline{{ }^{0}} T^{\prime \prime}$. For this CR structure with the supplement vector field $\left\{\left(M,{ }^{0} T^{\prime \prime}\right), \zeta\right\}$, we can introduce a mixed Hodge structure which should correspond to the mixed Hodge structure on a tubular neighborhood $U$ of $M$ in $V$. Here, we assume that there is a real vector field $\zeta$ satisfying:

$$
\begin{gather*}
\zeta_{p} \not{ }^{0} T_{p}^{\prime \prime}+{ }^{0} T_{p}^{\prime}  \tag{1}\\
{\left[\zeta, \Gamma\left(M,{ }^{0} T^{\prime \prime}\right)\right] \subset \Gamma\left(M,{ }^{0} T^{\prime \prime}\right)} \tag{2}
\end{gather*}
$$

While, during our studying deformation theory of CR structures, we learn that: for Calabi-Yau manifolds, the Kuranishi family is unobstructed. So, in order to obtain the analogy to isolated singularities, $Z^{1}$ space is found(see [3]).

$$
\begin{equation*}
Z^{1}=\left\{u: u \in F^{n-1,1}, d^{\prime \prime} u=0, d^{\prime} u=0\right\} \tag{3}
\end{equation*}
$$

In the case complex manifolds, $Z^{1}$ might be translated as follows. For a tubular neighborhood $U$ of $M$ in $V$, we set

$$
\begin{equation*}
\left\{u: u \in \Gamma\left(U, \wedge^{n-1}\left(T^{\prime} U\right)^{*} \wedge\left(T^{\prime \prime} U\right)^{*}\right), \bar{\partial} u=0, \partial u=0\right\} \tag{4}
\end{equation*}
$$

If $X^{(n)}$ is a compact n -dimensional Kaehler manifold, then

$$
\begin{equation*}
\left\{u: u \in \Gamma\left(X^{(n)}, \wedge^{n-1}\left(T^{\prime} X^{(n)}\right)^{*} \wedge\left(T^{\prime \prime} X^{(n)}\right)^{*}\right), \bar{\partial} u=0, \partial u=0\right\} \tag{5}
\end{equation*}
$$

includes the $\bar{\partial}$-harmonic space consisting of ( $n-1,1$ ) forms. While, here, we are treating an open manifold $U$ (tubular neighborhood of $M$ ). So even if the ( $n-1,1$ ) Kohn-Rossi cohomology does not vanish(the existence of a non- trivial $\bar{\partial}$-harmonic space consisting of ( $n-1,1$ ) forms), the above space might be 0 . Here we give a program to obtain a nontrivial element of (4) from a non- trivial simultaneous deformation.

Let $\tilde{V}$ be the resolution of the isolated singularity with complex dimension $n$ in $C^{N}, V$, and $\pi$ is the resolution map $\pi ; \tilde{V} \rightarrow V$. And consider non-trivial deformations of isolated singularity ( $V, o$ ) with this resolution. Namely, $\pi_{t}$ is a resolution map of $V_{t}$ in $C^{N}, \pi_{t} ; \quad \tilde{V}_{t} \rightarrow V_{t}$, $t \in T$, where $V_{t}$ is a deformation of $V, \tilde{V}_{t}$ is a deformation of $\tilde{V}, T$ is an analytic space with the origin, and at the origin, $\pi_{o}=\pi, \tilde{V}_{o}=$ $\tilde{V}, V_{o}=V$. Furthermore, we assume that $V_{t} \subset C^{N}$. Now we take a $C^{\infty}$ trivialization $i_{t}$ : a tubular neighborhood of $M_{o} \rightarrow$ a tubular neighborhood of $M_{t}$, which satisfies $i_{t}\left(M_{o}\right)=M_{t}$. In this setting, our program is as follows.

- (First Step) By using the simultaneous resolution, we construct a non-trivial $(n, 0)$ form $\omega_{t}$, which is not $d$ exact on $\tilde{V}_{t}$ for a generic $t$, and depends on $t$ complex analytically. In general, "to give an $(n, 0)$ form, satisfying a certain condition", might be easier than "to give an ( $n-1,1$ ) form with the corresponding condition".
- (Second Step) By choosing a proper $C^{\infty}$ trivialization of the simultaneous deformation, $i_{t}$,

$$
i_{t}^{*} \omega_{t}=\omega_{0}+\omega_{1} t+\cdots, \quad(\text { expansion with respect to } t) .
$$

- (Third Step) From $d \omega_{t}=0$, it follows that: $d \omega_{1}=0$. By the definition, $\omega_{1}$ is a form of type $(n, 0)+(n-1,1)$ on $\tilde{V}_{o}-\pi^{-1}(o)$, we write it by;

$$
\omega_{1}=\omega_{1}^{(n, 0)}+\omega_{1}^{(n-1,1)} .
$$

As $d \omega_{1}=0$, this is equivalent to

$$
\begin{gathered}
\bar{\partial} \omega_{1}^{(n-1,1)}=0, \\
\bar{\partial} \omega_{1}^{(n, 0)}+\partial \omega_{1}^{(n-1,1)}=0 .
\end{gathered}
$$

The $\bar{\partial}$-cohomology class, determined by $\omega_{1}^{(n-1,1)}$, is the induced one by the Kodaira-Spencer class of deformations. So, this must be non-trivial. In this setting, we would like to construct a nontrivial element of (4), associated with the given simultaneous deformation.
For the Third Step, we have to comment on a crucial point. The naive answer is that:

$$
\partial \omega_{1}^{(n-1,1)}=0
$$

This is too strong. There is an ambiguity to choose the $C^{\infty}$ trivialization, $i_{t}$. By changing the $C^{\infty}$ trivialization, $\omega_{1}\left(\right.$ resp. $\omega_{1}^{(n-1,1)}$ ) is replaced by $\omega_{1}-d u$ (resp. $\omega_{1}^{(n-1,1)}-\bar{\partial} u$ ), where $u$ is an $(n-1,1)$ form. Hence our problem(to obtain a non-trivial element of (4)) is reduced to that; is there any $C^{\infty}(n-1,1)$ form $u$, satisfying: $\bar{\partial} \omega_{1}^{(n-1,1)}-\partial \bar{\partial} u=0$ ? This is so called " $\partial \bar{\partial}$ lemma". For a compact Kaehler manifold, by taking the harmonic part, this is always solvable. However, for an open manifold, this is not an easy problem. One of our conjecture is that; if $\omega_{1}^{(n-1,1)}$ is induced by the simultaneous deformation, then this might be solvable. In the next section, we study this conjecture in $A_{l}$ singularities.

## §2. $A_{l}$ singularities

Let

$$
X=\left\{\left(z_{1}, \ldots, z_{n+1}\right):\left(z_{1}, \ldots, z_{n+1}\right) \in C^{n+1}, z_{1}^{2}+\cdots+z_{n+1}^{l+1}=0\right\}
$$

where $l$ is a positive integer. We call this isolated singularity $A_{l}$ singularity. Consider a family of deformations of $X$,

$$
X_{t}=\left\{\left(z_{1}, \ldots, z_{n+1}\right):\left(z_{1}, \ldots, z_{n+1}\right) \in C^{n+1}, z_{1}^{2}+\cdots+z_{n+1}^{l+1}=t\right\}
$$

Let $M=X \cap\left\{\left(z_{1}, \ldots, z_{n+1}\right):\left|z_{1}\right|^{2}+\cdots+\left|z_{n+1}\right|^{2}=1\right\}$. And consider a $C^{\infty}$ trivialization of this deformation over a neighborhood of $M$ in $X$.
Let $i_{t}:\left(z_{1}, \ldots, z_{n+1}\right) \rightarrow\left(z_{1}(t), \ldots, z_{n+1}(t)\right)$, where

$$
\begin{aligned}
z_{1}(t) & =z_{1}+\frac{1}{2 k(z, \bar{z})} \bar{z}_{1}\left(1+\left|z_{n+1}\right|^{2}+\cdots+\left|z_{n+1}\right|^{2 l}\right) t \\
\ldots & \\
z_{n}(t) & =z_{n}+\frac{1}{2 k(z, \bar{z})} \bar{z}_{n}\left(1+\left|z_{n+1}\right|^{2}+\cdots+\left|z_{n+1}\right|^{2 l}\right) t \\
z_{n+1}(t) & =z_{n+1}+\frac{1}{(l+1) k(z, \bar{z})} \bar{z}_{n+1}^{l} t
\end{aligned}
$$

Here

$$
k(z, \bar{z})=\left(1+\left|z_{n+1}\right|^{2}+\cdots+\left|z_{n+1}\right|^{2(l-1)}\right)\left(\left|z_{1}\right|^{2}+\cdots+\left|z_{n}\right|^{2}\right)+\left|z_{n+1}\right|^{2} l .
$$

So, on $M$, because of $\left|z_{1}\right|^{2}+\cdots+\left|z_{n}\right|^{2}=1-\left|z_{n+1}\right|^{2}, k(z, \bar{z})=1$ holds. And,

$$
\begin{aligned}
z_{1}(t)^{2}+ & \cdots+ \\
= & z_{n}(t)^{2}+z_{n+1}(t)^{l+1} \\
= & z_{1}^{2}+\cdots+z_{n}^{2}+z_{n+1}^{l+1} \\
& +\frac{1}{k(z, \bar{z})}\left\{\left(1+\left|z_{n+1}\right|^{2}+\cdots+\left|z_{n+1}\right|^{2(l-1)}\right)\left(\left|z_{1}\right|^{2}+\cdots+\left|z_{n}\right|^{2}\right)\right. \\
& \left.+\left|z_{n+1}\right|^{2 l}\right\} t+\text { higher order term of } t \\
\equiv & t \quad \bmod t^{2}
\end{aligned}
$$

By adjusting higher order term, we have a $C^{\infty}$ trivialization $i_{t}: X \rightarrow X_{t}$ over a neighborhood of $M$. However, in this paper, we discuss only differential forms of type $(n-1,1)$. So the above map is enough.

## §3. An approach to the First Step

In this section, we give a non-trivial holomorphic ( $n, 0$ ) form on $X_{t} \cap\left(\right.$ a neighborhood of $M$ in $\left.C^{n+1}\right)$, which depends on $t$, complex
analytically. Let $f=z_{1}^{2}+\cdots+z_{n}^{2}+z_{n+1}^{l+1}$. Like in [2], we, first, set a type $(1,0)$ vector field $Z_{f}$, defined on a neighborhood of $M$ in the $C^{n+1}$, as follows. Let $\Omega$ be the standard symplectic form.

$$
\Omega=\sum_{i=1}^{n+1} \sqrt{-1} d z_{i} \wedge d \bar{z}_{i}
$$

By using this metric, we define a $(1,0)$ vector field $Z_{f}$ on a neighborhood of $M$ by;

$$
d f(X)=\Omega\left(X, \bar{Z}_{f}\right), \quad \text { for all }(1,0) \text { vector field } X
$$

This $Z_{f}$ is easily written down as follows.

$$
\begin{aligned}
Z_{f} & =\sqrt{-1} \sum_{i=1}^{n+1} \overline{\left(\frac{\partial f}{\partial z_{i}}\right)} \frac{\partial}{\partial z_{i}} \\
& =\sqrt{-1}\left\{\sum_{i=1}^{n} 2 \bar{z}_{i} \frac{\partial}{\partial z_{i}}+(l+1) \bar{z}_{n+1}^{l} \frac{\partial}{\partial z_{n+1}}\right\}
\end{aligned}
$$

So,

$$
\begin{aligned}
Z_{f}(f) & =\sqrt{-1}\left(2^{2} \sum_{i=1}^{n}\left|z_{i}\right|^{2}+(l+1)^{2}\left|z_{n+1}\right|^{2 l}\right) \\
& \neq 0 \text { on a neighborhood of } M
\end{aligned}
$$

Let $\omega=d z_{1} \wedge \cdots \wedge d z_{n+1}$. For $X_{t}$, we set a holomorphic $(n, 0)$ form $\omega^{\prime}(t)$, which depends on $t$, complex analytically by ;

$$
\left.\omega^{\prime}(t)=Z_{f}\right\rfloor \omega \quad \text { on } \quad X_{t} \text { (inner product with vector field } Z_{f} \text { ). }
$$

And set

$$
\omega_{t}^{\prime}=\frac{1}{\sum_{i=1}^{n} 2^{2}\left|z_{i}\right|^{2}+(l+1)^{2}\left|z_{n+1}\right|^{2 l}} \omega^{\prime}(t)
$$

By the type of $\omega$, our $\omega_{t}^{\prime}$ is of type $(n, 0)$ on $X_{t}$. We must show that our $\omega_{t}^{\prime}$ is holomorphic on $X_{t}$. For this, we recall the following lemma.

Lemma 3.1. $\quad \omega=-\sqrt{-1} d f \wedge \omega_{t}^{\prime}$ on a neighborhood of $M$.
We sketch the proof of this lemma. For a point $p$ of a neighborhood of $M$ in $C^{n+1}, T_{p}^{\prime} C^{n+1}$ is spanned by $Z_{f}$ and $\left\{X_{i}(p)\right\}_{1 \leq i \leq n}$, which satisfy $X_{i}(p) f=0$. So, with these vector fields, just by a direct computation, we have our lemma.

By this lemma, on $X_{t}$,

$$
d \omega_{t}^{\prime}=0
$$

We have to see that our $\omega_{o}^{\prime}$ is not a d-exactn on $X_{o}=X$. But if we restric $\omega_{t}$ to

$$
\left\{\left(z_{1}, \cdots, z_{n}, z_{n+1}\right): z_{1}^{2}+\cdots+z_{n}^{2}+z_{n+1}^{l+1}=0, z_{n+1}=0\right\}
$$

a complex $n-1$ dimensional $A_{1}$ singularity, then it gives a non-trivial $n-1$ dimensional cohomology(by the definition of our $\omega_{t}^{\prime}$, it coincides with nontrivial element, constructed in [2]). So, we have a non trivial form.

## §4. An approach to the Third Step

By the $C^{\infty}$ trivialization of the simultaneous deformations, $i_{t}$, constructed in Section 2, on a tubular neighborhood of $M$,

$$
i_{t}^{*} \omega_{t}=\omega_{0}+\omega_{1} t+\cdots, \quad(\text { expansion with respect to } t)
$$

We explain a difficulty about this part. For example, we take $A_{1}$ singularity (in our notations, $l=1$ ). Then, in the $C^{\infty}$ isomorphism map, $i_{t}$, as a denominator, $k(z, \bar{z})$ appears. Only on the boundary case(CR case)

$$
k(z, \bar{z})=1 \text { on the boundary. }
$$

But we are treating the tubular neighborhood case. So, it is not so valid that there is no extra non-trivial $(n, 0)$ term of $\omega_{1}$ (we write it by $\omega_{1}^{(n, 0)}$ ). Fortunately, for the case $l=1$ ( the case of an ordinary double point $),(n, 0)$ term doesn't appear(this means that it is not necessary to change the $C^{\infty}$ trivialization $i_{t}$, constructed in Section 2). So, in this case, $d \omega_{1}=0$ means that; $\partial \omega_{1}=0$ and $\bar{\partial} \omega_{1}=0$. For the other $l$, we have to control the difficulty which arises from the term $k(z, \overline{)}$. In another paper, we discuss the other case.

For the case $l=1$, the $C^{\infty}$ isomorphism map is as follows.

$$
z_{i}(t)=z_{i}+\frac{1}{2 \sum_{i=1}^{n+1}\left|z_{i}\right|^{2}} \bar{z}_{i} t, \quad i=1, \ldots, n+1
$$

And

$$
Z_{f}=2\left(\sum_{i=1}^{n+1} \bar{z}_{i} \frac{\partial}{\partial z_{i}}\right)
$$

In order to simplify the sketch, we assume $n=2$. Then,

$$
Z_{f}=2\left(\bar{z}_{1} \frac{\partial}{\partial z_{1}}+\bar{z}_{2} \frac{\partial}{\partial z_{2}}+\bar{z}_{3} \frac{\partial}{\partial z_{3}}\right)
$$

And so,

$$
\begin{aligned}
& \left.Z_{f}\right\rfloor \omega=2\left(\bar{z}_{1} d z_{2} \wedge d z_{3}-\bar{z}_{2} d z_{1} \wedge d z_{3}+\bar{z}_{3} d z_{1} \wedge d z_{2}\right) \\
& \qquad \begin{aligned}
Z_{f}(f) & =4\left(\left|z_{1}\right|^{2}+\left|z_{2}\right|^{2}+\left|z_{3}\right|^{2}\right) \\
& =4 r^{2}
\end{aligned}
\end{aligned}
$$

Here $r^{2}=\left|z_{1}\right|^{2}+\left|z_{2}\right|^{2}+\left|z_{3}\right|^{2}$. And

$$
\begin{aligned}
z_{1}(t) & =z_{1}+\frac{1}{2} \frac{1}{r^{2}} \bar{z}_{1} t \\
z_{2}(t) & =z_{2}+\frac{1}{2} \frac{1}{r^{2}} \bar{z}_{2} t \\
z_{3}(t) & =z_{3}+\frac{1}{2} \frac{1}{r^{2}} \bar{z}_{3} t
\end{aligned}
$$

Now we compute $\omega_{1}$.

$$
\begin{aligned}
\left.i_{t}^{*}\left(\frac{1}{4 r^{2}} Z_{f}\right\rfloor \omega\right)= & \frac{1}{2} i_{t}^{*}\left(\frac{1}{r^{2}}\left(\bar{z}_{1} d z_{2} \wedge d z_{3}-\bar{z}_{2} d z_{1} \wedge d z_{3}+\bar{z}_{3} d z_{1} \wedge d z_{2}\right)\right) \\
& =\frac{1}{2}\left(\frac{\left.\bar{z}_{1}(t) d z_{2}(t) \wedge d z_{3}(t)-\bar{z}_{2}(t) d z_{1}(t) \wedge d z_{3}(t)+\bar{z}_{3}(t) d z_{1}(t) \wedge d z_{2}(t)\right)}{z_{1}(t) \bar{z}_{1}(t)+z_{2}(t) \bar{z}_{2}(t)+z_{3}(t) \bar{z}_{3}(t)}\right) \\
& \equiv \frac{1}{2}\left(\frac{\left.\bar{z}_{1} d z_{2}(t) \wedge d z_{3}(t)-\bar{z}_{2} d z_{1}(t) \wedge d z_{3}(t)+\bar{z}_{3} d z_{1}(t) \wedge d z_{2}(t)\right)}{z_{1}(t) \bar{z}_{1}+z_{2}(t) \bar{z}_{2}+z_{3}(t) \bar{z}_{3}}\right) \bmod \left(t^{2}, \bar{t}\right) \\
= & \frac{1}{2}\left(\frac{\left.\bar{z}_{1} d z_{2}(t) \wedge d z_{3}(t)-\bar{z}_{2} d z_{1}(t) \wedge d z_{3}(t)+\bar{z}_{3} d z_{1}(t) \wedge d z_{2}(t)\right)}{z_{1} \bar{z}_{1}+z_{2} \bar{z}_{2}+z_{3} \bar{z}_{3}}\right) \\
& \quad \text { because of } z_{1}^{2}+z_{2}^{2}+z_{3}^{2}=0 .
\end{aligned}
$$

While

$$
\begin{aligned}
\bar{z}_{1} d z_{2}(t) \wedge d z_{3}(t) & =\bar{z}_{1}\left(d z_{2}+\frac{1}{2}\left(d\left(\frac{1}{r^{2}}\right)\right) \bar{z}_{2} t+\frac{1}{2} \frac{1}{r^{2}} d \bar{z}_{2} t\right) \wedge\left(d z_{3}+\frac{1}{2}\left(d\left(\frac{1}{r^{2}}\right)\right) \bar{z}_{3} t+\frac{1}{2} \frac{1}{r^{2}} d \bar{z}_{3} t\right) \\
& \equiv \bar{z}_{1} d z_{2} \wedge d z_{3}+\left\{\bar{z}_{1} \frac{1}{2}\left(d\left(\frac{1}{r^{2}}\right)\right) \bar{z}_{2} \wedge d z_{3}+\bar{z}_{1} \frac{1}{2} \frac{1}{r^{2}} d \bar{z}_{2} \wedge d z_{3}\right. \\
& \left.+\bar{z}_{1} d z_{2} \wedge \frac{1}{2}\left(d\left(\frac{1}{r^{2}}\right)\right) \bar{z}_{3}+\bar{z}_{1} d z_{2} \frac{1}{2} \frac{1}{r^{2}} d \bar{z}_{3}\right\} t \quad \bmod t^{2}
\end{aligned}
$$

Therefore from this term, $(2,0)$ part is

$$
\frac{1}{2} \bar{z}_{1} \bar{z}_{2} \partial\left(\frac{1}{r^{2}}\right) \wedge d z_{3}+\frac{1}{2} \bar{z}_{1} \bar{z}_{3} d z_{2} \wedge \partial\left(\frac{1}{r^{2}}\right) .
$$

By the same way, from $-\bar{z}_{2} d z_{1}(t) \wedge d z_{3}(t)$, as a $(2,0)$ part,

$$
-\frac{1}{2} \bar{z}_{1} \bar{z}_{2} \partial\left(\frac{1}{r^{2}}\right) \wedge d z_{3}-\frac{1}{2} \bar{z}_{2} \bar{z}_{3} d z_{1} \wedge \partial\left(\frac{1}{r^{2}}\right)
$$

And from $\bar{z}_{3} d z_{1}(t) \wedge d z_{2}(t),(2,0)$ part is

$$
\frac{1}{2} \bar{z}_{1} \bar{z}_{3} \partial\left(\frac{1}{r^{2}}\right) \wedge d z_{2}+\frac{1}{2} \bar{z}_{2} \bar{z}_{3} d z_{1} \wedge \partial\left(\frac{1}{r^{2}}\right)
$$

So summing up these three terms, in this case, we see that $(2,0)$ part does not appear.

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