# 49. Studies on Holonomic Quantum Fields. IV 

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(Communicated by Kôsaku Yosida, M. J. A., Nov. 12, 1977)

This is a continuation of our previous notes [1], [2] and together with the latter constitutes the second part of the work referred to in [1]. We use the same notation as in [1], [2], [3].

1. First we shall show that the wave function $\boldsymbol{w}_{F, n}={ }^{t}\left(\hat{w}_{F, n}^{1}(x)\right.$, $\cdots, \hat{w}_{F, n}^{n}(x)$ ) constructed in [2] forms a basis of $W_{a_{1}, \ldots, a_{n}}^{\text {strict, }}$. By (30) the local expansion of $\boldsymbol{w}_{F, n}$ in the sense of (10) in [1] takes the form

$$
\begin{equation*}
\boldsymbol{w}_{F, n} \sim \frac{i}{2}\left[\sum_{l=0}^{\infty} C_{F, l}[A]_{l} \boldsymbol{w}-\sum_{l=0}^{\infty} \bar{C}_{F, l} w_{l}^{*}[A]\right] \tag{42}
\end{equation*}
$$

where ( $\left.i / 2) C_{F, l}={ }^{t}{ }^{t} \boldsymbol{c}_{l}\left(\hat{w}_{F, n}^{1}\right), \cdots,{ }^{t} \boldsymbol{c}_{l}\left(\hat{w}_{F, n}^{n}\right)\right)$. From (31) it follows that if we write $C_{F, 0}=1-T, T$ is purely imaginary and hermitian: $T=-\bar{T}$ $=-{ }^{t} T$. Since $\boldsymbol{w} \mathscr{R}$ is a basis of $W_{a_{1}, \ldots, a_{n}}^{\text {strict, }}$, there exists a real $n \times n$ matrix $C$ satisfying $\boldsymbol{w}_{F, n}=C \boldsymbol{w} \mathcal{R}$. Comparing the 0 -th coefficients of their local expansions we have $(i / 2) C_{F, 0}=C C_{\mathcal{R}, 0}$ or equivalently $1-T=2 C e^{-H}$. Taking the complex conjugate we have $1+T=2 C e^{H}$, and hence

$$
\begin{equation*}
C=(2 \cosh H)^{-1}, \quad T=\tanh H=(1-G)(1+G)^{-1} . \tag{43}
\end{equation*}
$$

Hence $\boldsymbol{w}_{F, n}$ is also a basis of $W_{a_{1}, \ldots, a_{n}}^{\text {strict }, ~}$.
The relation between $\boldsymbol{w}_{F}$ and $\boldsymbol{w} \mathcal{R}$ enables us to express the coefficients $B, E$ appearing in the system (12) in [1] satisfied by $\boldsymbol{w} \mathcal{R}$, in terms of $\tau_{F, n}$ and $\tau_{F, n}^{\mu \nu}$. From (11), (40), (41) and (43) we have

$$
\begin{array}{ll}
F=\left[U^{-1} V, m A\right], & G=U\left(2 \tau_{F, n}-U\right)^{-1}  \tag{44}\\
B=\sqrt{G} m A \sqrt{G^{-1}}, & E=\sqrt{G} F \sqrt{G^{-1}}
\end{array}
$$

where
(45)

$$
\begin{aligned}
& U=\tau_{F, n}(1-T)=\left(\begin{array}{cccc}
\tau_{F, n} & i \tau_{F, n}^{12} & \cdots & i \tau_{F, n}^{1 n} \\
-i \tau_{F, n}^{12} & \tau_{F, n} & \cdots & i \tau_{F, n}^{2 n} \\
\vdots & \vdots & \ddots & \vdots \\
-i \tau_{F, n}^{1 n} & -i \tau_{F, n}^{2 n} & \cdots & \tau_{F, n}
\end{array}\right) \\
& V=2\left(\begin{array}{cclc}
m^{-1} \partial_{-a_{1}} \tau_{F, n} & i m^{-1} \partial_{-a_{2}} \tau_{F, n}^{12} & \cdots & i m^{-1} \partial_{-a_{n}} \tau_{F, n}^{1 n} \\
-i m^{-1} \partial_{-a_{1}} \tau_{F, n}^{12} & m^{-1} \partial_{-a_{2}} \tau_{F, n} & \cdots & i m^{-1} \partial_{-a_{n}} \tau_{F, n}^{2 n} \\
\vdots & \vdots & \ddots & \vdots \\
-i m^{-1} \partial_{-a_{1}-} \tau_{F, n}^{1 n} & -i m^{-1} \partial_{-a_{2}-} \tau_{F, n}^{2 n} & \cdots & m^{-1} \partial_{-a_{n}} \tau_{F, n}
\end{array}\right] .
\end{aligned}
$$

Thus we have constructed, in terms of $\psi, \varphi_{F}$ and $\varphi^{F}$, not only a solution to the extended holonomic system (12) but also one to the system of total differtntial equations (18).
2. Now we will give a closed expression for $\tau_{F, n}$ by means of solution matrices to the total differential equations (18) in [1]. From (40) and (41) we see that

$$
\begin{equation*}
\omega=d \log \tau_{F, n}=\frac{1}{2}\left(\operatorname{tr} C_{F, 1} m d A+\operatorname{tr} \bar{C}_{F, 1} m d \bar{A}\right) . \tag{46}
\end{equation*}
$$

From (13) in [1] and (43), after a little computation we rewrite (46) in the following form.

$$
\begin{align*}
\omega= & \frac{1}{2} \operatorname{tr}\left[\frac{1}{2} T \Theta-\frac{1}{2} F \Theta+m^{2}\left(-{ }^{t} G \bar{A} G+\bar{A}\right) d A\right]  \tag{47}\\
& + \text { complex conjugate. }
\end{align*}
$$

We note that the 1 -form in the right hand side of (47) is shown to be a closed 1 -form and is invariant under the Euclidean motion group even for an arbitrary solution to (18) in [1].
$\hat{w}_{F, n}^{\nu_{1}, \cdots, \nu_{m}}(x)$ is written as a linear combination of the components of $\boldsymbol{w}_{F, n}$ as follows.

Comparing the local expansion of both sides of (48) we have

$$
\begin{aligned}
& =\text { Pfaffian }\left(i(\tanh H)_{\nu, \nu^{\prime}}\right)_{\nu, \nu^{\prime}}=\nu_{1}, \cdots, \nu_{m} .
\end{aligned}
$$

More generally we have
(50) $\hat{w}_{F, n}^{\nu_{1}, \cdots, \nu_{m}}\left(x_{1}, \cdots, x_{k}\right)=$ Pfaffian

Erratum in Sato-Miwa-Jimbo [3]. The expressions in paragraphs $\S 3$ and §4 should be corrected as follows.
p. 7, line 5 from the bottom:

$$
\left\langle w, w^{\prime}\right\rangle=\frac{1}{2} \int_{-\infty}^{+\infty} m d x^{1}\left(w_{+}(x) w_{+}^{\prime}(x)+w_{-}(x) w_{-}^{\prime}(x)\right)
$$

lines 4-3 from the bottom:

$$
\frac{1}{2} \int_{-\infty}^{+\infty} m d x^{1}\left(w_{+}(x) \psi_{+}(x)+w_{-}(x) \psi_{-}(x)\right)
$$

p. 8, line 13 from the bottom:

$$
\begin{aligned}
\phi_{ \pm}(u)= & \varepsilon(u) \lim _{t \rightarrow \pm \infty} \frac{i}{2} \int_{x^{0}=t} d x^{1}\left(e^{i m(x-u+x+u-1)}\left(\partial / \partial x^{0}\right) \varphi^{F}(x)\right. \\
& \left.-\varphi^{F}(x)\left(\partial / \partial x^{0}\right) e^{i m(x-u+x+u-1)}\right) .
\end{aligned}
$$

## References

[1] M. Sato, T. Miwa, and M. Jimbo: Proc. Japan Acad., 53A, 147-152 (1977).
[ 2 ] ——: ibid., 153-158 (1977).
[ 3] -: ibid., 6-10 (1977).

