## PAPERS COMMUNICATED

## 113. Concircular Geometry IV. Theory of Subspaces.

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In three previous papers, ${ }^{1)}$ we have considered the Concircular Geometry, that is to say, the geometry in which one seeks for the properties of Riemannian spaces invariant under the conformal transformations of the metric

$$
\bar{g}_{\mu \nu}=\rho^{2} g_{\mu \nu} \quad(\lambda, \mu, \nu, \ldots=1,2,3, \ldots, n),
$$

with functions $\rho$ satisfying the following partial differential equations

$$
\rho_{\mu \nu} \equiv \frac{\partial \rho_{\mu}}{\partial u^{\nu}}-\rho_{\lambda}\left\{\lambda_{\mu \nu}^{\lambda}\right\}-\rho_{\mu} \rho_{\nu}+\frac{1}{2} g^{\alpha \beta} \rho_{a} \rho_{\beta} \rho_{\mu \nu}=\phi g_{\mu \nu} \quad\left(\rho_{\mu}=\frac{\partial \log \rho}{\partial u^{\mu}}\right) .
$$

In the present paper, we shall deal with the theory of subspaces in the concircular geometry.
§1. Let us consider a subspace $V_{m}$ immersed in a Riemannian space $V_{n}$ whose parametric representation is

$$
\begin{equation*}
u^{\lambda}=u^{\lambda}\left(u^{\dot{i}}, u^{\dot{2}}, \ldots, u^{\dot{m}}\right) \tag{1.1}
\end{equation*}
$$

where $\left(u^{\lambda}\right)$ and ( $u^{i}$ ) $(i, j, k, \ldots=\dot{1}, \dot{2}, \ldots, \dot{m})$ denote the coordinate systems of $V_{n}$ and $V_{m}$ respectively. A conformal transformation

$$
\begin{equation*}
\bar{g}_{\mu \nu}=\rho^{2} g_{\mu \nu} \tag{1.2}
\end{equation*}
$$

of the fundamental tensor of $V_{n}$, being a concircular one with the function $\rho$ satisfying the equations

$$
\begin{equation*}
\rho_{\mu \nu} \equiv \rho_{\mu ; \nu}-\rho_{\mu} \rho_{\nu}+\frac{1}{2} g^{\alpha \beta} \rho_{a} \rho_{\beta} g_{\mu \nu}=\phi g_{\mu \nu} \tag{1.3}
\end{equation*}
$$

where the semi-colon denotes the covariant differentiation with respect to the Christoffel symbols $\left\{\begin{array}{l}\lambda \nu \nu\end{array}\right\}$ formed with $g_{\mu \nu}$, the induced conformal transformation

$$
\begin{equation*}
g_{j k}=\rho^{2} g_{j k} \tag{1.4}
\end{equation*}
$$

of the fundamental tensor

$$
\begin{equation*}
g_{j k}=g_{\mu \nu} B_{j}^{\mu} B_{k_{k}}^{\nu} \quad\left(B_{j}^{\mu}=\frac{\partial u^{\mu}}{\partial u^{j}}\right) \tag{1.5}
\end{equation*}
$$

of the subspace is not in general a concircular one.

[^0]We shall, first, seek for the subspace $V_{m}$ for which the induced conformal transformation is also a concircular one.

Putting

$$
\begin{equation*}
\rho_{j k} \equiv \rho_{j ; k}-\rho_{j} \rho_{k}+\frac{1}{2} g^{a b} \rho_{a} \rho_{b} g_{j k}, \tag{1.6}
\end{equation*}
$$

where

$$
\begin{equation*}
\rho_{j}=\frac{\partial \log \rho}{\partial u^{j}}=\rho_{\mu} B_{j}^{\mu}, \tag{1.7}
\end{equation*}
$$

and $\rho_{j ; k}$ denotes the covariant derivative of $\rho_{j}$ with respect to the threeindex symbols of Christoffel $\left\{{ }_{j}^{i} k\right\}$ formed with $g_{j k}$, we obtain
or

$$
\rho_{j k}=\rho_{\mu \nu} B_{j}^{\cdot \mu} B_{k}^{\cdot \nu}+\rho_{\mu} H_{\dot{j} \dot{k}}^{\mu}-\frac{1}{2} \rho_{a} \rho_{\beta} B_{A}^{\cdot a} B_{A}^{\dot{\beta}} g_{j k}
$$

where $B_{A^{a}}(A, B, \ldots=\dot{m}+\dot{1}, \ldots \ldots, \dot{n})$ are $n-m$ mutually orthogonal unit vectors normal to $V_{m}$ and

$$
\begin{equation*}
H_{\dot{j} \dot{k}^{\mu}}=\frac{\partial B_{j}^{\cdot \mu}}{\partial u^{k}}+B_{j}^{\cdot a} B_{k}^{; \beta}\left\{\alpha_{a \beta}^{\mu}\right\}-B_{a}^{-\mu}\left\{g_{j k}\right\} . \tag{1.10}
\end{equation*}
$$

The conformal transformation (1.2) being a concircular one, we have

$$
\rho_{\mu \nu}=\phi g_{\mu \nu}
$$

Substituting these equations in (1.9), we have

$$
\begin{equation*}
\rho_{j k}=\rho_{\mu} H_{\ddot{j}_{k}{ }^{\mu}}+\left(\phi-\frac{1}{2} \rho_{a} \rho_{\beta} B_{A}^{-a} B_{A}^{\beta}\right) g_{j k} . \tag{1.11}
\end{equation*}
$$

If we suppose that the induced conformal transformation (1.4) is also concircular, we must have the equations of the form

$$
\begin{equation*}
\rho_{\mu} \boldsymbol{M}_{\ddot{\boldsymbol{j}} \boldsymbol{k}}{ }^{\mu}=0 \tag{1.12}
\end{equation*}
$$

where

$$
\begin{equation*}
M_{\ddot{j} \ddot{z}^{\mu}}=H_{\dot{j} \ddot{j}^{\mu}}-\frac{1}{m} g^{a b} H_{\dot{a} \dot{b}}{ }^{\mu} g_{j k} . \tag{1.13}
\end{equation*}
$$

Conversely, if the equation (1.12) is satisfied, it is easily seen that the conformal transformation (1.4) is a concircular one.

Thus we have the following theorems:
Theorem I. The necessary and sufficient condition that a concircular transformation of the fundamental tensor of a Riemannian space induce a concircular transformation on a subspace is that the function $\rho$ satisfy the equations $\rho_{\mu} M_{\dot{j} i_{i}}{ }^{\mu}=0$ as well as (1.3).

Theorem II. The conformal transformation induced on a totally umbilical subspace by a concircular transformation is always a concircular one.
§2. We have seen, in a previous paper, ${ }^{1)}$ that the curvature tensor of $V_{n}$ defined by

$$
\begin{equation*}
Z_{\mu \nu \omega}^{\lambda}=R_{\mu \nu \omega}^{\lambda}-\frac{R}{n(n-1)}\left(g_{\mu \nu} \delta_{\omega}^{\lambda}-g_{\mu \omega} \delta_{\nu}^{\lambda}\right) \tag{2.1}
\end{equation*}
$$

is a concircular invariant. When the subspace $V_{m}$ is not a totally umbilical one, the curvature tensor of $V_{m}$

$$
\begin{equation*}
Z_{j k h}^{i}=R_{j k h}^{i}-\frac{g^{a b} R_{a b}}{m(m-1)}\left(g_{j k} \delta_{h}^{i}-g_{j h} \delta_{k}^{i}\right) \tag{2.2}
\end{equation*}
$$

where

$$
R_{j k h}^{i}=\frac{\partial\left\{\begin{array}{l}
i  \tag{2.3}\\
j k
\end{array}\right\}}{\partial u^{h}}-\frac{\partial\left\{\begin{array}{l}
i \\
j h
\end{array}\right\}}{\partial u^{k}}+\left\{\begin{array}{l}
a \\
j k k
\end{array}\right\}\left\{\begin{array}{c}
i \\
a h
\end{array}\right\}-\left\{\begin{array}{c}
a \\
j h
\end{array}\right\}\left\{\begin{array}{c}
i \\
a k
\end{array}\right\}
$$

is not in general a concircular invariant.
But the Weyl conformal curvature tensor

$$
\begin{align*}
C_{j k h}^{i}=R_{j k h}^{i} & -\frac{1}{m-2}\left(R_{j k} \delta_{h}^{i}-R_{j h} \delta_{k}^{i}+g_{j k} R_{\cdot h}^{i}-g_{j h} R_{\cdot k}^{i}\right)  \tag{2.4}\\
& +\frac{g^{a b} R_{a b}}{(m-1)(m-2)}\left(g_{j k} \delta_{h}^{i}-g_{j h} \delta_{k}^{i}\right)
\end{align*}
$$

is, of course, a concircular invariant. This conformal curvature tensor $C_{j k h}^{i}$ may be expressed by means of $Z_{j k h}^{i}$ and $Z_{j k}=Z_{j k i}^{i}$ as follows:

$$
\begin{equation*}
C_{j k h}^{i}=Z_{j k h}^{i}-\frac{1}{m-2}\left(Z_{j k} \delta_{h}^{i}-Z_{j h} \delta_{k}^{i}+g_{j k} Z_{\cdot h}^{i}-g_{j h} Z_{\cdot k}^{i}\right) \tag{2.5}
\end{equation*}
$$

where

$$
Z_{i_{h}}^{i}=g^{i k} Z_{k h}
$$

We shall, in the following, establish the relations between the concircular curvature tensor $Z_{\mu \nu \omega}^{\lambda}$ and the conformal curvature tensor $C_{j k h}^{i}$. The equations of Gauss of $V_{m}$ in $V_{n}$ are

$$
\begin{equation*}
R_{j k h}^{i}=B_{j j k h}^{i \mu \nu \omega} R_{\mu \nu \omega}^{\lambda}+H_{\dot{j} \dot{j}^{\lambda}}^{\lambda^{\prime}} H_{\cdot h \lambda}^{i}-H_{\dot{j} h}^{{ }^{\lambda}} H_{\cdot k \lambda}^{i} \tag{2.6}
\end{equation*}
$$

where

$$
B_{j i k h}^{i \mu \nu}=B_{{ }_{\lambda}}^{i} B_{j}^{\cdot \mu} B_{k}^{\cdot \nu} B_{h}^{\cdot \omega}, \quad B_{\cdot \lambda}^{i}=g^{i j} g_{\lambda \mu} B_{j}^{\cdot \mu} \quad \text { and } \quad H_{\cdot h \lambda}^{i}=g^{i k} g_{\lambda \mu} H_{\dot{k h}}{ }^{\mu} .
$$

Contracting (2.1) with $B_{i j k h, ~ w e ~ h a v e ~}^{i \mu j \omega}$

$$
B_{i j k h}^{i \mu \nu} Z_{\mu \nu \omega}^{\lambda}=B_{i j k h}^{i \mu \omega} R_{\mu \nu \omega}^{\lambda}-\frac{R}{n(n-1)}\left(g_{j k} \delta_{h}^{i}-g_{j h} \delta_{k}^{i}\right)
$$

Then substituting these equations in (2.6), we obtain

$$
\begin{equation*}
R_{j k h}^{i}=B_{j i k h}^{i \mu \nu} Z_{\mu \nu \omega}^{\lambda}+H_{\ddot{j k}}^{\lambda} H_{\cdot h \lambda}^{i}-H_{\dot{j h}}^{{ }^{\lambda}} H_{\cdot k \lambda}^{i}+\frac{R}{n(n-1)}\left(g_{j k} \delta_{h}^{i}-g_{j h} \delta_{k}^{i}\right) \tag{2.7}
\end{equation*}
$$

From (2.7), we find by contration

1) K. Yano: Concircular geometry I. loc. cit.

$$
\begin{equation*}
R_{j k}=B_{j \dot{j} k}^{\mu \nu} B_{\lambda}^{\omega} Z_{\mu \nu \omega}^{\lambda}+H_{\ddot{j} \ddot{m}^{\lambda}}^{H_{\cdot b \lambda}^{b}}-H_{\ddot{j} \dot{\imath}}^{\lambda} H_{\cdot k \lambda}^{b}+\frac{(m-1)}{n(n-1)} R g_{j k} \tag{2.8}
\end{equation*}
$$

and

$$
\begin{equation*}
g^{j k} R_{j k}=B^{\mu \nu} B_{\lambda}^{\omega} Z_{\mu \nu \omega}^{\lambda}+H_{\cdot a}^{a \cdot \lambda} H_{\cdot b \lambda}^{b}-H_{\cdot b}^{a \cdot \lambda} H_{\cdot a \lambda}^{b}+\frac{m(m-1)}{n(n-1)} R \tag{2.9}
\end{equation*}
$$

where

$$
B_{\lambda}^{\omega}=B_{\cdot \lambda}^{i} B_{i}^{\omega \omega}, \quad B_{j k}^{\mu \nu}=B_{j}^{\mu} B_{k}^{\cdot \nu}, \quad B^{\mu \nu}=B_{a b}^{\mu \nu} g^{a b}, \quad \text { and } \quad H_{: k}^{i \cdot \lambda}=g^{i j} H_{\dot{j} \ddot{l}_{k}^{\lambda}} .
$$

The equations (2.7), (2.8) and (2.9) give us

$$
\begin{align*}
Z_{j k h}^{i} & =B_{\lambda j j k h}^{i \mu \omega} Z_{\mu \nu \omega}^{\lambda}-\frac{B^{\mu \nu} B_{\lambda}^{\omega} Z_{\mu \nu \omega}^{\lambda}}{m(m-1)}\left(g_{j k} \delta_{\dot{\hbar}}^{i}-g_{j h} \delta_{k}^{i}\right)  \tag{2.10}\\
& +H_{\dot{j}{ }^{\lambda} H_{\cdot h \lambda}^{i}-H_{j \dot{\prime}}^{\lambda} H_{\cdot k \lambda}^{i}-\frac{H_{a}^{a \cdot \lambda} H_{\cdot b \lambda}^{b}}{m(m-1)}\left(g_{j k} \delta_{h}^{i}-g_{j h} \partial_{k}^{i}\right)} \\
& +\frac{H_{\cdot b}^{a \cdot \lambda} H_{\cdot a \lambda}^{b}}{m(m-1)}\left(g_{j k} \delta_{h}^{i}-g_{j h} \delta_{k}^{i}\right)
\end{align*}
$$

and

$$
\begin{align*}
Z_{j k} & =B_{j k}^{\mu \nu} B_{\lambda}^{\omega} Z_{\mu \nu \omega}^{\lambda}-\frac{1}{m} B^{\mu \nu} B_{\lambda}^{\omega} Z_{\mu \nu \omega}^{\lambda} g_{j k}+H_{\dot{j} k}^{\cdot \lambda} H_{\cdot b \lambda}^{b}  \tag{2.11}\\
& -H_{j b}^{\cdot} H_{\cdot k \lambda}^{b}-\frac{1}{m} H_{\cdot a}^{a \cdot \lambda} H_{\cdot{ }_{b \lambda}}^{b} g_{j k}+\frac{1}{m} H_{\cdot b}^{a \cdot \lambda} H_{\cdot \alpha}^{b} g_{j k}
\end{align*}
$$

Substituting the equations (2.10) and (2.11) in (2.4), we obtain

$$
\begin{align*}
& C_{j k h}^{i}=Z_{j k h}^{i}-\frac{1}{m-2}\left(Z_{j k} \delta_{h}^{i}-Z_{j h} \delta_{k}^{i}+g_{j k} Z^{i}{ }_{h}-g_{j h} Z^{i}{ }_{k}\right)  \tag{2.12}\\
& =B_{i j k h}^{i \mu j \omega} Z_{\mu \nu \omega}^{\lambda}-\frac{1}{m-2}\left(B_{j k}^{\mu \nu} B_{\lambda}^{\omega} Z_{\mu \nu \omega}^{\lambda} \delta_{h}^{i}-B_{j h}^{\mu \nu} B_{\lambda}^{\omega} Z_{\mu \nu \omega}^{\lambda} \delta_{k}^{i}\right. \\
& \left.+g_{j k} B_{a h}^{\mu \nu} B_{\lambda}^{\omega} g^{a i} Z_{\mu \nu \omega}^{\lambda}-g_{j h} B_{a k}^{\mu \nu} B_{\lambda}^{\omega} g^{\alpha i} Z_{\mu \nu \omega}^{\lambda}\right) \\
& +\frac{B^{\mu \nu} B_{\lambda}^{\omega} Z_{\mu \nu \omega}^{\lambda}}{(m-1)(m-2)}\left(g_{j k} \partial_{h}^{i}-g_{j h} \delta_{\dot{i}}^{i}\right)+H_{\dot{j} k}^{\wedge} H_{\cdot h \lambda}^{i}-H_{\dot{j} h}{ }^{\lambda} H_{\cdot k \lambda}^{i} \\
& -\frac{H_{a}^{a \cdot \lambda} H_{\cdot b \lambda}^{b}}{m(m-1)}\left(g_{j k} \delta_{h}^{i}-g_{j h} \delta_{k}^{i}\right)+\frac{H_{\cdot b}^{a \cdot \lambda} H_{a \lambda}^{b}}{m(m-1)}\left(g_{j k} \delta_{h}^{i}-g_{j h} \delta_{k}^{i}\right)
\end{align*}
$$

$$
\begin{aligned}
& +g_{j h} H_{b}^{i \cdot \lambda} H^{b}{ }_{k \lambda}-\frac{2 H_{a}^{a \cdot \lambda} H_{b \lambda}^{b}}{m}\left(g_{j k} \delta_{h}^{i}-g_{j h} \delta_{k}^{i}\right) \\
& \left.+\frac{2 H_{b}^{a \cdot \lambda} H_{a \lambda}^{b}}{m}\left(g_{j k} \delta_{h}^{i}-g_{j h} \delta_{k}^{i}\right)\right] .
\end{aligned}
$$

We have, on the other hand,

$$
\begin{aligned}
& H_{\dot{j} \ddot{k}^{\lambda}}=M_{\dot{j} k}^{\lambda}+\frac{1}{m} H_{\cdot a}^{a \cdot \lambda} g_{j k}, \\
& H_{\cdot h \lambda}^{i}=M_{\cdot h \lambda}^{i}+\frac{1}{m} H_{\cdot a \lambda}^{a} \delta_{h k}^{i}, \quad\left(M_{\cdot h \lambda}^{i}=g^{i j} g_{\lambda \mu \mu} M_{\dot{j} \ddot{h}^{\mu}}\right) .
\end{aligned}
$$

Substituting these equations in (2.12), we obtain finally

$$
\begin{align*}
& Z_{j k h}^{i}-\frac{1}{m-2}\left(Z_{j k} \delta_{h}^{i}-Z_{j h} \delta_{k}^{i}+g_{j k} Z_{\cdot h}^{i}-g_{j h} Z^{i}{ }_{\cdot k}\right)  \tag{2.13}\\
& =B_{\lambda j k h}^{i \mu j \omega} Z_{\mu \nu \omega}^{\lambda}-\frac{1}{m-2}\left(B_{j k}^{\mu \nu} B_{\lambda}^{\omega} Z_{\mu \nu \omega}^{\lambda} \delta_{h}^{i}-B_{j h}^{\mu \nu} B_{\lambda}^{\omega} Z_{\mu \nu \omega}^{\lambda} \delta_{k c}^{i}\right. \\
& \left.+g_{j k} B_{a h}^{\mu \nu} B_{\lambda}^{\omega} g^{a i} Z_{\mu \nu \omega}^{\lambda}-g_{j h} B_{a k}^{\mu \nu} B_{\lambda}^{\omega} g^{a i} Z_{\mu \nu \omega}^{\lambda}\right) \\
& +\frac{B^{\mu \nu} B_{\lambda}^{\omega} Z_{\mu \nu \omega}^{\lambda}}{(m-1)(m-2)}\left(g_{j k} \delta_{h}^{i}-g_{j h} \delta_{\dot{k}}^{i}\right)+M_{\ddot{j} \ddot{j}^{\lambda}} M_{\cdot h \lambda}^{i}-M_{\ddot{j} \dot{H}^{\lambda}} M_{\cdot k \lambda}^{i} \\
& +\frac{1}{m-2}\left[\left(g_{j k} M_{\cdot a}^{i \cdot \lambda} M_{\cdot h \lambda}^{a}+M_{\ddot{j a}}{ }^{\lambda} M_{\cdot k \lambda}^{a} \delta_{h}^{i}\right)\right. \\
& -\left(g_{j h} M_{\cdot a}^{i \cdot \lambda} M_{\cdot k \lambda}^{a}+M_{\dot{j a}}^{\left.\left.\cdot \ddot{-}^{\lambda} M_{\cdot h \lambda}^{a} \partial_{\bar{k}}^{i}\right)\right]}\right. \\
& -\frac{M_{: b}^{a \cdot \lambda} M_{\cdot a \lambda}^{b}}{(m-1)(m-2)}\left(g_{j k} \partial_{h}^{i}-g_{j h} \delta_{k}^{i}\right) .
\end{align*}
$$

The left member of (2.13) representing the Weyl conformal curvature tensor, the equations (2.13) are the equations of Gauss of $V_{m}$ in $V_{n}$ in our concircular geometry.
§ 3. Let $B_{A}^{\cdot \lambda}(A, B, C, \ldots=\dot{m}+\dot{1}, \dot{m}+\dot{2}, \ldots, \dot{n})$ be $n-m$ mutually orthogonal unit vectors normal to the subspace $V_{m}$, then the equations of Weingarten may be written in the form

$$
\begin{equation*}
B_{\dot{A} ; j}^{\lambda}=-B_{a}^{\cdot \lambda} H_{\cdot j A}^{a}+L_{A B j} B_{\dot{B}}^{\lambda} \tag{3.1}
\end{equation*}
$$

where we have put

$$
H_{\cdot j A}^{a}=g_{\lambda \mu} H_{\cdot j}^{a \cdot \lambda} B_{\dot{A}}^{\mu} \quad \text { and } \quad L_{A B j}=g_{\lambda \mu} B_{\dot{A}^{\prime} ; j} B_{\dot{B}}{ }^{\mu}
$$

From (3.1) we can derive the equations of Codazzi

$$
\begin{equation*}
B_{\lambda A j k}^{i \mu \omega} R_{\mu \nu \omega}^{\lambda}=-H_{{ }_{j A} ; k}^{i}+H_{{ }^{i}{ }_{k A ; j}}^{i}+H_{\cdot j B}^{i} L_{A B k}-H_{{ }^{i}{ }_{k B}}^{i_{A B j}} L_{A B j}, \tag{3.2}
\end{equation*}
$$

where

$$
B_{\lambda A j k}^{i \mu \nu \omega}=B_{\cdot}^{i}{ }_{\lambda} B_{A}{ }^{\mu} B_{j}^{\cdot \nu} B_{k}^{*}{ }^{\omega} .
$$

Multiplying (2.1) by $B_{\alpha A j k e}^{i \mu \nu \omega}$ and contracting with respect to the indices $\lambda, \mu, \nu$ and $\omega$, we find $B_{\lambda A j k}^{i \mu \nu \omega} Z_{\mu \nu \omega}^{\lambda}=B_{\lambda A j k}^{i \mu \omega} R_{\mu \nu \omega}^{\lambda}$, consequently

$$
\begin{equation*}
B_{A A j k}^{i \mu \nu} Z_{\mu \nu \omega}^{\lambda}=-H_{{ }_{j A} ; k}^{i}+H_{{ }_{k A ; j}}^{i}+H_{\cdot j B}^{i} L_{A B k}-H_{\cdot k B}^{i} L_{A B j} . \tag{3.3}
\end{equation*}
$$

These are the equations of Codazzi in our concircular geometry.
From (3.3) we can conclude that the tensor whose compoments are

$$
\begin{equation*}
-H_{\cdot j A ; k}^{i}+H_{{ }_{k A} ; j}^{i}+H_{{ }_{j B}}^{i} L_{A B k}-H_{{ }_{k B}}^{i} L_{A B j} \tag{3.4}
\end{equation*}
$$

is a semi-concircular tensor, that is to say, it will be multiplied by a power of $\rho$ by the concircular transformation.

If we consider a hypersurface $V_{n-1}$ in $V_{n}$ and denote by $B^{\lambda}$ the unit vector normal to $V_{n-1}$, we have

$$
\begin{equation*}
H_{i j}^{-\lambda}=H_{i j} B^{\lambda}, \quad L_{A B j}=0 \tag{3.5}
\end{equation*}
$$

and the equations (3.3) reduce to

$$
B_{\lambda}^{i} B^{\mu} B_{j k}^{\nu \omega} Z_{\mu \nu \omega}^{\lambda}=-H_{\cdot j ; k}^{i}+H_{\cdot k ; j}^{i}
$$

or

$$
\begin{equation*}
B_{\cdot \lambda}^{i} B^{\mu} B_{j k}^{\nu \omega} Z_{\mu \nu \omega}^{\lambda}=-M_{\cdot j ; k}^{i}+M_{\cdot k ; j}^{i}-\frac{1}{n-1} H_{\cdot a ; k}^{a} \delta_{j}^{i}+\frac{1}{n-1} H_{\cdot a ; j}^{a} \delta_{k}^{i} \tag{3.6}
\end{equation*}
$$

where

$$
H_{\cdot k}^{i}=g_{i j}^{i j} H_{j k}, \quad M_{i \cdot}^{\lambda}=M_{i j} B^{\lambda} \quad \text { and } \quad M_{\cdot k}^{i}=g^{i j} M_{j k}
$$

§ 4. In this paragraph, we state some of theorems which may be easily deduced from the formulae proved in three preceding paragraphs. They are all well known theorems, but it may not be of no use to emphasize here that they are theorems which may be considered in the concircular geometry.

Theorem III. A totally umbilical subspace in a concircularly flat space is also concircularly flat.

Proof. For a totally umbilical subspace, we have $H_{\dot{j} k}^{\lambda}=\frac{1}{m} H_{a}^{a \cdot \lambda} g_{j k}$. In such a case equations (2.10) become

$$
\begin{equation*}
Z_{j k h}^{i}=B_{j j k h}^{i \mu \nu \omega} Z_{\mu \nu \omega}^{\lambda}-\frac{B^{\mu \nu} B_{\lambda}^{\omega} Z_{\mu \nu \omega}^{\lambda}}{m(m-1)}\left(g_{j k} \delta_{h}^{i}-g_{j h} \delta_{k}^{i}\right) \tag{4.1}
\end{equation*}
$$

Thus we can see that if the enveloping space $V_{n}$ is a concircularly flat one ( $Z_{\mu \nu \omega}^{\lambda}=0$ ), the subspace is also concircularly flat one ( $Z_{j k h}^{i}=0$ ).

Theorem IV. The mean curvature of totally umbilical hypersurface in a concircularly flat space is constant.

Proof. The conditions $Z_{\mu \nu \omega}^{\lambda}=0, M_{i j}=0$ and equations (3.6) give us

$$
H_{a ; k}^{a} \delta_{j}^{i}-H_{a ; j}^{a} \delta_{k}^{i}=0
$$

from which we have

$$
\begin{equation*}
H_{a ; k}^{a}=0 \tag{4.2}
\end{equation*}
$$

Thus the theorem is proved.
Theorem V. If there exists always a totally umbilical hypersurface of constant mean curvature touching an arbitrary hyperplane passing through any point of the enveloping space, then the enveloping space is concircularly flat.

Proof. $M_{{ }_{j}}^{i}$ and $H_{a ; j}^{a}$ being zero, we have from (3.6)
(4.3)

$$
B_{\lambda}^{i} B^{\mu} B_{j k}^{\nu \omega} Z_{\mu \nu \omega}^{\lambda}=0,
$$

which must be satisfied for any $B_{i}^{\lambda}$ and $B^{\lambda}$ satisfying

$$
g_{\lambda \mu} B_{i}^{\prime \lambda} B^{\mu}=0
$$

Consequently we have from (4.3)

$$
Z_{\mu \nu \omega}^{\lambda}=0 .
$$

This proves the theorem.


[^0]:    1) K. Yano, Concircular Geometry I, Proc. 16 (1940), 195-200, II, Proc. 16 (1940), 359-360 and III, 16 (1940), 442-458.
