SURVIVAL OF NEAREST-PARTICLE SYSTEMS WITH LOW BIRTH RATE¹

By MAURY BRAMSON

University of Minnesota

Nearest-particle systems form a class of continuous-time interacting particle systems on \mathbb{Z} . The birth rate $\beta(l,r)$ at a given site depends on the distances l and r to the nearest occupied sites on the left and right; deaths occur at rate 1. Assume that $b(n) = \sum_{l+r=n} \beta(l,r)$, $2 \le n < \infty$, $b(\infty) = \sum_{l=1}^{\infty} \beta(l,\infty) + \sum_{r=1}^{\infty} \beta(\infty,r)$, is constant. In Liggett [6] the question was posed whether for $b(n) \equiv 1 + \varepsilon$, $2 \le n \le \infty$, with $0 < \varepsilon \le 1$, there are such systems which survive for all t. Here, we answer affirmatively for all such ε and construct a class of examples.

1. Introduction. Nearest-particle systems form a familiar class of continuous-time interacting particle systems on \mathbb{Z} . As time evolves, particles are born and die. At most one particle is permitted at each site at a given time; states can therefore be identified with occupied subsets $A \subset \mathbb{Z}$. The process is Markov with birth rates $\beta(l, r)$ and death rate 1:

(1a)
$$A \to A \cup \{x\}$$
 at rate $\beta(l_A(x), r_A(x))$ for each $x \notin A$

and

(1b)
$$A \to A \setminus \{x\}$$
 at rate 1 for each $x \in A$.

Here

$$l_A(x) = x - \max\{y: y \le x \text{ and } y \in A\},$$

$$r_A(x) = \min\{y: y \ge x \text{ and } y \in A\} - x$$

and $l_A(x)$ or $r_A(x)$ is $+\infty$ if the maximum or minimum is not defined. It is also typically assumed that

$$\beta(\infty,\infty)=0$$

(so that \emptyset is a trap), and

(3)
$$\beta(l,r) = \beta(r,l) \quad \text{for all } 1 \le l, r \le \infty,$$
$$\sum_{l=1}^{\infty} \beta(l,\infty) < \infty.$$

Here, ξ_t will denote such a process. $\xi_t(x) = 0,1$ will denote the state at x, and ξ_t^A the process with initial state A. (The superscript will often be suppressed.) As is typically the case for interacting particle systems, one is interested in

Received July 1987; revised June 1988.

¹Research supported in part by NSF Grant DMS-83-01080.

AMS 1980 subject classification. 60K35.

Key words and phrases. Nearest-particle system, low birth rate, survival.

formulating conditions on $\beta(l, r)$ under which the process survives, that is,

$$P[\xi_t^A \neq \emptyset \text{ for all } t] > 0.$$

One is also frequently interested in the stronger statement

$$\liminf_{t\to\infty} P\big[x\in\xi_t^A\big]>0.$$

If the process does not survive, we will say that it dies out, or becomes extinct. Various classes of nearest-particle systems have been studied; see Liggett [7] for a general survey. Notable examples include reversible nearest-particle systems, first studied by Spitzer [8], and the contact process. The latter has rates

(4)
$$\beta(l, r) = \lambda \quad \text{if } l = r = 1, \\ = \lambda/2 \quad \text{if } l = 1, r > 1 \text{ or } l > 1, r = 1, \\ = 0 \quad \text{otherwise.}$$

This process has been the subject of a considerable amount of study. It was shown in Holley and Liggett [5] that the contact process survives for $\lambda \geq 4$; one can show that for $\lambda \leq 2.36$, the process dies out (Harris [4]).

Set

(5a)
$$b(n) = \sum_{l+r=n} \beta(l,r), \qquad 2 \le n < \infty$$

and

(5b)
$$b(\infty) = \sum_{l=1}^{\infty} \beta(l, \infty) + \sum_{r=1}^{\infty} \beta(\infty, r).$$

b(n) is the total birth rate on an interval of length n between two occupied sites and $b(\infty)$ the birth rate on the union of the two infinite intervals. One can ask how survival depends on $\{b(n)\}$. By using an argument based on that of [5], it was shown in [6] that for $0 < |A| < \infty$, ξ_t^A survives irrespective of the specific rule $\beta(l,r)$ as long as $b(n) \ge 4$ for $2 \le n \le \infty$. On the other hand, it is easy to show that if $b(n) \le 1$, then the process dies out. (One can dominate $|\xi_t^A|$ by the critical binary branching process.) Reversible nearest-particle systems provide examples for survival with $b(n) \equiv b$, b > 2 (Griffeath and Liggett [3] and Liggett [6]). In [6], the question was asked whether there are examples of nearest-particle systems with $b \le 2$ and which survive.

Here, we construct such examples. Specifically, let ${}^{M}\xi_{t}$ denote the nearest-particle system with birth rates given by

(6)
$${}^{M}\beta(l,r) = (1+\varepsilon)/((n-1)\wedge 2M) \quad \text{for } l\wedge r \leq M,$$
$$= 0 \quad \text{for } l\wedge r > M.$$

for $2 \le n \le \infty$ where n = l + r. If $n \le 2M + 1$, then the birth rate is evenly distributed over the interval between occupied sites, whereas if $n \ge 2M + 2$, then the birth rate is $(1 + \varepsilon)/2M$ at sites within distance M of an endpoint and 0 otherwise. Clearly, $b(n) \equiv 1 + \varepsilon$ for $2 \le n \le \infty$. We will show

THEOREM 1. For given $\varepsilon > 0$ and large enough M (depending on ε),

(7)
$$\liminf_{t \to \infty} P\left[x \in {}^{M}\xi_{t}^{A}\right] > 0$$

if $A \neq \emptyset$ and $x \in \mathbb{Z}$.

One may wish to compare ${}^{M}\xi_{t}$ with the uniform long-range contact process in Bramson and Gray [1]. The birth rate there is given by the first line of (6) but with the denominator not truncated by 2M; the initial state is \mathbb{Z} . The analog of (7) is shown for $\varepsilon > 3$.

2. Proof of Theorem 1.

Basic ideas. We begin by introducing several auxiliary processes of ${}^{M}\xi_{t}$. Let ${}^{M}\xi_{t}$ denote the nearest-particle system on $\mathbb Z$ with birth rates

(8)
$${}^{M}\tilde{\beta}(l,r) = {}^{M}\beta(l,r) \quad \text{for } l,r < \infty,$$

$$= 0 \quad \text{otherwise.}$$

The process ${}^{M}\tilde{\xi}_{t}$ has the same birth (and death) rates as ${}^{M}\xi_{t}$, except that births are now only permitted between particles, that is, $b(\infty) = 0$. The leftmost occupied site is therefore an increasing function of t and the rightmost site a decreasing function.

Set

(9)
$${}^{M}I(z) = [2Mz, 2M(z+1)),$$

where $z \in \mathbb{Z}$. We define ${}^{M}\hat{\xi}_{t}$ so that

(10)
$${}^{M}\hat{\xi}_{t} = \bigcup_{z=-\infty}^{\infty} {}^{M}_{z}\hat{\xi}_{t},$$

where $z^{M} \hat{\xi}_{t}$, $t \in [Tn, T(n+1))$, is the nearest-particle system with

and which has birth rates ${}^M\!\tilde{\beta}(l,r)$. (T>0) is large and fixed, and will be specified later. ${}^M_z\xi_t$ has the same birth and death rates as ${}^M\!\xi_t$, but is periodically restarted at times Tn on alternating intervals ${}^M\!I(z)$ with distributions ${}^M\!\xi_{Tn}\cap {}^M\!I(z)$. It is easy to see that

$$_{z}^{M}\hat{\xi}_{t}\subset^{M}I(z)$$

and that for $Tn \leq t_1 \leq t_2 < T(n+1)$,

$$_{z}^{M}\hat{\xi}_{t_{1}} = \varnothing \Rightarrow_{z}^{M}\hat{\xi}_{t_{2}} = \varnothing$$
.

If

$${}^{M}\xi_{0}\cap{}^{M}I(0)\neq\varnothing,$$

then of course ${}^{M}\xi_{0} \cap {}^{M}I(0) \neq \emptyset$. Since the transition mechanism for ${}^{M}\xi_{t}$ is translation invariant, we can assume wlog in proving (7) that (12) holds. We also

point out that ${}^M\!\xi_t$ and ${}^M\!\xi_t$ are attractive. (See [7] for background.) This becomes obvious if one notes that ${}^M\!\beta(l,r)$ is increasing in l and decreasing in r. On account of (8), (10) and (11), if ${}^M\!\xi_t = {}^M\!\xi_t$ at given t, then the birth rate for the first process at this time is at most that of the second process. Together with the attractiveness of ${}^M\!\xi_t$ (or of ${}^M\!\xi_t$), this implies one can couple ${}^M\!\xi_t$ and ${}^M\!\xi_t$ with

$$^{M}\hat{\xi}_{t}\subset ^{M}\xi_{t}.$$

(Refer to [7], page 127, for more detail.)

Define the discrete time process ${}^{M}\zeta_{n}$ on \mathbb{Z} , with ${}^{M}\zeta_{n}(z)=0,1,$ so that ${}^{M}\zeta_{n}(z)=1$ if

(14)
$${}^{M}\xi_{T(n+1)-} \neq \varnothing.$$

That is, ${}^M\zeta_n(z)=1$ if descendents under ${}^M\xi_t$ of particles at ${}^M\xi_{Tn}\cap {}^MI(z)$ are still alive by time T(n+1) – . Note that on account of (11), ${}^M\zeta_n(z)=0$ for $z\neq n \mod 2$.

The basic idea behind the proof of Theorem 1 is to compare ${}^{M}\zeta_{n}$ with a supercritical discrete-time contact process ${}^{p}\eta_{n}$ on (n,z) with $z=n \mod 2$. ${}^{p}\eta_{n}$ is defined so the probability that ${}^{p}\eta_{n}(z)=1$ is $1-(1-p)^{L}$, where L=0,1,2, is the number of occupied sites of ${}^{p}\eta_{n-1}$ among $\{z-1,z+1\}$. $({}^{p}\eta_{n}$ can also be formulated in terms of oriented percolation.) For $p>p_{0}$, $p_{0}<1$ chosen appropriately,

(15)
$$\liminf_{n \to \infty} P\left[z \in {}^{p}\eta_{2n}^{A}\right] > 0$$

if $z = n \mod 2$ and $A \cap 2\mathbb{Z} \neq \emptyset$ (Durrett [2]).

The remainder of this article is devoted to showing that for fixed p < 1, M may be chosen large enough (with appropriate T) so that ${}^p\eta_n$ and ${}^M\zeta_n$ can be coupled together with

$${}^{p}\eta_{n}\subset {}^{M}\zeta_{n}.$$

Theorem 1 is then an easy consequence of (a) (15) and of (13), (14) and (16), which compare ${}^{M}\xi_{t}$ to ${}^{M}\hat{\xi}_{t}$, ${}^{M}\hat{\xi}_{t}$ to ${}^{M}\zeta_{n}$ and ${}^{M}\zeta_{n}$ to ${}^{p}\eta_{n}$, and (b) the observation that

$$\inf P \Big[x \in {}^{M} \xi_{t} | {}^{M} \xi_{T(n+1)} = B \Big] > 0,$$

where inf is taken over all B satisfying $B\cap^M I(z)\neq\varnothing$, all $t\in [T(n+2),T(n+4)),\ x\in^M I(z)\cup^M I(z+1)$ and all n. Part (b) follows from the attractiveness of ${}^M\xi_t$, and implies that if $z\in^p\eta_n$ (and hence ${}^M\xi_{T(n+1)}\cap^M I(z)\neq\varnothing$), then $x\in^M I(z)\cup^M I(z+1)$ is likely to be occupied after allowing time T for the process ${}^M\xi_t$ to spread.

Notation. We justified the coupling ${}^M\!\xi_t\subset{}^M\!\xi_t$ in (13) by using the attractiveness of ${}^M\!\xi_t$ and the lower birth rate of ${}^M\!\xi_t$. One may also explicitly construct ${}^M\!\xi_t$ from ${}^M\!\xi_t$ as follows: At each site $(t,x)\in\mathbb{R}^+\times\mathbb{Z}$ where a birth for ${}^M\!\xi_t$ occurs, introduce an independent [0,1] uniformly distributed random variable $W_{t,x}$. If the birth rate for ${}^M\!\xi_t$ at (t,x) is greater than $W_{t,x}$, we say that a birth for ${}^M\!\xi_t$ occurs at (t,x); otherwise, no birth occurs there. Deaths occur at rate 1 for each

process, so we say that deaths occur at occupied sites of ${}^M\!\xi_t$ whenever they occur for ${}^M\!\xi_t$. This procedure may be carried out inductively since $|{}^M\!\xi_t| < \infty$. In this setting, ${}^M\!\xi_t \subset {}^M\!\xi_t$ clearly holds. We denote by ${}^M\!\mathcal{F}_t$ the σ -algebra generated by ${}^M\!\xi_s$ and $W_{s,x}$ up to time t and by the independent procedure used to construct ${}_t V_t$ in (22). The construction has the advantage of allowing us to construct simultaneously other subprocesses of ${}^M\!\xi_t$ [notably ${}_t \xi_t$ in (22)] by employing $W_{t,x}$. They will all be measurable with respect to ${}^M\!\mathcal{F}_t$ and satisfy couplings analogous to (13). Let L_t and R_t be the leftmost and rightmost particles of the process ${}^M\!\xi_t$ from (8) constructed in this manner. On account of (8), ${}^M\!\xi_t$ evolves independently of the births and deaths for ${}^M\!\xi_t$ and of values of $W_{t,x}$ which occur outside of $[L_t, R_t]$.

We require the use of several σ -algebras. (To simplify the notation a little, we omit here and later on the superscript M, which will be implicit.) Set

(17)
$$\mathscr{G}_n = \sigma(\zeta_m, 0 \le m \le n),$$

and for $t \in [Tn, T(n+1))$, set

(18)
$$\hat{\mathcal{F}}_t = \sigma(\xi_s, s \le Tn) \lor \sigma(\hat{\xi}_s, Tn \le s \le t),$$

$$\hat{\mathcal{F}}^t = \sigma(\hat{\xi}_s, t \le s < T(n+1)).$$

It is easy to check that

$$\mathscr{G}_n \subset \hat{\mathscr{F}}_t \vee \hat{\mathscr{F}}^t$$

for $t \in [Tn, T(n+1))$.

We will also use the following notation. Set v = (k, n, z, i), where k = 0, 1, ..., K - 1, $z = n \mod 2$ and $i = \pm 1$; the symbol v will henceforth be reserved for this 4-tuple. Introduce the times

(19)
$$S_{k,n} = Tn + (1 + t_h)k + 1, T_{k,n} = Tn + (1 + t_h)(k + 1);$$

 $t_h>0$ will be specified in (36), K will be specified after (49) and T will satisfy $T=K(1+t_h).$ Note that

$$Tn < S_{0,n} < T_{0,n} < S_{1,n} < \cdots < T_{K-1,n} = T(n+1).$$

We define

(20)
$$_{v}D = \{\omega : \exists y_{1}, y_{2} \text{ with } y_{i} \in \xi_{S_{h,n}} \cap I(z+i), y_{2} - y_{1} \geq M/2 \}.$$

 $_vD$ consists of those realizations which at $t=S_{k,n}$ have a pair of particles at sites y_1, y_2 in I(z+i) which are "far apart." Consider tuples $v_j=(k_j,n,z_j,i_j)$, where we equate $(k_j,n,z_j,1)$ and $(k_j,n,z_j+2,-1)$. One can for appropriate $\varepsilon_1>0$ choose subsets $_vE\subset_vD$ so that for distinct $v_i,\ j=1,\ldots,J$,

(21)
$$P\left[\bigcap_{j=1}^{J} v_{j} E | \mathscr{G}_{n}\right](\omega) = \varepsilon_{1}^{J}$$

if $\zeta_n(z_j) = 1$. That is, $v_1 E, \ldots, v_J E$ are independent under $\zeta_n(z_1) = \cdots = \zeta_n(z_J) = 1$ and \mathscr{G}_n . For $\omega \in E$ denote by v_1, v_2 occupied positions v_1, v_2 such as

in (20). $_vY_1$, $_vY_2$ may be chosen so that $_vY_2-_vY_1$ is uniformly distributed over [M/2,M] and is independent of everything else. We will justify this and (21) in Proposition 1. Let $_v\tilde{\xi}_t$, $t\in[S_{k,\,n},T(n+2))$, denote the process defined on $_vE$ with

(22)
$${}_{v}\tilde{\xi}_{S_{k-1}} = \left\{ {}_{v}Y_{1}, {}_{v}Y_{2} \right\}$$

and birth rate given by (8). Recall that $b(\infty) = 0$ here. Using the construction given before (17), ξ_t is defined on the same space as ξ_t with

(23)
$${}_{v}\tilde{\xi}_{T(n+1)} \subset \xi_{T(n+1)} \cap I(z+i) = {}_{z+i}\hat{\xi}_{T(n+1)}.$$

Since the birth rates of both $_{0}\tilde{\xi}_{t}$ and $_{z+i}\hat{\xi}_{t}$ are defined by (8), one has

(24)
$${}_{v}\tilde{\xi}_{t} \subset_{z+i}\hat{\xi}_{t} \text{ for } t \in [T(n+1), T(n+2)).$$

[Note that $\hat{\xi}_t \cap I(z+i) = \emptyset$ for $t \in [Tn, T(n+1))$.] We will (after Lemma 1) choose M large enough so that $x, y \in {}_v \hat{\xi}_t$ with |x-y| = 1 typically does not occur until $t > S_{k,n} + 2T$.

Define

(25)
$${}_{v}F = \left\{ \omega \colon {}_{v}\tilde{\xi}_{T_{k-1}} \neq \varnothing \right\}$$

and

$${}_{u}F = \bigcup_{k=0}^{K-1} {}_{v}F,$$

where u=(n,z,i). (The symbol u will henceforth be reserved for this 3-tuple.) $_vF$ is the subset of $_vE$ over which the process $_v\tilde{\xi}_t$ started at time $S_{k,n}$ is still alive by $T_{k,n}$. $_uF$ may be thought of as the set of those realizations where the process ξ_t has managed to "take hold" in I(z+i) over one of its "trial periods" $[S_{k,n},T_{k,n}]$. (Lemma 2 will make this more precise.) It will follow from Proposition 2 that $P[_uF|\mathcal{G}_n] \sim 1$ for $\zeta_n(z) = 1$ if K is chosen appropriately large.

Finally, let

(27)
$$\kappa(u) = \min\{k \colon \omega \in F\}$$

 $[\kappa(u) = \infty \text{ if } \omega \notin {}_{u}F] \text{ and }$

(28)
$$v_{\kappa} = (\kappa(n, z, i), n, z, i).$$

Then define

(29)
$${}_{u}G = \left\{\omega \colon {}_{v_{s}}\tilde{\xi}_{T(n+2)} \neq \varnothing\right\}.$$

 $[S_{\kappa,\,n},\,T_{\kappa,\,n}]$ is the interval over which ξ_t first "takes hold" and ${}_uG$ is the subset of ${}_uF$ over which ${}_{v_\kappa}\tilde{\xi}_t$ is still alive by time T(n+2) – . [Recall that $Tn < T_{k,\,n} \le T(n+1)$.] It will follow from Proposition 3 that $P[{}_uG|\mathscr{G}_n \lor \sigma({}_uF)] \sim 1$ for $\omega \in {}_uF$.

Note that the processes $v_j \tilde{\xi}_t$ used in (25) to construct $v_j F$ are defined over disjoint sets

$$I(z_j + i_j) \times [S_{k_i, n}, T_{k_i, n}]$$

in space-time for distinct v_j , $j=1,\ldots,J$. On account of (21) and the independence of $v_jY_2-v_jY_1$, this implies that distinct v_jF , $j=1,\ldots,J$, are independent under $\zeta_n(z_j)=1$ and \mathscr{G}_n . Similarly, u_jF and u_jG are defined over disjoint sets $I(z_j+i_j)$ for distinct u_j , $u_j=(n,z_j,i_j)$, and so u_jF , resp. u_jG , $j=1,\ldots,J$, are also independent under $\zeta_n(z_j)=1$ and \mathscr{G}_n .

Demonstration of (16). Our approach will be to show that for fixed p, if M is large enough (and T is chosen appropriately), then

(30)
$$P[{}_{u}G|\mathscr{G}_{n}](\omega) \geq 1 - (1-p)^{2}$$

for $z \in \zeta_n$, where u = (n, z, i). By (24) and (14),

(31)
$${}_{u}G \subset \left\{\omega \colon_{z+i} \hat{\xi}_{T(n+2)-} \neq \varnothing\right\} = \left\{\omega \colon z+i \in \zeta_{n+1}\right\}.$$

Recall that \mathscr{G}_n gives the entire history of ζ_m up to time n. Note also that the process ${}^p\eta_n$ defined above (15) is attractive. Using (30)–(31) and the independence of ${}_{u_j}G$, $j=1,\ldots,J$, under $\zeta_n(z_j)=1$ and \mathscr{G}_n , one can therefore couple ${}^p\eta_n$ and ζ_n as in (16) with ${}^p\eta_n \subset \zeta_n$.

To obtain (30), we compute the probabilities of $_vE$ (Proposition 1), $_vF|_vE$ and $_uF|_vE$ (Proposition 2) and $_uG|_uF$ (Proposition 3) under appropriate conditioning. In order to ensure independence where needed, some care must be exercised in the choice of the σ -algebras.

In Proposition 1, we compute a lower bound for the conditional probability of $_vE$ under $z\in\zeta_n$. The basic point is that for $z\in\zeta_n$, there is a positive probability that ξ_t spreads to I(z+i) over the time interval $[S_{k,\,n}-1,S_{k,\,n}]$ as in the right-hand side of (20), and that this probability and the manner in which ξ_t spreads do not depend on $\hat{\mathscr{F}}_{T(n+1)-}$. (Recall that $T_{k,\,n-1}=S_{k,\,n}-1$.) Since I(z) has length 2M, which is on the same scale as the birth rates for ξ_t given in (6), this probability also does not depend on M. Here

$$_{z}\hat{R}_{t} = \max\{x \colon x \in_{z}\hat{\xi}_{t}\},$$

$$_{z}\hat{L}_{t} = \min\{x \colon x \in_{z}\hat{\xi}_{t}\}$$

and

$$_{n,z}B = \left\{\omega \colon {}_{z}\hat{\xi}_{T(n+1)} \neq \varnothing\right\} = \left\{\omega \colon z \in \zeta_{n}\right\}.$$

PROPOSITION 1. For $\omega \in_{n,z} B$ and i = -1,1,

(32)
$$P[_{v}D|\mathscr{G}_{n}](\omega) \geq \varepsilon_{1}$$

for appropriate $\varepsilon_1 > 0$. Subsets ${}_vE \subset {}_vD$ may be chosen so that (a) if $\omega \in {}_{n,z_j}B$ for distinct $v_j = (k_j, n, z_j, i_j), j = 1, \ldots, J$, then

(33)
$$P\left[\bigcap_{j=1}^{J} v_{j} E | \mathscr{G}_{n}\right](\omega) = \varepsilon_{1}^{J}$$

and (b) on $_vE \cap_{n,z}B$, $\xi_{S_{k,n}} \cap I(z+i)$ contains occupied positions $_vY_1,_vY_2$ for which $_vY_2 - _vY_1$ is uniformly distributed on [M/2, M] under \mathscr{G}_n , and for which

on $\bigcap_{j=1}^{J} (v_j E \cap_{n, z_j} B)$, $v_j Y_2 - v_j Y_1$, $j = 1, \ldots, J$, are mutually independent under \mathscr{G}_n .

PROOF. By (6), the rate at which a birth of ξ_t occurs in $[{}_z\hat{R}_t + M/2, {}_z\hat{R}_t + M]$ is greater than 1/2. (For bookkeeping purposes, allow births to occur at already occupied sites.) Such a birth, later births to the right of this point and deaths in $(\hat{R}_{T_{k-1,n}}, (2z+4)M)$ occur at rates independent of $\hat{\mathscr{F}}_{T(n+1)-}$. Also note that $\omega \in {}_{n,z}B$ implies

$$_{z}\hat{\xi}_{T_{k-1,n}}\neq\emptyset$$
.

It is therefore not difficult to check that if $\omega \in_{n,z} B$, then for appropriate $c_1 > 0$,

(34)
$$P\Big[\exists y_j, j = 1, ..., J: y_j \in \xi_{S_{k,n}}, y_1 \in [2zM, (2z+3)M), \\ y_{j+1} - y_j \in [M/2, M] | \hat{\mathscr{F}}_{T(n+1)-} \Big] \ge (c_1/J)^J.$$

For J=5, there are at least two indices $j_1, j_2, j_2=j_1+1$, for which $y_{j_1}, y_{j_2} \in I(z+1)$. Since $\mathscr{G}_n \subset \hat{\mathscr{F}}_{T(n+1)-}$, this shows (32) for i=1 with $\varepsilon_1=(c_1/5)^5$. The reasoning for i=-1 is analogous.

Recall that ξ_t is attractive; one can therefore choose y_{j_1} , y_{j_2} so that their difference is uniformly distributed over [M/2, M] under $\hat{\mathscr{F}}_{T(n+1)-}$ (and hence \mathscr{G}_n). A subset $E \subset D$ can be chosen with

$$P[_{v}E|\mathscr{G}_{n}](\omega)=\varepsilon_{1}$$

on $_{n,\,z}B$, so that this difference remains uniformly distributed on $_vE\cap_{n,\,z}B$. Now assume that v_j are distinct [where we equate $(k_j,\,n,\,z_j,1)$ and $(k_j,\,n,\,z_j+2,\,-1)$]. The births as in (34) needed for $_{v_j}D$ occur over disjoint intervals of the form

$$(R_{S_{k,n}},(2z+4)M) \times [T_{k-1,n},S_{k,n}]$$

and

$$\left[\left(2z-2\right)\!M,L_{S_{k,n}}\right)\times\left[\left.T_{k-1,\,n},S_{k,\,n}\right.\right]$$

in space-time. A little thought therefore shows that the subsets $_vE\subset_vD$ can be chosen so that under $\omega\in_{n,\,z_j}B,\ j=1,\ldots,\,J,$ (33) holds and the corresponding differences $_{v_j}Y_2-_{v_j}Y_1$ are independent under $\mathscr{F}_{T(n+1)-}$, and hence \mathscr{G}_n . \square

Once particles spread from I(z) to neighboring blocks I(z-1) and I(z+1), they start to reproduce; we need to know that their descendents survive with some probability. To make this more precise, we introduce the continuous-time birth-death process X_t on $0, 1, 2, \ldots$, with transitions

(35)
$$k \to k+1 \quad \text{at rate } (1+\varepsilon)(k-1)^+,$$
$$\to k-1 \quad \text{at rate } k,$$

and $X_0 = 2$. Comparison with a supercritical binary branching process shows that:

LEMMA 1. $P[X_t > 0 \text{ for all } t] = \varepsilon_2 > 0.$

Choose t_h so that

$$(36) P[X_t > 0 \forall t | X_{t_h} > 0] \ge h,$$

where $h \ge 1/2$. We will need $h \sim 1$ to obtain (30) with $p \sim 1$. Also, introduce the stopping time

(37)
$$\tau^{A} = \inf\{t: |x - y| = 1 \text{ for some } x, y \in {}^{M} \xi_{t}^{A}\},$$

where ${}^{M}\tilde{\xi}_{t}^{A}$ is the process defined by (8). It follows easily from (6) that for $A = \{0, m\}$ and fixed t,

(38)
$$\sup_{m \ge M/2} P[\tau^A \le t] \to 0 \quad \text{as } M \to \infty.$$

Choose M_h so that for $2m \ge M \ge M_h$,

(39)
$$P[\tau^A \le 2T] \le \varepsilon_2(1-h).$$

[T will be specified later on; recall that $(1 + t_h)$ divides T.] $|\tilde{\xi}_t^A|$ can be coupled with X_t so that

$$|\tilde{\xi}_t^A| = X_t \text{ for } t \leq T^A.$$

One therefore obtains from Lemma 1, (36) and (39) that:

LEMMA 2. For $A = \{0, m\}$ with $2m \ge M \ge M_h$,

(a)
$$P\big[\tilde{\xi}_t^A \neq \varnothing \text{ for } t \leq t_h\big] \geq \varepsilon_2/2,$$

(b)
$$P\Big[\tilde{\xi}_t^A \neq \varnothing \text{ for } t \leq 2T | \tilde{\xi}_{t_h}^A \neq \varnothing \Big] \geq 2h - 1.$$

By utilizing Proposition 1 and Lemma 2, we now compute lower bounds on the conditional probability of $_{\nu}F$.

Proposition 2. For $z \in \zeta_n$,

$$(40) P\lceil_{n} F|\mathscr{G}_{n}\rceil \geq 1 - (1 - \varepsilon_{1} \varepsilon_{2}/2)^{K},$$

where \mathscr{G}_n and $_uF$ are as in (17) and (26) and $K = T/(1 + t_h)$.

Proof. By Proposition 1,

$$(41) P[_{v}E|\mathscr{G}_{n}](\omega) \geq \varepsilon_{1}$$

for $\omega \in_{n,z} B$. Also, by the Markov property and Lemma 2(a),

(42)
$$P[{}_{v}F|\mathscr{F}_{S_{k,n}}](\omega) = P[{}_{v}F|\sigma({}_{v}Y_{2} - {}_{v}Y_{1})](\omega) \\ \geq \varepsilon_{2}/2$$

for $\omega \in {}_{v}E$. Because of (8), it is not difficult to see that the random variable on the left-hand side of (42) is independent of $\hat{\mathscr{F}}^{S_{k,n}}$. Therefore,

(43)
$$P\left[{}_{v}F|\mathscr{F}_{S_{k,n}}\vee\hat{\mathscr{F}}^{S_{k,n}}\right](\omega)\geq\varepsilon_{2}/2$$

for $\omega \in E$. Since

$$\mathscr{G}_n \subset \mathscr{F}_{S_{k,n}} \vee \hat{\mathscr{F}}^{S_{k,n}},$$

it follows from (41) and (43) that

(44)
$$P[_{n}F|\mathscr{G}_{n}](\omega) \geq \varepsilon_{1}\varepsilon_{2}/2$$

on $\{\omega: z \in \zeta_n\} = {}_{n,z}B$. But, ${}_{(k,\,n,\,z,\,i)}F$ are independent for $k=0,1,\ldots,\,K-1$. Therefore,

(45)
$$P[_{n}F|\mathscr{G}_{n}](\omega) \geq 1 - (1 - \varepsilon_{1}\varepsilon_{2}/2)^{K}$$

for $z \in \zeta_n$. \square

Using Lemma 2, we can also compute lower bounds for ${}_{u}G|_{u}F$.

Proposition 3. For $\omega \in {}_{\mu}F$,

(46)
$$P[_{u}G|\mathscr{G}_{n}\vee\sigma(_{u}F)](\omega)\geq 2h-1.$$

PROOF. By the Markov property and Lemma 2(b),

(47)
$$P\left[{}_{u}G|\mathscr{F}_{T_{k,n}}\vee\sigma(\kappa)\right](\omega) = P\left[{}_{u}G\cap\{\kappa=k\}|\mathscr{F}_{T_{k,n}}\vee\sigma(\kappa)\right](\omega) \geq 2h-1$$

for

$$(48) \qquad \qquad \omega \in {}_{\nu}F \cap \{\kappa = k\} = {}_{\nu}F.$$

The random variable in (47) depends only on events in I(z+i) and is therefore independent of $\mathscr{F}^{T_{k,n}}$. Since

$$\mathscr{G}_n \subset \mathscr{F}_{T_{k,n}} \vee \hat{\mathscr{F}}^{T_{k,n}},$$

this implies that

$$P[{}_{u}G|\mathscr{G}_{n}\vee\sigma(\kappa)](\omega)\geq 2h-1$$

for ω as in (48). Since the inequality does not depend on k (as long as $k < \infty$) and since $_{u}F = \bigcup_{k=0}^{K-1} \{\kappa = k\}$, (46) follows. \square

By combining Propositions 2 and 3, one sees that for $z \in \zeta_n$,

(49)
$$P[{}_{u}G|\mathscr{G}_{n}](\omega) \geq (2h-1)[1-(1-\varepsilon_{1}\varepsilon_{2}/2)^{K}].$$

Choosing h close enough to 1 and K large enough, one obtains

(50)
$$P[{}_{\nu}G|\mathscr{G}_{n}](\omega) \geq 1 - (1-p)^{2}$$

for fixed p < 1. One can now specify t_h [using (36)] and T, with $T = K(1 + t_h)$.

(39) tells us how to choose M. (50) therefore gives us the desired inequality (30). This completes the proof of Theorem 1.

We close with a few observations. First note that the assumption $b(n) \equiv 1 + \varepsilon$ is not essential. One can instead demonstrate Theorem 1 for processes satisfying $\liminf_{n \to \infty} b(n) \ge 1 + \varepsilon$ and $b(\infty) \ge \varepsilon$. [Replacing $1 + \varepsilon$ by ε for $b(\infty)$ only changes the constant ε_1 in Proposition 1.] The basic proof goes through, although one needs to condition more carefully when the process is not attractive. Also, note that for $|A| = \infty$, the comparison we have made with the contact process shows that for each t_0 ,

(51)
$$P\left[x \in {}^{M}\xi_{t}^{A} \text{ for some } t \geq t_{0}\right] = 1.$$

One can instead condition on nonextinction rather than assuming $|A| = \infty$.

If one wishes, one can with somewhat less effort than needed for Theorem 1 exhibit processes ξ_t^A , with $b(n) \equiv 1 + \varepsilon$, for which

(52)
$$P[\xi_t^A \neq \emptyset \text{ for all } t] > 0,$$

 $0<|A|<\infty$. For these processes, $\beta(l,\infty)=\beta(\infty,l)$ has a "very fat tail" (with not even fractional moments) and $\beta(l,r)=0$ except near l=r. The basic idea for demonstrating (52) is to use the tail behavior of $\beta(l,\infty)$ to obtain births at R_t+m far to the right of the rightmost particle R_t of ξ_t , with the rate only decreasing slowly as $m\to\infty$. As in Lemma 2, for large m the number of progeny between particles at R_t and R_t+m can be compared with the birth-death process X_t for long times. As m increases and this time increases, one has more time to look for births much further to the right of R_t+m . Iterating in this manner, one can obtain (52). These processes, while providing examples of survival, are not as convincing as M_{ξ_t} since (1) survival is only shown on sets $A_n \subset M_{\xi_{t,n}}$ min $\{x: x \in A_n\} \to \infty$ as $t_n \to \infty$, and (2) these processes, with their fat tails, are scarcely representative members of the class with $b(n) \equiv 1 + \varepsilon$.

REFERENCES

- BRAMSON, M. and GRAY, L. (1981). A note on the survival of the long-range contact process. Ann. Probab. 9 885-890.
- [2] DURRETT, R. (1984). Oriented percolation in two dimensions. Ann. Probab. 12 999-1040.
- [3] GRIFFEATH, D. and LIGGETT, T. M. (1982). Critical phenomena for Spitzer's reversible nearest particle systems. Ann. Probab. 10 881-895.
- [4] HARRIS, T. E. (1974). Contact interactions on a lattice. Ann. Probab. 2 969-988.
- [5] HOLLEY, R. and LIGGETT, T. M. (1978). The survival of contact processes. Ann. Probab. 6 198-206.
- [6] LIGGETT, T. M. (1984). Finite nearest particle systems. Z. Wahrsch. verw. Gebiete 68 65-73.
- [7] LIGGETT, T. M. (1985). Interacting Particle Systems. Springer, New York.
- [8] SPITZER, F. (1977). Stochastic time evolution of one-dimensional infinite-particle systems. Bull. Amer. Math. Soc. 83 880-890.

DEPARTMENT OF MATHEMATICS UNIVERSITY OF WISCONSIN MADISON, WISCONSIN 53706