## 107. Interaction of Two Nonlinear Waves at the Boundary

By Toru Sasaki<br>Department of Mathematics, University of Tokyo<br>(Communicated by Kôsaku Yosida, m. J. A., Dec. 14, 1987)

§ 0. Introduction. In this paper, we will consider the following nonlinear mixed probrem :

$$
\begin{aligned}
& \left(\frac{\partial^{2}}{\partial t^{2}}-\frac{\partial^{2}}{\partial x^{2}}-\frac{\partial^{2}}{\partial y^{2}}\right) u=f\left(\frac{\partial u}{\partial y}\right) \\
& \quad \text { in } \Omega_{M}^{+}=\left\{(t, x, y) \in R^{3} ; 0<y,|t|<M\right\} \\
& \left.u\right|_{y=0}=0
\end{aligned}
$$

where $f$ is a smooth function which will be chosen later and $M$ a positive number.

On the semilinear Cauchy problem, Rauch and Reed [5] have shown by a simple example that anomalous singularities arise when three characteristic hypersurfaces $\Sigma_{1}, \Sigma_{2}$ and $\Sigma_{3}$, carrying progressing waves intersect. On the other hand, Bony [2], [3] and Melrose and Ritter [4] have shown that the phenomena of interactions of singularities do not occur when two hypersurfaces $\Sigma_{1}$ and $\Sigma_{2}$ intersect.

In the case of nonlinear mixed problem, Beals and Métivier [1] have shown that when single characteristic hypersurface hits the boundary transversally, then the solution will be conormal with respect to the union of the surface and the reflected characteristic hypersurface.

We will apply the method of [5] to the nonlinear mixed problem and show by an example that anomalous singularities arise when even two hypersurfaces hit the boundary at the same time.

The author expresses his sincere gratitude to Prof. H. Komatsu, Dr. M. Yamazaki, and Dr. N. Tose for valuable suggestions.
§ 1. Singularities of the solution to a linear problem. In this section, we will estimate from below the singularities of the solution $V$ of the equations

$$
\begin{aligned}
& \left(\frac{\partial^{2}}{\partial t^{2}}-\frac{\partial^{2}}{\partial x^{2}}-\frac{\partial^{2}}{\partial y^{2}}\right) V=\chi_{\Gamma^{+}} \\
& \left.V\right|_{y=0}=0 \\
& \left.V\right|_{t=0}=\left.\frac{\partial V}{\partial t}\right|_{t=0}=0 .
\end{aligned}
$$

Here $\Gamma^{+}=\left\{(t, x, y) \in R^{3} ; y>0, t>0, y<-x+\sqrt{2} t, y<x+\sqrt{2} t\right\}$ and $\chi_{\Gamma^{+}}$is its characteristic function.

Proposition 1. Sing supp $V$ contains the forward light cone $C_{0}^{+}$with vertex at origin.

Proof. We will consider sing supp $V \cap\{t=1\}$. For general $t$, one can
show the similar assertion.
Case 1. We show that $V$ is not smooth across $C_{0}^{+} \cap\{y>\sqrt{2} / 2\}$. For simplicity we assume $x=0$.

Let

$$
Z_{1}(t, x, y)=\int_{0}^{t} \iint_{R^{2}} E(t-\tau, x-\xi, y-\eta) \cdot \chi_{\Delta}(\tau, \xi, \eta) d \xi d \eta d \tau
$$

where $E$ denotes the fundamental solution of $\partial^{2} / \partial t^{2}-\partial^{2} / \partial x^{2}-\partial^{2} / \partial y^{2}$ and $\Delta=\{(t, x, y) ; y<-x+\sqrt{2} t, y<x+\sqrt{2} t\}$. Hörmander's theorem implies that $Z_{1}$ is smooth across $C_{0}^{+} \cap\{y>\sqrt{2} / 2\}$.

Put $W_{1}=Z_{1}-V$. Then $W_{1}=0$ in $\Gamma^{+} \cap\left\{y>\sqrt{1-x^{2}},|x|<1\right\}$ and $W_{1}$ is written as

$$
W_{1}=\int_{0}^{t} \iint_{R^{2}} E(t-\tau, x-\xi, y-\eta) \cdot\left(\chi_{\Delta \backslash \Gamma^{+}}+\chi_{\Gamma^{-}}\right)(\tau, \xi, \eta) d \xi d \eta d \tau .
$$

Here $\Gamma^{-}=\{(t, x, y) ; t>0, y<0, y>-x-\sqrt{2} t, y>x-\sqrt{2} t\}$. Assume ( $1, x, y$ ) is in the interior region of $C_{0}^{+}$and $\varepsilon$ is the distance of $(x, y)$ and $C_{0}^{+} \cap\{t=1\}$. Then as in [5], we can show that

$$
\begin{aligned}
W_{1} & \geqq \int_{0}^{t} \iint_{R^{2}} E(t-\tau, x-\xi, y-\eta) \cdot \chi_{\Delta \Gamma^{+}}(\tau, \xi, \eta) d \xi d \eta d \tau \\
& \geqq(\mathrm{const}) \varepsilon^{7 / 2} .
\end{aligned}
$$

This shows that $W_{1}$ is not smooth across $C_{0}^{+} \cap\{y>\sqrt{2} / 2\}$. Thus $V$ is not smooth across $C_{0}^{+} \cap\{y>\sqrt{2} / 2\}$.

Case 2. We show that $V$ is not smooth across $C_{0}^{+} \cap\{0<y<\sqrt{2} / 2\}$. Put $\Lambda=\{(t, x, y) ; y<x+\sqrt{2} t, y>-x-\sqrt{2} t\}, \Lambda^{ \pm}=\Lambda \cap\{ \pm y>0\}$, and

$$
Z_{2}=\int_{0}^{t} \iint_{R^{2}} E(t-\tau, x-\xi, y-\eta) \cdot\left(\chi_{A^{+}}-\chi_{A^{-}}\right)(\tau, \xi, \eta) d \xi d \eta d \tau
$$

As above, $Z_{2}$ is smooth near the light cone $C_{0}^{+} \cap\{0<y<\sqrt{2}\}$.
Let

$$
\begin{aligned}
W_{2}(t, x, y) & =Z_{2}(t, x, y)-V(t, x, y) \\
& =\int_{0}^{t} \iint_{R^{2}} E(t-\tau, x-\xi, y-\eta) \cdot\left(\chi_{\Lambda^{+} \backslash \Gamma^{+}}-\chi_{\Lambda^{-} \backslash \Gamma^{-}}\right)(\tau, \xi, \eta) d \xi d \eta d \tau .
\end{aligned}
$$

The integral is integrated over the intersection of $\left(\Lambda^{+} \backslash \Gamma^{+}\right) \cup\left(\Lambda^{-} \backslash \Gamma^{-}\right)$and $\tilde{C}_{(t, x, y)}$, where $\tilde{C}_{(t, x, y)}$ denotes the interior region of the backward light cone with vertex at $(t, x, y)$.

It is sufficient to estimate the value of $W_{2}$ from below when $(1, x, y)$ is in $\tilde{C}_{0}^{+}$. Suppose $x<0,(1, x, y) \in \tilde{C}_{0}^{+}$and put $\varepsilon=\sqrt{1-x^{2}}-y$, then $\varepsilon$ is positive. Put

$$
T_{1}=\left\{\tau \in(0,1) ; C_{(t, x, y)}^{-} \cap\left(\Lambda^{-} \backslash \Gamma^{-}\right) \neq \varnothing\right\}
$$

and

$$
T_{2}=\left\{\tau \in(0,1) ; C_{(t, x, y)}^{-} \cap\left(\Lambda^{-} \backslash \Gamma^{-}\right)=\varnothing, C_{(t, x, y)}^{-} \cap\left(\Lambda^{+} \backslash \Gamma^{+}\right) \neq \varnothing\right\}
$$

Then we have

$$
\begin{aligned}
W_{2} & =\int_{T_{1}} \iint_{R^{2}} E(t-\tau, x-\xi, y-\eta) \cdot\left(\chi_{A^{+} \backslash \Gamma^{+}}-\chi_{A^{-} \backslash \Gamma^{-}}\right)(\tau, \xi, \eta) d \xi d \eta d \tau \\
& +\int_{T_{2}} \iint_{R^{2}} E(t-\tau, x-\xi, y-\eta) \cdot\left(\chi_{A^{+} \backslash \Gamma^{+}}-\chi_{A^{-} \backslash \Gamma^{-}}\right)(\tau, \xi, \eta) d \xi d \eta d \tau \\
& =W_{2}^{1}+W_{2}^{2} .
\end{aligned}
$$

$W_{2}^{1}(t, x, y)$ is positive because $E(t-\tau, x-\xi, y-\eta)$ takes constant value along the arc $\left\{(\xi, \eta) ;(x-\xi)^{2}+(y-\eta)^{2}=\tau_{0}\right\}\left(0<\tau_{0}<1-\tau\right)$, and for each arc the part on which the integrand is positive is longer than the other part. Thus we have

$$
W_{2}(t, x, y) \geqq W_{2}^{2}(t, x, y) \geqq(\text { const }) \varepsilon^{7 / 2} .
$$

This completes the proof of Proposition 1.
Proposition 2. For each $M>0$, there exists $c_{M}>0$ such that

$$
\left|\frac{\partial V}{\partial y}\right|<c_{M} \quad \text { in } \Omega_{M}^{+} .
$$

Proof. Let $\Omega_{M, R}^{+}=\Omega_{M}^{+} \cap\{(t, x, y) ;|x|<R,|y|<R\}$. It is sufficient to show that there exists $c_{m, R}>0$ such that

$$
\left|\frac{\partial}{\partial y} V\right|<c_{M, R} \quad \text { in } \Omega_{M, R}^{+}
$$

for all $R>0$. For supp $V \cap \Omega_{M}^{+}$is contained in $\Omega_{M, R}^{+}$if $R$ is sufficiently large. We write $V$ as follows:

$$
\begin{aligned}
V(t, x, y)= & \int_{0}^{t} \iint_{R^{2}} E(t-\tau, x-\xi, y-\eta) \cdot \chi_{\Gamma^{+}}(\tau, \xi, \eta) d \xi d \eta d \tau \\
& -\int_{0}^{t} \iint_{R^{2}} E(t-\tau, x-\xi, y-\eta) \cdot \chi_{\Gamma^{-}}(\tau, \xi, \eta) d \xi d \eta d \tau \\
= & V_{+}-V_{-} .
\end{aligned}
$$

We will estimate $(\partial / \partial y) V_{+}$. It is easy to see that

$$
\begin{aligned}
\frac{\partial}{\partial y} V_{+}(t, x, y)= & \frac{1}{2 \pi} \int_{0}^{t} \int_{-\infty}^{\infty} \frac{Y\left(t-\tau-\sqrt{(x-\xi)^{2}+y^{2}}\right)}{\sqrt{(t-\tau)^{2}-(x-\xi)^{2}-y^{2}}} d \tau d \xi \\
& -\frac{1}{2 \pi} \int_{0}^{t} \int_{0}^{\infty} \frac{Y\left(t-\tau-\sqrt{(x-\xi)^{2}+(y+\xi-\sqrt{2} t)^{2}}\right)}{\sqrt{(t-\tau)^{2}-(x-\xi)^{2}-(y+\xi-\sqrt{2} t)^{2}}} d \tau d \xi \\
& -\frac{1}{2 \pi} \int_{0}^{t} \int_{-\infty}^{0} \frac{Y\left(t-\tau-\sqrt{(x-\xi)^{2}+(y-\xi-\sqrt{2} t)^{2}}\right)}{\sqrt{(t-\tau)^{2}-(x-\xi)^{2}-(y-\xi-\sqrt{2} t)^{2}}} d \tau d \xi .
\end{aligned}
$$

Here $Y$ denotes the Heaviside function. If $(t, x, y)$ is in $\Omega_{M, R}^{+}$, then the domain of integration of each term is contained in a compact set. Thus each term is bounded in $\Omega_{M, R}^{+}$. This shows that $(\partial / \partial y) V_{-}$is bounded in $\Omega_{M, R}^{+}$by the same way.
§ 2. Constructing of a nonlinear equation with interacting reflection of singularities. Using $V$ in the last section, we shall oonstruct the solution of the boundary value problem

$$
\begin{aligned}
& \left(\frac{\partial^{2}}{\partial t^{2}}-\frac{\partial^{2}}{\partial x^{2}}-\frac{\partial^{2}}{\partial y^{2}}\right) u=f\left(\frac{\partial u}{\partial y}\right) \quad \text { in } \Omega_{M}^{+} \\
& \left.u\right|_{y=0}=0 .
\end{aligned}
$$

First of all, put $v_{1}(t, x, y)=(y-x+\sqrt{2} t)_{+}-(-y-x+\sqrt{2} t)_{+}$, and $v_{2}(t, x, y)$ $=(y+x+\sqrt{2} t)_{+}-(-y+x+\sqrt{2} t)_{+}$. Here

$$
X_{+}= \begin{cases}X & (X \geqq 0) \\ 0 & (X<0)\end{cases}
$$

Then $v_{1}$ and $v_{2}$ satisfy

$$
\begin{aligned}
& \left(\frac{\partial^{2}}{\partial t^{2}}-\frac{\partial^{2}}{\partial x^{2}}-\frac{\partial^{2}}{\partial y^{2}}\right) v=0 \\
& \left.v\right|_{y=0}=0
\end{aligned}
$$

Moreover for any positive $\delta, u_{\dot{\delta}}=v_{1}+v_{2}+\delta V$ satisfies the equation

$$
\begin{aligned}
& \left(\frac{\partial^{2}}{\partial t^{2}}-\frac{\partial^{2}}{\partial x^{2}}-\frac{\partial^{2}}{\partial y^{2}}\right) u_{\delta}=\delta \chi_{\Gamma^{+}} \\
& \left.u_{\delta}\right|_{y=0}=0 .
\end{aligned}
$$

On the other hand, $(\partial / \partial y)\left(v_{1}+v_{2}\right) \leqq 2$ if $t$ is negative. When $t$ is positive, $(\partial / \partial y)\left(v_{1}+v_{2}\right)=4$ in $\Gamma^{+}$and $(\partial / \partial y)\left(v_{1}+v_{2}\right) \leqq 3$ outside $\Gamma^{+}$.

Thus we can see from Proposition 2 that ( $\partial / \partial y$ ) $u_{\delta} \geqq 4-\delta c_{M}$ in $\Gamma^{+} \cap \Omega_{M}^{+}$ and $(\partial / \partial y) u_{i} \leqq 3+\delta c_{M}$ in $\Omega_{M}^{+} \backslash \Gamma^{+}$. We take $\delta$ so small that $\delta c_{M}<1 / 3$ and then take smooth function $f_{\delta}$ as

$$
f_{\dot{\delta}}(X)= \begin{cases}\delta & (X>4-1 / 3) \\ 0 & (X<3+1 / 3)\end{cases}
$$

Then $f_{\delta}\left((\partial / \partial y) u_{j}\right)=\delta \chi_{\Gamma^{+}}$in $\Omega_{M}^{+}$, and $u_{\delta}$ satisfies the equations

$$
\begin{aligned}
& \left(\frac{\partial^{2}}{\partial t^{2}}-\frac{\partial^{2}}{\partial x^{2}}-\frac{\partial^{2}}{\partial y^{2}}\right) u=f_{\delta}\left(\frac{\partial}{\partial y} u\right) \quad \text { in } \Omega_{M}^{+} \\
& \left.u\right|_{y=0}
\end{aligned}
$$

Finally we consider the singularities of the solution $u_{\dot{j}}$. For negative $t$, sing supp $u_{\delta}$ is contained only four characteristic planes. But Proposition 1 assures that for positive $t$, sing supp $u_{\delta}$ contains $C_{0}^{+}$besides those planes.

## References

[1] M. Beals and G. Métivier: Progressing wave solutions to certain nonlinear mixed problems. Duke Math. J., 51, 125-137 (1986).
[2] J.-M. Bony: Interaction des singularités pour les équations aux dérivées partielles nonlinéaires. Sém. Goulaouic-Meyer-Schwartz, $\mathrm{n}^{\circ} 10$ (1981/82).
[3] -: Propagation et interaction des singularités par les solutions des équations aux dérivées partielles non linéaires. Proc. Int. Cong. Math. Warzawa, 11331147 (1983).
[4] R. Melrose and N. Ritter: Interaction of non-linear progressing waves for semilinear wave equations. Ann. of Math., 121, 187-213 (1985).
[5] J. Rauch and M. Reed: Singularities produced by the nonlinear interaction of three progressing waves; examples. Comm. in P. D. E., 7, 1117-1133 (1982).

