# Ergodic decomposition of probability measures on the configuration space 

Dedicated to Professor Takeshi Hirai on his 60th birthday

## By

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## Introduction

Let $X$ be a locally compact space which satisfies the second countable axiom. Any locally finite subset of $X$ is called a configuration in $X$, that is a subset $\gamma \subset X$ such that $\gamma \cap K$ is finite for any compact set $K \subset X$. Let us denote by $\Delta_{X}$ the space of all infinite and by $B_{X}$ the space of all finite configurations in $X$, and set $\Gamma_{X}:=\Delta_{X} \cup B_{X}$. We introduce a measurable structure $\mathscr{C}$ on $\Gamma_{X}$ such that $\mathscr{C}$ is a minimal $\sigma$-algebra with which all the functions, $\gamma \in \Gamma_{X} \rightarrow|\gamma \cap B| \in \mathbf{R}$ are measurable, where $B$ runs through all the Borel sets in $X$ and $|\gamma \cap B|$ is the number of the set $\gamma \cap B$. It is known that $\left(\Gamma_{X}, \mathscr{C}\right)$ is a standard space (See, theorem 1.2 in [3]) and hence any probability measure $\mu$ on ( $\left.\Gamma_{X}, \mathscr{C}\right)$ is decomposed into conditional probability measures with resect to any sub- $\sigma$-field of $\mathscr{C}$. The subject of this paper are two kinds of measures on ( $\Gamma_{X}, \mathscr{C}$ ) with well known properties and their ergodic decompositions. The first one is a $\mathrm{Diff}_{0} X$-quasiinvariant probability measure $\mu$, where $X$ is a connected para-compact $C^{\infty}$-manifold and $\operatorname{Diff}_{0} X:=\{\psi \mid \psi:$ diffeomorphism on $X$ with compact support $\}$. In 1975, Vershick-Gel'fand-Graev introduced elementary representations $U_{\mu}$ generated by these $\mu$ 's and discussed fully their interesting properties in [5]. In particular they showed that $U_{\mu}$ is irreducible if and only if $\mu$ is ergodic. Thus our subject correspondes to an irreducible decomposition of $U_{\mu}$. It will be shown in section 1 that an ergodic decomposition of $\mathrm{Diff}_{0} X$-quasi-invariant probability measure is actually possible.
The second one is a consideration of Gibbs measures $\mu$ having been discussed in great detail in statistical mechanics. An ergodic decomposition of such $\mu$ relative to the tail- $\sigma$-field leads us to a remarkable fact that there exist typical extremal measures which are regarded as a base on a convex set formed by such $\mu$ 's. These contents will be discussed in section 2. In both of section 1 and section 2, we denote a $\sigma$-finite non atomic Borel measure on $X$ by $m$. The direct product $m^{n}$ of $n$ copies of $m$ is naturally regarded as a measure on $\tilde{X}^{n}:=\left\{\left(x_{1}, \cdots, x_{n}\right) \in X^{n} \mid x_{i} \neq x_{j}\right.$ for all $\left.i \neq j\right\}$ and thus an image measure $p_{n} m^{n}$ is obtained by the natural map $p_{n}:\left(x_{1}, \cdots, x_{n}\right) \in \tilde{X}^{n} \rightarrow\left\{x_{1}, \cdots, x_{n}\right\} \in B_{x}^{n}:=\left\{\gamma \in \Gamma_{X} \| \gamma \mid=n\right\}$. We denote it by $m_{X, n}$.

## 1. Ergodic decomposition of Diff $_{0} \mathrm{X}$-quasi-invariant measures

1.1. Basic notion and result. As before let $X$ be a d-dimensional $C^{\infty}$-manifold and $m$ be a locally Euclidean Borel measure on $X$ with smooth densities. And we associate with each $\psi \in \operatorname{Diff}_{0} X$ a transformation $T_{\psi}$ on $\Gamma_{X}$ such that $T_{\psi}\left\{x_{1}, \cdots, x_{n}, \cdots\right\}=\left\{\psi\left(x_{1}\right), \cdots, \psi\left(x_{n}\right), \cdots\right\}$. A probability measure $\mu$ on $\left(\Gamma_{X}, \mathscr{C}\right)$ is said to be Diff $_{0} X$-quasi-invariant, if and only if $T_{\psi} \mu \simeq \mu$ for all $\psi \in \operatorname{Diff}_{0} X$, where the symbol $\simeq$ means the equivalence relation of measures. Moreover $\mu$ is said to be $\operatorname{Diff}_{0} X$-ergodic, if $\mu(A)=1$ or 0 provided that $\mu\left(T_{\psi}(A) \ominus A\right)=0$ for all $\psi \in \operatorname{Diff}_{0} X$. It is an aim of the present section that after suitably setting a measure space $(\Lambda, \lambda)$ we decompose $\mu$ such as $\mu=\int_{\Lambda} \mu_{l} \lambda(\mathrm{~d} l)$ with Diff $_{0} X$-ergodic measures $\mu_{l}$. Besides it is to be desired that $\mu_{l}$ 's are mutually singular.
Now let $\mu$ be a $\operatorname{Diff}_{0} X$-quasi-invariant probability measure and put $\alpha:=\mu\left(B_{X}\right)$ and $\beta:=\mu\left(\Delta_{X}\right)$. Then we have $\mu=\alpha \mu_{1}+\beta \mu_{2}$, where $\mu_{1}(E):=\mu\left(E \cap B_{X}\right) / \alpha$ and $\mu_{2}(E):=\mu\left(E \cap \Delta_{X}\right) / \beta$ for all $E \in \mathscr{C}$. Furthermore we put $\alpha_{n}:=\mu_{1}\left(B_{X}^{n}\right)$. Then $\mu_{1}$ is decomposed as

$$
\begin{equation*}
\mu_{1}=\sum_{n=0}^{\infty} \alpha_{n} \mu_{1, n}, \tag{1.1}
\end{equation*}
$$

where $\mu_{1, n}(E)=\mu_{1}\left(E \cap B_{X}^{n}\right) / \alpha_{n}$ for all $\mathrm{E} \in \mathscr{C}$. Since $B_{X}^{n}(n=0,1, \cdots)$ is a $\operatorname{Diff}_{0} X$ invariant set, so $\mu_{1, n}$ is a $\operatorname{Diff}_{0} X$-quasi-invariant measure. Here we give the following theorem.

Theorem 1.1 ([5]). Any non zero $\sigma$-finite Diff $_{0} X$-quasi-invariant measure on $B_{X}^{n}$ is equivalent to $m_{X, n}$.

The proof will be seen in a discussion for the proof of Lemma 1.2 which will be stated later on.
In anyway, it follows immediately from the above theorem that any non zero $\sigma$-finite Diff $_{0} X$-quasi-invariant measure on $B_{X}^{n}$ is Diff $_{0} X$-ergodic. Hence $\mu_{1, n}$ is ergodic and (1.1) is actually an ergodic decomposition of $\mu_{1}$. Next let us observe $\mu_{2}$, so we shall assume that $\mu\left(\Delta_{X}\right)=1$ from now on. Here we introduce a set $\tilde{X}^{\infty}:=\left\{\left(x_{1}, \cdots, x_{n}, \cdots\right) \in X^{\infty} \mid x_{i} \neq x_{j}\right.$ for all $i \neq j$ and the set $\left\{x_{1}, \cdots, x_{n}, \cdots\right\}$ has no accumulation points $\}$ and consider a cross section $s$ of the natural map

$$
p:\left(x_{1}, \cdots, x_{n}, \cdots\right) \in \tilde{X}^{\infty} \rightarrow\left\{x_{1}, \cdots, x_{n}, \cdots\right\} \in \Delta_{X} .
$$

Let us take and fix an increasing sequence $\left\{Y_{n}\right\}$ of connected open sets with compact closure such that $\bar{Y}_{n} \subset Y_{n+1}$ and $Y_{n} \uparrow X$. Then there exists a measurable section $s$ possessing the following property ( P ) with this $\left\{Y_{n}\right\}$.
(P) If we have $\left|\gamma \cap Y_{1}\right|=k_{1},\left|\gamma \cap\left(Y_{2} \backslash Y_{1}\right)\right|=k_{2}, \cdots,\left|\gamma \cap\left(Y_{n} \backslash Y_{n-1}\right)\right|=k_{n}, \cdots$ for $\gamma \in \Delta_{X}$, then the first $k_{1}$ elements of $s(\gamma)$ are in $\gamma \cap Y_{1}$, the next $k_{2}$ element of $s(\gamma)$ are in $\gamma \cap\left(Y_{2} \backslash Y_{1}\right)$ and so on.

We call it to be admissible. Notice that $s(E)$ is a Borel set in $\tilde{X}^{\infty}$ for any $E \in \mathscr{C} \cap \Delta_{X}$, because $s$ is one to one and measurable, and the space $\left(\Gamma_{X}, \mathscr{C}\right)$ is standard.
For measures on the natural measurable space ( $\tilde{X}^{\infty}, \mathfrak{B}\left(\tilde{X}^{\infty}\right)$ ) we also obtain the notion of $\operatorname{Diff}_{0} X$-quasi-invariance and ergodicity with maps $\tilde{T}_{\psi}:\left(x_{1}, \cdots x_{n}, \cdots\right) \in \tilde{X}^{\infty}$ $\rightarrow\left(\psi\left(x_{1}\right), \cdots, \psi\left(x_{n}\right), \cdots\right) \in \tilde{X}^{\infty}$ for $\psi \in \operatorname{Diff}_{0} X$. Here we shall define a new measure $\tilde{\mu}$ on $\mathfrak{B}\left(\tilde{X}^{\infty}\right)$ from a given probability measure $\mu$ taking the above measurable admissible cross section $s$ :

$$
\begin{equation*}
\tilde{\mu}(E):=\sum_{\sigma \in \Theta_{\infty}} c(\sigma)(s \mu) \sigma(E) \tag{1.2}
\end{equation*}
$$

for all $E \in \mathfrak{B}\left(\tilde{X}^{\infty}\right)$, where $\Theta_{\infty}$ is the set of all finite permutations on $\mathbf{N},\{c(\sigma)\}_{\sigma \in \mathcal{S}_{\infty}}$ is a fixed positive sequence such that $\sum_{\sigma \in \Theta_{\infty}} c(\sigma)=1$ and $(s \mu) \sigma$ is a image measure of $\mu$ by the map,

$$
\gamma \in \Delta_{X} \xrightarrow{s} s(\gamma)=\left(x_{1}, \cdots, x_{n}, \cdots\right) \xrightarrow{\sigma} s(\gamma) \sigma:=\left(x_{\sigma(1)}, \cdots, x_{\sigma(n)}, \cdots\right) \in \tilde{X}^{\infty} .
$$

Theorem 1.2 (section 2 in [5]). Under the above notations,
(a) $\mu$ is $\operatorname{Diff}_{0} X$-quasi-invariant if and only if so is $\tilde{\mu}$.
(b) $\mu$ is $\mathrm{Diff}_{0} X$-ergodic if and only if so is $\tilde{\mu}$.
(c) If a Borel probability measure $\mu_{1}$ on $\tilde{X}^{\infty}$ is $\widehat{\varsigma}_{\infty}$-quasi-invariant and $\mu_{1}\left(\bigcup_{\sigma \in \mathfrak{E}_{\infty}} s\left(\Gamma_{X}\right) \sigma\right)=1$, then $p \mu_{1}$ is equivalent to $\mu_{1}$.
(d) $A$ Diff $_{0} X$-quasi-invariant probability measure $\tilde{\mu}$ on $\left(\tilde{X}^{\infty}, \mathfrak{B}\left(\tilde{X}^{\infty}\right)\right)$ is $\operatorname{Diff}_{0} X$ ergodic if and only if $\tilde{\mu}(A)=1$ or 0 for any $A \in \mathfrak{B}_{\infty}$, where $\mathfrak{B}_{\infty}=\bigcap_{n=1}^{\infty} q_{n}^{-1}\left(\mathcal{B}\left(\tilde{X}^{\infty}\right)\right)$ and $q_{n}:\left(x_{1}, \cdots, x_{n}, \cdots\right) \in \tilde{X}^{\infty} \rightarrow\left(x_{n+1}, \cdots, x_{m}, \cdots\right) \in \tilde{X}^{\infty} . \mathfrak{B}_{\infty}$ is called the tail- $\sigma$-field.
1.2. Diff $\boldsymbol{Y}$-quasi-invariant measure on $\boldsymbol{B}_{Y}$ and one parameter group of $\mathrm{Diff}_{0} X$. In this paragraph the letter $Y$ stands for connected open subset in $X$ with compact closure, and we observe subgroups of diffeomorphisms on $X$ whose support is contained in $Y$ which will be denoted by $\operatorname{Diff}_{0} Y$. Using the sequence $\left\{Y_{n}\right\}$ already stated in the admissible cross section, we have,

$$
\begin{equation*}
\operatorname{Diff}_{0} X=\bigcup_{n=1}^{\infty} \operatorname{Diff}_{0} Y_{n} . \tag{1.3}
\end{equation*}
$$

Now from a trivial equality, $\gamma=(\gamma \cap Y) \cup\left(\gamma \cap Y^{c}\right)$ we can identify $\Gamma_{X}$ with a product space $B_{Y}$ and $\Gamma_{Y c}$. Put $\pi_{Y}: \gamma \in \Gamma_{X} \rightarrow \gamma \cap Y \in B_{Y}$ and $\pi_{Y c}: \gamma \in \Gamma_{X} \rightarrow \gamma \cap Y^{c} \in \Gamma_{Y c}$. Since $B_{Y}$ and $\Gamma_{Y c}$ are naturally regarded as subspaces of $\Gamma_{X}$ so the measurable structure $\mathscr{C}_{Y}$ and $\mathscr{C}_{Y c}$ are induced from $\mathscr{C}$ respectively. It is easy to see that the above identification $\Gamma_{X} \simeq B_{Y} \times \Gamma_{Y}$ is an isomorphism with the measurable structure $\mathscr{C}$ and $\mathscr{C}_{Y} \times \mathscr{C}_{Y}$. By the way probability measures $v$ on $\left(B_{Y}, \mathscr{C}_{Y}\right)$ naturally arises, if we decompose Diff $_{0} Y$-quasi-invariant probability measures with respect to sub- $\sigma$-field $\pi_{Y_{c}}^{-1}\left(\mathscr{C}_{Y c}\right)$. So we shall observe such $v^{\prime}$, especially with $v\left(B_{Y}^{n}\right)=1$ for a while. As before we define

$$
\begin{equation*}
\tilde{v}(E)=\sum_{\sigma \in \Theta_{n}}\left(s_{n} v\right) \sigma(E) \tag{1.4}
\end{equation*}
$$

for all $E \in \mathfrak{B}\left(\tilde{Y}^{n}\right)$, taking a measurable cross section $s_{n}$ of the natural map
$p_{n}:\left(y_{1}, \cdots, y_{n}\right) \in \tilde{Y}^{n} \rightarrow\left\{y_{1}, \cdots, y_{n}\right\} \in B_{Y}^{n}$. Then $v$ and $\tilde{v}$ have the same kind of quasi-invariance. Now $\tilde{Y}^{n}$ is covered by countable sets of the form $O_{1}^{m} \times \cdots \times O_{n}^{m}$ ( $m=1, \cdots$ ), where $O_{i}^{m}$ is an open set with compact closure which is diffeomorphic to $\mathbf{R}^{d}$ by a map $\psi_{i}^{m}$ and $O_{i}^{m} \cap O_{j}^{m}=\phi$ for $i \neq j$. Since Diff ${ }_{0} Y$ acts on $\tilde{Y}^{n}$ transitively, it follows that there exists an at most countable set $\left\{\varphi_{k}\right\}$ such that $\tilde{Y}^{n}=\bigcup_{k=1}^{\infty} \varphi_{k}\left(O_{1}^{m} \times \cdots \times O_{n}^{m}\right)$ for each $m$. Hence $\tilde{v}\left(O_{1}^{m} \times \cdots \times O_{n}^{m}\right)>0$, if $\tilde{v}$ is quasi-invariant under a group $H_{n}$ generated by such $\varphi_{k}$ 's. Here we shall prepare some basic lemma.

Lemma 1.1. There exists a one parameter group $\pi_{i}^{l}(i=1, \cdots, d, l=1, \cdots)$ of Diff $\mathbf{R}^{d}$ which satisfies $\pi_{i}^{l}(t)(\xi)=\left(\xi_{1}, \cdots, \xi_{i}+t, \cdots, \xi_{d}\right)$ for all $\xi=\left(\xi_{1}, \cdots, \xi_{d}\right) \in \mathbf{R}^{d}$ and $t \in \mathbf{R}$ such that $\operatorname{Max}_{1 \leq t \leq d}\left|\xi_{i}\right|<l$ and $|t|<l$.

Proof. For the existence of such one parameter group, we solve the following differential equation (1.5) with a function $f_{l}$ of $C^{\infty}$-class on $\mathbf{R}$ such that $f_{l}(s)=1$ on $|s| \leq 2 l$ and $f_{l}(s)=0$ on $|s| \geq 3 l$.

$$
\left\{\begin{array}{l}
\frac{d \mathbf{x}}{d t}=f_{l}\left(x_{1}\right) \cdots f_{l}\left(x_{d}\right) \mathbf{e}_{i}  \tag{1.5}\\
\mathbf{x}(0)=\xi
\end{array}\right.
$$

where $\mathrm{e}_{i}=(0, \cdots, \stackrel{i}{1}, \cdots, 0)$. Then the solution $\mathbf{x}(t, \xi)$ of (1.5) gives directly a desired diffeomorphism.

Let $G_{d}^{0}$ be a group generated by the one parameter groups $\pi_{i}^{l}(i=1, \cdots, d$, $l=1, \cdots)$. Then it is easily seen that

Proposition 1.1. For any $l \in \mathbf{N}$ and for any $\tau=\left(t_{1}, \cdots, t_{d}\right) \in \mathbf{R}^{d}$, there exists $\psi \in G_{d}^{0}$ such that

$$
\psi(\xi)=\xi+\tau \text { for all } \xi \text { with } \operatorname{Max}_{1 \leq i \leq d}\left|\xi_{i}\right|<l .
$$

Proposition 1.2. Any $\sigma$-finite $G_{d}^{0}$-quasi-invariant Borel measure on $\mathbf{R}^{d}$ is equivalent to the Lebesgue measure.

Now let us pull back each element of $G_{d}^{0}$ by the maps $\psi_{i}^{m}(i=1, \cdots, d, m=1, \cdots$ ) and extend it to the element of $\operatorname{Diff}_{0} Y$. Then considering the restriction of $\tilde{v}$ to $O_{1}^{m} \times \cdots \times O_{n}^{m}$, we deduce that

Lemma 1.2. In Diff $_{0} Y$ there exist one parameter groups $\pi_{i, n}(i=1, \cdots)$ and $a$ countable subgroup $H_{n}$ such that the following are equivalent for any Borel probability measure $v$ on $B_{Y}^{n}$.
(a) $v$ is quasi-invariant under the groups $\pi_{i, n}(i=1, \cdots)$ and $H_{n}$.
(b) $v$ is equivalent to $m_{X, n}$.
(c) $v$ is Diff $_{0} Y$-quasi-invariant.

The following lemma is an immediate consequence of the above lemma by letting $n$ run from 0 to $\infty$.

Lemma 1.3. In Diff $_{0} Y$ there exists one parameter subgroups $\pi_{i, Y}(i=1, \cdots)$ and a countable group $H_{Y}$ such that the following are equivalent for any Borel probability measure $v$ on $B_{Y}$.
(a) $v$ is quasi-invariant under the groups $\pi_{i, Y}(i=1, \cdots)$ and $H_{Y}$.
(b) $v$ is Diff $_{0} Y$-quasi-invariant.

Next let us decompose a probability measure $\mu$ on $\left(\Gamma_{X}, \mathscr{C}\right)$ into the regular conditional probability measures $\left\{\mu^{\nu}\right\}_{y \in \Gamma_{Y c}}$ on $\left(\mathfrak{B}_{Y}, \mathscr{C}_{Y}\right)$ with respect to the map $\pi_{Y^{c}}$ which satisfy
(1.6) $\mu^{\nu}(A)$ is a $\mathscr{C}_{Y c}$-measurable function of $\gamma \in \Gamma_{Y^{c}}$ for each fixed $A \in B_{Y}$, and

$$
\begin{equation*}
\mu(A \times B)=\int_{B} \mu^{\nu}(A) \pi_{Y c} \mu(d \gamma) \tag{1.7}
\end{equation*}
$$

for all $A \in \mathscr{C}_{Y}$ and $B \in \mathscr{C}_{Y c}$.
Lemma 1.4. Under the above notations, the following are equivalent.
(a) $\mu$ is Diff $_{0} Y$-quasi-invariant.
(b) $\mu^{\gamma}$ is Diff $_{0} Y$-quasi-invariant for $\pi_{Y c} \mu$-a.e. $\gamma$.

Proof. There is nothing to prove "(b) implies (a)". Let us see the converse relation. For this we calculate $T_{\psi} \mu(A \times B), \psi \in \operatorname{Diff}_{0} Y$ in two ways. The first one is,

$$
\begin{equation*}
T_{\psi} \mu(A \times B)=\mu\left(T_{\psi}^{-1}(A) \times B\right)=\int_{B} T_{\psi} \mu^{\nu}(A) \pi_{\mathrm{Y}} \mu(d \gamma) . \tag{1.8}
\end{equation*}
$$

And the other one is,

$$
\begin{equation*}
T_{\psi} \mu(A \times B)=\int_{B} \int_{A} \frac{d T_{\psi} \mu}{d \mu}\left(\gamma^{\prime}\right) \mu^{\gamma}\left(d \gamma^{\prime}\right) \pi_{Y с} \mu(d \gamma) . \tag{1.9}
\end{equation*}
$$

It follows from (1.8) and (1.9) that

$$
\begin{equation*}
T_{\psi} \mu^{\nu}(\cdot)=\int_{(\cdot)} \frac{d T_{\psi} \mu}{d \mu}\left(\gamma^{\prime}\right) \mu^{\nu}\left(d \gamma^{\prime}\right) \tag{1.10}
\end{equation*}
$$

for $\pi_{Y c} \mu$-a.e. $\gamma$, and thus we have

$$
\begin{equation*}
T_{\psi} \mu^{\gamma} \simeq \mu^{\nu} \tag{1.11}
\end{equation*}
$$

for $\pi_{\mathrm{Yc}} \mu$-a.e. $\gamma$. Here we take an arbitrary one parameter group $\left\{\psi_{t}\right\}_{t \in \mathbf{R}}$ of $\operatorname{Diff}_{0} Y$ and set

$$
\begin{aligned}
& \Pi:=\left\{(t, \gamma) \in \mathbf{R} \times \Gamma_{\mathbf{Y}^{c}} \mid T_{\psi_{t}} \mu^{\nu} \simeq \mu^{\nu}\right\} \text { and } \\
& \Pi_{0}:=\left\{(t, \gamma) \in \mathbf{R} \times \Gamma_{Y_{c}} \left\lvert\, T_{\psi_{1}} \mu^{\nu}(\cdot)=\int_{(\cdot)} \frac{d T_{\psi_{1}} \mu}{d \mu}\left(\gamma^{\prime}\right) \mu^{\gamma}\left(d \gamma^{\prime}\right)\right.\right\} .
\end{aligned}
$$

Then $\Pi_{0}$ is a $\mathfrak{B}(\mathbf{R}) \times \mathscr{C}_{Y}$-measurable subset of $\Pi$ and for any fixed $t$ the $\mathbf{R}$-section $\Pi_{0}^{t}$ determined by $t$ has full measure for $\pi_{Y^{c}} \mu$. Thus by virtue of Fubini's theorem the $\Gamma_{Y^{c}}$-section $\Pi_{0}^{\gamma}$ determined by $\gamma$ has full Lebesgue measure for $\pi_{Y c} \mu$-a.e. $\gamma$. So the $\Gamma_{\mathrm{Yc}}$-section $\Pi^{\nu}$ is Lebesgue measurable and it is a subgroup of $\mathbf{R}$ with positive measure for $\pi_{Y c} \mu$-a.e. $\gamma$. This implies that $\Pi^{\gamma}=\mathbf{R}$ for $\pi_{Y c} \mu$-a.e. $\gamma$. Now consider groups $\pi_{i, Y}(i=1, \cdots)$ and $H_{Y}$ stated in Lemma 1.3. Applying the above arguments to these subgroups, we conclude that $\mu^{\nu}$ is $\operatorname{Diff}_{0} Y$-quasi-invariant for $\pi_{Y c} \mu$-a.e. $\gamma$.

From (1.3), Lemma 1.3 and Lemma 1.4 we have the following theorem.
Theorem 1.3. In Diff $_{0} X$, there exist one parameter groups $\pi_{i}(i=1, \cdots)$ which are subgroups of $\operatorname{Diff}_{0}\left(Y_{k_{i}}\right)$ and a countable group $G_{0}$ such that the following are equivalent for any probability measure $\mu$ on $\left(\Gamma_{X}, \mathscr{C}\right)$.
(a) $\mu$ is Diff $_{0} X$-quasi-invariant.
(b) $\mu$ is quasi-invariant under the groups $\pi_{i}(i=1, \cdots)$ and $G_{0}$.
1.3. Ergodic decomposition of Diff $_{0} X$-quasi-invariant measure. Let $\mu$ be a Diff $_{0} X$-quasi-invariant probability measure on $\left(\Gamma_{X}, \mathscr{C}\right)$ with $\mu\left(\Delta_{X}\right)=1$ and $\tilde{\mu}$ be the Borel measure on $\tilde{X}^{\infty}$ defined by (1.2). We decompose $\tilde{\mu}$ into conditional probability measures $\left\{\tilde{\mu}^{x}\right\}_{x \in \tilde{X}^{\infty}}$ with respect to the tail- $\sigma$-field $\mathfrak{B}_{\infty}$. Namely,

$$
\begin{align*}
& \tilde{\mu}^{x}(B) \text { is a } \mathfrak{B}_{\infty} \text {-measurable function of } x \in \tilde{X}^{\infty} \text { for each fixed } B \in \mathfrak{B}\left(\tilde{X}^{\infty}\right) \text {, and }  \tag{1.12}\\
& \tilde{\mu}(A \cap B)=\int_{A} \tilde{\mu}^{x}(B) \tilde{\mu}(d x) \tag{1.13}
\end{align*}
$$

for all $A \in \mathfrak{B}_{\infty}$ and $B \in \mathfrak{B}\left(\tilde{X}^{\infty}\right)$. Since the measurable space ( $\tilde{X}^{\infty}, \mathfrak{B}\left(\tilde{X}^{\infty}\right)$ ) is standard, and $\mathfrak{B}_{\infty}$ is an intersection of a decreasing sequence of the countably generated $\sigma$-fields $q_{n}^{-1}\left(\boldsymbol{B}\left(\tilde{X}^{\infty}\right)\right.$ ), so by the well known fact, (For example see theorem 2.3 in [2])

$$
\begin{equation*}
{ }^{\exists} A_{1} \in \mathfrak{B}_{\infty} \text { with } \tilde{\mu}\left(A_{1}\right)=1 \text { s.t., }{ }^{\forall} x \in A_{1}, \tilde{\mu}^{x}(\cdot)=1 \text { or } 0 \text { on } \mathfrak{B}_{\infty} \text {. } \tag{1.14}
\end{equation*}
$$

Furthermore it follows easily from the construction of $\tilde{\mu}$,

$$
\begin{align*}
& { }^{\exists} A_{2} \in \mathfrak{B}_{\infty} \text { with } \tilde{\mu}\left(A_{2}\right)=1 \text { s.t., } \forall x \in A_{2}, \tilde{\mu}^{x} \text { is } \Im_{\infty} \text {-quasi-invariant and }  \tag{1.15}\\
& \tilde{\mu}^{x}\left(\cup_{\sigma \in \Theta_{\infty}} s\left(\Gamma_{X}\right) \sigma\right)=1 \text {. }
\end{align*}
$$

Consequently putting $\mu^{[x]}:=p \tilde{\mu}^{x}$, we have $\mu^{[x]} \simeq \tilde{\mu}^{x}$ for all $x \in A_{2}$ by virtue of (c) in Theorem 1.2. Next we have for each fixed $\psi \in \operatorname{Diff}_{0} X, T_{\psi} \tilde{\mu}^{x} \simeq \tilde{\mu}^{x}$ for $\tilde{\mu}$-a.e. $x$, because every set in $\mathfrak{B}_{\infty}$ is Diff $X$-invariant. So using Theorem 1.3 and proceeding similar manner with the proof of Lemma 1.4, we deduce that

$$
\begin{equation*}
{ }^{\exists} A_{3} \in \mathfrak{B}_{\infty} \text { with } \tilde{\mu}\left(A_{3}\right)=1 \text { s.t., }{ }^{\forall} x \in A_{3}, \mu^{[x]} \text { is } \text { Diff }_{0} X \text {-quasi-invariant. } \tag{1.16}
\end{equation*}
$$

Thus we have,

$$
\begin{equation*}
{ }^{\forall} x \in A_{1} \cap A_{2} \cap A_{3}, \mu^{[x]} \text { is } \text { Diff }_{0} X \text {-ergodic, } \tag{1.17}
\end{equation*}
$$

by virtue of (c) and (d) in Theorem 1.2. Since

$$
\mu\left(s^{-1}\left(A_{1} \cap A_{2} \cap A_{3}\right)\right)=\sum_{\sigma \in \mathfrak{G}_{\infty}} c(\sigma)(s \mu) \sigma\left(A_{1} \cap A_{2} \cap A_{3}\right)=\tilde{\mu}\left(A_{1} \cap A_{2} \cap A_{3}\right)=1
$$

so the following result is obtained.
Theorem 1.4. Let $\mu$ be a Diff $_{0} X$-quasi-invariant probability measure on $\left(\Gamma_{X}, \mathscr{C}\right)$ with $\mu\left(\Delta_{X}\right)=1$. Then there exists a family of $\mathrm{Diff}_{0} X$-ergodic probability measures $\left\{\mu^{\nu}\right\}_{\gamma \in \Delta(X)}$ such that
(a) $\mu^{\gamma}(B)$ is a $s^{-1}\left(\mathfrak{B}_{\infty}\right)$-measurable function of $\gamma \in \Delta_{X}$ for each fixed $B \in \mathscr{C}$, and
(b) $\mu\left(B \cap s^{-1}(A)\right)=\int_{s^{-1}(A)} \mu^{\gamma}(B) \mu(d \gamma)$ for all $B \in \mathscr{C}$ and $A \in \mathfrak{B}_{\infty}$.

Proof. For it we have only to put $\mu^{\gamma}:=\mu^{[s(\gamma)]}$ if $\gamma \in s^{-1}\left(A_{1} \cap A_{2} \cap A_{3}\right)$ and $\mu^{\nu}:=\theta$, otherwise, where $\theta$ is some definite Diff $_{0} X$-ergodic probability measure on ( $\Gamma_{X}, \mathscr{C}$ ).

We wish to rewrite this decomposition in a somewhat elegant style which is independent of the admissible sections. For this let us put

$$
\mathfrak{A}_{\infty}:=\left\{B \in \mathscr{C} \mid T_{\psi} B=B \text { for all } \psi \in \operatorname{Diff}_{0} X\right\} .
$$

Then we have $s^{-1}\left(\mathfrak{B}_{\infty}\right) \subset \mathfrak{Q H}_{\infty}$, as is easily seen. Moreover for the $\mu$ in Theorem 1.4

$$
\mu(A \ominus \tilde{A})=0 \text { for all } A \in \mathfrak{H}_{\infty}, \text { where } \tilde{A}:=\left\{\gamma \in \Delta_{\boldsymbol{X}} \mid \mu^{\nu}(A)=1\right\} .
$$

Because,

$$
\mu(B \cap \tilde{A})=\int_{\tilde{A}} \mu^{\nu}(B) \mu(d \mu)
$$

for all $B \in \mathscr{C}$ by virtue of Theorem 1.4, while
by virtue of Theorem 1.4, while

$$
\mu(B \cap A)=\int_{\Delta_{x}} \mu^{\nu}(B) \mu^{\nu}(A) \mu(d \gamma)=\int_{\tilde{A}} \mu^{\nu}(B) \mu(d \gamma) .
$$

Theorem 1.5. Let $\mu$ be a Diff ${ }_{0} X$-quasi-invariant probability measure on $\left(\Gamma_{X}, \mathscr{C}\right)$. Then there exists a family of probability measures $\left\{\mu^{\nu}\right\}_{\gamma \in \Gamma_{X}}$ on $\left(\Gamma_{X}, \mathscr{C}\right)$ such that
(a) $\mu^{\gamma}$ is Diff $_{0} X$-ergodic for each $\gamma \in \Gamma_{X}$,
(b) $\mu^{\nu}(B)$ is an $\mathfrak{\mathfrak { r }}_{\infty}$-measurable function of $\gamma$ for each fixed $B \in \mathscr{C}$
(c) $\mu(A \cap B)=\int_{A} \mu^{\nu}(B) \mu(d \gamma)$ for all $A \in \mathfrak{N}_{\infty}$ and $B \in \mathscr{C}$.

Proof. First we divide $\mu$ into $\mu_{1}$ and $\mu_{2}$ as in the first place of this section and decompose $\mu_{1}$ into $\mu_{1, n}$ according to (1.1). Further we decompose $\mu_{2}$ into $\left\{\mu_{2}^{\nu}\right\}_{y \in A_{x}}$ as in Theorem 1.4. Next we define $\left\{\mu^{\nu}\right\}_{\gamma \in \Gamma_{X}}$ such that $\mu^{\gamma}=\mu_{2}^{\gamma}$ for $\gamma \in \Delta_{X}$ and $\mu^{\gamma}=\mu_{1, n}$ for $\gamma \in B_{x}^{n}$. Then the result easily follows from what we stated.

Lemma 1.5. For any $\operatorname{Diff}_{0} X$-quasi-invariant probability measures $\mu$ and $v$, the following are equivalent.
(a) $v$ is absolutely continuous with $\mu$.
(b) There exists $A \in \mathfrak{H}_{\infty}$ such that " $v(B)=0$ if and only if $\mu(A \cap B)=0$ ".

Proof. We have only to check the implication "(a) implies (b)". From the assumption there exists $A_{0} \in \mathscr{C}$ such that

$$
" v(B)=0 \text { if and only if } \mu\left(A_{0} \cap B\right)=0 "
$$

Thus $A_{0}$ must satisfy $\mu\left(A_{0} \ominus T_{\psi}\left(A_{0}\right)\right)=0$ for all $\psi \in \operatorname{Diff}_{0} X$. It follows from the above theorem that $\mu^{\nu}\left(A_{0} \ominus T_{\psi}\left(A_{0}\right)\right)=0$ for $\mu$-a.e. $\gamma$. Here let us take an arbitrary one parameter group $\left\{\psi_{t}\right\}_{t \in \mathbf{R}}$ contained in some $\operatorname{Diff}_{0} Y$. Then $\mu^{\nu}\left(A_{0} \ominus T_{\psi_{t}}\left(A_{0}\right)\right)$ is a $\mathfrak{B}(\mathbf{R}) \times \mathfrak{A}_{\infty}$-measurable function of $(t, \gamma) \in \mathbf{R} \times \Gamma_{X}$, which is easily checked, so by virtue of Fubini's theorem the Lebesgue measure of $Q_{\gamma}:=\left\{t \in \mathbf{R} \mid \mu^{\gamma}\left(A_{0} \ominus T_{\psi_{t}}\left(A_{0}\right)\right)=\right.$ $0\}$ is full for $\mu$-a.e. $\gamma$. As $Q_{\gamma}$ is a group, so $Q_{\gamma}=\mathbf{R}$ for $\mu$-a.e. $\gamma$. It follows from Theorem 1.3 that a measure $v^{\gamma}$ defined by the restriction of $\mu^{\nu}$ to the set $A_{0}$ is Diff $_{0} X$-quasi-invariant for $\mu$-a.e. $\gamma$. Since $\mu^{\gamma}$ is Diff $_{0} X$-ergodic, so $\mu^{\nu} \simeq \nu^{\nu}$ unless $\mu^{\nu}\left(A_{0}\right)=0$. That is $\mu^{\nu}\left(A_{0}\right)=1$ or 0 for $\mu$-a.e. $\gamma$. Thus we have $\mu\left(A_{0} \ominus A\right)=0$ for an $A$ defined by $A:=\left\{\gamma \in \Gamma_{X} \mid \mu^{\nu}\left(A_{0}\right)=1\right\} \in \mathfrak{H}_{\infty}$.

Theorem 1.6. For any Diff $_{0} X$-quasi-invariant probability measures $\mu$ and $v$, (a) $v \lesssim \mu$ if and only if $v \lesssim \mu$ on $\mathfrak{A}_{\infty}$.
(b) $\mu$ is Diff $_{0} X$-ergodic if and only if $\mu(\cdot)=1$ or 0 on $\mathfrak{H}_{\infty}$.
(c) If $\mu$ and $v$ are Diff $_{0} X$-ergodic, then $\mu \simeq v$ or $\mu \perp v$.

Proof. (a) Suppose that $v \lesssim \mu$ on $\mathfrak{A}_{\infty}$ and put $\lambda=(\mu+v) / 2$. Then by virtue of the above lemma, there exists $A \in \mathfrak{H}_{\infty}$ such that " $\mu(B)=0$ if and only if $\lambda(A \cap B)=0$ ". Especially we have $\mu\left(A^{c}\right)=0$ and thus $v\left(A^{c}\right)=0$. Consequently, $v(B)=v(B \cap A) \leq 2 \lambda(A \cap B)$, which implies $v(B)=0$ if $\mu(B)=0$. The converse relation is obvious. (b) and (c) easily follow from (a).

If we wish to be that factor measures $\left\{\mu^{\nu}\right\}_{\gamma \in \Gamma_{x}}$ appearing in Theorem 1.5 are mutually singular, then the following technique will be useful. First notice that a minimal $\sigma$-algebra $\mathscr{D}$ with which all the functions, $\gamma \in \Gamma_{X} \rightarrow \mu^{\nu}(B) \in \mathbf{R}$, where $B$ runs through $\mathscr{C}$, are measurable is countably generated and thus $\mathscr{D}=g^{-1}(\mathscr{B}(\mathbf{R}))$ with a suitable map $g: \Gamma_{X} \rightarrow \mathbf{R}$. It is not difficult to see,

$$
\begin{equation*}
g(\gamma)=g\left(\gamma^{\prime}\right), \text { if and only if } \mu^{\nu}=\mu^{\gamma^{\prime}} \text {. } \tag{1.18}
\end{equation*}
$$

Further by virtue of (c) in Theorem 1.5 we have,

$$
\begin{equation*}
\int_{g^{-1}(\boldsymbol{F})} \mu^{\gamma}\left(g^{-1}(E)\right) \mu(d \gamma)=\int_{g^{-1}(F)} \chi_{E}(g(\gamma)) \mu(d \gamma) \tag{1.19}
\end{equation*}
$$

for all $E, F \in \mathfrak{B}(\mathbf{R})$, and hence for $\mu$-a.e. $\gamma$,

$$
\begin{equation*}
\mu^{\nu}\left(g^{-1}(E)\right)=\chi_{E}(g(\gamma)) \tag{1.20}
\end{equation*}
$$

for all $E \in \mathfrak{B}(\mathbf{R})$. Especially we have,

$$
\begin{equation*}
\mu^{\gamma}\left(g^{-1}\{g(\gamma)\}\right)=1 \tag{1.21}
\end{equation*}
$$

for $\mu$-a.e. $\gamma$. Now define $\mu_{t}^{\prime}=\mu^{\gamma}$, if $t=g(\gamma)$ and $\mu_{t}^{\prime}=\theta$, otherwise, where $\theta$ is some definite $\mathrm{Diff}_{0} X$-ergodic probability measure on $\left(\Gamma_{X}, \mathscr{C}\right)$. Then we have
(1.22) For each fixed $B \in \mathscr{C}, \mu_{t}^{\prime}(B)$ is a universally measurable function of $t \in \mathbf{R}$.

Because $g\left\{\gamma \mid \mu^{\gamma}(B) \leq a\right\}$ is an analytic set for every $a \in \mathbf{R}$. And further we have,

$$
\begin{equation*}
\mu\left(B \cap g^{-1}(E)\right)=\int_{E} \mu_{t}^{\prime}(B) g \mu(d t) \tag{1.23}
\end{equation*}
$$

for all $B \in \mathscr{C}$ and $E \in \mathfrak{B}(\mathbf{R})$. Compairing $\mu_{t}^{\prime}$ with regular conditional probability measure $p(t, \cdot)$ given $g=t$, we deduce that

$$
\begin{equation*}
{ }^{\exists} T_{1} \in \mathfrak{B}(\mathbf{R}) \text { with } g \mu\left(T_{1}\right)=1 \text { such that }{ }^{\forall} t \in T_{1}, \mu_{t}^{\prime}=p(t, \cdot) \text {. } \tag{1.24}
\end{equation*}
$$

Finally we put $\mu_{t}(\cdot)=p(t, \cdot)$, if $t \in T_{1}$ and $\mu_{t}=\theta$, otherwise. Then
Theorem 1.7. Let $\mu$ be a Diff $_{0} X$-quasi-invariant probability measure. Then there exist a map $g$ and a family of probability measures $\left\{\mu_{t}\right\}_{t \in \mathbf{R}}$ on $\left(\Gamma_{X}, \mathscr{C}\right)$ such that
(a) $g$ is a measurable map from $\left(\Gamma_{X}, \mathfrak{A}_{\infty}\right)$ to $(\mathbf{R}, \mathfrak{B}(\mathbf{R}))$,
(b) $\mu_{t}$ is Diff $_{0} X$-ergodic for every $t \in \mathbf{R}$,
(c) $\mu_{t}(B)$ is a Borel measurable function of $t \in \mathbf{R}$ for each fixed $B \in \mathscr{C}$,
(d) there exists a Borel set $T_{0}$ with $g \mu\left(T_{0}\right)=1$ such that $\mu_{t}\left(g^{-1}\{t\}\right)=1$ for all $t \in T_{0}$, especially $\mu_{t}\left(t \in T_{0}\right)$ are mutually singular, and
(e) $\mu\left(B \cap g^{-1}(E)\right)=\int_{E} \mu_{t}(B) g \mu(d t)$ for all $B \in \mathscr{C}$ and $E \in \mathfrak{B}(\mathbf{R})$.

## 2. Ergodic decomposition of Gibbs measure

2.1. Basic properties. In this section $X$ is a general locally compact topological space which satisfies the second countable axiom and $m$ stands for non atomic Radon measures on $\mathfrak{B}(X)$ which is the natural Borel $\sigma$-field on $X$. A function $U(x \mid \gamma) \in(-\infty, \infty]$ defined on $(x, \gamma) \in X \times \Gamma_{X}$ is said to be a potential if it satisfies

$$
\begin{equation*}
U(x \mid \gamma) \text { is a } \mathfrak{B}(\mathbf{R}) \times \mathscr{C} \text {-measurable function, and } \tag{2.1}
\end{equation*}
$$

$$
\begin{equation*}
U(x \mid \gamma \cup\{y\})+U(y \mid \gamma)=U(y \mid \gamma \cup\{x\})+U(x \mid \gamma) \tag{2.2}
\end{equation*}
$$

for all $x, y \in X$ and $\gamma \in \Gamma_{X}$. We shall extend the domain of definition of the potential to $B_{x}^{n}$ such that
$U(\phi \mid \gamma):=0$ for $n=0, U\left(\left\{x_{1}, x_{2}\right\} \mid \gamma\right):=U\left(x_{1} \mid \gamma \cup\left\{x_{2}\right\}\right)+U\left(x_{2} \mid \gamma\right)$ for $n=2$, and $U(\underline{x} \mid \gamma)=U\left(\left\{x_{1}, \cdots, x_{n-1}\right\} \mid \gamma \cup\left\{x_{n}\right\}\right)+U\left(x_{n} \mid \gamma\right)$ for $\underline{x}:=\left\{x_{1}, \cdots, x_{n}\right\} \in B_{X}^{n}$ inductively. These are well defined by the property (2.2).
Now let $\mu$ be a probability measure on $\left(\Gamma_{X}, \mathscr{C}\right)$ and denote the conditional expectation of a $\mathscr{G}$-measurable function $f$ on $\Gamma_{X}$ with respect to the $\sigma$-field $\pi_{Y_{c}}^{-1}\left(\mathscr{C}_{Y c}\right)$ by $\operatorname{Exp}\left(f \mid \mathscr{C}_{Y c}\right)$. Let us proceed to the definition of Gibbs measure. A
probability measure $\mu$ on $\left(\Gamma_{X}, \mathscr{C}\right)$ is said to be $(U, m)$-Gibbsian or simply Gibbsian (in a sense of Dobrushin, Ruelle, Lanford) for a potential $U$ and a measure $m$ if and only if it satisfies,

$$
\begin{equation*}
\Xi_{K}:=\sum_{n=0}^{\infty} n!^{-1} \int_{B_{K}^{\prime \prime}} \exp \left(-U\left(\underline{x} \mid \gamma \cap K^{c}\right)\right) m_{K, n}(d \underline{x})<\infty \tag{2.3}
\end{equation*}
$$

for $\mu$-a.e. $\gamma$, and

$$
\begin{align*}
& \operatorname{Exp}\left(f \mid \mathscr{C}_{K}\right)(\gamma)=  \tag{2.4}\\
& \quad \Xi_{K}(\gamma)^{-1} \sum_{n=0}^{\infty} n!^{-1} \int_{B_{K}^{u}} \exp \left(-U\left(\underline{x} \mid \gamma \cap K^{c}\right)\right) f\left(\underline{x} \cup\left(\gamma \cap K^{c}\right)\right) m_{K, n}(d \underline{x})
\end{align*}
$$

for each non negative bounded $\mathscr{C}$-measurable function $f$ on $\Gamma_{X}$. Notice that we always have $\Xi_{K}(\gamma) \geq 1$. And it is fairly easy to see that a set of all $(U, m)$-Gibbsian measure is closed under the convex combination. From now on we shall write

$$
\int^{K} \exp \left(-U\left(\underline{x} \mid \gamma \cap K^{c}\right)\right) f\left(\underline{x} \cdot \gamma \cap K^{c}\right) m(d \underline{x})
$$

instead of

$$
\sum_{n=0}^{\infty} n!^{-1} \int_{B_{K}^{\prime}} \exp \left(-U\left(\underline{x} \mid \gamma \cap K^{c}\right)\right) f\left(\underline{x} \cup\left(\gamma \cap K^{c}\right)\right) m_{K, n}(d \underline{x})
$$

according to [4].
Lemma 2.1. (2.3) and (2.4) is equivalent to the following condition (2.5).

$$
\begin{equation*}
\int_{\Gamma_{X}} f(\gamma) \mu(d \gamma)=\int_{\{\gamma| | \gamma \cap K \mid=0\}} \mu(d \gamma) \int^{K} \exp (-U(\underline{x} \mid \gamma)) f(\underline{x} \cdot \gamma) m(d \underline{x}) \tag{2.5}
\end{equation*}
$$

for each compact set $K$ and non negative bounded $\mathscr{C}$-measurable function $f$.
Proof. Suppose that (2.3) and (2.4) are satisfied and let $\chi_{N_{K}}$ be the indicator function of the set $N_{K}:=\{\gamma \| \gamma \cap K \mid=0\}$. Then for $f=\chi_{N_{K}}(2.4)$ gives

$$
\begin{equation*}
\operatorname{Exp}\left(\chi_{N_{K}} \mid \mathscr{E}_{K^{c}}\right)(\gamma)=\Xi_{K}(\gamma)^{-1} . \tag{2.6}
\end{equation*}
$$

Thus,

$$
\begin{aligned}
& \int_{\Gamma_{X}} f(\gamma) \mu(d \gamma)=\int_{\Gamma_{X}} \operatorname{Exp}\left(f \mid \mathscr{C}_{K^{c}}\right)(\gamma) \mu(d \gamma) \\
= & \left.\int_{\Gamma_{X}} \mu(d \gamma) \operatorname{Exp}\left(\chi_{N_{K}} \mid \mathscr{C}_{K^{c}}\right)(\gamma) \int^{K} \exp \left(-U\left(\underline{x} \mid \gamma \cap K^{c}\right)\right) f\left(\underline{x} \cdot \gamma \cap K^{c}\right)\right) m(d \underline{x}) \\
= & \left.\int_{N_{K}} \mu(d \gamma)\right)^{K} \exp (-U(\underline{x} \mid \gamma)) f(\underline{x} \cdot \gamma) m(d \underline{x}) .
\end{aligned}
$$

Conversely, Put $F_{K}:=\left\{\gamma \mid \Xi_{K}(\gamma)<\infty\right\}$ and substitute the indicator function $\chi_{F_{K}}$ for $f$ in (2.5). Then it yields that

$$
\mu\left(F_{K^{c}}\right)=\int_{N_{K}} \chi_{F_{K}^{c}}(\gamma) \Xi_{K}(\gamma) \mu(d \gamma),
$$

and thus $\mu\left(F_{K}^{c}\right)=0$. The rest of the proof easily follows from (2.6) which is easily derived from (2.5).

Let us look quickly how the Gibbsian property implies Diff $_{0} X$-quasiinvariance. So let $\mu$ be a Gibbs measure and $Y$ be any open subset with compact closure. Then as is easily seen, (2.4) also holds for such $Y$ provided that $m(\bar{Y} \backslash Y)=0$. Thus the conditional probability measure $\mu^{\nu}$ with respect to $\pi_{Y c}^{-1}\left(\mathscr{C}_{Y c}\right)$ is,

$$
\begin{equation*}
\mu^{\nu}(A)=\Xi_{Y}(\gamma)^{-1} \sum_{n=0}^{\infty} n!^{-1} \int_{A \cap B_{Y}^{u}} \exp \left(-U\left(\underline{x} \mid \gamma \cap Y^{c}\right)\right) m_{Y, n}(d \underline{x}) \tag{2.7}
\end{equation*}
$$

for all $A \in \mathscr{C}_{Y}$. Hence we have,
Theorem 2.1. Let $X$ be a connected $\sigma$-compact $C^{\infty}$-manifold and $m$ be a locally Euclidean Radon measure on $X$. Then under the assumption that the potential function $U(x \mid \gamma)$ is always finite, any $(U, m)$-Gibbs measure $\mu$ is Diff ${ }_{0} X$-quasi-invariant.

Proof. Take a sequence $\left\{Y_{n}\right\}$ of connected open sets with compact closure such that $m\left(\bar{Y}_{n} \backslash Y_{n}\right)=0$ and apply Lemma 1.4.

Let $\left\{U_{n}\right\}$ be a countable open base of $X$ such that $\bar{U}_{n}$ is compact for all $n, \mathscr{K}$ be a collection of all the sets being finite union of $U_{n}(n=1, \cdots)$, and $\mathscr{F}$ be a countable field generatig $\mathscr{C}$.

Lemma 2.2. In order that a probability measure $\mu$ on $\left(\Gamma_{X}, \mathscr{C}\right)$ is Gibbsian, it is necessary and sufficient that (2.5) is satisfied for all $K \in \mathscr{K}$ and $\chi_{B}(=f)$, where $\chi_{B}$ is the indicator function of a set $B \in \mathscr{F}$.

Proof. We have only to check the sufficiency. Now it is immediate from the assumption that (2.5) holds for all $K \in \mathscr{K}$ and for all non negative bounded $\mathscr{C}$-measurable functions. Hence proceeding in the same way with the proof of Lemma 2.1, we have for each $K \in \mathscr{K}$

$$
\begin{equation*}
\Xi_{K}(\gamma)<\infty \tag{2.8}
\end{equation*}
$$

for $\mu$-a.e. $\gamma$, and for each $K \in \mathscr{K}$

$$
\begin{equation*}
\Xi_{K}^{-1}(\gamma)=\operatorname{Exp}\left(\chi_{N_{K}} \mid \mathscr{C}_{K^{c}}\right)(\gamma) \tag{2.9}
\end{equation*}
$$

for $\mu$-a.e. $\gamma$. Take any compact set $K$. Then there exists a sequence $\left\{K_{n}\right\} \subset \mathscr{K}$ such that $K_{n} \downarrow K$. It gives that

$$
\left\{\gamma \|\left|\gamma \cap K_{n}\right|=0\right\} \uparrow\{\gamma \| \gamma \cap K \mid=0\} \text { and }\left\{\gamma \mid \gamma \cap K_{n}^{c}=\gamma \cap K^{c}\right\} \uparrow \Gamma_{X} .
$$

Here we notice that if $\gamma \cap K^{c}=\gamma \cap K_{n}^{c}$ for some $n$, then

$$
\Xi_{K}(\gamma) \leq \int^{K_{n}} \exp \left(-U\left(\underline{x} \mid \gamma \cap K_{n}^{c}\right)\right) m(d \underline{x})=\Xi_{K_{n}}(\gamma) .
$$

Thus (2.8) and the above relation show that (2.3) holds for all compact sets $K$. By the assumption (2.5) holds for all $K_{n}$, so by virtue of Lebesgue-Fatou's lemma we have,

$$
\begin{equation*}
\left.\int_{\Gamma \mathbb{x}} f(\gamma) \mu(d \gamma) \geq \int_{N_{K}} \mu(d \gamma) \int^{K} \exp \left(-U\left(\underline{x} \mid \gamma \cap K^{c}\right)\right) f\left(\underline{x} \cdot \gamma \cap K^{c}\right)\right) m(d \underline{x}) \tag{2.10}
\end{equation*}
$$

And hence,

$$
\begin{align*}
& \operatorname{Exp}\left(f \mid \mathscr{C}_{\mathbf{K}^{c}}\right)(\gamma) \geq  \tag{2.11}\\
& \quad \operatorname{Exp}\left(\chi_{N_{K}} \mid \mathscr{C}_{\mathbf{K}^{c}}\right)(\gamma) \cdot \int^{K} \exp \left(-U\left(\underline{x} \mid \gamma \cap K^{c}\right)\right) f\left(\underline{x} \cdot \gamma \cap K^{c}\right) m(d \underline{x})
\end{align*}
$$

for $\mu$-a.e. $\gamma$. Especially,

$$
\begin{equation*}
\Xi_{K}(\gamma)^{-1} \geq \operatorname{Exp}\left(\chi_{N_{K}} \mid \mathscr{C}_{K^{c}}\right)(\gamma) \tag{2.12}
\end{equation*}
$$

for $\mu$-a.e. $\gamma$. Now let us consider the relation (2.9) for $K=K_{n}$. As is easily seen, $\pi_{K_{n}^{c}}^{-1}\left(\mathscr{C}_{K_{n}^{c}}\right) \uparrow \pi_{K^{c}}^{-1}\left(\mathscr{C}_{K^{c}}\right)$, so the right hand side of (2.9) is

$$
\left\{\begin{array}{l}
\operatorname{Exp}\left(\chi_{N_{K_{n}}} \mid \mathscr{C}_{K_{n}^{c}}\right)(\gamma) \leq \operatorname{Exp}\left(\chi_{N_{K}} \mid \mathscr{C}_{K_{n}^{c}}\right)(\gamma)  \tag{2.13}\\
\operatorname{Exp}\left(\chi_{N_{K}} \mid \mathscr{C}_{K_{n}^{c}}\right)(\gamma) \rightarrow \operatorname{Exp}\left(\chi_{K} \mid \mathscr{C}_{K^{c}}\right)(\gamma) \text { as } n \rightarrow \infty
\end{array}\right.
$$

While for the left hand side, first we put $F_{\infty}:=\bigcap_{n=1}^{\infty} F_{K_{n}}$. Then (2.8) gives $\mu\left(F_{\infty}\right)=1$. And if $\gamma \in F_{\infty}$ and $\gamma \cap K_{N}^{c}=\gamma \cap K^{c}$ for some $N$, then for all $n \geq N$,

$$
\Xi_{K_{n}}(\gamma)=\sum_{l=0}^{L} l!^{-1} \int_{B_{K_{n}}^{\prime}} \exp \left(-U\left(\underline{x} \mid \gamma \cap K^{c}\right)\right) m_{K_{n}, l}(d \underline{x})+\varepsilon_{L, n},
$$

where

$$
\begin{aligned}
\varepsilon_{L, n}:= & \sum_{l=L+1}^{\infty} l!^{-1} \int_{B_{K_{n}}^{\prime}} \exp \left(-U\left(\underline{x} \mid \gamma \cap K^{c}\right)\right) m_{K_{n}, l}(d \underline{x}) \leq \\
& \sum_{l=L+1}^{\infty} l!^{-1} \int_{B_{K_{N}}^{\prime}} \exp \left(-U\left(\underline{x} \mid \gamma \cap K^{c}\right)\right) m_{K_{N}, l}(d \underline{x}) .
\end{aligned}
$$

And if we take a sufficiently large $L$, the last term becomes smaller than $\varepsilon$ for a given $\varepsilon>0$. Consequently for such an $L$,

$$
\overline{\lim }_{n} \Xi_{K_{n}}(\gamma) \leq \varepsilon+\sum_{l=0}^{L} l!^{-1} \int_{B_{K}^{\prime}} \exp \left(-U\left(\underline{x} \mid \gamma \cap K^{c}\right)\right) m_{K, l}(d \underline{x}) \leq \varepsilon+\Xi_{K}(\gamma) .
$$

So we have,

$$
\begin{equation*}
\lim _{n} \Xi_{K_{n}}(\gamma) \leq \Xi_{K}(\gamma) \tag{2.14}
\end{equation*}
$$

for $\mu$-a.e. $\gamma$. It follows from (2.13) and easy calculations that

$$
\begin{equation*}
\Xi_{K}(\gamma)^{-1} \leq \operatorname{Exp}\left(\chi_{N_{K}} \mid \mathscr{C}_{K^{c}}\right)(\gamma) \tag{2.15}
\end{equation*}
$$

for $\mu$-a.e. $\gamma$. This and (2.12) show that (2.6) holds for all compact sets $K$. Now the inequality (2.11) becomes,

$$
\begin{equation*}
\operatorname{Exp}\left(f \mid \mathscr{C}_{K^{c}}\right)(\gamma) \geq \Xi_{K}(\gamma)^{-1} \int^{K} \exp \left(-U\left(\underline{x} \mid \gamma \cap K^{c}\right)\right) f\left(\underline{x} \cdot \gamma \cap K^{c}\right) m(d \underline{x}) \tag{2.16}
\end{equation*}
$$

for $\mu$-a.e. $\gamma$. By the way (2.16) becomes an equality for $f=$ const, thus it is actually an equality for any $f \geq 0$.

Let us take and fix an increasing sequence $\left\{K_{n}\right\}$ of compact sets such that $K_{n} \uparrow X$, and consider the tail $\sigma$-field $\mathscr{C}_{\infty}:=\bigcap_{n=1}^{\infty} \pi_{K_{n}^{c}}^{-1}\left(\mathscr{C}_{K_{n}^{c}}\right)$. $\mathscr{C}_{\infty}$ does not depend on a particular choice of $\left\{K_{n}\right\}$.

Theorem 2.1. Let $\mu$ be a $(U, m)$-Gibbs measure and $\left\{\mu_{\infty}^{\nu}\right\}_{\nu \in \Gamma_{x}}$ be a family of conditional probability measure of $\mu$ with respect to $\mathscr{C}_{\infty}$. Then $\mu_{\infty}^{\gamma}$ is $(U, m)$ Gibbsian for $\mu$-a.e. $\gamma$.

Proof. For $A \in \mathscr{C}_{\infty}$ and $B \in \mathscr{F}$ we calculate $\mu(A \cap B)$ in two ways. The first one is,

$$
\mu(A \cap B)=\int_{A} \mu_{\infty}^{\gamma}(B) \mu(d \gamma)
$$

and the other one is,

$$
\mu(A \cap B)=\int_{N_{K}} \mu(d \gamma) \chi_{A}(\gamma) \int^{K} \exp \left(-U\left(\underline{x} \mid \gamma \cap K^{c}\right)\right) \chi_{B}\left(\underline{x} \cdot \gamma \cap K^{c}\right) m(d \underline{x}),
$$

where $K$ is taken from $\mathscr{K}$. These show that

$$
\begin{equation*}
\mu_{\infty}^{\delta}(B)=\int_{N_{K}} \mu_{\infty}^{\delta}(d \gamma) \int^{K} \exp \left(-U\left(\underline{x} \mid \gamma \cap K^{c}\right)\right) \chi_{B}\left(\underline{x} \cdot \gamma \cap K^{c}\right) m(d \underline{x}) \tag{2.17}
\end{equation*}
$$

for $\mu$-a.e. $\delta$. Since $\mathscr{K}$ and $\mathscr{F}$ are countable, so the assertion directly follows from Lemma 2.2.

Here we introduce a notion of ergodicity. A Gibbs measure $\mu$ on $\left(\Gamma_{X}, \mathscr{C}\right)$ is said to be ergodic if and only if $\mu(\cdot)=1$ or 0 on $\mathscr{C}_{\infty}$. As $\left(\Gamma_{X}, \mathscr{C}\right)$ is a standard space and $\mathscr{C}_{\infty}$ is an intersection of a decreasing sequence of countably generated $\sigma$-fields $\pi_{K_{n}^{-1}}^{-1} \mathscr{C}_{K_{n}^{c}}$, so $\mu_{\infty}^{\gamma}$ is ergodic for almost all $\gamma$ by a well known result (For example see theorem 2.3 in [2]) and the ergodic decomposition seems
to be settled.
2.2. Ergodic measures as a base. However we will have a stronger result that factor measures $\mu_{\infty}^{\nu}\left(\gamma \in \Gamma_{X}\right)$ do not depend on each $\mu$. From now on we take and fix a countable field $\mathscr{F}_{0}$ generating $\mathscr{C}$ such that any finitely additive finite measure on $\mathscr{F}_{0}$ has a $\sigma$-additive extension on $\mathscr{C}$. The existence of such $\mathscr{F}_{0}$ is assured by the fact that $\left(\Gamma_{X}, \mathscr{C}\right)$ is standard. Now set

$$
\begin{aligned}
& \Omega_{1}:=\left\{\gamma \mid \Xi_{K_{n}}(\gamma)<\infty \text { holds except finitely many } n ’ \mathrm{~s}\right\} \\
& \Omega_{2}:=\left\{\gamma \in \Omega_{1} \mid \lim _{n} \Xi_{K_{n}}(\gamma)^{-1} \int^{K_{n}} \exp \left(-U\left(\underline{x} \mid \gamma \cap K_{n}^{c}\right)\right) \chi_{B}\left(\underline{x} \cdot \gamma \cap K_{n}^{c}\right) m(d \underline{x})\right. \\
& \left.\quad \text { exists for every } B \in \mathscr{F}_{0}\right\} .
\end{aligned}
$$

Then $\Omega_{1}, \Omega_{2} \in \mathscr{C}_{\infty}$ and for any ( $U, m$ )-Gibbs measure $\mu, \mu\left(\Omega_{2}\right)=1$ by virtue of the martingale convergence theorem. And by the nice property of $\mathscr{F}_{0}$, we can define a probability measure $\omega_{\gamma}^{0}\left(\gamma \in \Omega_{2}\right)$ on $\left(\Gamma_{X}, \mathscr{C}\right)$ as the extension of a finitely additive measure:

$$
B \in \mathscr{F}_{0} \rightarrow \lim _{n} \Xi_{K_{n}}(\gamma)^{-1} \int^{K_{n}} \exp \left(-U\left(\underline{x} \mid \gamma \cap K_{n}^{c}\right)\right) \chi_{B}\left(\underline{x} \cdot \gamma \cap K_{n}^{c}\right) m(d \underline{x}) .
$$

Let us make up a definition $\omega_{\gamma}^{0}$ as $\omega_{\gamma}^{0}=\zeta$ for $\gamma \in \Omega_{2}^{c}$, where $\zeta$ is some definite $(U, m)$-Gibbs ergodic measure. Then

$$
\begin{equation*}
\omega_{\gamma}^{0}(B) \text { is a } \mathscr{C}_{\infty} \text {-measurable function of } \gamma \in \Gamma_{X} \text { for each fixed } B \in \mathscr{C} . \tag{2.18}
\end{equation*}
$$

Further by virtue of the martingale convergence theorem we have for any Gibbs measure $\mu$,

$$
\begin{equation*}
\mu(A \cap B)=\int_{A} \omega_{\gamma}^{0}(B) \mu(d \gamma) \tag{2.19}
\end{equation*}
$$

for all $A \in \mathscr{C}_{\infty}$ and $B \in \mathscr{C}$. Because (2.19) is first valid for $B \in \mathscr{F}_{0}$ and holds in general by the extension property. It follows from (2.19) that

$$
\begin{equation*}
\omega_{\gamma}^{0}=\mu_{\infty}^{\gamma} \tag{2.20}
\end{equation*}
$$

for $\mu$-a.e. $\gamma$. Here we shall put

$$
\begin{array}{r}
\Omega_{3}:=\left\{\delta \in \Gamma_{X} \mid \omega_{\delta}^{0}(B)=\int_{N_{K}} \omega_{\delta}^{0}(d \gamma) \int^{K} \exp \left(-U\left(\underline{x} \mid \gamma \cap K^{c}\right)\right) \chi_{B}\left(\underline{x} \cdot \gamma \cap K^{c}\right) m(d \underline{x})\right. \\
\text { for all } \left.B \in \mathscr{F}_{0} \text { and } K \in \mathscr{K}\right\} .
\end{array}
$$

Then $\Omega_{3} \in \mathscr{C}_{\infty}$ and (2.20) gives $\mu\left(\Omega_{3}\right)=1$ for any Gibbs measure $\mu$. And it follows from Lemma 2.2 that

$$
\begin{equation*}
\omega_{\delta}^{0} \text { is }(U, m) \text {-Gibbsian for each } \delta \in \Omega_{3} . \tag{2.21}
\end{equation*}
$$

Thus (2.19) derives that for $\delta \in \Omega_{3}$

$$
\begin{equation*}
\omega_{\delta}^{0}(A \cap B)=\int_{A} \omega_{\gamma}^{0}(B) \omega_{\delta}^{0}(d \gamma) \tag{2.22}
\end{equation*}
$$

for all $A \in \mathscr{C}_{\infty}$ and $B \in \mathscr{C}$. Finally we put

$$
\Omega_{4}:=\left\{\delta \in \Omega_{3} \mid \int_{\Gamma x}\left\{\omega_{\gamma}^{0}(B)-\omega_{\delta}^{0}(B)\right\}^{2} \omega_{\delta}^{0}(d \gamma)=0 \quad \text { for all } B \in \mathscr{F}_{0}\right\} .
$$

Then we have $\Omega_{4} \in \mathscr{C}_{\infty}$ and for any ( $U, m$ )-Gibbs measure $\mu$,

$$
\int_{\Gamma_{x}} \int_{\Gamma_{x}}\left\{\omega_{\gamma}^{0}(B)-\omega_{\delta}^{0}(B)\right\}^{2} \omega_{\delta}^{0}(d \gamma) \mu(d \gamma)=\int_{\Gamma_{x}} \int_{\Gamma_{x}}\left\{\mu_{\infty}^{\gamma}(B)-\mu_{\infty}^{\delta}(B)\right\}^{2} \mu_{\infty}^{\delta}(d \gamma) \mu(d \gamma)=0
$$

Thus we have $\mu\left(\Omega_{4}\right)=1$. Moreover it follows from (2.22) that $\omega_{\delta}^{0}(A \cap B)=$ $\omega_{\delta}^{0}(A) \omega_{\delta}^{0}(B)$ for each $\delta \in \Omega_{4}$, which implies $\omega_{\delta}^{0}(A)=1$ or 0 for all $A \in \mathscr{C}_{\infty}$. Thus

$$
\begin{equation*}
\omega_{\delta}^{0} \text { is ergodic for each } \delta \in \Omega_{4} \text {. } \tag{2.23}
\end{equation*}
$$

Define $\omega_{\delta}=\omega_{\delta}^{0}$, if $\delta \in \Omega_{4}$ and $\omega_{\delta}=\zeta$, otherwise. Then we have,
Theorem 2.2. As for a convex set formed by all $(U, m)$-Gibbs measures, there exists a family of probability measures $\left\{\omega_{\gamma}\right\}_{\gamma \in \Gamma_{X}}$ on $\left(\Gamma_{X}, \mathscr{C}\right)$ such that
(a) $\omega_{\gamma}$ is a $(U, m)$-Gibbs ergodic measure for each $\gamma \in \Gamma_{X}$,
(b) $\omega_{\gamma}(B)$ is a $\mathscr{C}_{\infty}$-measurable function of $\gamma \in \Gamma_{X}$ for each fixed $B \in \mathscr{C}$ and
(c) for any $(U, m)$-Gibbs measure $\mu$

$$
\mu(A \cap B)=\int_{A} \omega_{\gamma}(B) \mu(d \gamma) \text { for all } A \in \mathscr{C}_{\infty} \text { and } B \in \mathscr{C} .
$$

Corollary. For any $(U, m)$-Gibbs measures $\mu$ and $v$,
(a) $\mu=v$ if and only if $\mu=v$ on $\mathscr{C}_{\infty}$.
(b) $v \lesssim \mu$ if and only if $v \lesssim \mu$ on $\mathscr{C}_{\infty}$.
(c) If $\mu$ and $v$ are ergodic, then $\mu=v$ or $\mu \perp v$.

Let us take and fix an above family $\left\{\omega_{\gamma}\right\}_{\gamma \in \Gamma_{x}}$ and consider a minimal $\sigma$-field $\mathscr{C}_{\infty, \omega}$ with which all the functions, $\gamma \rightarrow \omega_{\gamma}(B)(B \in \mathscr{C})$ are measurable. Since $\mathscr{C}_{\infty, \omega}$ is countably generated, so there exists a map $h: \Gamma_{X} \rightarrow \mathbf{R}$ such that $\mathscr{C}_{\infty, \omega}=$ $h^{-1}(\mathfrak{B}(\mathbf{R}))$. As before it is easily checked that

$$
\begin{equation*}
\omega_{\gamma}=\omega_{\gamma^{\prime}} \text { if and only if } h(\gamma)=h\left(\gamma^{\prime}\right) . \tag{2.24}
\end{equation*}
$$

Further we claim that

$$
\begin{equation*}
h \omega_{\gamma}=\delta_{h(\gamma)} \tag{2.25}
\end{equation*}
$$

for all $\gamma \in \Gamma_{X}$, where $\delta_{s}$ is the Dirac measure at $s \in \mathbf{R}$. For, put $S:=\left\{\gamma \in \Gamma_{X} \mid\right.$ $\left.h \omega_{\gamma}=\delta_{h(y)\}}\right\}$. Then we have $S \in \mathscr{C}_{\infty, \omega}$ and for any Gibbs measure $\mu$,

$$
\mu\left(h^{-1}(E \cap F)\right)=\int_{h^{-1}(F)} \omega_{\gamma}\left(h^{-1}(E)\right) \mu(d \gamma)=\int_{h^{-1}(F)} \chi_{E}(h(\gamma)) \mu(d \gamma)
$$

for all $E, F \in \mathfrak{B}(\mathbf{R})$. Since both the integrants are $\mathscr{C}_{\infty, \omega}$-measurable, so it follows that $\omega_{\gamma}\left(h^{-1}(E)\right)=\chi_{E}(h(\gamma))$ for $\mu$-a.e. $\gamma$, and thus $\mu(S)=1$. Especially we have,

$$
\begin{equation*}
\omega_{\gamma}(S)=1 \tag{2.26}
\end{equation*}
$$

for all $\gamma \in \Gamma_{\boldsymbol{X}}$. Now for any fixed $\gamma \in \Gamma_{X}$ let us take $\sigma \in\left\{\delta \mid \omega_{\delta}=\omega_{\gamma}\right\} \cap S$. Notice that the last set is not empty because we have $\omega_{\delta}=\omega_{\gamma}$ for $\omega_{\gamma}$-a.e. $\delta$ by virtue of the ergodicity of $\omega_{\gamma}$. Then,

$$
\omega_{\gamma}\left(h^{-1}(E)\right)=\omega_{\sigma}\left(h^{-1}(E)\right)=\chi_{E}(h(\sigma))=\chi_{E}(h(\gamma)) .
$$

We settle these arguments as the following theorem.
Theorem 2.3. Under the notation in Theorem 2.2, let $\mathscr{C}_{\infty, \omega}$ be the minimal $\sigma$-field with which all the functions, $\gamma \rightarrow \omega_{\gamma}(B)(B \in \mathscr{C})$ are measurable. Then there exists a map $h: \Gamma_{X} \rightarrow \mathbf{R}$ such that $\mathscr{C}_{\infty, \omega}=h^{-1}(\mathfrak{B}(\mathbf{R}))$ and we have
(a) $\omega_{\gamma}=\omega_{\gamma^{\prime}}$ if and only if $h(\gamma)=h\left(\gamma^{\prime}\right)$, and
(b) $h \omega_{\gamma}=\delta_{h(\gamma)}$ for all $\gamma \in \Gamma_{X}$.

Remark 2.1. As for the uniqueness of such a family $\left\{\omega_{\gamma}\right\}_{\gamma \in \Gamma_{x}}$, it is desirable to state it without exceptional set instead of with exceptional set of measure 0 . The following is an answer for this question. Namely, in order that such families $\left\{\omega_{\gamma}\right\}_{\gamma \in \Gamma_{x}}$ and $\left\{\omega_{\gamma}^{\prime}\right\}_{\gamma \in \Gamma_{x}}$ coincide, it is necessary and sufficient that the $\sigma$-fields generated by them are the same one, i.e., $\mathscr{C}_{\infty, \omega}=\mathscr{C}_{\infty, \omega}$.

Theorem 2.4. A map : $\lambda(\cdot) \rightarrow \int_{\Gamma_{x}} \omega_{\gamma}(\cdot) \lambda(d \gamma)$ is a bijection from a space of all probability measures on $\left(\Gamma_{X}, \mathscr{C}_{\infty, \omega}\right)$ to the space of all $(U, m)$-Gibbs measures.

The proof is obvious from what we have stated.
We conclude this paragraph with the following theorem.
Theorem 2.5. As for the convex set formed by all $(U, m)$-Gibbs measures, there exist a map $h: \Gamma_{X} \rightarrow \mathbf{R}$ and a family of probability measures $\left\{\beta_{\tau}\right\}_{\tau \in h\left(\Gamma_{x}\right)}$ on $\left(\Gamma_{X}, \mathscr{C}\right)$ such that
(a) $h$ is a measurable map from $\left(\Gamma_{X}, \mathscr{C}_{\infty}\right)$ to $(\mathbf{R}, \mathfrak{B}(\mathbf{R}))$,
(b) $\beta_{\tau}$ is a $(U, m)$-Gibbs ergodic measure for each $\tau \in h\left(\Gamma_{X}\right)$,
(c) $\beta_{\tau}(B)$ is a $h\left(\Gamma_{X}\right) \cap \mathfrak{B}(\mathbf{R})$-measurable function of $\tau \in h\left(\Gamma_{X}\right)$ for each fixed $B \in \mathscr{C}$, and hence it is universally measurable,
(d) $\beta_{\tau}\left(h^{-1}(\tau)\right)=1$ for all $\tau \in h\left(\Gamma_{X}\right)$, especially $\beta_{\tau}\left(\tau \in h\left(\Gamma_{X}\right)\right)$ are mutually singular, and
(e) for any $(U, m)$-Gibbs measure $\mu$,

$$
\mu\left(B \cap h^{-1}(E)\right)=\int_{E} \beta_{\tau}(B) h \mu(d \tau) \text { for all } B \in \mathscr{C} \text { and } E \in \mathfrak{B}(\mathbf{R}) \text {. }
$$

Proof. Let us put $\beta_{\tau}=\omega_{\gamma}$, if $\tau=h(\gamma)$. Then the well-definedness and (d) come from Theorem 2.3. Next we shall show (c). Notice that for each fixed
$a \in \mathbf{R}$, there exists some $E_{a} \in \mathfrak{B}(\mathbf{R})$ such that $\left\{\gamma \in \Gamma_{X} \mid \omega_{\gamma}(B) \leq a\right\}=h^{-1}\left(E_{a}\right)$. Hence

$$
\left\{\tau \in h\left(\Gamma_{X}\right) \mid \beta_{\tau}(B) \leq a\right\}=h\left\{\gamma \in \Gamma_{X} \mid \omega_{\gamma}(B) \leq a\right\}=E_{a} \cap h\left(\Gamma_{X}\right) .
$$

Since an image of a Borel set in a standard space by a Borel map is an analytic set, so $h\left(\Gamma_{X}\right)$ is universally measurable. The rest of the proof easily follows from Theorem 2.2.

We remark that $\beta_{\tau}\left(\tau \in h\left(\Gamma_{X}\right)\right)$ runs all over the set of all $(U, m)$-Gibbs ergodic measures.
2.3. Specific Gibbs measures. First we shall characterize Gibbs measures with total mass on $B_{X}$.

Theorem 2.6. If $a(U, m)$-Gibbs measure $\mu$ on $\left(\Gamma_{X}, \mathscr{C}\right)$ have total mass on $B_{X}$, then it follows that
(a) $S:=\int^{x} \exp (-U(\underline{x} \mid \phi)) m(d \underline{x})<\infty$,
and the explicit form of $\mu$ is given by
(b) $\int_{\Gamma_{x}} f(\gamma) \mu(d \gamma)=S^{-1} \int^{X} \exp (-U(\underline{x} \mid \phi)) f(\underline{x}) m(d \underline{x})$ for all non negative bounded measurable function $f$.
Conversely if (a) holds, then a measure $\mu$ given by $(b)$ is $(U, m)$-Gibbsian with total mass on $B_{X}$.

Proof. As is easily seen, $B_{X}$ is an atom of $\mathscr{C}_{\infty}$, so the measure $\mu$ with total mass on $B_{X}$ must be ergodic. It follows from the martingale convergence theorem that

$$
\begin{equation*}
\int_{B_{X}} f(\delta) \mu(d \delta)=\lim _{n} \frac{\int_{K_{n}}^{K_{n}} \exp \left(-U\left(\underline{x} \mid \gamma \cap K_{n}^{c}\right)\right) f\left(\underline{x} \cdot \gamma \cap K_{n}^{c}\right) m(d \underline{x})}{\int^{K_{n}} \exp \left(-U\left(\underline{x} \mid \gamma \cap K_{n}^{c}\right)\right) m(d \underline{x})} \tag{2.27}
\end{equation*}
$$

for $\mu$-a.e. $\gamma$. However $\gamma \in B_{X}$ implies $\gamma \cap K_{n}^{c}=\phi$ for sufficiently large $n$, so (2.27) is actually,

$$
\begin{equation*}
\int_{B_{X}} f(\gamma) \mu(d \mu)=\lim _{n} \frac{\int^{K_{n}} \exp (-U(\underline{x} \mid \phi)) f(\underline{x}) m(d \underline{x})}{\int^{K_{n}} \exp (-U(\underline{x} \mid \phi)) m(d \underline{x})} . \tag{2.28}
\end{equation*}
$$

By the assumption we have $\mu\left(B_{X, k, l}\right)>0$ for some $k$ and $l$, where $B_{X, k, l}=$ $\left\{\gamma \in B_{X}^{k} \mid \gamma \subset K_{l}\right\}$. Let us put the indicator function of $B_{X, k, l}$ for $f$. Then the numerator under the limit sign in (2.28) becomes

$$
k!^{-1} \int \cdots{ }_{K_{l}^{k}} \int \exp \left(-U\left(\left\{x_{1}, \cdots, x_{l}\right\} \mid \phi\right)\right) m^{l}(d x)
$$

for all $n \geq l$ which is independent of $n$. So we have

$$
\lim _{n} \int^{K_{n}} \exp (-U(\underline{x} \mid \phi)) m(d \underline{x})=S<\infty
$$

and (b) follows directly from (2.28).
Conversely, we claim that (2.5) holds for $\mu$ defined by (b) under the assumption (a). For it we have only to check it for functions $f(\gamma) \chi_{B_{X}^{n}}(\gamma)(n=0,1, \cdots)$. It is obvious for $n=0$, so let $n>0$. Then,

$$
\begin{aligned}
& \int_{N_{K}} \mu(d \gamma) \int^{K} \exp \left(-U(\underline{x} \mid \gamma) f(\underline{x} \cdot \gamma) \chi_{B_{\underline{x}}^{n}}(\underline{x} \cdot \gamma) m(d \underline{x})\right. \\
& =\sum_{l=0}^{n} \int_{\left\{\gamma| | \gamma \cap K\left|=0,\left|\gamma \cap K^{c}\right|=l\right\}\right.} \mu(d \gamma)(n-l)!^{-1} \int_{x \in K^{n-1}} \int \exp \left(-U\left(\left\{x_{1}, \cdots, x_{n-l}\right\}\right.\right. \\
& \left.\left.\mid \gamma \cap K^{c}\right)\right) \cdot f\left(\left\{x_{1}, \cdots, x_{n-l}\right\} \cup\left(\gamma \cap K^{c}\right)\right) m^{n-l}(d x) \\
& =S^{-1} \sum_{l=0}^{n}\{l!(n-l)!\}^{-1} \int_{x \in K^{n-1}} \iint_{y \in\left(K^{c}\right)^{\prime}} \int \exp \left(-U\left(\left\{x_{1}, \cdots, x_{n-l}\right\} \mid\left\{y_{1}, \cdots, y_{l}\right\}\right)\right. \\
& \cdot \exp \left(-U\left(\left\{y_{1}, \cdots, y_{l}\right\} \mid \phi\right)\right) f\left(\left\{x_{1}, \cdots, x_{n-l}, y_{1}, \cdots, y_{l}\right\}\right) m^{n-l}(d x) m^{l}(d y) \\
& =S^{-1} n!^{-1} \sum_{l=0}^{n} C_{l} \int_{x \in K^{n-1}} \iint_{y \in\left(K^{c}\right)^{l}} \int \exp \left(-U\left(\left\{x_{1}, \cdots, x_{n-l}, y_{1}, \cdots, y_{l}\right\} \mid \phi\right)\right) \text {. } \\
& f\left(\left\{x_{1}, \cdots, x_{n-l}, y_{1}, \cdots, y_{l}\right\}\right) m^{n-l}(d x) m^{l}(d y) \\
& =S^{-1} n!^{-1} \int \cdots_{x^{n}} \int \exp \left(-U\left(\left\{z_{1}, \cdots, z_{n}\right\} \mid \phi\right)\right) f\left(\left\{z_{1}, \cdots, z_{n}\right\}\right) m^{n}(d z) \\
& =\int_{B_{\boldsymbol{X}}} f(\gamma) \chi_{B_{\mathbf{X}}^{n}}(\gamma) \mu(d \gamma) .
\end{aligned}
$$

Theorem 2.7. If the potential $U$ is constant, say $U(x \mid \gamma)=-\log a$, then the convex set of all $(U, m)$-Gibbs measures consists of only a Poisson measure $P_{a m}$ with intesity am. ( $P_{a m}$ is of course ergodic by virtue of $0-1$ law.)

Proof. Let $\mu$ be any $(U, m)$-Gibbs measure. Then for each compact set $K$ we have,

$$
\begin{aligned}
\pi_{K} \mu(B) & =\int_{N_{K}} \mu(d \gamma) \int^{K} \exp (-U(\underline{x} \mid \gamma)) \chi_{B}(\underline{x}) m(d \underline{x}) \\
& =\mu\left(N_{K}\right) \sum_{n=0}^{\infty} n!^{-1} a^{n} m_{K, n}\left(B \cap B_{K}^{n}\right)
\end{aligned}
$$

for all $B \in \mathscr{C}_{\boldsymbol{K}}$. So by virtue of Obata's result [1], there exists a Borel probability measure $\lambda$ on $[0, \infty)$ such that $\mu=\int_{0}^{\infty} P_{c m} \lambda(d c)$ in the case $m(X)=\infty$, or $\mu$ is a convex combination of $m_{X, n} / m(X)^{n}$ in the case $m(X)<\infty$. First we shall consider the infinite case. So let us take a set $E \in \mathfrak{B}(\mathbf{R})$ and a function

$$
\rho(\gamma)= \begin{cases}\lim _{n} \frac{1}{n} \sum_{l=1}^{n} \frac{\left|\gamma \cap\left(K_{l+1} \backslash K_{l}\right)\right|}{m\left(K_{l+1} \backslash K_{l}\right)}, & \text { if the limit exists. }  \tag{2.29}\\ 0, & \text { otherwise. }\end{cases}
$$

And we calculate $\int_{\Gamma x} \chi_{E}(\rho(\gamma)) \mu(d \gamma)$ in two ways, noting that $\rho(\gamma)=c$ for $P_{c m}$-a.e. $\gamma$. the first one is,

$$
\int_{\Gamma_{x}} \chi_{E}(\rho(\gamma)) \mu(d \gamma)=\int_{\Gamma_{x}} \chi_{E}(\rho(\gamma)) P_{c m}(d \gamma) \lambda(d c)=\lambda(E),
$$

and the other one is,

$$
\begin{aligned}
& \int_{\Gamma_{X}} \chi_{E}(\rho(\gamma)) \mu(d \gamma)=\int_{0}^{\infty} \lambda(d c) \int_{N_{K}} P_{c m}(d \gamma) \int^{K} a^{|\underline{x}|} \chi_{E}\left(\rho\left(\underline{x} \cdot \gamma \cap K^{c}\right)\right) m(d \underline{x}) \\
= & \exp (a m(K)) \int_{E} \exp (-c m(K)) \lambda(d c) .
\end{aligned}
$$

These show that $c=a$ for $\lambda$-a.e. $c$ and hence $\mu=P_{a m}$. Next we shall consider the finite case. So there exists a non negative sequence $\left\{c_{n}\right\}$ with $\sum_{n=0}^{\infty} c_{n}=1$ such that

$$
\begin{equation*}
\mu=\sum_{n=0}^{\infty} c_{n} m(X)^{-n} m_{X, n} . \tag{2.30}
\end{equation*}
$$

Then for each compact set $K$,

$$
\begin{aligned}
c_{n} & =\mu\left(B_{X}^{n}\right)=\int_{N_{K}} \mu(d \gamma) \int^{K} a^{|\underline{x}|} \chi_{B_{X}^{n}}\left(\underline{x} \cdot \gamma \cap K^{c}\right) m(d \underline{x}) \\
& =\sum_{l=0}^{\infty} c_{l} m(X)^{-l} \int_{N_{K}} m_{X, l}(d \gamma) \int^{K} a^{|\underline{x}|} \chi_{B_{x}^{n}}\left(\underline{x} \cdot \gamma \cap K^{c}\right) m(d \underline{x}) .
\end{aligned}
$$

So we have,

$$
\begin{equation*}
c_{n}=\sum_{l=0}^{n}\left\{(n-l)!m(X)^{l}\right\}^{-1} a^{n-l} m(K)^{n-l} m\left(K^{c}\right)^{l} c_{l} . \tag{2.31}
\end{equation*}
$$

It follows from the mathematical induction for $n$ that

$$
c_{n}=n!^{-1} a^{n} m(X)^{n} c_{0}, \quad \text { and } \quad c_{0}=\exp (-a m(X)),
$$

which follows from the normalizing condition. Thus we have

$$
\mu=\exp (-a m(X)) \sum_{n=0}^{\infty} n!^{-1} a^{n} m_{X, n}=P_{a m} .
$$

Corollary. Let $U$ be a potential defined by $U(x \mid \gamma)=-\log \rho(\gamma)$, where $\rho$ is a function defined by (2.29). If a Borel Radon measure $m$ on $X$ is infinite, then the extremal points of the convex set of all $(U, m)$-Gibbs measures consists of $\left\{P_{c m}\right\}_{c \in[0, \infty)}$. That is, for any $(U, m)$-Gibbs measure $\mu$,

$$
\mu\left(B \cap \rho^{-1}(E)\right)=\int_{E} P_{c m}(B) \rho \mu(d c) \text { for all } B \in \mathscr{C} \text { and } E \in \mathfrak{B}([0, \infty)) \text {. }
$$

Proof. Let $\left\{\omega_{\tau}\right\}_{\tau \epsilon \Gamma_{x}}$ be as in Theorem 2.2. Then the $\mathscr{C}_{\infty_{\infty}}$-measurabilty of $\rho$ implies that for each fixed $\gamma \in \Gamma_{X}$, there exists a constant $a(\gamma)$ such that

$$
\begin{equation*}
\rho(\sigma)=a(\gamma) \tag{2.32}
\end{equation*}
$$

for $\omega_{\gamma}$-a.e. $\sigma$. Thus by the above theorem we have

$$
\begin{equation*}
\omega_{\gamma}=P_{a(\gamma) m} \tag{2.33}
\end{equation*}
$$

for all $\gamma \in \Gamma_{X}$. Especially for each compact set $K$,

$$
\omega_{\gamma}\left(N_{K}\right)=P_{a(\gamma) m}\left(N_{K}\right)=\exp (-a(\gamma) m(K))
$$

It follows that $a(\gamma)$ is also a $\mathscr{C}_{\infty}$-measurable function of $\gamma$, and therefore (2.32) implies that for any $(U, m)$-Gibbs measure $\mu$,

$$
\begin{equation*}
\rho(\gamma)=a(\gamma) \tag{2.34}
\end{equation*}
$$

for $\mu$-a.e. $\gamma$. Thus we have,

$$
\mu\left(B \cap \rho^{-1}(E)\right)=\int_{\rho^{-1}(E)} \omega_{\gamma}(B) \mu(d \gamma)=\int_{\rho^{-1}(E)} P_{\rho(\gamma) m}(B) \mu(d \gamma)=\int_{E} P_{c m}(B) \rho \mu(d c) .
$$

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