

## Stationary Phase Method with Estimate of Remainder Term over a Space of Large Dimension

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### Abstract.

Let  $r_d(\nu)$  denote the remainder term of the stationary phase method over  $R^d$ . Then an estimate of  $\nu^{d/2+1}r_d(\nu)$ , as  $d \rightarrow \infty$ , is given under certain assumptions, which are tolerable for application to Feynman path integrals.

### §1. Stationary phase method

Stationary phase method is a method to evaluate asymptotically, as  $\nu \rightarrow \infty$ , oscillatory integrals over  $R^d$  of the following form:

$$I(S, a, \nu) = \int_{R^d} e^{-i\nu S(x)} a(x) dx,$$

where  $S(x)$  is a real valued  $C^\infty$  function called the phase function,  $a(x)$  is a  $C^\infty$  function called the amplitude and  $\nu$  is a large positive parameter. In the simplest case that  $a(x) \in C_0^\infty(R^d)$  and that  $S(x)$  has only one critical point  $x^*$ , where  $\text{Hess } S(x^*)$  is non-degenerate, it gives

$$I(S, a, \nu) = \left( \frac{2\pi}{i\nu} \right)^{d/2} [\det\{\text{Hess } S(x^*)\}]^{-1/2} (e^{-i\nu S(x^*)} a(x^*) + r_d(\nu))$$

and an estimate of the remainder term

$$r_d(\nu) = O(\nu^{-d/2-1}).$$

If support of  $a(x)$  is not compact, we have to require some additional assumption that control the behaviour of  $a(x)$  at the infinity. For instance (cf. [1]), the same conclusion holds if we assume the following

**Hypothesis (H.0).** (i)  $\sup_x |\partial_x^\alpha S(x)| < \infty$  for any multi-index  $\alpha$  with  $|\alpha| \geq 2$ . (ii) There exists a constant  $\delta > 0$  such that  $|\det \text{Hess } S(x)| \geq \delta$ . (iii) For any multi-index  $\alpha$ ,  $\sup_x |\partial_x^\alpha a(x)| < \infty$ .

Since the stationary phase method is closely related to the mathematical theory of Feynman path integrals (cf. [3], [4], [5] and [6]), we wish to investigate the following

**Question.** Can one control  $\nu^{d/2+1} r_d(\nu)$  as  $d \rightarrow \infty$  ?

We give a positive answer to this question. Detailed discussions can be found in [2]. Applications are discussed in [4], [5] and [6].

## §2. Statement of results

We shall treat the following oscillatory integral over  $L - 1$  dimensional space:

$$I(\{t_j\}, S, a, \nu)(x_L, x_0) = \prod_{j=1}^L \left( \frac{\nu i}{2\pi t_j} \right)^{1/2} \int_{R^{L-1}} e^{-i\nu S(x_L, \dots, x_0)} a(x_L, \dots, x_0) \prod_{j=1}^{L-1} dx_j,$$

with large positive parameter  $\nu$  and small positive parameters  $\{t_j\}$ . Our hypothesis for the phase function is

**Hypothesis (H.1).**  $S(x_L, \dots, x_0)$  is of the form

$$S(x_L, \dots, x_0) = \sum_{j=1}^L S_j(t_j, x_j, x_{j-1}),$$

where

$$S_j(t_j, x_j, x_{j-1}) = \frac{|x_j - x_{j-1}|^2}{2t_j} + t_j \omega_j(t_j, x_j, x_{j-1}).$$

For any  $m \geq 2$  there exists a positive constant  $\kappa_m$  such that

$$\sup_{x_j, x_{j-1}} |\partial_{x_j}^\alpha \partial_{x_{j-1}}^\beta \omega_j(t_j, x_j, x_{j-1})| \leq \kappa_m$$

if  $2 \leq \alpha + \beta \leq m$ .

We will give two examples of phase functions satisfying hypothesis (H.1).

**Example 1.** Let  $L(\xi, x) = \frac{1}{2}\xi^2 - V(x)$ ,  $(\xi, x) \in \mathbb{R}^2$ , be a Lagrangian with a potential  $V(x)$ . Assume that the potential  $V(x)$  is a real-valued  $C^\infty$ -function satisfying estimates:

$$\sup_x |V^{(k)}(x)| < \infty \quad \text{for any } k \geq 2.$$

Then for a small  $T > 0$ , there exists a unique classical orbit  $\gamma^{cl}(t)$  such that  $\gamma^{cl}(0) = y, \gamma^{cl}(T) = x$ . Let

$$S^{cl}(T, x, y) = \int_0^T L(\dot{\gamma}^{cl}(t), \gamma^{cl}(t)) dt$$

be the classical action. Then  $S^{cl}(T, x, y)$  is of the form

$$S^{cl}(T, x, y) = \frac{|x - y|^2}{2T} + T\phi^{cl}(T, x, y)$$

and for any  $m \geq 2$  there exists a constant  $C_m$  such that

$$\sup_x |\partial_x^\alpha \partial_y^\beta \phi^{cl}(T, x, y)| \leq C_m$$

if  $2 \leq \alpha + \beta \leq m$ . Therefore,  $S(x_L, \dots, x_0) = \sum_{j=1}^L S(t_j, x_j, x_{j-1})$  satisfies the hypothesis (H.1).

**Example 2.** Let  $L(\xi, x)$  be the same lagrangian. Let  $\gamma^{ln}(t)$  be the straight line connecting  $(0, y)$  and  $(T, x)$  in the time-space, i.e.,

$$\gamma^{ln}(t) = \frac{t}{T}x + \frac{T-t}{T}y.$$

Let

$$S^{ln}(T, x, y) = \int_0^T L(\dot{\gamma}^{ln}(t), \gamma^{ln}(t)) dt.$$

Then function  $S^{ln}(T, x, y)$  is of the form

$$S^{ln}(T, x, y) = \frac{|x - y|^2}{2T} + T\phi^{ln}(T, x, y)$$

and for any  $m \geq 2$  there exists a positive constant  $C_m$  such that

$$\sup_x |\partial_x^\alpha \partial_y^\beta \phi^{ln}(T, x, y)| \leq C_m$$

if  $2 \leq \alpha + \beta \leq m$ . Therefore,  $S^{ln}(x_L, \dots, x_0) = \sum_{j=1}^L S^{ln}(t_j, x_j, x_{j-1})$  satisfies the hypothesis (H.1).

Under hypothesis (H.1) the critical point of the function  $(x_{L-1}, \dots, x_1) \rightarrow S(x_L, x_{L-1}, \dots, x_1, x_0)$  is unique if  $T_L = \sum_{j=1}^L t_j$  is small. We denote it by  $(x_{L-1}^*, \dots, x_1^*)$ . We abbreviate  $S(x_L, x_{L-1}^*, \dots, x_1^*, x_0)$  as  $S(\overline{x_L}, \overline{x_0})$ . We can write the Hessian of  $S$  at the critical point as  $H + W$ , where

$$H = \begin{pmatrix} \frac{1}{t_1} + \frac{1}{t_2} & -\frac{1}{t_2} & 0 & 0 & \dots \\ -\frac{1}{t_2} & \frac{1}{t_2} + \frac{1}{t_3} & -\frac{1}{t_3} & 0 & \dots \\ 0 & -\frac{1}{t_3} & \frac{1}{t_3} + \frac{1}{t_4} & -\frac{1}{t_4} & \dots \\ \vdots & \vdots & \vdots & \vdots & \vdots \end{pmatrix}$$

and

$$W = \begin{pmatrix} t_1 \partial_{x_1}^2 \omega_1 + t_2 \partial_{x_1}^2 \omega_2 & t_2 \partial_{x_1} \partial_{x_2} \omega_2 & 0 & \dots \\ t_2 \partial_{x_1} \partial_{x_2} \omega_2 & t_2 \partial_{x_2}^2 \omega_2 + t_3 \partial_{x_2}^2 \omega_3 & t_3 \partial_{x_2} \partial_{x_3} \omega_3 & \dots \\ 0 & t_3 \partial_{x_2} \partial_{x_3} \omega_3 & t_3 \partial_{x_3}^2 \omega_3 + t_4 \partial_{x_3}^2 \omega_4 & \dots \\ \vdots & \vdots & \vdots & \vdots \end{pmatrix}.$$

It is clear that

$$\det H = \frac{T_L}{t_1 t_2 \dots t_L} \neq 0.$$

We can state our first result.

**Theorem 1.** *Under the hypothesis (H.1) there exists a positive constant  $\delta_1$  independent of  $L$  such that if  $T_L = t_1 + \dots + t_L \leq \delta_1$  then*

$$\begin{aligned} & I(\{t_j\}, S, 1, \nu)(x_L, x_0) \\ &= \left( \frac{\nu i}{2\pi T_L} \right)^{1/2} e^{-i\nu S(\overline{x_L}, \overline{x_0})} [\det(I + H^{-1}W)]^{-1/2} (1 + r(\nu, x_L, x_0)), \end{aligned}$$

where the remainder term  $r(\nu, x_L, x_0)$  satisfies the estimate: For any  $K \geq 0$  there exists positive constants  $C_K$  such that if  $|\alpha_0|, |\alpha_L| \leq K$

$$|\partial_{x_0}^{\alpha_0} \partial_{x_L}^{\alpha_L} r(\nu, x_L, x_0)| \leq C_K T_L^3 \nu^{-1}.$$

*Remark.*  $\delta_1$  and  $C_K$  are independent of  $L$  as far as  $T_L$  is bounded. Therefore, we can control  $r(\nu, x_L, x_0)$  even when  $L$  tends to  $\infty$ .

In order to state the result for general integral with amplitude  $a(x)$ , we require a little more preparations. Let  $1 \leq k \leq l \leq L$ . Then the

critical point of the function  $(x_{l-1}, \dots, x_{k+1}) \rightarrow \sum_{j=k+1}^l S_j(t_j, x_j, x_{j-1})$  is unique if  $t_{k+1} + \dots + t_l$  is small. Let  $(x_{l-1}^*, \dots, x_{k+1}^*)$  denote the critical point, which is a function of  $x_l$  and  $x_k$ . We abbreviate  $a(x_L, \dots, x_l, x_{l-1}^*, \dots, x_{k+1}^*, x_k, \dots, x_0)$  to  $a(x_L, \dots, x_{l+1}, \overline{x_l, x_k}, x_{k-1}, \dots, x_0)$ .

Our hypothesis concerning the amplitude function is the following:

**Hypothesis (H.2).** *For any integer  $K \geq 0$  there exists a positive constant  $A_K$  with the following properties: (i) If  $|\alpha_j| \leq K$  for  $j = 0, 1, \dots, L$ , then*

$$\left| \prod_{j=0}^L \partial_{x_j}^{\alpha_j} a(x_L, \dots, x_0) \right| \leq A_K.$$

(ii) *For any sequence of positive integers  $\{j_1, \dots, j_s\}$  satisfying*

$$0 = j_0 < j_1 - 1 < j_1 < j_2 - 1 < \dots < j_s - 1 < j_s < L$$

*we have*

$$\left| \partial_{x_0}^{\alpha_0} \partial_{x_L}^{\alpha_L} \prod_{k=1}^s \partial_{x_{j_k-1}}^{\alpha_{j_k-1}} \partial_{x_{j_k}}^{\alpha_{j_k}} a(\overline{x_L, x_{j_s}, x_{j_s-1}, x_{j_{s-1}}, \dots, x_{j_1-1}, x_{j_0}) \right| \leq A_K,$$

*as far as  $|\alpha_j| \leq K$  for  $j = 0, j_1 - 1, j_1, \dots, j_s - 1, j_s, L$ .*

Before stating our second theorem, we give an example of amplitude functions satisfying hypothesis (H.2).

**Example.** Let  $b_j(x_j, x_{j-1})$ ,  $j = 1, \dots, L$ , be functions bounded together with their derivatives of all order, i.e., for any positive integer  $K$  there exists  $C_K$  such that

$$\sup_x \left| \partial_{x_j}^{\alpha_j} \partial_{x_{j-1}}^{\alpha_{j-1}} b_j(x_j, x_{j-1}) \right| \leq C_K \quad 0 \leq \alpha_j, \alpha_{j-1} \leq K.$$

Then  $a(x_L, \dots, x_0) = e^{(\sum_{j=1}^L t_j b_j(x_j, x_{j-1}))}$  satisfies hypothesis (H.2) above.

Now we can state our main

**Theorem 2.** *Under the hypotheses (H.1) and (H.2) there exists a positive constant  $\delta_1$  such that if  $0 < T_L \leq \delta_1$*

$$\begin{aligned} & I(\{t_j\}, S, a, \nu)(x_L, x_0) \\ &= \left( \frac{\nu i}{2\pi T_L} \right)^{1/2} e^{-i\nu S(\overline{x_L, x_0})} [\det(I + H^{-1}W)]^{-1/2} \\ & \quad \times (a(\overline{x_L, x_0}) + r(\nu, x_L, x_0)), \end{aligned}$$

where  $r(\nu, x_L, x_0)$  satisfies the estimate: For any  $K \geq 0$  there exists positive constants  $C_K$  and  $M(K)$  such that if  $|\alpha_0|, |\alpha_L| \leq K$  we have

$$|\partial_{x_0}^{\alpha_0} \partial_{x_L}^{\alpha_L} r(\nu, x_L, x_0)| \leq C_K T_L \nu^{-1} A_{M(K)}.$$

*Remark.*  $\delta_1$ ,  $C_K$  and  $M(K)$  are independent of  $L$  as far as  $T_L$  is bounded. Therefore, we can control  $r(\nu, x_L, x_0)$  even when  $L$  tends to  $\infty$ .

### §3. Sketch of the proof

We begin with our key lemma, which is valid under hypothesis (H.3) weaker than (H.2) and is interesting in its own sake.

**Hypothesis (H.3).** For any integer  $K \geq 0$  there exists a positive constant  $A_K$  such that if  $|\alpha_j| \leq K$  for  $j = 0, 1, \dots, L$ ,

$$\left| \prod_{j=0}^L \partial_{x_j}^{\alpha_j} a(x_L, \dots, x_0) \right| \leq A_K.$$

We can state

**Key Lemma.** Under the hypotheses (H.1) and (H.3) there exists a positive constant  $\delta_0$  such that if  $T_L \leq \delta_0$  we have

$$\begin{aligned} I(\{t_j\}, S, a, \nu)(x_L, x_0) \\ = \left( \frac{\nu i}{2\pi T_L} \right)^{1/2} e^{-i\nu S(\overline{x_L}, \overline{x_0})} [\det(I + H^{-1}W)]^{-1/2} b(\nu, x_L, x_0), \end{aligned}$$

where  $b(\nu, x_L, x_0)$  satisfies the estimate: For any  $K \geq 0$  there exists positive constants  $C_1(K)$  and  $M(K)$  such that if  $|\alpha_0|, |\alpha_L| \leq K$  we have

$$|\partial_{x_0}^{\alpha_0} \partial_{x_L}^{\alpha_L} b(\nu, x_L, x_0)| \leq C_1(K)^L A_{M(K)}.$$

*Remark.*  $C(K)$  and  $M(K)$  are independent of  $\{t_j\}$ ,  $L$ ,  $(x_L, x_0)$  and  $\nu$  as long as  $T_L \leq \delta_0$ .

Above Lemma can be proved by modifying the proof of Theorem 6.8 in Chapt. 10 of Kumano-go [7].

Omitting the proof of lemma we proceed to the proof of Theorem 2. To make notations simpler we denote  $\frac{\nu i}{2\pi}$  by  $E$ . With this notation we can write

$$I(\{t_j\}, S, a, \nu)(x_L, x_0) = \prod_{j=1}^L \left( \frac{E}{t_j} \right)^{1/2} \int_{R^{L-1}} e^{-i\nu S(x_L, \dots, x_0)} a(x_L, \dots, x_0) \prod_{j=1}^{L-1} dx_j.$$

We perform integration over  $x_1$ -space. Using stationary phase method, we have

$$\begin{aligned} & \prod_{j=1}^2 \left( \frac{E}{t_j} \right)^{1/2} \int_R e^{-i\nu \{S_2(t_2, x_2, x_1) + S_1(t_1, x_1, x_0)\}} a(x_L, \dots, x_2, x_1, x_0) dx_1 \\ &= \left( \frac{E}{T(2, 1)} \right)^{1/2} e^{-i\nu S_{21}^*(x_2, x_0)} (P_1 a(x_L, \dots, x_2, x_0) + R_1 a(x_L, \dots, x_2, x_0)). \end{aligned}$$

Here  $T(2, 1) = t_2 + t_1$ ,  $S_{21}^*(x_2, x_0)$  denotes the critical value of  $S_2(t_2, x_2, x_1) + S_1(t_1, x_1, x_0)$  with respect to the variable  $x_1$ ,  $P_1 a$  is the main part and  $R_1 a$  is the remainder term of the stationary phase method.

*Remark.* (A) Clearly, we have

$$P_1(a)(x_L, \dots, x_2, x_0) = a(x_L, x_{L-1}, \dots, \overline{x_2, x_0}) D(S_1 + S_2; x_2, x_0)^{-1/2}$$

here

$$D(S_1 + S_2; x_2, x_0) = 1 + \frac{t_1 t_2}{t_1 + t_2} (t_2 \partial_{x_1}^2 \omega_2(t_2, x_2, x_1^*) + t_1 \partial_{x_1}^2 \omega_1(t_1, x_1^*, x_0)).$$

(B) The remainder term  $R_1 a$  is a very complicated function with respect to  $x_2$  but is simple with respect to the variable  $(x_L, \dots, x_3, x_0)$ . In fact, we have  $\partial_{x_j}(R_1 a) = R_1 \partial_{x_j} a$  for  $j = 0$  and  $3 \leq j \leq L$ . And  $R_1 a$  is small in the following sense: For any integer  $K \geq 0$  there exists a constant  $C_K$  such that

$$\begin{aligned} & |\partial_{x_0}^{\alpha_0} \partial_{x_2}^{\alpha_2} \dots \partial_{x_L}^{\alpha_L} R_1 a(x_L, \dots, x_2, x_0)| \\ & \leq C_K \nu^{-1} \frac{t_1 t_2}{t_1 + t_2} \max_{x_1} \sup |\partial_{x_0}^{\alpha_0} \partial_{x_1}^{\beta_1} \partial_{x_2}^{\beta_2} \partial_{x_3}^{\alpha_3} \dots \partial_{x_L}^{\alpha_L} a(x_L, \dots, x_2, x_1, x_0)|. \end{aligned}$$

Here max is taken with respect to  $\beta_1, \beta_2$  for  $\beta_1 \leq \alpha_2 + 4, \beta_2 \leq \alpha_2$ .

Next we integrate the term  $P_1a$  over  $x_2$ -space and apply the stationary phase method. We obtain

$$\begin{aligned} & \left(\frac{E}{t_3}\right)^{1/2} \left(\frac{E}{T(2,1)}\right)^{1/2} \\ & \int_R e^{-i\nu\{S_3(t_3, x_3, x_2) + S_{21}^*(x_2, x_0)\}} P_1a(x_L, \dots, x_2, x_0) dx_2 \\ & = \left(\frac{E}{T(3,1)}\right)^{1/2} e^{-i\nu S_{31}^*(x_3, x_0)} \\ & \quad (P_2P_1a(x_L, \dots, x_3, x_0) + R_2P_1a(x_L, \dots, x_3, x_0)). \end{aligned}$$

Here  $S_{31}^*(x_3, x_0)$  denotes the critical value of the function  $x_2 \rightarrow S_3(t_3, x_3, x_2) + S_{21}^*(x_2, x_0)$ ,  $P_2P_1a$  is the main term and  $R_2P_1a$  is the remainder. Since  $P_2P_1a$  is a simple function of  $x_3$ , we integrate it over  $x_3$ -space and apply the stationary phase method. The main term includes  $P_3P_2P_1a$  and the remainder includes  $R_3P_2P_1a$ .

Repeating this procedure  $L - 1$  times, we obtain

$$A_0(x_L, x_0) = \left(\frac{E}{T(L,1)}\right)^{1/2} e^{-i\nu S_{L1}^*(x_L, x_0)} P_{L-1} \dots P_1a(x_L, x_0),$$

which is nothing but the main term of Theorem 2.

Now we must treat the remainder term. Since  $R_1a$  is a complicated function of  $x_2$ , we skip integration over  $x_2$  space and perform integration over  $x_3$ -space. Then we obtain

$$\begin{aligned} & \left(\frac{E}{t_4}\right)^{1/2} \left(\frac{E}{t_3}\right)^{1/2} \left(\frac{E}{T(2,1)}\right)^{1/2} \\ & \int_R e^{-i\nu\{S_4(t_4, x_4, x_3) + S_3(t_3, x_3, x_2) + S_{21}^*(x_2, x_0)\}} R_1a(x_L, \dots, x_4, x_3, x_2, x_0) dx_3 \\ & = \left(\frac{E}{T(4,3)}\right)^{1/2} \left(\frac{E}{T(2,1)}\right)^{1/2} e^{-i\nu\{S_{43}^*(x_4, x_2) + S_{21}^*(x_2, x_0)\}} \\ & \quad (P_3R_1a(x_L, \dots, x_4, x_2, x_0) + R_3R_1a(x_L, \dots, x_4, x_2, x_0)). \end{aligned}$$

Here  $S_{43}^*(x_4, x_2)$  denotes the critical value of the function  $x_3 \rightarrow S_4(t_4, x_4, x_3) + S_3(t_3, x_3, x_2)$ ,  $P_3R_1a$  denotes the main term and  $R_3R_1a$  is the remainder.  $P_3R_1a$  is a simple function of the variable  $x_4$  but  $R_3R_1a$  is not. We integrate  $P_3R_1a$  over  $x_4$ -space but we skip integration of  $R_3R_1a$  over  $x_4$ -space.

Similarly, we skip integration of  $R_2P_1a$  over  $x_3$ -space and integrate it over  $x_4$ -space. We obtain

$$\begin{aligned} & \left(\frac{E}{t_5}\right)^{1/2} \left(\frac{E}{t_4}\right)^{1/2} \left(\frac{E}{T(3,1)}\right)^{1/2} \\ & \int_R e^{-i\nu\{S_5(t_5, x_5, x_4) + S_4(t_4, x_4, x_3) + S_{31}^*(x_3, x_0)\}} R_2P_1a(x_L, \dots, x_4, x_3, x_0) dx_4 \\ & = \left(\frac{E}{T(5,4)}\right)^{1/2} \left(\frac{E}{T(3,1)}\right)^{1/2} e^{-i\nu\{S_{54}^*(x_5, x_3) + S_{31}^*(x_3, x_0)\}} \\ & \quad (P_4R_2P_1a(x_L, \dots, x_5, x_3, x_0) + R_4R_2P_1a(x_L, \dots, x_5, x_3, x_0)). \end{aligned}$$

We continue this process. *The rule is that we apply the stationary phase method when we integrate over  $x_k$ -space and if  $R_k$  appears then we skip integration over  $x_{k+1}$ -space.* We finally obtain the following expression:

$$I(\{t_j\}, S, a, \nu)(x_L, x_0) = A_0(x_L, x_0) + \sum^* A_{j_s j_{s-1} \dots j_1}(x_L, x_0),$$

where  $\sum^*$  denotes summation with respect to indices  $(j_s, \dots, j_1)$  satisfying

$$1 < j_1 < j_2 - 1 < j_2 < j_3 - 1 < \dots < j_s - 1 < j_s,$$

and each term is an oscillatory integral

$$\begin{aligned} & A_{j_1 j_2 \dots j_s}(x_L, x_0) \\ & = \prod_{m=1}^s \left(\frac{E}{T(j_m, j_m - 1)}\right)^{1/2} \\ & \quad \int_{R^s} e^{-i\nu S_{j_s \dots j_1}(x_L, x_{j_s}, \dots, x_{j_1}, x_0)} b_{j_s \dots j_1}(x_L, x_{j_s}, \dots, x_{j_1}, x_0) \prod_{m=1}^s dx_{j_m}, \end{aligned}$$

whose phase function is

$$\begin{aligned} & S_{j_s \dots j_1}(x_L, x_{j_s}, \dots, x_{j_1}, x_0) \\ & = S_{L j_s}^*(x_L, x_{j_s}) + S_{j_s j_{s-1}}^*(x_{j_s}, x_{j_{s-1}}) + \dots + S_{j_1 0}^*(x_{j_1}, x_0) \end{aligned}$$

and the amplitude is

$$b_{j_s \dots j_1}(x_L, x_{j_s}, \dots, x_{j_1}, x_0) = Q_{L-1} Q_{L-2} \dots Q_1 a(x_L, x_{j_s}, \dots, x_{j_1}, x_0),$$

with

$$Q_j = \begin{cases} Id, & \text{for } j = j_s, j_{s-1}, \dots, j_1, \\ R_j, & \text{for } j = j_s - 1, j_{s-1} - 1, \dots, j_1 - 1, \\ P_j, & \text{otherwise.} \end{cases}$$

Furthermore, we can prove that  $b_{j_s \dots j_1}(x_L, x_{j_s}, \dots, x_{j_1}, x_0)$  satisfies hypothesis (H.3).

**Proposition.** *For any integer  $K \geq 0$  there exist positive constants  $C_2(K)$  and integer  $m(K)$  such that*

$$\begin{aligned} & \left| \partial_{x_L}^{\alpha_L} \partial_{x_{j_s}}^{\alpha_{j_s}} \dots \partial_{x_{j_1}}^{\alpha_{j_1}} \partial_{x_0}^{\alpha_0} b_{j_s \dots j_1}(x_L, x_{j_s}, \dots, x_{j_1}, x_0) \right| \\ & \leq C_2(K)^s A_{m(K)} \prod_{k=1}^s \nu^{-1} t_{j_k}. \end{aligned}$$

Now we apply our key lemma to  $A_{j_s j_{s-1} \dots j_1}(x_L, x_0)$  and use the proposition above. Then we obtain

$$A_{j_s j_{s-1} \dots j_1}(x_L, x_0) = \left( \frac{E}{T_{L,1}} \right)^{1/2} e^{-i\nu S(\overline{x_L, x_0})} a_{j_s j_{s-1} \dots j_1}(x_L, x_0),$$

where the function  $a_{j_s j_{s-1} \dots j_1}(x_L, x_0)$  satisfies the following estimates: For any integer  $K \geq 0$  we have

$$\left| \partial_{x_L}^{\alpha_L} \partial_{x_0}^{\alpha_0} a_{j_s j_{s-1} \dots j_1}(x_L, x_0) \right| \leq C_1(K)^s C_2(M(K))^s A_{m(M(K))} \prod_{k=1}^s \nu^{-1} t_{j_k}.$$

This implies that the remainder term  $r(\nu, x_L, x_0)$  can be written as

$$r(\nu, x_L, x_0) = \sum^* a_{j_s j_{s-1} \dots j_1}(x_L, x_0).$$

If  $\alpha_0, \alpha_L \leq K$  we have

$$\begin{aligned} \left| \partial_{x_L}^{\alpha_L} \partial_{x_0}^{\alpha_0} r(\nu, x_L, x_0) \right| & \leq \sum^* \left| \partial_{x_L}^{\alpha_L} \partial_{x_0}^{\alpha_0} a_{j_s j_{s-1} \dots j_1}(x_L, x_0) \right| \\ & \leq \sum^* C_3(K)^s A_{m(M(K))} \prod_{k=1}^s \nu^{-1} t_{j_k} \\ & \leq A_{m(M(K))} \left( \prod_{j=1}^L (1 + C_3(K) \nu^{-1} t_j) - 1 \right), \end{aligned}$$

where we abbreviated  $C_1(K)C_2(M(K))$  as  $C_3(K)$ . This proves Theorem 2.

Theorem 1 can be proved similarly.

More detailed discussions are given by [2].

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