# STRICT CONVEXITY OF THE LIMIT SHAPE IN FIRST-PASSAGE PERCOLATION 

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## Abstract

Sufficient conditions are given for the strict convexity of the limit shape in standard firstpassage percolation. These conditions involve (1) asymptotic "straightness" of the geodesics, and (2) existence of mean-zero limit distributions for the first-passage times.

## 1 Introduction

The limit shape of the infected set in standard first-passage percolation (see [5] and [7] for an introduction to the subject) is well known to be convex. It is widely believed that, with few exceptions (perhaps only when the distribution $F$ of the passage time of an edge is supported by two points), the limit set $\mathcal{B}$ is strictly convex in the sense that line segments connecting any two points of $\mathcal{B}$ are entirely contained, except for their endpoints, in the interior of $\mathcal{B}$. Strict convexity of the limit set plays an important role in certain questions regarding the evolution of multi-type stochastic competition models - see, for instance, [8]. Nevertheless, strict convexity has not yet been established for any distribution $F$, not even the exponential distribution.
It has been observed by Durrett and Liggett [2] that when the edge passage time distribution $F$ is supported by two points, the limit shape need not be strictly convex: its boundary may include nontrivial sections of hyperplanes. This situation seems to be atypical in at least two other important respects. First, time-minimizing paths from the origin to distant points need not be (to first order) straight. Second, the first-passage distributions are, for distant points in directions along the "flat spots", highly concentrated: In particular, if $T(x)$ is the first passage to the point of $\mathbb{Z}^{d}$ nearest to $x \in \mathbb{R}^{d}$ and $a$ is the smaller point in the support of $F$, then for certain directions $u$, as $n \rightarrow \infty$,

$$
\begin{equation*}
T(n u)-n a \xrightarrow{\mathcal{D}} G \tag{1}
\end{equation*}
$$

for a certain distribution $G$ independent of the direction $u$.

[^0]That the latter two properties are indeed closely related to the strict convexity of the the limit shape is further suggested by more recent results of Newman, Piza, and Licea [13], [12], [10], [9]. They have shown that, under certain additional hypotheses on $F$, if the limit shape is not only strictly convex but uniformly curved then (a) the variance of $T(n u)$ must grow at least logarithmically in $n$ (see also [14] for another proof of this in the case where $F$ is an exponential distribution), and (b) geodesic segments (time-minimizing paths) between distant points of $\mathbb{Z}^{d}$ must be (to first order) straight line segments.
The purpose of this note is to show that properties closely related to (a) and (b) above are in fact sufficient for the strict convexity of the limit shape $\mathcal{B}$. Assume that the passage times for different edges of the integer lattice $\mathbb{Z}^{d}$ are independent and identically distributed, and that their common distribution $F$ is supported by $[0, \infty)$, is nonatomic, and has finite exponential moment $\int e^{\gamma t} d F(t)$ for some $\gamma>0$. These assumptions ensure that time-minimizing paths between points $x, y$ of $\mathbb{Z}^{d}$ are unique, and that the Shape Theorem holds, with a compact, convex limit shape $\mathcal{B}$ containing the origin in its interior - see Theorems 1.7 and 1.15 of [7]. For each $x \in \mathbb{R}^{d} \backslash\{0\}$ define

$$
\begin{equation*}
\mu(x)=\lim _{n \rightarrow \infty} \frac{E T(n x)}{n} \tag{2}
\end{equation*}
$$

to be the inverse infection speed in direction $x$. The limit shape $\mathcal{B}$ consists of the origin and those $x \neq 0$ such that $\mu(x) \leq 1$.
Let $u$ be a fixed nonzero vector in $\mathbb{R}^{d}$, and let $\mathcal{L}_{u}$ be the ray through $u$ emanating from the origin.

Hypothesis 1 For any convex cone $\mathcal{A}$ of $\mathbb{R}^{d}$ containing the vector $u$ in its interior, and for each $\delta>0$ there exists $R=R(\delta, \mathcal{A})<\infty$ such that the following is true: For each point $v \in \mathbb{Z}^{d} \cap \mathcal{A}$ at distance $\leq 2$ from the line $\mathcal{L}_{u}$, the probability that the time-minimizing path from the origin to $v$ is contained in $\mathcal{A} \cup\{x:\|x\| \leq R\}$ is at least $1-\delta$.

In the following, the only convex cones considered will be those whose intersections with the unit sphere $\Sigma_{d-1} \subset \mathbb{R}^{d}$ are spherical caps centered at the point $\mathcal{L}_{u} \cap \Sigma_{d-1}$; such cones will be called spherical cones. Two spherical cones will be said to have the same aperture if the spherical caps that determine them are congruent.

Hypothesis 2 There exist a mean-zero probability distribution $G_{u} n$ the real line and a scalar sequence $a(n) \rightarrow \infty$ such that as $n \rightarrow \infty$,

$$
\begin{equation*}
\frac{T(n u)-n \mu(u)}{a(n)} \xrightarrow{\mathcal{D}} G_{u} . \tag{3}
\end{equation*}
$$

Theorem 1 Let $u$, v be linearly independent vectors in $\mathbb{R}^{d}$. If Hypotheses 1 - 2 hold for both $u$ and $v$ then for each $\alpha \in(0,1)$,

$$
\begin{equation*}
\mu(\alpha u+(1-\alpha) v)<\alpha \mu(u)+(1-\alpha) \mu(v) \tag{4}
\end{equation*}
$$

Theorem 1 clearly implies that if the hypotheses 1 and 2 hold for all $u$ in a dense set of unit vectors then the limit shape $\mathcal{B}$ is strictly convex. Unfortunately, the convergence in law of the rescaled first-passage times $T(n u)$ has not been established for any distribution $F$, nor is it at all clear that even if such convergence holds the limit distribution $G_{u}$ should have mean zero. It has been conjectured that the variance of $T(n u)$ is of order $n^{2 / 3}$ at least for continuous distributions $F$ with finite exponential moments, and it is now known [1] that at
least for certain distributions the variance of $T(n u)$ grows sublinearly with $n$. This suggests that a limiting distribution, if it exists, may not be Gaussian, but rather a Tracy-Widom distribution. See Johansson [6] for results in this direction on what seems to be a closely related model.

## 2 Proof of Theorem 1

The proof, like that of [2], makes use of a derived oriented percolation process. This oriented percolation is always 2 -dimensional, regardless of the ambient dimension of the first-passage percolation, and, in the general case, is a combined edge+site percolation: Bernoulli- $p_{v}$ random variables are attached to the vertical edges, Bernoulli- $p_{h}$ random variables are attached to the horizontal edges, and Bernoulli- $p_{s}$ random variables are attached to the vertices of the lattice $\mathbb{Z}^{2}$; these random variables are mutually independent. An open path in the oriented percolation is a finite alternating sequence $e_{i}, v_{i}$ of edges and vertices, all of which are "open" (that is, the attached Bernoulli random variables all take the value 1 ), with each vertex $v_{i}$ incident to the edges $e_{i}$ and $e_{i+1}$, and such that the sequence of vertices visited by successive edges in the path is nondecreasing with respect to the usual partial order on $\mathbb{Z}^{d}$. If there is an infinite open path starting at the origin, then percolation is said to occur. When percolation occurs, an infinite open path starting at the origin may visit the diagonal infinitely often; denote this event by $\Omega$.

Proposition 3 There is a critical value $p_{c}<1$ such that if $p_{v}, p_{h}$ and $p_{s}$ all exceed $p_{c}$ then $P(\Omega)>0$.

See [3] for the proof in the case where $p_{v}=p_{h}$ and $p_{s}=1$; the extension stated here may be proved in much the same manner.
Assume now that the hypotheses of Theorem 1 are in force. For any path $\gamma$ in the integer lattice, let $\tau(\gamma)$ denote the time required to traverse $\gamma$ (that is, the sum of the edge passage times over all edges in $\gamma$ ). The derived oriented percolation process on $\mathbb{Z}^{2}$ will be defined by making an appropriate embedding of $\mathbb{Z}^{2}$ into the ambient space $\mathbb{Z}^{d}$ and using the travel times $\tau(\gamma)$ along appropriate paths to determine which edges and vertices of $\mathbb{Z}^{2}$ will be open. The embedding $\varphi: \mathbb{Z}^{2} \rightarrow \mathbb{Z}^{d}$ will be a linear mapping that sends $(1,0)$ to $x$ and $(0,1)$ to $y$ for certain linearly independent vectors $x, y \in \mathbb{Z}^{d}$. Along with this embedding there is a natural associated mapping, also denoted by $\varphi$, of the edges of the lattice $\mathbb{Z}^{2}$ to line segments in $\mathbb{R}^{d}$. The mapping $\varphi$ of vertices and edges will be augmented by the following geometric objects: (A) Each embedded vertex $\varphi(v)$ will be surrounded by a closed ball $B_{v}$ of radius $R$ centered at $\varphi(v)$; the constant $R$ will be the same for all vertices, and small enough that no two distinct balls $B_{v}, B_{v^{\prime}}$ intersect. (B) Each embedded edge $\varphi(e)$ will be enclosed in a region $D_{e}$ that consists of the intersection of two congruent spherical cones, one based at each endpoint of $\varphi(v)$. The cones used in this construction will all have the same aperture, and the aperture will be sufficiently small that no two of the regions $D_{e}, D_{e^{\prime}}$ overlap except possibly at the vertices $\varphi(v)$. See Figure 1 for a depiction of a portion of the decorated lattice.
Bernoulli random variables $X_{e}, X_{v}$ indexed by the edges $e$ and vertices $v$ are now defined as follows: For each edge $e$ of $\mathbb{Z}^{2}$, denote by $\psi(e)$ the segment of the edge $\varphi(e)$ exterior to the balls $B_{v}$ and $B_{v^{\prime}}$, where $v, v^{\prime}$ are the vertices incident to $e$, and let $z=z_{e}, z^{\prime}=z_{e}^{\prime}$ be the lattice points closest to the endpoints of the segment $\psi(e)$. Define $Y_{e}$ to be the minimum traversal


Figure 1:
time $\tau(\gamma)$ among all paths $\gamma$ from $z$ to $z^{\prime}$ that lie entirely in the region $D_{e}$. Set

$$
\begin{array}{lll}
X_{e}=1 & \text { if } \quad Y_{e} \leq t_{H} & \text { and } e \text { is horizontal; } \\
X_{e}=1 & \text { if } \quad Y_{e} \leq t_{V} & \text { and } e \text { is vertical } \tag{6}
\end{array}
$$

and $X_{e}=0$ otherwise. The constants $t_{H}, t_{V}$ will be specified later. Observe that because the regions $D_{e}$ are nonoverlapping, the random variables $Y_{e}$, and hence also $X_{e}$, are mutually independent. For each vertex $v \in \mathbb{Z}^{2}$, set

$$
\begin{equation*}
X_{v}=1 \tag{7}
\end{equation*}
$$

if the passage times of all edges in the ball $B_{v}$ are bounded above by $t_{S}$, a constant to be specified later; otherwise, $X_{v}=0$.
Note that if the edge passage time distribution $F$ has support contained in $\left[0, t_{S}\right]$ then $X_{v}=1$ almost surely for each vertex $v$; in this case, the percolation event $\Omega$ is determined completely by the edge r.v.s $X_{e}$. In general, if the support of $F$ is unbounded, then it will not be the case that $X_{v}=1$ a.s.; moreover, the random variables $X_{e}$ and $X_{v}$ will not necessarily be mutually independent, and so Proposition 3 will not apply directly. A modification of the construction to remedy this difficulty will be proposed later. However, the random variables $X_{e}$ indexed by edges $e$ will be mutually independent, and $P\left\{X_{e}=1\right\}$ will depend only on whether the edge $e$ is vertical or horizontal: denote the two possible values by $p_{V}, p_{H}$, respectively. Similarly, the distribution of the travel time $Y_{e}$ depends only on whether $e$ is vertical or horizontal.
To complete the construction, it remains to specify the constants $R, t_{H}, t_{V}, t_{S}$, the vectors $x, y$, and the aperture of the spherical cone that determines the regions $D_{e}$. Here Hypotheses $1-2$ come into play: Hypothesis 1 will be used to determine $R$ and the aperture of the cone,
and Hypothesis 2 the remaining parameters. Let $u, v$ be linearly independent vectors in $\mathbb{R}^{d}$ for which Hypothesis 1 holds. Given this choice $u, v$, choose an aperture sufficiently small that

$$
\begin{equation*}
\mathcal{A}_{u} \cap \mathcal{A}_{v}=\{0\} \tag{8}
\end{equation*}
$$

where $\mathcal{A}_{u}, \mathcal{A}_{v}$ are the closed spherical cones of this aperture centered at the points where the rays $\mathcal{L}_{u}, \mathcal{L}_{v}$ intersect the unit sphere $\Sigma_{d-1}$. By Hypothesis 1 , for any $\delta>0$ there exists $1<R<\infty$ such that for any points $\xi, \zeta \in \mathbb{Z}^{d}$ both of norm at least $R / 2$ and both at distance less than 2 from the ray $\mathcal{L}_{u}$ (alternatively, the ray $\mathcal{L}_{v}$ ), there is probability at least $1-\delta$ that the geodesic connecting the points $\xi, \zeta$ lie entirely in the cone $\mathcal{A}_{u}$ (alternatively, the cone $\mathcal{A}_{v}$ ). Consequently, if $R$ is used as the radius of the balls $B_{v}$ and if each region $D_{e}$ is the intersection of two spherical cones each congruent to $\mathcal{A}_{u}$ (and therefore also $\mathcal{A}_{v}$ ), then for each edge $e$ the probability that the geodesic segment connecting the lattice points $z_{e}$ and $z_{e}^{\prime}$ lies entirely in $D_{e}$ is at least $1-2 \delta$.
Assume now that the vectors $u, v$ satisfy both Hypotheses 1 and 2 . Let $R$ and the cone aperture be chosen as in the preceding paragraph (these aso depend on $\delta>0$, which will be specified below). According to Hypothesis 2, there are scalar sequences $a_{u}(n), a_{v}(n)$ so that the travel times $T(n u), T(n v)$, normalized as in (3), have mean-zero limit distributions $G_{u}, G_{v}$. Because these limit distributions have mean zero, if $g_{u}, g_{v}$ are $(1-\epsilon)$ th quantiles of the distributions $G_{u}, G_{v}$, for any $\epsilon>0$, then

$$
\nu_{u}:=\frac{\int_{-\infty}^{g_{u}} t d G_{u}(t)}{1-G_{u}\left(g_{u}\right)}<0 \quad \text { and } \quad \nu_{v}:=\frac{\int_{-\infty}^{g_{v}} t d G_{v}(t)}{1-G_{v}\left(g_{v}\right)}<0
$$

This is in fact the crucial point of the proof. The plan now is to let $x, y$ be the points of $\mathbb{Z}^{d}$ closest to $n u$ and $n v$ for a large integer $n$, and for a suitably small value $\epsilon>0$ to let

$$
t_{H}=n \mu(u)+a_{u}(n) g_{u} \quad \text { and } \quad t_{V}=n \mu(v)+a_{v}(n) g_{v}
$$

The integer $n$ is chosen so that $R^{d}$ is negligible compared to the minimum of $n, a_{u}(n)$, and $a_{v}(n)$ - this is possible because of the hypothesis that the sequences $a_{u}(n), a_{v}(n)$ go to $\infty$ and so that the approximation (3) to the distributions of $T(n u), T(n v)$ by $G_{u}, G_{v}$ (in particular, to their $(1-\epsilon)$ th quantiles and their conditional expectations) is sufficiently accurate. For such $n$, if $\delta>0$ is sufficiently small then for all edges $e$,

$$
\begin{equation*}
P\left\{X_{e}=1\right\}>1-2 \delta-2 \epsilon, \tag{9}
\end{equation*}
$$

and

$$
\begin{equation*}
E\left(Y_{e}-n \mu(u) \mid X_{e}=1\right)<-4 C R^{d} t_{S} \tag{10}
\end{equation*}
$$

where $C$ is the volume of the $d$-dimensional unit ball (thus, $C R^{d}$ is, approximately, the number of edges of the lattice that intersect the ball $B_{v}$ ). Observe that (9) follows from the choice of $t_{H}, t_{V}$, as $g_{u}, g_{v}$ are $(1-\epsilon)$ th quantiles of $G_{u}, G_{v}$, and (10) from the fact that $\nu_{u}, \nu_{v}<0$. Finally, choose $t_{S}$ so large that $P\left\{X_{v}=1\right\}>1-\delta$ for each vertex $v$ (or, if the passage time distribution $F$ has finite support, so that $P\left\{X_{v}=1\right\}=1$ ). Note that if $X_{v}=1$ then the travel time between any two vertices in $B_{v}$ is bounded by $2 C R^{d} t_{S}$. In particular, if $X_{v}=X_{v^{\prime}}=1$ for both endpoints $v, v^{\prime}$ of an edge $e$, then the travel time from $\varphi(v)$ to $\varphi\left(v^{\prime}\right)$ is at most $4 C R^{d} t_{S}+Y_{e}$.
The proof of Theorem 1 in the special case where the edge passage time distribution $F$ has bounded support may now be completed. Assume that $\epsilon, \delta>0$ are small enough that $1-2 \epsilon-2 \delta$
is no smaller than the critical value $p_{c}$ for standard oriented bond percolation on the twodimensional integer lattice. Then by (9) and Proposition 3, the percolation event $\Omega$ has positive probability. On this event, there is an infinite open oriented path $\gamma$ starting at the origin that crosses the diagonal infinitely often. Consider now the image in $\mathbb{Z}^{d}$ of this path by the embedding $\varphi$ : For each edge $e \in \gamma$, the traversal time $\tau(\gamma)$ is no larger than $Y_{e}+4 C R^{d} t_{S}$. Moreover, conditional on the realization $\left\{X_{e}\right\}$, the random variables $Y_{e}$, for $e \in \gamma$, are independent, with conditional distribution depending only on the orientation of $e$ (vertices or horizontal), and with conditional expectations satisfying (10) above. Thus, by the SLLN, for $(m, m)$ on the path $\gamma$ the traversal time $T_{m}$ of the segment of $\varphi(\gamma)$ from the origin to $\varphi(m, m)$ will satisfy

$$
\begin{equation*}
\liminf _{m \rightarrow \infty} T_{m} / m-n \mu(u)-n \mu(v)<0 \tag{11}
\end{equation*}
$$

Moreover, the liminf, when divided by $n$, remains bounded away from 0 as $n \rightarrow \infty$. As $n \rightarrow \infty$, the limiting directions of $x \approx n u$ and $y \approx n v$ approach $u$ and $v$, respectively. Consequently, by the homogeneity and continuity of the inverse speed function $\mu(\cdot)$, the relation (4) follows for $\alpha=1 / 2$.

Note. To prove the relation (4) in general, it suffices to prove it for $\alpha=1 / 2$. This follows from the homogeneity of $\mu$ and the fact that any convex combination $\alpha u+(1-\alpha) v$ may be realized as $\left(u^{\prime}+v^{\prime}\right) / 2$ for suitable positive scalar multiples $u^{\prime}=c u$ and $v^{\prime}=c^{\prime} v^{\prime}$.

As has already been noted, if the edge passage time distribution $F$ does not have bounded support then a modification of the argument is necessary, because in this case the Bernoulli random variables $X_{e}$ and $X_{v}$ need not be independent, and so Proposition 3 is not directly applicable. In this case, the random variables $Y_{e}$ should be replaced by $Y_{e}^{\prime}$, where $Y_{e}^{\prime}$ is defined as follows: Reset the edge passage time to $t_{S}$ for all edges $e$ that intersect one of the balls $B_{v}$, and let all other edge passage times be chosen from the distribution $F$. Now let $Y_{e}^{\prime}$ be the minimum traversal time $\tau^{\prime}(\gamma)$ among all paths from $z$ to $z^{\prime}$ that lie entirely in $D_{e}$, as before, and define

$$
\begin{array}{lll}
X_{e}^{\prime}=1 & \text { if } \quad Y_{e}^{\prime} \leq t_{H} \quad \text { and } e \text { is horizontal } \\
X_{e}^{\prime}=1 \quad \text { if } \quad Y_{e}^{\prime} \leq t_{V} \quad \text { and } e \text { is vertical } \tag{13}
\end{array}
$$

and $X_{e}^{\prime}=0$ otherwise. The random variables $X_{e}^{\prime}, X_{v}$ are now mutually independent, and so Proposition 3 applies. Moreover, the random variables $Y_{e}^{\prime}$ are, conditional on the realization $X_{e}^{\prime}, X_{v}$ of the modified oriented percolation process, independent, with distribution depending only on the orientation of the edge $e$, and conditional means satisfying (10), at least when $n$ is sufficiently large (since the modification of the edge passage times in the balls $B_{v}$ will be swamped by the renormalization factors $a_{u}(n), a_{v}(n)$ for large $\left.n\right)$. The argument may now be completed in the same manner as in the special case discussed above.

## References

[1] Benjamini, I., Kalai, G., and Schramm, O. (2003) First passage percolation has sublinear distance variance. To appear in Annals of Probability.
[2] Durrett, R., and Liggett, T. (1981) The shape of the limit set in Richardson's growth model. Annals of Probability 9 186-193.
[3] Durrett, R. (1984) Oriented percolation in two dimensions. Annals of Probability 12999 - 1040.
[4] Grimmett, G. (1989) Percolation. Springer-Verlag, New York.
[5] Hammersley, J. M. and Welsh, D. J. A. (1965) First-passage percolation, subadditive processes, stochastic networks, and generalized renewal theory. Proc. Internat. Res. Sem., Statist. Lab., Univ. California, Berkeley, Calif. 61-110.
[6] Johansson, K. (2000) Transversal fluctuations for increasing subsequences on the plane. Probability Theory and Related Fields 116445 - 456.
[7] Kesten (1986) Aspects of first passage percolation. École d'été de probabilitéde SaintFlour, XIV-1984 125-264.
[8] Kordzakhia, G. (2003) The problem of coexistence in multi-type competition models. University of Chicago Ph. D. dissertation.
[9] Licea, C.; Newman, C. M.; Piza, M. S. T. (1996) Superdiffusivity in first-passage percolation. Probability Theory and Related Fields 106 559-591.
[10] Licea, C.; Newman, C. M. (1996) Geodesics in two-dimensional first-passage percolation. Annals of Probability 24399 - 410.
[11] Liggett, T. (1985) Interacting Particle Systems. Springer-Verlag, New York
[12] Newman, C. M.; Piza, M. S. T. (1995) Divergence of shape fluctuations in two dimensions. Annals of Probability 23 977-1005.
[13] Newman, C. M. (1995) A surface view of first-passage percolation. Proceedings of the International Congress of Mathematicians, Vol. 1, 2 (Zurich, 1994) 1017-1023.
[14] Pemantle, R.; Peres, Y. (1994) Planar first-passage percolation times are not tight. Probability and phase transition (Cambridge, 1993) $261-264$.
[15] Russo, L. (1981) On the critical percolation probabilities. Z. Warscheinlichkeitstheorie 56 229-237.


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