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# Derivatives of sup-functionals of fractional Brownian motion evaluated at $H=\frac{1}{2}^{*}$ 

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#### Abstract

We consider a family of sup-functionals of (drifted) fractional Brownian motion with Hurst parameter $H \in(0,1)$. This family includes, but is not limited to: expected value of the supremum, expected workload, Wills functional, and Piterbarg-Pickands constant. Explicit formulas for the derivatives of these functionals as functions of Hurst parameter evaluated at $H=\frac{1}{2}$ are established. In order to derive these formulas, we develop the concept of derivatives of fractional $\alpha$-stable fields introduced by Stoev \& Taqqu (2004) and propose Paley-Wiener-Zygmund representation of fractional Brownian motion.


Keywords: fractional Brownian motion; expected workload; Wills functional; Pickands constant; Piterbarg constant.
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## 1 Introduction

Sup-functionals of the supremum of fractional Brownian motion play crucial role in many problems arising both in theoretical probability and its applications, as in e.g. statistics, financial mathematics, risk theory, queueing theory, see e.g. [2, 11, 26, 30]. Unfortunately, despite intensive research on their properties, apart from some particular cases - reduced mostly to standard Brownian motion - the exact value of such functionals is not known.

[^0]Let $\left\{B_{H}(t): t \in \mathbb{R}_{+}\right\}$be fractional Brownian motion (fBm) with Hurst parameter $H \in(0,1)$, that is a centered Gaussian process with continuous sample paths a.s. and

$$
\begin{equation*}
\operatorname{Cov}\left(B_{H}(s), B_{H}(t)\right)=\frac{1}{2}\left(|s|^{2 H}+|t|^{2 H}-|s-t|^{2 H}\right) \tag{1.1}
\end{equation*}
$$

for all $s, t \in \mathbb{R}_{+}$. In this manuscript, we consider the following two families of functionals

$$
\begin{equation*}
\mathscr{M}_{H}(T, a):=\mathbb{E}\left(\sup _{t \in[0, T]} B_{H}(t)-a t\right), \quad \mathscr{P}_{H}(T, a):=\mathbb{E}\left(\exp \left\{\sup _{t \in[0, T]} \sqrt{2} B_{H}(t)-a t^{2 H}\right\}\right) \tag{1.2}
\end{equation*}
$$

where $a \in \mathbb{R}$ is the intensity of the drift and $T \in \mathbb{R}_{+} \cup\{\infty\}$ is the time horizon. These functionals cover a range of interesting quantities in the extreme value theory of Gaussian processes. In particular:
(i) For $a \in \mathbb{R}$ and $T \in \mathbb{R}_{+}$, the quantity $\mathscr{M}_{H}(T, a)$ is the expected value of the workload in the fluid queue with an fBm input at time $T$ under the assumption that at time 0 the system starts off empty. Analogously, if $a>0$ and $T=\infty$, then $\mathscr{M}_{H}(\infty, a)$ is the expected stationary workload of a queue with an fBm input, see e.g. [26, 11];
(ii) for $T \in \mathbb{R}_{+}$, the quantity $\mathscr{H}_{H}(T):=\mathscr{P}_{H}(T, 1)$ is known as Wills functional (truncated Pickands constant), see e.g. [39, 9] and references therein;
(iii) for $a>1$, the quantity $\mathscr{P}_{H}(\infty, a)$ is known as Piterbarg constant; see e.g. [31, 4] and references therein;
(iv) the quantity $\mathscr{H}_{H}:=\lim _{T \rightarrow \infty} \frac{1}{T} \mathscr{H}_{H}(T)$ is known as Pickands constant; see e.g. [30, 31].

The values of functionals $\mathscr{M}_{H}(T, a)$ and $\mathscr{P}_{H}(T, a)$ are notoriously difficult to find in cases other than $H \in\left\{\frac{1}{2}, 1\right\}$. Thus, most of the work is focused on finding upper and lower bounds for these quantities (see, e.g. [34, 12, 10, 4, 7, 8, 5]) or determining their asymptotic behavior in various settings (as $H$ goes to $0, H$ goes to $1, T$ grows large, $a$ goes to 0 etc.); see, e.g., $[21,5]$. We note that in many cases simulation methods do not help in estimation of $\mathscr{P}_{H}(T, a)$. For example, the second moment of $\mathscr{P}_{H}(\infty, a)$ does not exist when $a<2$, which makes it very difficult to assess error bounds in this case. When $H \rightarrow 0$, then the approximation error resulted from simulation becomes overwhelming, see [18, 17].

In recent years, Delorme, Wiese and others studied the behavior of fBm and its time for $H \approx \frac{1}{2}$ using the perturbation theory $[13,15,14,16]$. Our work was initially inspired by their result in [13], where the following expansion for Pickands constant (see item (iv) above) was derived

$$
\begin{equation*}
\mathscr{H}_{H}=1-2\left(H-\frac{1}{2}\right) \gamma_{\mathrm{E}}+O\left(\left(H-\frac{1}{2}\right)^{2}\right), \quad H \rightarrow \frac{1}{2} \tag{1.3}
\end{equation*}
$$

where $\gamma_{\mathrm{E}}$ is the Euler-Mascheroni constant.
The main goal of this contribution is to develop tools for researching expansions similar to (1.3) for functionals introduced in Eq. (1.2). In particular, we find explicit formulas for the derivatives of these functionals evaluated at $H=\frac{1}{2}$, i.e.

$$
\begin{equation*}
\mathscr{M}_{1 / 2}^{\prime}(T, a):=\left.\frac{\partial}{\partial H} \mathscr{M}_{H}(T, a)\right|_{H=1 / 2} \quad \text { and } \quad \mathscr{P}_{1 / 2}^{\prime}(T, a):=\left.\frac{\partial}{\partial H} \mathscr{P}_{H}(T, a)\right|_{H=1 / 2} \tag{1.4}
\end{equation*}
$$

For these purposes, we consider a probability space on which fBms with all values of $H \in(0,1)$ are coupled in a non-trivial way. In particular, we consider the Mandelbrot \& van Ness' (MvN) fractional Brownian motion introduced in the seminal work [25]; see Eq. (2.1) below for precise definition. Following [29] the realizations of MvN field can be viewed as two-dimensional random surfaces $(H, t) \mapsto B_{H}(t)$, as opposed to
one-dimensional trajectories $t \mapsto B_{H}(t)$ for each fixed $H \in(0,1)$. While these random surfaces are non-differentiable in the time direction (with respect to $t$ ), they turn out to be smooth functions of the Hurst parameter (with respect to $H$ ). Therefore, following Stoev and Taqqu $[36,37]$ we define the derivatives of the MvN field with respect to the Hurst parameter $H$. Intuitively speaking, the concept of $H$-derivative of fBm allows us to rigorously write $\frac{\partial}{\partial H} \mathbb{E} B_{H}(\tau)=\mathbb{E} \frac{\partial}{\partial H} B_{H}(\tau)$, where $\tau$ is some (well-behaved) random time and $\left\{\frac{\partial}{\partial H} B_{H}(t), t \in \mathbb{R}_{+}\right\}$is a certain explicitly defined Gaussian process. It will turn out that in our context, the latter expression (i.e. $\mathbb{E} \frac{\partial}{\partial H} B_{H}(\tau)$ ) is tractable (explicitly computable).

We note that Stoev and Taqqu considered a broader class of fractional $\alpha$-stable fields, while for the purposes of this paper we limit ourselves to the Gaussian case, i.e. $\alpha=2$, which corresponds to the fBm and, in particular, the MvN field. See also [37, Eq. (1.9)] with $\left(a_{+}, a_{-}\right)=(1,0)$ and [32, Chapter 7.4] for more information on fractional $\alpha$-stable fields. Focusing on the Gaussian case, we strengthen some of the results derived by the authors. In particular, we show sample path continuity of the derivative fields (Proposition 4.2) and strengthen [36, Lemma 4.1] to almost sure convergence for all $(H, t) \in(0,1) \times \mathbb{R}_{+}$(Proposition 4.3). These propositions are then used in the proofs of main results.

Finally, we propose a Paley-Wiener-Zygmund (PWZ) representation of MvN field and its derivatives. We show that PWZ field is a continuous modification of MvN field (Proposition 4.1) and the difference quotients of $n$-th derivative converge everywhere to ( $n+1$ )-st derivative almost surely (Proposition 4.3). The PWZ representation is defined in terms of Riemann integrals, as opposed to stochastic integrals in the original Mandelbrot \& van Ness' definition. In the context of this manuscript this representation is more tractable and allows us to express the main quantities of interest, cf. (1.4), as integrals of elementary functions, which we then calculate explicitly in special cases, see Section 3.

The manuscript is organized as follows. In Section 2 we define the fractional Brownian motion through its MvN and PWZ representations. We also recall the facts related to the joint distribution of the supremum of drifted Brownian motion and its argmax. In particular, we introduce the conditional distribution of Brownian motion, conditioned on the value of the supremum and its argmax, which follows the law of the generalized 3-dimensional Bessel bridge. In Section 3 we state main results of this paper. In Theorem 1 we give a formula for $\mathscr{M}_{1 / 2}^{\prime}(T, a)$ and in Theorem 2 for $\mathscr{P}_{1 / 2}^{\prime}(T, a)$ in terms of integrals of explicit elementary functions. More explicit formulae for $\mathscr{M}_{1 / 2}^{\prime}(T, 0), \mathscr{M}_{1 / 2}^{\prime}(\infty, a)$ and $\mathscr{P}_{1 / 2}^{\prime}(\infty, a)$ are given in Corollary 3.3. Additionally, in Corollary 3.4 we show that Piterbarg constants $\mathscr{P}_{H}(\infty, a)$ are monotone as functions of Hurst parameter. While this result is not directly related to our topic, it is a direct consequence of Proposition C. 1 and it might be of independent interest. In Section 4 we define and examine the $H$-derivatives of the fBm both with the use of MvN and PWZ representations. In Section 5 we give the proofs of main theorems. Since they require quite a lot of different results, throughout this section we introduce various preliminary results. Some of them, for the argmax of $B_{H}(t)-a t$ and $B_{H}(t)-a t^{2 H}$ are presented in Proposition 5.1 and Proposition 5.2 respectively. More technical results are deferred to appendices. In Appendix A we show the equivalence between $\mathcal{L}^{2}$ and Paley-Wiener-Zygmund stochastic integrals. In Appendix B we present proofs of results from Section 2 and Section 4. In Appendix C we state results needed in the proof of Theorem 3.1 and Theorem 3.2, which can be also of independent interest such as monotonicity of Piterbarg constants and bounds on moments of the supremum of a Gaussian process at a random time. Finally, in Appendix D we write down all the calculations needed to prove Corollary 3.3.

## 2 Preliminaries

Let $\{B(t): t \in \mathbb{R}\}$ be a standard, two-sided Wiener process and let $\mathbb{R}_{+}:=[0, \infty)$. In their seminal paper, [25], introduced a family of processes $\left\{X_{H}(t), t \in \mathbb{R}_{+}\right\}$for $H \in(0,1)$, where

$$
\begin{equation*}
X_{H}(t)=\int_{-\infty}^{0}\left[(t-s)^{H-\frac{1}{2}}-(-s)^{H-\frac{1}{2}}\right] \mathrm{d} B(s)+\int_{0}^{t}(t-s)^{H-\frac{1}{2}} \mathrm{~d} B(s) \tag{2.1}
\end{equation*}
$$

For each $H \in(0,1)$, $\left\{X_{H}(t): t \in \mathbb{R}_{+}\right\}$is a centered Gaussian process with

$$
\operatorname{Cov}\left(X_{H}(s), X_{H}(t)\right)=V(H) \cdot \operatorname{Cov}\left(B_{H}(s), B_{H}(t)\right)
$$

where, $B_{H}(t)$ is a fractional Brownian motion introduced in (1.1), and

$$
V(H):=\int_{0}^{\infty}\left((1+s)^{H-\frac{1}{2}}-s^{H-\frac{1}{2}}\right)^{2} \mathrm{~d} s+\int_{0}^{1} s^{2 H-1} \mathrm{~d} s=\frac{\Gamma\left(\frac{1}{2}+H\right) \Gamma(2-2 H)}{2 H \Gamma\left(\frac{3}{2}-H\right)}
$$

see e.g. [27, Appendix A], where the explicit formula for $V(H)$ is derived. This shows that, up to the scaling factor, for each $H \in(0,1)$, process $\left\{X_{H}(t): t \in \mathbb{R}_{+}\right\}$is a fractional Brownian motion. Therefore, we call $\left\{X_{H}(t):(H, t) \in(0,1) \times \mathbb{R}_{+}\right\}$the Mandelbrot \& van Ness (MvN) fractional Brownian field. At the same time, there exists another representation $\left\{\widetilde{X}_{H}(t):(H, t) \in(0,1) \times \mathbb{R}_{+}\right\}$of MvN field,

$$
\begin{align*}
\widetilde{X}_{H}(t): & =\left(H-\frac{1}{2}\right) \int_{-\infty}^{0}\left[(t-s)^{H-\frac{3}{2}}-(-s)^{H-\frac{3}{2}}\right] B(s) \mathrm{d} s+t^{H-\frac{1}{2}} B(t) \\
& -\left(H-\frac{1}{2}\right) \int_{0}^{t}(t-s)^{H-\frac{3}{2}}(B(t)-B(s)) \mathrm{d} s \tag{2.2}
\end{align*}
$$

which we call the Paley-Wiener-Zygmund (PWZ) representation. In Section 4 it is shown that field $\left\{\widetilde{X}_{H}(t):(H, t) \in(0,1) \times \mathbb{R}_{+}\right\}$is a modification of MvN field, whose sample paths $(H, t) \mapsto \widetilde{X}_{H}(t)$ are continuous almost surely; see Proposition 4.1.

From now on, the fractional Brownian motion is defined through the process $X_{H}(\cdot)$, i.e. with $D(H):=(V(H))^{-1 / 2}$, we put

$$
\begin{equation*}
B_{H}(t):=D(H) \cdot X_{H}(t) \tag{2.3}
\end{equation*}
$$

For any $a \in \mathbb{R}$ we define

$$
Y_{H}(t ; a):=B_{H}(t)-a t, \quad \text { and } \quad Y_{H}^{*}(t ; a):=B_{H}(t)-a t^{2 H}
$$

Notice that $Y_{1 / 2}^{*}=Y_{1 / 2}$ for all $a \in \mathbb{R}$. Furthermore, we define suprema and their locations for these processes

$$
\begin{array}{ll}
\bar{Y}_{H}(T, a):=\sup _{t \in[0, T]} Y_{H}(t ; a), & \tau_{H}(T, a):=\arg \max \left\{t \in[0, T]: Y_{H}(t ; a)\right\}, \\
\bar{Y}_{H}^{*}(T, a):=\sup _{t \in[0, T]} Y_{H}^{*}(t ; a), \quad \tau_{H}^{*}(T, a):=\arg \max \left\{t \in[0, T]: Y_{H}^{*}(t ; a)\right\} . \tag{2.4}
\end{array}
$$

It is known that when $a \in \mathbb{R}$ and $T \in \mathbb{R}_{+}$, or $a>0$ and $T=\infty$ then $\bar{Y}, \bar{Y}^{*}$ are almost surely finite and $\tau_{H}$ and $\tau_{H}^{*}$ are well-defined (unique) and almost surely finite; see [19] citing [24].

Next, we recall the formulae for the joint density of the supremum of (drifted) Brownian motion over $[0, T]$ and its argmax. In the following, for $z \in \mathbb{R}$ we define the error function and the complementary error function respectively:

$$
\operatorname{erf}(z):=\frac{2}{\sqrt{\pi}} \int_{0}^{z} e^{-t^{2}} \mathrm{~d} t, \quad \operatorname{erfc}(z):=1-\operatorname{erf}(z)
$$

For brevity of exposition, in this section we write $Y(t):=Y_{1 / 2}(t ; a), \bar{Y}(T):=\bar{Y}_{1 / 2}(T, a)$, and $\tau(T):=\tau_{1 / 2}(T, a)$. When $a \in \mathbb{R}$ and $T \in \mathbb{R}_{+}$, according to [35] we have:

$$
\frac{\mathbb{P}(\tau(T) \in \mathrm{d} t, \bar{Y}(T) \in \mathrm{d} y, Y(T) \in \mathrm{d} x)}{\mathrm{d} t \mathrm{~d} y \mathrm{~d} x}=\frac{1}{\pi} \frac{y(y-x)}{t^{3 / 2}(T-t)^{3 / 2}} \exp \left\{-\frac{y^{2}}{2 t}-\frac{(y-x)^{2}}{2(T-t)}\right\} e^{-a x-a^{2} T / 2}
$$

After the integration of the above density with respect to $x$ over the domain $x<y$ we recover the joint density of the pair $(\tau(T), \bar{Y}(T))$

$$
\begin{align*}
p(t, y ; T, a) & :=\frac{\mathbb{P}(\tau(T) \in \mathrm{d} t, \bar{Y}(T) \in \mathrm{d} y)}{\mathrm{d} t \mathrm{~d} y} \\
& =\frac{y e^{-\frac{(y+t a)^{2}}{2 t}}}{\sqrt{\pi} t^{3 / 2}} \cdot\left(\frac{e^{-a^{2}(T-t) / 2}}{\sqrt{\pi(T-t)}}+\frac{a}{\sqrt{2}} \cdot \operatorname{erfc}\left(-a \sqrt{\frac{T-t}{2}}\right)\right) \tag{2.5}
\end{align*}
$$

for $(t, y) \in(0, T) \times \mathbb{R}_{+}$; see also [6, 2.1.13.4]. When $a>0$, then the pair $(\tau(\infty), \bar{Y}(\infty))$ is well-defined, with density

$$
\begin{equation*}
p(t, y ; \infty, a)=\frac{\sqrt{2} a y e^{-\frac{(y+t a)^{2}}{2 t}}}{\sqrt{\pi} t^{3 / 2}} \tag{2.6}
\end{equation*}
$$

for $(t, y) \in \mathbb{R}_{+}^{2}$; see e.g. [6, 2.1.13.4(1)].
When $a \in \mathbb{R}$ and $T \in \mathbb{R}_{+}$or $a>0$ and $T=\infty$ then, conditionally on $\bar{Y}(T)=y, \tau(T)=t$, the process $\{Y(t)-Y(t-s)\}_{s \in[0, t]}$ has the law of a 3-dimensional Bessel bridge from $(0,0)$ to $(t, y)$, see e.g. [3, Prop. 2]. The law of this process does not depend on the value of the drift $a$ nor on the time horizon $T$. Moreover, the density of the marginal distribution of this process is known. In the following, for $t, y>0$ and $s \in(0, t)$ we define

$$
g(x, s ; t, y):=\frac{\mathbb{P}(Y(t)-Y(t-s) \in \mathrm{d} x \mid \tau(T)=t, \bar{Y}(T)=y)}{\mathrm{d} x}
$$

Then, according to e.g. [23, Theorem 1(ii)] we have

$$
\begin{equation*}
g(x, s ; t, y)=\frac{\frac{x}{s^{3} / 2} \exp \left\{-\frac{x^{2}}{2 s}\right\}}{\frac{y}{t^{3 / 2}} \exp \left\{-\frac{y^{2}}{2 t}\right\}} \cdot \frac{1}{\sqrt{2 \pi(t-s)}}\left[e^{-\frac{(x-y)^{2}}{2(t-s)}}-e^{-\frac{(x+y)^{2}}{2(t-s)}}\right] 1_{\{x<y\}} \tag{2.7}
\end{equation*}
$$

The following functional of the 3-dimensional Bessel bridge from $(0,0)$ to $(t, y)$ will be important later on

$$
\begin{equation*}
I(t, y):=\mathbb{E}\left(\left.\int_{0}^{t} \frac{Y(t)-Y(t-s)}{s} \mathrm{~d} s \right\rvert\, \tau(T)=t, \bar{Y}(T)=y\right) \tag{2.8}
\end{equation*}
$$

Using the fact that we have explicit formula for the density $g(\cdot)$ in Eq. (2.7), we can express $I(t, y)$ as a double integral, see the proposition below, whose proof is given in Appendix B.
Proposition 2.1. For any $t, y>0$ we have

$$
I(t, y):=\sqrt{\frac{2}{\pi}} \cdot t y^{-1} \int_{0}^{\infty} \int_{0}^{\infty} \frac{x^{2}}{q\left(1+q^{2}\right)^{2}}\left(e^{-(x-y q / \sqrt{t})^{2} / 2}-e^{-(x+y q / \sqrt{t})^{2} / 2}\right) \mathrm{d} x \mathrm{~d} q
$$

## 3 Main results

Recall the definitions of the functionals $\mathscr{M}_{H}$ and $\mathscr{P}_{H}$ introduced in the Introduction in Eq. (1.2). We can rewrite them in terms of the expectation of random variables $\bar{Y}$ and $\bar{Y}^{*}$, i.e

$$
\mathscr{M}_{H}(T, a)=\mathbb{E}\left(\bar{Y}_{H}(T, a)\right), \quad \mathscr{P}_{H}(T, a)=\mathbb{E}\left(\exp \left\{\sqrt{2} \cdot \bar{Y}_{H}^{*}\left(T, \frac{a}{\sqrt{2}}\right)\right\}\right)
$$

We shall derive formulas for the derivative of functions $H \mapsto \mathscr{M}_{H}(T, a)$ and $H \mapsto \mathscr{P}_{H}(T, a)$ evaluated at $H=\frac{1}{2}$. Following Eq. (1.4), in what follows let

$$
\mathscr{M}_{1 / 2}^{\prime}(T, a):=\lim _{H \rightarrow 1 / 2} \frac{\mathscr{M}_{H}(T, a)-\mathscr{M}_{1 / 2}(T, a)}{H-\frac{1}{2}}
$$

and

$$
\mathscr{P}_{1 / 2}^{\prime}(T, a):=\lim _{H \rightarrow 1 / 2} \frac{\mathscr{P}_{H}(T, a)-\mathscr{P}_{1 / 2}(T, a)}{H-\frac{1}{2}} .
$$

Theorem 3.1. If $a \in \mathbb{R}$ and $T \in \mathbb{R}_{+}$or $a>0$ and $T=\infty$, then

$$
\begin{equation*}
\mathscr{M}_{1 / 2}^{\prime}(T, a)=\int_{0}^{T} \int_{0}^{\infty}(y(1+\log (t))+a t \log (t)-I(t, y)) p(t, y ; T, a) \mathrm{d} y \mathrm{~d} t . \tag{3.1}
\end{equation*}
$$

Theorem 3.2. If $a \in \mathbb{R}$ and $T \in \mathbb{R}_{+}$or $a>1$ and $T=\infty$, then

$$
\begin{equation*}
\mathscr{P}_{1 / 2}^{\prime}(T, a)=\int_{0}^{T} \int_{0}^{\infty} \sqrt{2}\left(y(1+\log (t))-\frac{a}{\sqrt{2}} t \log (t)-I(t, y)\right) e^{\sqrt{2} y} p\left(t, y ; T, \frac{a}{\sqrt{2}}\right) \mathrm{d} y \mathrm{~d} t . \tag{3.2}
\end{equation*}
$$

In order to gain a better intuitive understanding of the main results, we provide an outline of the proof of Theorem 3.1 below; the proof of Theorem 3.2 will be similar. Full proofs of these theorems are given in Section 5.

Outline of the proof of Theorem 3.1. The main part of the proof is to show that

$$
\left.\frac{\partial}{\partial H}\left[\mathbb{E} Y_{H}\left(\tau_{H}(T, a), a\right)\right]\right|_{H=1 / 2}=\mathbb{E}\left[\left.\frac{\partial}{\partial H} Y_{H}\left(\tau_{1 / 2}(T, a), a\right)\right|_{H=1 / 2}\right]
$$

where the expression on the left is, by definition, equal to $\mathscr{M}_{1 / 2}^{\prime}(T, a)$. To explain in words, we may swap the order of taking the expected value and differentiation in the definition of $\mathscr{M}_{1 / 2}^{\prime}(T, a)$ and swap $\tau_{H}(T, a)$ with $\tau_{1 / 2}(T, a)$ above. The derivative $\frac{\partial}{\partial H} Y_{H}(t, a)$ is understood point-wise, for every fixed $t \in \mathbb{R}_{+}$. In proofs we will need an $H$-calculus, which is formally introduced and worked out in Section 4. Here we need only the first derivative
$X_{H}^{(1)}(t)=\int_{-\infty}^{0}\left[\log (t-s)(t-s)^{H-\frac{1}{2}}-\log (-s)(-s)^{H-\frac{1}{2}}\right] \mathrm{d} B(s)+\int_{0}^{t} \log (t-s)(t-s)^{H-\frac{1}{2}} \mathrm{~d} B(s)$.
In $H$-calculus, the Leibniz formula is valid and the $H$-derivative at $H=1 / 2$ of the fBm $\frac{\partial}{\partial H} B_{H}(t)$ is $X(t)+X^{(1)}(t)$. This is derived later in Section 4, see (4.9). As soon as the equation above is established, we find that

$$
\left.\frac{\partial}{\partial H} Y_{H}\left(\tau_{1 / 2}(T, a), a\right)\right|_{H=1 / 2}=X(\tau)+X^{(1)}(\tau)
$$

where $\tau:=\tau_{1 / 2}(T, a)$. Finally, $\mathscr{M}_{1 / 2}^{\prime}(T, a)$ is equal to the expectation of the expression above and it can be expressed as the definite integral in Eq. (3.1) using PWZ representations of $X$ and $X^{(1)}$ and the fact that distribution of Brownian motion conditioned on its supremum and time of supremum is known; see Section 2 for more information.

We note that the derivatives $\mathscr{M}_{1 / 2}^{\prime}(T, a)$ and $\mathscr{P}_{1 / 2}^{\prime}(T, a)$ in (3.1) and (3.2) are expressed as definite integrals. Thus, they can be computed numerically for any drift $a$ and time horizon $T$, which satisfy the assumptions of Theorem 3.1 and Theorem 3.2 respectively. In addition, we were able to calculate these derivatives explicitly in special cases; see the corollary below.
Corollary 3.3. It holds that
(i) if $T \in \mathbb{R}_{+}$, then $\mathscr{M}_{1 / 2}^{\prime}(T, 0)=\sqrt{\frac{2 T}{\pi}} \cdot(\log (T)-2)$,
(ii) if $a>0$, then $\mathscr{M}_{1 / 2}^{\prime}(\infty, a)=-\frac{1}{a}\left(\gamma_{\mathrm{E}}+\log \left(2 a^{2}\right)\right)$,
(iii) if $a>1$, then $\mathscr{P}_{1 / 2}^{\prime}(\infty, a)=-\frac{2 a}{a-1}\left(1+(a-2) \log \left(\frac{a-1}{a}\right)\right)$.

The calculations leading to Corollary 3.3 are deferred to Appendix D. From Corollary 3.3(iii), it straightforwardly follows that for any $a>1$, the function $H \rightarrow \mathscr{P}_{H}(\infty, a)$ is decreasing in the neighborhood of $H=\frac{1}{2}$ (because $\mathscr{P}_{1 / 2}^{\prime}(\infty, a)<0$ ). The following corollary extends this observation to the whole domain $H \in(0,1)$; its proof follows straightforwardly from Proposition C. 1 given in Appendix C.
Corollary 3.4. Suppose that $0<H_{1}<H_{2}<1$.
(i) If $a \in \mathbb{R}$ and $T \in \mathbb{R}_{+}$, then

$$
\mathscr{P}_{H_{2}}(T, a) \leq \mathscr{P}_{H_{1}}\left(T^{H_{2} / H_{1}}, a\right)
$$

(ii) If $a>1$, then

$$
\mathscr{P}_{H_{2}}(\infty, a) \leq \mathscr{P}_{H_{1}}(\infty, a)
$$

Remark 3.5. Using Mathematica software [40] and applying certain simplifications, we are able to calculate the following limit

$$
\begin{equation*}
\lim _{T \rightarrow \infty} \lim _{H \rightarrow 1 / 2} \frac{\mathscr{H}_{H}(T) / T-\mathscr{H}_{1 / 2}(T) / T}{H-\frac{1}{2}}=\lim _{T \rightarrow \infty} \frac{\mathscr{P}_{1 / 2}^{\prime}(T, 1)}{T}=-2 \gamma_{\mathrm{E}} \tag{3.3}
\end{equation*}
$$

It is noted that the result above is very much related to the result established in [13] on the derivative of Pickands constant at $H=\frac{1}{2}$, that is

$$
\begin{equation*}
\lim _{H \rightarrow 1 / 2} \lim _{T \rightarrow \infty} \frac{\mathscr{H}_{H}(T) / T-\mathscr{H}_{1 / 2}(T) / T}{H-\frac{1}{2}}=\lim _{H \rightarrow 1 / 2} \frac{\mathscr{H}_{H}-\mathscr{H}_{1 / 2}}{H-\frac{1}{2}}=-2 \gamma_{\mathrm{E}} \tag{3.4}
\end{equation*}
$$

see also Eq. (1.3). The difference between Eq. (3.3) and Eq. (3.4) is the order of the limit operations.

## 4 Mandelbrot \& van Ness' fractional Brownian field

In this section we provide some properties of Mandelbrot \& van Ness' field, which play important role in the proofs of results given in Section 3.

In contrast to the definition of $B_{H}(t)$, which we have given at the beginning of Section 1, the definition in (2.3) provides an additional coupling between fBms with different values of $H$. In fact, we will view the process $\left\{X_{H}(t),(H, t) \in(0,1) \times \mathbb{R}_{+}\right\}$ as a centered Gaussian field and refer to it as Mandelbrot \& van Ness' field (MvN). Realizations of the MvN field are trajectories (surfaces) $(H, t) \mapsto X_{H}(t)$; using the standard rules of $\mathcal{L}^{2}$ theory of stochastic integrals, we find that

$$
\begin{align*}
\operatorname{Cov}\left(X_{H}(t), X_{H^{\prime}}\left(t^{\prime}\right)\right)= & \int_{-\infty}^{0}\left[(t-s)^{H-\frac{1}{2}}-(-s)^{H-\frac{1}{2}}\right]\left[\left(t^{\prime}-s\right)^{H^{\prime}-\frac{1}{2}}-(-s)^{H^{\prime}-\frac{1}{2}}\right] \mathrm{d} s  \tag{4.1}\\
& +\int_{0}^{t \wedge t^{\prime}}(t-s)^{H-\frac{1}{2}}\left(t^{\prime}-s\right)^{H^{\prime}-\frac{1}{2}} \mathrm{~d} s
\end{align*}
$$

for any $H, H^{\prime} \in(0,1)$ and $t, t^{\prime} \in \mathbb{R}_{+}$. Explicit value of covariance function above was found in [38, Theorem 4.1]. While it is a well-known fact that fBm is self-similar, the same holds true for the MvN field, i.e. for any $c>0$ we have

$$
\begin{equation*}
\left\{B_{H}(c t):(H, t) \in(0,1) \times \mathbb{R}_{+}\right\} \stackrel{\mathrm{d}}{=}\left\{c^{H} B_{H}(t):(H, t) \in(0,1) \times \mathbb{R}_{+}\right\} \tag{4.2}
\end{equation*}
$$

where ' ${ }^{\text {d }}$ ' stands for the equality of finite-dimensional distributions, see e.g. [36, Theorem 2.1(c)]; this can also be seen by a direct calculation using (4.1). It has been shown that there exists a continuous modification of the MvN field, see [29, Theorem 4].

We remark that for the purposes of this paper the domain of MvN field is $(H, t) \in$ $(0,1) \times \mathbb{R}_{+}$, however, we note that it can be extended to $(H, t) \in(0,1) \times \mathbb{R}$.

Next, for any $n \in \mathbb{Z}_{+}$, where $\mathbb{Z}_{+}$is the set of non-negative integers, we define the $n$th derivative of MvN field (with respect to the Hurst parameter), to be the stochastic process $\left\{X_{H}^{(n)}(t):(H, t) \in(0,1) \times \mathbb{R}_{+}\right\}$, where

$$
\begin{equation*}
X_{H}^{(n)}(t)=\int_{-\infty}^{0} f_{H}^{(n)}(t, s) \mathrm{d} B(s)+\int_{0}^{t} g_{H}^{(n)}(t, s) \mathrm{d} B(s) \tag{4.3}
\end{equation*}
$$

for any $(H, t) \in(0,1) \times \mathbb{R}_{+}$, with $f_{H}(t, s):=(t-s)^{H-\frac{1}{2}}-(-s)^{H-\frac{1}{2}}, g_{H}(t, s):=(t-s)^{H-\frac{1}{2}}$, and

$$
\begin{align*}
f_{H}^{(n)}(t, s) & :=\frac{\partial^{n}}{\partial H^{n}} f_{H}(t, s)=\log ^{n}(t-s)(t-s)^{H-\frac{1}{2}}-\log ^{n}(-s)(-s)^{H-\frac{1}{2}},  \tag{4.4}\\
g_{H}^{(n)}(t, s) & :=\frac{\partial^{n}}{\partial H^{n}} g_{H}(t, s)=\log ^{n}(t-s)(t-s)^{H-\frac{1}{2}} .
\end{align*}
$$

We follow the convention that $\frac{\partial^{0}}{\partial x^{0}} h(x):=h(x)$ for any function $h(x)$. In particular, for $n=0$, the definition (4.3) is equivalent to the original MvN field (2.1), i.e. $X_{H}^{(0)}=X_{H}$. Since the case $H=\frac{1}{2}$ will be used particularly often throughout this manuscript, we write $X^{(n)}(\cdot):=X_{1 / 2}^{(n)}(\cdot)$ for brevity. Definition (4.3) can be found in [37].

Let us emphasize that all derivatives of the MvN field live on the same probability space, and in fact they are jointly Gaussian. The covariance between random variables $X_{H}^{(n)}(t)$, and $X_{H^{\prime}}^{\left(n^{\prime}\right)}\left(t^{\prime}\right)$ can be found in analogous way to (4.1).

Similarly to the previous section, the derivatives of MvN field also have their PWZ representation. For $n \in \mathbb{Z}_{+}$we define a Paley-Wiener-Zygmund (PWZ) representation of the MvN field and its derivatives $\left\{\widetilde{X}_{H}^{(n)}(t):(H, t) \in(0,1) \times \mathbb{R}_{+}\right\}$, where

$$
\begin{align*}
\widetilde{X}_{H}^{(n)}(t): & =-\int_{-\infty}^{0}\left[\frac{\partial}{\partial s} f_{H}^{(n)}(t, s)\right] B(s) \mathrm{d} s+g_{H}^{(n)}(t, 0) B(t) \\
& +\int_{0}^{t}\left[\frac{\partial}{\partial s} g_{H}^{(n)}(t, s)\right](B(t)-B(s)) \mathrm{d} s, \tag{4.5}
\end{align*}
$$

for any $(H, t) \in(0,1) \times \mathbb{R}_{+}$, where

$$
\begin{align*}
\frac{\partial}{\partial s} f_{H}^{(n)}(t, s) & =\frac{\partial}{\partial s} g_{H}^{(n)}(t, s)-\frac{\partial}{\partial s} g_{H}^{(n)}(0, s) \\
\frac{\partial}{\partial s} g_{H}^{(n)}(t, s) & =-n \log ^{n-1}(t-s)(t-s)^{H-\frac{3}{2}}-\left(H-\frac{1}{2}\right) \log ^{n}(t-s)(t-s)^{H-\frac{3}{2}} \tag{4.6}
\end{align*}
$$

Again, we recognize that when $n=0$, the definition (4.5) is equivalent to (2.2), i.e. $\widetilde{X}_{H}^{(0)}=\widetilde{X}_{H}$.

The following proposition justifies calling the PWZ, the equivalent representation of MvN field.
Proposition 4.1. For all $n \in \mathbb{Z}_{+}$, the field $\left\{\widetilde{X}_{H}^{(n)}(t):(H, t) \in(0,1) \times \mathbb{R}_{+}\right\}$is a continuous modification of $\left\{X_{H}^{(n)}(t):(H, t) \in(0,1) \times \mathbb{R}_{+}\right\}$.

The result in Proposition 4.1 is a direct consequence of the considerations in Appendix A (Lemma A. 1 in particular) and the proposition below.

Proposition 4.2. For all $n \in \mathbb{Z}_{+}$,

$$
\mathbb{P}\left((H, t) \mapsto \widetilde{X}_{H}^{(n)}(t) \text { is a continuous mapping for all }(H, t) \in(0,1) \times \mathbb{R}_{+}\right)=1
$$

We note that PWZ representation of fractional Brownian motion is defined in terms of Riemann integrals, as opposed to the MvN representation, which is defined through $\mathcal{L}^{2}$ stochastic integrals; see Appendix A.

The rest of this section is devoted to showing various fundamental properties of the PWZ field; due to Proposition 4.1, these results carry over to the MvN field and its derivatives. All the proofs are deferred to Appendix B.

In a special case $n=0$, Proposition 4.2 was proven by [29]. The following result justifies calling processes $X_{H}^{(n)}(t)$ for $n>0$ the derivatives with respect to the Hurst parameter.
Proposition 4.3. For all $n \in \mathbb{Z}_{+}$,

$$
\mathbb{P}\left(\lim _{\Delta \rightarrow 0} \frac{\widetilde{X}_{H+\Delta}^{(n)}(t)-\widetilde{X}_{H}^{(n)}(t)}{\Delta}=\widetilde{X}_{H}^{(n+1)}(t) \text { for all }(H, t) \in(0,1) \times \mathbb{R}_{+}\right)=1
$$

In is noted that weaker versions of Proposition 4.2 and 4.3 were derived in [37].
Thanks to Proposition 4.3, it makes sense to write $\frac{\partial^{n}}{\partial H^{n}} X_{H}(t)=X_{H}^{(n)}(t)$. In the following, for any $k \in \mathbb{N}$ let: $a_{1}, \ldots, a_{k} \in \mathbb{R}, n_{1}, \ldots, n_{k} \in \mathbb{Z}_{+}, H_{1}, \ldots, H_{k} \in(0,1)$ and define

$$
\begin{equation*}
\eta(t):=\sum_{i=1}^{k} a_{i} \cdot X_{H_{i}}^{\left(n_{i}\right)}(t) . \tag{4.7}
\end{equation*}
$$

Proposition 4.4. The Gaussian process $\left\{\eta(t): t \in \mathbb{R}_{+}\right\}$defined in (4.7) is centered and it has stationary increments.

For $N \in \mathbb{Z}_{+}, H \in(0,1)$ we define the Taylor sum remainder

$$
\begin{equation*}
R_{H}(t ; N):=X_{H}(t)-\sum_{n=0}^{N-1} \frac{X^{(n)}(t)}{n!} \cdot\left(H-\frac{1}{2}\right)^{n} \tag{4.8}
\end{equation*}
$$

It is noted that $R_{H}(t ; 0)=X_{H}(t)$ and, according to Proposition 4.4, $\left\{R_{H}(t ; N) ; t \in \mathbb{R}_{+}\right\}$is a centered Gaussian process with stationary increments. Its variance function can also be proven to be Hölder continuous; see Lemma 4.5 below.
Lemma 4.5. Let $0<\underline{H}<\bar{H}<1$ be such that $\frac{1}{2} \in[\underline{H}, \bar{H}]$. Then, for any $\varepsilon>0$ and $N \in \mathbb{Z}_{+}$, there exists a constant $C$ such that

$$
\mathbb{E}\left|R_{H}(t ; N)-R_{H}(s ; N)\right|^{2} \leq C\left(H-\frac{1}{2}\right)^{2 N} \cdot\left(|t-s|^{2 \underline{H}-\varepsilon}+|t-s|^{2 \bar{H}+\varepsilon}\right)
$$

for all $t, s>0$ and $H \in[\underline{H}, \bar{H}]$.
Remark 4.6 (Extension to normalized MvN field). Let $B_{H}(t)=B_{H}^{(0)}(t)=D(H) \widetilde{X}_{H}(t)$ and $B_{H}^{(n)}(t):=\frac{\partial}{\partial H^{n}} B_{H}(t)$. Noticing that $D(H)$ is a smooth function, for all $n \in \mathbb{Z}_{+}$we have

$$
\mathbb{P}\left(\lim _{\Delta \rightarrow 0} \frac{B_{H+\Delta}^{(n)}(t)-B_{H}^{(n)}(t)}{\Delta}=B_{H}^{(n+1)}(t) \text { for all }(H, t) \in(0,1) \times \mathbb{R}_{+}\right)=1
$$

Moreover, by a simple application of Leibniz formula for the $n$th derivative of product of functions we find that

$$
\begin{equation*}
B_{H}^{(n)}(t)=\sum_{k=0}^{n}\binom{n}{k} D^{(n-k)}(H) \cdot \widetilde{X}_{H}^{(k)}(t) \tag{4.9}
\end{equation*}
$$

The value of $D^{(n)}(H)$ at $H=\frac{1}{2}$ for all $n \in \mathbb{Z}_{+}$can be found by direct calculation; for example

$$
\begin{equation*}
D^{(0)}\left(\frac{1}{2}\right)=1, \quad D^{(1)}\left(\frac{1}{2}\right)=1, \quad D^{(2)}\left(\frac{1}{2}\right)=-1-\frac{\pi^{2}}{3}, \quad D^{(3)}\left(\frac{1}{2}\right)=3-\pi^{2}-6 \zeta(3), \tag{4.10}
\end{equation*}
$$

where $\zeta(\cdot)$ is the Riemann-zeta function.

## 5 Proof of Theorem 3.1 and Theorem 3.2

Before proving Theorem 3.1 and Theorem 3.2 we need to develop some preliminary results.

Recall the definition of $\tau_{H}(T, a)$ and $\tau_{H}^{*}(T, a)$ from Eq. (2.4). In the following proposition we establish that the all-time suprema locations $\tau_{H}(\infty, a)$ and $\tau_{H}^{*}(\infty, a)$ on the MvN field are uniformly bounded with probability growing to 1.
Proposition 5.1. Let $0<\underline{H}<\bar{H}<1$ and $a>0$. Then there exist $C, \gamma>0$ such that
(i) $\mathbb{P}\left(\sup _{H \in[\underline{H}, \bar{H}]} \tau_{H}(\infty, a)>T\right) \leq C e^{-\gamma T^{2-2 \bar{H}}}$,
(ii) $\mathbb{P}\left(\sup _{H \in[\underline{H}, \bar{H}]} \tau_{H}^{*}(\infty, a)>T\right) \leq C e^{-\gamma T^{2 \underline{H}}}$
for all $T$ sufficiently large.
Proof. Since the proofs of (i) and (ii) follow by the same idea, we focus only on (ii). Observe that, for $T \geq 1$,

$$
\begin{align*}
\mathbb{P}\left(\sup _{H \in[\underline{H}, \bar{H}]} \tau_{H}^{*}(\infty, a)>T\right) & \leq \mathbb{P}\left(\sup _{H \in[\underline{H}, \bar{H}, t \geq T} B_{H}(t)-a t^{2 H}>0\right) \\
& =\mathbb{P}\left(\sup _{H \in[\underline{H}, \bar{H}], t \geq T} \frac{B_{H}(t)}{t^{2 H}}>a\right) \\
& \leq \sum_{k=0}^{\infty} \mathbb{P}\left(\sup _{H \in\left[\underline{H}, \bar{H}, t \in\left[2^{k} T, 2^{k+1} T\right]\right.} \frac{B_{H}(t)}{t^{2 H}}>a\right) \\
& =\sum_{k=0}^{\infty} \mathbb{P}\left(\sup _{H \in[\underline{H}, \bar{H}], t \in[1,2]} \frac{B_{H}\left(2^{k} T t\right)}{\left(2^{k} T t\right)^{2 H}}>a\right) \\
& =\sum_{k=0}^{\infty} \mathbb{P}\left(\sup _{H \in[\underline{H}, \bar{H}], t \in[1,2]} \frac{2^{H k} T^{H}}{2^{2 H k} T^{2 H}} \frac{B_{H}(t)}{t^{2 H}}>a\right)  \tag{5.1}\\
& \leq \sum_{k=0}^{\infty} \mathbb{P}\left(\sup _{H \in[\underline{H}, \bar{H}], t \in[1,2]} \frac{B_{H}(t)}{t^{2 H}}>a 2^{\underline{H} k} T^{H}\right),
\end{align*}
$$

where (5.1) follows from self-similarity of the MvN field, cf. (4.2). By the continuity of the MvN field, in view of Borell inequality (see, e.g. [1, Theorem 2.1]), we know that

$$
\mathbb{E}\left(\sup _{H \in[\underline{H}, \bar{H}], t \in[1,2]} \frac{B_{H}(t)}{t^{2 H}}\right)<\infty
$$

and for sufficiently large $T$, using that $\sup _{H \in[\underline{H}, \bar{H}], t \in[1,2]} \operatorname{Var}\left(\frac{B_{H}(t)}{t^{2 H}}\right)=1$,

$$
\mathbb{P}\left(\sup _{H \in[\underline{H}, \bar{H}], t \in[1,2]} \frac{B_{H}(t)}{t^{2 H}}>a 2^{\underline{H} k} T^{\underline{H}}\right) \leq 2 \exp \left(-\frac{a^{2} 2^{2 \underline{H} k} T^{2} \underline{H}}{8}\right)
$$

Thus, there exists $C>0$ such that for sufficiently large $T$

$$
\sum_{k=0}^{\infty} \mathbb{P}\left(\sup _{H \in[\underline{H}, \bar{H}], t \in[1,2]} \frac{B_{H}(t)}{t^{2 H}}>a 2^{\underline{H} k} T^{\underline{H}}\right) \leq C \exp \left(-\frac{a^{2}}{8} T^{2 \underline{H}}\right)
$$

This completes the proof.
Proposition 5.2. If $a \in \mathbb{R}$ and $T \in \mathbb{R}_{+}$or $a>0$ and $T=\infty$, then
(i) $\lim _{H \rightarrow H^{\prime}} \tau_{H}(T, a)=\tau_{H^{\prime}}(T, a)$, a.s.
(ii) $\lim _{H \rightarrow H^{\prime}} \tau_{H}^{*}(T, a)=\tau_{H^{\prime}}^{*}(T, a)$, a.s.
for any $H^{\prime} \in(0,1)$.
Proof. In view of Proposition 4.2, the random bivariate function $(H, t) \mapsto X_{H}(t)$ is continuous almost surely, this implies that also, for any fixed $a \in \mathbb{R}$, the functions $(H, t) \mapsto Y_{H}(t ; a)$, and $(H, t) \mapsto Y_{H}^{*}(t ; a)$ are continuous. Now, consider the case $a \in$ $\mathbb{R}, T \in \mathbb{R}_{+}$- in this case, (i) and (ii) follow from the fact that $Y_{H}(t ; a) \rightarrow Y_{H^{\prime}}(t ; a)$, $Y_{H}^{*}(t ; a) \rightarrow Y_{H^{\prime}}^{*}(t ; a)$ converge uniformly on $t \in[0, T]$, so the argmax functionals must also converge, see e.g. [33, Lemma 2.9].

We now show that (i) holds also when $a>0$ and $T=\infty$. Let $\mathcal{A} \subset(0,1)$ be any compact interval containing $H^{\prime}$ and for $n \in \mathbb{N}$ let $A_{n}:=\left\{\sup _{H \in \mathcal{A}} \tau_{H}(\infty, a)<n\right\}$. Moreover, we have $\cup_{n} A_{n}=\left\{\sup _{H \in \mathcal{A}} \tau_{H}(\infty, a)<\infty\right\}$, which according to Proposition 5.1, is a set of full measure. We thus can write

$$
\begin{aligned}
\mathbb{P}\left(\lim _{H \rightarrow H^{\prime}} \tau_{H}(\infty, a)=\tau_{H^{\prime}}(\infty, a)\right) & =\mathbb{P}\left(\lim _{H \rightarrow H^{\prime}} \tau_{H}(\infty, a)=\tau_{H^{\prime}}(\infty, a) ; \cup_{n=1}^{\infty} A_{n}\right) \\
& =\mathbb{P}\left(\bigcup_{n=1}^{\infty}\left\{\lim _{H \rightarrow H^{\prime}} \tau_{H}(\infty, a)=\tau_{H^{\prime}}(\infty, a) ; A_{n}\right\}\right) \\
& =\lim _{n \rightarrow \infty} \mathbb{P}\left(\lim _{H \rightarrow H^{\prime}} \tau_{H}(\infty, a)=\tau_{H^{\prime}}(\infty, a) ; A_{n}\right),
\end{aligned}
$$

where in the last line we used the continuity property of probability measures for increasing sets. Finally, we notice that on the event $A_{n}$ we have $\tau_{H}(\infty, a)=\tau_{H}(n, a)$ for all $H \in \mathcal{A}$, thus

$$
\mathbb{P}\left(\lim _{H \rightarrow H^{\prime}} \tau_{H}(\infty, a)=\tau_{H^{\prime}}(\infty, a) ; A_{n}\right)=\mathbb{P}\left(\lim _{H \rightarrow H^{\prime}} \tau_{H}(n, a)=\tau_{H^{\prime}}(n, a) ; A_{n}\right)=\mathbb{P}\left(A_{n}\right)
$$

because we have already established that $\tau_{H}(T, a) \rightarrow \tau_{H^{\prime}}(T, a)$ a.s. for any fixed $T \in \mathbb{R}_{+}$. This concludes the proof of (i) because $\mathbb{P}\left(A_{n}\right) \rightarrow 1$, as $n \rightarrow \infty$. The proof of (ii) is analogous.

Corollary 5.3. Let $0<\underline{H}<\bar{H}<1, N \in \mathbb{Z}_{+}$. If $a \in \mathbb{R}$ and $T \in \mathbb{R}_{+}$or $a>0$ and $T=\infty$, then for any $n \in \mathbb{N}$,
(i) $\sup _{H \in[\underline{H}, \bar{H}]} \mathbb{E}\left|\frac{R_{H}\left(\tau_{H}(T, a) ; N\right)}{\left(H-\frac{1}{2}\right)^{N}}\right|^{n}<\infty$, and
(ii) $\sup _{H \in[\underline{H}, \bar{H}]} \mathbb{E}\left|\frac{R_{H}\left(\tau_{H}^{*}(T, a) ; N\right) \mid}{\left(H-\frac{1}{2}\right)^{N}}\right|^{n}<\infty$.

Proof of Corollary 5.3. From Lemma 4.5 we know that for every $\varepsilon>0$ there exists $C>0$ such that for all $t \in[0,1]$,

$$
\sup _{H \in[H, \bar{H}]} \mathbb{E}\left(R_{H}(t ; N)\right)^{2} \leq C\left(H-\frac{1}{2}\right)^{2 N} \cdot t^{2(\underline{H}-\varepsilon)}
$$

where $\varepsilon>0$ is taken small enough so that $\underline{H}-\varepsilon>0$. Now, again, applying Lemma 4.5 to Lemma C. 3 we find that for every $\varepsilon>0$ and $n \in \mathbb{N}$ there exists a constant $C^{\prime}>0$ such that

$$
\sup _{H \in[\underline{H}, \bar{H}]} \mathbb{E} \sup _{t \in[0,1]}\left|R_{H}(t ; N)\right|^{n} \leq C^{\prime}\left(H-\frac{1}{2}\right)^{n N}\left(1+\mu^{n-1}(\underline{H}-\varepsilon)+\mu^{n}(\underline{H}-\varepsilon)\right),
$$

where $\mu(\underline{H}-\varepsilon)<\infty$. Finally, applying these two bounds and the result in Proposition 5.1(i) to Lemma C. 4 we find that there exists $C^{\prime \prime}>0$ such that

$$
\sup _{H \in[\underline{H}, \bar{H}]} \mathbb{E}\left|R_{H}\left(\tau_{H} ; N\right)\right|^{n} \leq C^{\prime \prime}\left(H-\frac{1}{2}\right)^{n N} .
$$

The same reasoning applies to $\tau_{H}^{*}$, which completes the proof.

### 5.1 Proof of Theorem 3.1

In the following, let $Y_{H}(t):=Y_{H}(t ; a), \tau_{H}:=\tau_{H}(T, a)$, and $Y(t):=Y_{1 / 2}(t), \tau:=\tau_{1 / 2}$. We have

$$
Y_{H}\left(\tau_{H}\right)-Y(\tau)=\left(Y_{H}\left(\tau_{H}\right)-Y_{H}(\tau)\right)+\left(Y_{H}(\tau)-Y(\tau)\right)
$$

We split the proof into three parts. In the first part of the proof we show that

$$
\begin{equation*}
\lim _{H \rightarrow 1 / 2} \mathbb{E}\left(\frac{Y_{H}\left(\tau_{H}\right)-Y_{H}(\tau)}{H-\frac{1}{2}}\right)=0 \tag{5.2}
\end{equation*}
$$

In the second part of the proof we show that

$$
\begin{equation*}
\lim _{H \rightarrow 1 / 2} \mathbb{E}\left(\frac{Y_{H}(\tau)-Y(\tau)}{H-\frac{1}{2}}\right)=\mathbb{E}\left(\left.\frac{\partial}{\partial H} Y_{H}(\tau)\right|_{H=1 / 2}\right) . \tag{5.3}
\end{equation*}
$$

Finally, in the last part of the proof we show that the claim in Theorem 3.1 holds.
Proof that Eq. (5.2) holds. By definition, $Y_{H}(t)=D(H) X_{H}(t)-a t$, so for any $t>0$, we have

$$
Y_{H}(t)=a t(D(H)-1)+D(H) Y(t)+D(H) X^{(1)}(t)\left(H-\frac{1}{2}\right)+D(H) R_{H}(t ; 2)
$$

where $R_{H}(t ; 2):=X_{H}(t)-\left[X(t)+X^{(1)}(t)\left(H-\frac{1}{2}\right)\right]$; see (4.8). Furthermore,

$$
\begin{aligned}
Y_{H}\left(\tau_{H}\right)-Y_{H}(\tau) & =a(D(H)-1)\left(\tau_{H}-\tau\right)+D(H)\left[Y\left(\tau_{H}\right)-Y(\tau)\right] \\
& +D(H)\left(H-\frac{1}{2}\right)\left[X^{(1)}\left(\tau_{H}\right)-X^{(1)}(\tau)\right]+D(H)\left[R_{H}\left(\tau_{H} ; 2\right)-R_{H}(\tau ; 2)\right]
\end{aligned}
$$

Since $Y(\tau)-Y\left(\tau_{H}\right) \geq 0$ and $Y_{H}\left(\tau_{H}\right)-Y_{H}(\tau) \geq 0$, then

$$
\begin{align*}
& 0 \leq \frac{Y_{H}\left(\tau_{H}\right)-Y_{H}(\tau)}{H-\frac{1}{2}} \leq a\left(\tau_{H}-\tau\right) \cdot \frac{D(H)-1}{H-\frac{1}{2}}+D(H)\left[X^{(1)}\left(\tau_{H}\right)-X^{(1)}(\tau)\right] \\
& +D(H) \cdot \frac{R_{H}\left(\tau_{H} ; 2\right)-R_{H}(\tau ; 2)}{H-\frac{1}{2}}=: J_{1}(H)+J_{2}(H)+J_{3}(H) \tag{5.4}
\end{align*}
$$

So we need to show that $\mathbb{E} J_{i}(H) \rightarrow 0$, as $H \rightarrow 0$ for $i \in\{1,2,3\}$. According to Proposition 5.2(i) we have $J_{1}(H) \rightarrow 0$ a.s. In view of the de la Vallèe Poussin theorem, in order to show $\mathbb{E} J_{1}(H) \rightarrow 0$, it suffices to show that $\sup _{H \in[1 / 3,2 / 3]} \mathbb{E} J_{1}^{2}<\infty$. According to Proposition 5.1, for any $0<\underline{H}<\frac{1}{2}<\bar{H}<1$ there exist positive constants $C, \gamma, \beta$ such that $\mathbb{P}\left(\tau_{H}>T\right) \leq C e^{-\gamma T^{\beta}}$, therefore

$$
\mathbb{E} \tau_{H}^{2}=\int_{0}^{\infty} 2 t \mathbb{P}\left(\tau_{H}>t\right) \mathrm{d} t \leq 2 C \int_{0}^{\infty} t e^{-\gamma t^{\beta}} \mathrm{d} t
$$

for all $H \in[\underline{H}, \bar{H}]$. Since the right-hand-side is finite and independent of $H$, then $\sup _{H \in[\underline{H}, \bar{H}]} \mathbb{E} \tau_{H}^{2}<\infty$ and $\mathbb{E} J_{1}(H) \rightarrow 0$.

Now, according to Proposition 4.2 and Proposition 5.2(i) we have $J_{2}(H) \rightarrow 0$ a.s. Moreover, combining Proposition 5.1(i) with Lemma C.4, and the fact that all moments of the supremum $\sup _{t \in[0,1]} X^{(1)}(t)$ exist, we find that the second moments, $\mathbb{E}\left[X^{(1)}\left(\tau_{H}\right)\right]^{2}$, $\mathbb{E}\left[X^{(1)}(\tau)\right]^{2}$ are uniformly bounded for all $H$ close enough to $\frac{1}{2}$; again, the de la Vallèe Poussin theorem implies that $\mathbb{E} J_{2}(H) \rightarrow 0$.

Finally, we observe that

$$
\mathbb{E}\left|J_{3}(H)\right| \leq \frac{\mathbb{E}\left|R_{H}\left(\tau_{H} ; 2\right)\right|}{H-\frac{1}{2}}+\frac{\mathbb{E}\left|R_{H}(\tau ; 2)\right|}{H-\frac{1}{2}}
$$

Due to Corollary 5.3, both terms above tend to 0 , as $H \rightarrow \frac{1}{2}$, therefore $\mathbb{E} J_{3}(H) \rightarrow 0$.
Proof that Eq. (5.3) holds. By virtue of Proposition 4.3, it is clear that

$$
\left.\frac{Y_{H}(\tau)-Y(\tau)}{H-\frac{1}{2}} \rightarrow \frac{\partial}{\partial H} Y_{H}(\tau)\right|_{H=1 / 2} \text { a.s. }
$$

Moreover, we recognize that $Y_{H}(\tau)-Y(\tau)=R_{H}(\tau ; 1)$, see the definition in (4.8), so in view of the de la Vallèe Poussin theorem, in order to show that (5.3) holds, it is enough to show that

$$
\sup _{H \in[1 / 3,2 / 3]} \mathbb{E}\left(\frac{R_{H}(\tau ; 1)}{H-\frac{1}{2}}\right)^{2}<\infty
$$

which follows from Corollary 5.3(i).
Proof that Eq. (3.1) holds. Notice that for any $t \in \mathbb{R}_{+}$it holds that

$$
\frac{\partial}{\partial H} Y_{H}(t)=\frac{\partial}{\partial H}\left(B_{H}(t)-a t\right)=B_{H}^{(1)}(t)
$$

Combining the results in (5.2) and (5.3), the Leibniz formula derived in (4.9), and the fact that $D^{\prime}(1 / 2)=1$ (cf. (4.10)) we obtain

$$
\begin{equation*}
\mathscr{M}_{1 / 2}^{\prime}(T, a)=\mathbb{E}\left[X(\tau)+X^{(1)}(\tau)\right] \tag{5.5}
\end{equation*}
$$

Now, using the PWZ representation for $X^{(1)}$ in (4.5) we find that

$$
X^{(1)}(\tau)=\int_{-\infty}^{0}\left((\tau-s)^{-1}-(-s)^{-1}\right) B(s) \mathrm{d} s+\log (\tau) B(\tau)-\int_{0}^{\tau} \frac{B(\tau)-B(s)}{\tau-s} \mathrm{~d} s
$$

Since $\tau$ is independent of $\{B(s): s<0\}$ then the expected value of the first integral is 0 , and

$$
\begin{aligned}
\mathbb{E} X^{(1)}(\tau) & =\mathbb{E}\left[\log (\tau) B(\tau)-\int_{0}^{\tau} \frac{B(\tau)-B(s)}{\tau-s} \mathrm{~d} s\right] \\
& =\mathbb{E}\left[\log (\tau)(B(\tau)-a \tau)+a \tau \log (\tau)-\int_{0}^{\tau} \frac{(B(\tau)-a \tau)-(B(s)-a s)}{\tau-s} \mathrm{~d} s-a \tau\right] \\
& =\mathbb{E}\left[\log (\tau) Y(\tau)+a \tau(\log (\tau)-1)-\int_{0}^{\tau} \frac{Y(\tau)-Y(s)}{\tau-s} \mathrm{~d} s\right]
\end{aligned}
$$

Now, we recognize that

$$
\begin{aligned}
\mathbb{E}\left(\int_{0}^{\tau} \frac{Y(\tau)-Y(s)}{\tau-s} \mathrm{~d} s\right) & =\mathbb{E}\left(\mathbb{E}\left(\left.\int_{0}^{\tau} \frac{Y(\tau)-Y(s)}{\tau-s} \mathrm{~d} s \right\rvert\,(\tau, Y(\tau))\right)\right) \\
& =\mathbb{E}(I(\tau, Y(\tau)))
\end{aligned}
$$

with $I(t, y)$ defined in (2.8). Since $\mathbb{E} X(\tau)=\mathbb{E}(Y(\tau)+a \tau)$, then from (5.5) we obtain

$$
\mathscr{M}_{1 / 2}^{\prime}(T, a)=\mathbb{E}[Y(\tau)(1+\log (\tau))+a \tau \log (\tau)-I(\tau, Y(\tau))] .
$$

The claim now follows because $(\tau, Y(\tau))$ has joint density $p(t, y ; T, a)$, cf. Eq. (2.5).

### 5.2 Proof of Theorem 3.2

In the following, for brevity let $Z_{H}(t):=\sqrt{2} B_{H}(t)-a t^{2 H}, \tau_{H}^{*}:=\tau_{H}^{*}\left(T, \frac{a}{\sqrt{2}}\right)$, so that $\sup _{t \in[0, T]} Z_{H}(t)=Z_{H}\left(\tau_{H}^{*}\right)$. Additionally we denote $Z(t):=Z_{1 / 2}(t), \tau^{*}:=\tau_{1 / 2}^{*}$. Notice that

$$
e^{Z_{H}\left(\tau_{H}^{*}\right)}-e^{Z\left(\tau^{*}\right)}=\left(e^{Z_{H}\left(\tau_{H}^{*}\right)}-e^{Z_{H}\left(\tau^{*}\right)}\right)+\left(e^{Z_{H}\left(\tau^{*}\right)}-e^{Z\left(\tau^{*}\right)}\right) .
$$

We split the proof into three parts. First part of the proof is to show that

$$
\begin{equation*}
\lim _{H \rightarrow 1 / 2} \mathbb{E}\left(\frac{e^{Z_{H}\left(\tau_{H}^{*}\right)}-e^{Z_{H}\left(\tau^{*}\right)}}{H-\frac{1}{2}}\right)=0 . \tag{5.6}
\end{equation*}
$$

In the second part of the proof we show that

$$
\begin{equation*}
\lim _{H \rightarrow 1 / 2} \mathbb{E}\left(\frac{e^{Z_{H}\left(\tau^{*}\right)}-e^{Z\left(\tau^{*}\right)}}{H-\frac{1}{2}}\right)=\mathbb{E}\left(\left.\frac{\partial}{\partial H} e^{Z_{H}\left(\tau^{*}\right)}\right|_{H=1 / 2}\right) . \tag{5.7}
\end{equation*}
$$

Finally, in the last part of the proof we show that the claim of Theorem 3.2 holds.
Proof that Eq. (5.6) holds. Since $Z_{H}\left(\tau_{H}^{*}\right)-Z_{H}\left(\tau^{*}\right) \geq 0$, then, using the mean value theorem we find that

$$
\begin{equation*}
\mathbb{E}\left[e^{Z_{H}\left(\tau_{H}^{*}\right)}-e^{Z_{H}\left(\tau^{*}\right)}\right] \leq \mathbb{E}\left[\left(Z_{H}\left(\tau_{H}^{*}\right)-Z_{H}\left(\tau^{*}\right)\right) \cdot e^{Z_{H}\left(\tau_{H}^{*}\right)}\right], \tag{5.8}
\end{equation*}
$$

so it suffices to show that the bound above converges to 0 .
By definition $Z_{H}(t)=\sqrt{2} D(H) X_{H}(t)-a t^{2 H}$, so for any $t>0$, we have

$$
Z_{H}(t)=a\left(D(H) t-t^{2 H}\right)+D(H) Z(t)+\sqrt{2} D(H)\left(X^{(1)}(t)\left(H-\frac{1}{2}\right)+R_{H}(t ; 2)\right)
$$

where $R_{H}(t ; 2):=X_{H}(t)-\left[X(t)+X^{(1)}(t)\left(H-\frac{1}{2}\right)\right]$ was defined in (4.8). Furthermore,

$$
\begin{aligned}
& Z_{H}\left(\tau_{H}^{*}\right)-Z_{H}\left(\tau^{*}\right)=a\left[D(H)\left(\tau_{H}^{*}-\tau^{*}\right)-\left(\left(\tau_{H}^{*}\right)^{2 H}-\left(\tau^{*}\right)^{2 H}\right)\right]+D(H)\left[Z\left(\tau_{H}^{*}\right)-Z\left(\tau^{*}\right)\right] \\
& +\sqrt{2} D(H)\left(H-\frac{1}{2}\right)\left[X^{(1)}\left(\tau_{H}^{*}\right)-X^{(1)}\left(\tau^{*}\right)\right]+\sqrt{2} D(H)\left[R_{H}\left(\tau_{H}^{*} ; 2\right)-R_{H}\left(\tau^{*} ; 2\right)\right]
\end{aligned}
$$

Now, for $H \in(0,1), t \in \mathbb{R}_{+}$we define

$$
U_{H}(t):=D(H) \cdot \frac{t-t^{2 H}}{H-\frac{1}{2}}+t^{2 H} \cdot \frac{D(H)-1}{H-\frac{1}{2}} .
$$

Since $Z\left(\tau^{*}\right)-Z\left(\tau_{H}^{*}\right) \geq 0$ and $Z_{H}\left(\tau_{H}^{*}\right)-Z_{H}\left(\tau^{*}\right) \geq 0$, then

$$
0 \leq \frac{Z_{H}\left(\tau_{H}^{*}\right)-Z_{H}\left(\tau^{*}\right)}{H-\frac{1}{2}} \leq a J_{1}(H)+\sqrt{2} D(H)\left(J_{2}(H)+J_{3}(H)\right)
$$

where

$$
J_{1}(H):=U_{H}\left(\tau_{H}^{*}\right)-U_{H}\left(\tau^{*}\right), \quad J_{2}(H):=X^{(1)}\left(\tau_{H}^{*}\right)-X^{(1)}\left(\tau^{*}\right)
$$

and

$$
J_{3}(H):=\frac{R_{H}\left(\tau_{H}^{*} ; 2\right)-R_{H}\left(\tau^{*} ; 2\right)}{H-\frac{1}{2}}
$$

The proofs that $\mathbb{E} J_{i}(H) \rightarrow 0, i \in\{2,3\}$, as $H \rightarrow \frac{1}{2}$ are analogous to the corresponding statements in the proof of Theorem 3.1, cf. Eq. (5.4). Thus, we show only that $\mathbb{E} J_{1}(H) \rightarrow$ 0 , as $H \rightarrow \frac{1}{2}$. We have

$$
\lim _{H \rightarrow 1 / 2} U_{H}(t)=-2 t \log (t)+t=: U(t)
$$

and the convergence is uniform for any compact subset of $\mathbb{R}_{+}$. Moreover,

$$
\left|U_{H}\left(\tau_{H}^{*}\right)-U_{H}\left(\tau^{*}\right)\right| \leq\left|U_{H}\left(\tau_{H}^{*}\right)-U\left(\tau_{H}^{*}\right)\right|+\left|U\left(\tau_{H}^{*}\right)-U\left(\tau^{*}\right)\right|+\left|U\left(\tau^{*}\right)-U_{H}\left(\tau^{*}\right)\right|,
$$

where each of the terms of the sum above converges to 0 due to uniform convergence $U_{H}(t) \rightarrow U_{H}(t)$, continuity of $U$ and the fact that $\tau_{H}^{*} \rightarrow \tau^{*}$ a.s. (see Proposition 5.2(ii)), so we have $J_{1}(H) \rightarrow 0$ a.s. Using the mean value theorem we find that, with $\underline{H}:=\min \left\{\frac{1}{2}, H\right\}$, $\bar{H}:=\max \left\{\frac{1}{2}, H\right\}$,

$$
\begin{equation*}
\left|\frac{t-t^{2 H}}{H-\frac{1}{2}}\right| \leq \sup _{\widetilde{H} \in[\underline{H}, \bar{H}]}\left|2 \log (t) t^{2 \widetilde{H}}\right| \leq 2|\log (t)|\left(1+t^{2 \bar{H}}\right) \tag{5.9}
\end{equation*}
$$

Now, we have that

$$
U_{H}^{2}(t) \leq 2 D^{2}(H)\left(\frac{t-t^{2 H}}{H-\frac{1}{2}}\right)^{2}+2 t^{4 H}\left(\frac{D(H)-1}{H-\frac{1}{2}}\right)^{2}
$$

so using the bound in (5.9) and Proposition 5.1(ii), we find that the second moments $\mathbb{E}\left[U_{H}\left(\tau_{H}^{*}\right)\right]^{2}, \mathbb{E}\left[U_{H}\left(\tau^{*}\right)\right]^{2}<\infty$ are uniformly bounded for all $H$ close enough to $\frac{1}{2}$. In view of the de la Vallèe Poussin theorem, this observation combined with the fact that $J_{1}(H) \rightarrow 0$ a.s. imply that $\mathbb{E} J_{1}(H) \rightarrow 0$.

Proof that Eq. (5.7) holds. By virtue of Proposition 4.3 and (5.2) combined, it is clear that

$$
\left.\frac{e^{Z_{H}\left(\tau^{*}\right)}-e^{Z\left(\tau^{*}\right)}}{H-\frac{1}{2}} \rightarrow \frac{\partial}{\partial H} e^{Z_{H}\left(\tau^{*}\right)}\right|_{H=1 / 2} \text { a.s. }
$$

In view of the de la Vallèe Poussin theorem, it suffices to show that some $1+\varepsilon$ with $\varepsilon>0$, the absolute moment of the pre-limit above is bounded for all $H$ close enough to $1 / 2$. Using the mean value theorem as in (5.8) and applying Hölder's inequality afterwards, we find that

$$
\begin{aligned}
& \left.\mathbb{E}\left(\frac{e^{Z_{H}\left(\tau^{*}\right)}-e^{Z\left(\tau^{*}\right)}}{H-\frac{1}{2}}\right)^{1+\varepsilon} \leq \mathbb{E}\left(\left|\frac{Z_{H}\left(\tau^{*}\right)-Z\left(\tau^{*}\right)}{H-\frac{1}{2}}\right| \cdot \sup _{H \in[\underline{H}, \bar{H}]} \exp \left\{Z_{H}\left(\tau^{*}\right)\right\}\right\}\right)^{1+\varepsilon} \\
& \quad \leq\left(\mathbb{E}\left|\frac{Z_{H}\left(\tau^{*}\right)-Z\left(\tau^{*}\right)}{H-\frac{1}{2}}\right|^{(1+\varepsilon) q}\right)^{1 / q} \cdot\left(\sup _{H \in[\underline{H}, \bar{H}]} \mathbb{E} \exp \left\{(1+\varepsilon) p Z_{H}\left(\tau^{*}\right)\right\}\right)^{1 / p}
\end{aligned}
$$

where $p, q>1,1 / p+1 / q=1$, and $H$ is assumed to be close enough to $\frac{1}{2}$ such that there exist $0<\underline{H}<\frac{1}{2}<\bar{H}<1$ with $H \in[\underline{H}, \bar{H}]$. Now, according to Proposition C.1, if we take $\varepsilon>0, p>1$ small enough to satisfy $(1+\varepsilon) p<a$, then the second term is finite. For the finiteness of the first term, note that $Z_{H}\left(\tau^{*}\right)-Z\left(\tau^{*}\right)=\sqrt{2} R_{H}\left(\tau^{*} ; 1\right)-a\left(\left(\tau^{*}\right)^{2 H}-\tau^{*}\right)$ and therefore, for some $C>0$, we can bound

$$
\left|\frac{Z_{H}\left(\tau^{*}\right)-Z\left(\tau^{*}\right)}{H-\frac{1}{2}}\right|^{(1+\varepsilon) q} \leq C\left(\left|\frac{R_{H}\left(\tau^{*} ; 1\right)}{H-\frac{1}{2}}\right|^{(1+\varepsilon) q}+\left|\frac{\left(\tau^{*}\right)^{2 H}-\tau^{*}}{H-\frac{1}{2}}\right|^{(1+\varepsilon) q}\right)
$$

Using the same argumentation as in the proof of Corollary 5.3(ii), for arbitrarily large $q>1$ we have

$$
\sup _{H \in[\underline{H}, \bar{H}]} \mathbb{E}\left|\frac{R_{H}\left(\tau^{*} ; 1\right)}{H-\frac{1}{2}}\right|^{(1+\varepsilon) q}<\infty .
$$

Using Eq. (5.9) we have

$$
\sup _{H \in[\underline{H}, \bar{H}]} \mathbb{E}\left|\frac{\left(\tau^{*}\right)^{2 H}-\tau^{*}}{H-\frac{1}{2}}\right|^{(1+\varepsilon) q} \leq \mathbb{E}\left(2\left|\log \left(\tau^{*}\right)\right|\left(1+\left(\tau^{*}\right)^{2 \bar{H}}\right)\right)^{(1+\varepsilon) q}<\infty
$$

where the finiteness for arbitrarily large $q>1$ is the result of Proposition 5.1(ii).
Proof that Eq. (3.2) holds. Notice that for any $t \in \mathbb{R}_{+}$it holds that

$$
\begin{aligned}
\frac{\partial}{\partial H} e^{Z_{H}(t)} & =\sqrt{2}\left[\frac{\partial}{\partial H}\left(B_{H}(t)-a t^{2 H}\right)\right] e^{Z_{H}(t)} \\
& =\sqrt{2}\left(B_{H}^{(1)}(t)-\frac{2 a}{\sqrt{2}} \cdot t^{2 H} \log (t)\right) \exp \left(\sqrt{2} B_{H}(t)-a t^{2 H}\right)
\end{aligned}
$$

Combining the results in (5.6) and (5.7), the Leibniz formula derived in (4.9), and the fact that $D^{\prime}\left(\frac{1}{2}\right)=1$, cf. (4.10), we obtain

$$
\mathscr{P}_{1 / 2}^{\prime}(T, a)=\mathbb{E}\left[\sqrt{2}\left(X(\tau)+X^{(1)}(\tau)-\frac{2 a}{\sqrt{2}} \cdot \tau \log (\tau)\right) \exp (\sqrt{2} X(\tau)-a \tau)\right]
$$

The rest of the proof is analogous to the third part of the proof of Theorem 3.1.

## A Paley-Wiener-Zygmund representation of $\mathcal{L}^{2}$ stochastic integrals

One of the core tools used in this contribution is the equivalence between MvN and PWZ representation of the fractional Brownian field and its derivatives introduced in Section 4. In this appendix we explain how to find the Paley-Wiener-Zygmund representation of $\mathcal{L}^{2}$ stochastic integral for a general class of processes.

Define formally a stochastic process $\left\{\xi(t): t \in \mathbb{R}_{+}\right\}$

$$
\begin{equation*}
\xi(t):=P_{1}(t)+P_{2}(t), \quad P_{1}(t):=\int_{-\infty}^{0} \psi(t, s) \mathrm{d} B(s), P_{2}(t):=\int_{0}^{t} \phi(t, s) \mathrm{d} B(s) . \tag{A.1}
\end{equation*}
$$

For a moment we may think about the integrals in $\mathcal{L}^{2}$ sense. The processes $P_{1}$ and $P_{2}$ are well defined if

$$
\begin{equation*}
\int_{-\infty}^{0} \psi(t, s)^{2} \mathrm{~d} s<\infty, \quad \int_{-\infty}^{t} \phi(t, s)^{2} d s<\infty, \quad \text { for all } t \in \mathbb{R}_{+} \tag{A.2}
\end{equation*}
$$

Suppose $\{B(t): t \in \mathbb{R}\}$ is a Brownian motion on a probability space $(\Omega, \mathcal{E}, \mathbb{P})$. Let $T<$ 0 . For $k=1,2, \ldots$ let $M_{k}:=\left|\left\lceil T 2^{k}\right\rceil\right|$. Define integral sums for $P_{1}(t, T):=\int_{T}^{0} \psi(t, s) \mathrm{d} B(s)$ by

$$
\begin{equation*}
P_{k, 1}(t, T):=\sum_{i=1}^{M_{k}-1} \psi\left(t,-\frac{i}{2^{k}}\right)\left(B\left(-\frac{i}{2^{k}}\right)-B\left(-\frac{i+1}{2^{k}}\right)\right) . \tag{A.3}
\end{equation*}
$$

Similarly, set $N_{k}=\left\lfloor t 2^{k}\right\rfloor$ and define integral sums for $P_{2}(t)$ by

$$
\begin{equation*}
P_{k, 2}(t)=\sum_{i=0}^{N_{k}-1} \phi\left(t, \frac{i+1}{2^{k}}\right)\left(B\left(\frac{i+1}{2^{k}}\right)-B\left(\frac{i}{2^{k}}\right)\right) . \tag{A.4}
\end{equation*}
$$

Clearly we have

$$
\begin{equation*}
P_{k, 1}(t, T) \xrightarrow{\mathcal{L}^{2}} P_{1}(t, T)=\int_{T}^{0} \psi(t, s) \mathrm{d} B(s), \quad \text { and } \quad P_{k, 2}(t) \xrightarrow{\mathcal{L}^{2}} P_{2}(t)=\int_{0}^{t} \phi(t, s) \mathrm{d} B(s) . \tag{A.5}
\end{equation*}
$$

We denote $\psi^{\prime}(t, s)=\frac{\partial}{\partial s} \psi(t, s)$ and $\phi^{\prime}(t, s)=\frac{\partial}{\partial s} \phi(t, s)$. In what follows, we introduce the following assumptions
(A) For each $t>0$, the function $s \mapsto \psi(t, s)$ belongs to $C^{1}(-\infty, 0)$ and there exist constants $C, \varepsilon>0$ such that

- $|\psi(t, s)| \leq C|s|^{-1 / 2+\epsilon},\left|\psi^{\prime}(t, s)\right| \leq C|s|^{-3 / 2+\epsilon}$, as $s \uparrow 0$,
- $|\psi(t, T)| \leq C|T|^{-1 / 2-\varepsilon},\left|\psi^{\prime}(t, T)\right| \leq C|T|^{-3 / 2-\varepsilon}$, as $T \downarrow-\infty$.
(B) For each $t>0$, the function $s \mapsto \phi(t, s)$ belongs to $C^{1}[0, t)$ and there exist constants $C, \varepsilon>0$ such that
- $|\phi(t, s)| \leq C|t-s|^{-1 / 2+\varepsilon},\left|\phi^{\prime}(t, s)\right| \leq C|t-s|^{-3 / 2+\epsilon}$, as $s \uparrow t$.

Note that conditions (A) and (B) imply (A.2). Conditions (A) and (B) were formulated having in mind $\psi=f_{H}^{(n)}$ and $\phi=g_{H}^{(n)}$, which were defined in (4.4).
Lemma A.1. Let functions $\psi(t, s), \phi(t, s)$ satisfy the assumptions (A) and (B) respectively. Then, with $\xi(t)$ defined in (A.1), for each $t \geq 0, \xi(t)=\widetilde{P}_{1}(t)+\widetilde{P}_{2}(t)$ a.s., where $\widetilde{P}_{1}(t)=$ $-\int_{-\infty}^{0} \psi^{\prime}(t, s) B(s) d s$ and $\widetilde{P}_{2}(t)=\phi(t, 0) B(t)+\int_{0}^{t} \phi^{\prime}(t, s)(B(t)-B(s)) d s$.

Proof. Set $N_{k}=\left\lfloor 2^{k} t\right\rfloor$. Then, the partial sum $P_{k, 2}(t)$ in (A.4) can be rewritten as follows

$$
\begin{aligned}
& \sum_{i=0}^{N_{k}-1} \phi\left(t, \frac{i+1}{2^{k}}\right)\left(B\left(\frac{i+1}{2^{k}}\right)-B\left(\frac{i}{2^{k}}\right)\right) \\
& =\sum_{i=1}^{N_{k}-1}\left(\phi\left(t, \frac{i}{2^{k}}\right)-\phi\left(t, \frac{i+1}{2^{k}}\right)\right) B\left(\frac{i}{2^{k}}\right)+\phi\left(t, \frac{N_{k}}{2^{k}}\right) B\left(\frac{N_{k}}{2^{k}}\right) \\
& =\sum_{i=1}^{N_{k}-1}\left(\phi\left(t, \frac{i}{2^{k}}\right)-\phi\left(t, \frac{i+1}{2^{k}}\right)\right)\left(B\left(\frac{i}{2^{k}}\right)-B(t)\right)+\phi\left(t, \frac{N_{k}}{2^{k}}\right) B\left(\frac{N_{k}}{2^{k}}\right) \\
& \quad+B(t) \sum_{i=1}^{N_{k}-1}\left(\phi\left(t, \frac{i}{2^{k}}\right)-\phi\left(t, \frac{i+1}{2^{k}}\right)\right) .
\end{aligned}
$$

Using the mean value theorem $\phi\left(t, \frac{i}{2^{k}}\right)-\phi\left(t, \frac{i+1}{2^{k}}\right)=-\frac{1}{2^{k}} \phi^{\prime}\left(t, \theta_{k, i}\right)$, for $\theta_{k, i} \in\left(\frac{i}{2^{k}}, \frac{i+1}{2^{k}}\right)$ and hence the above equals to

$$
\frac{1}{2^{k}} \sum_{i=1}^{N_{k}-1} \phi^{\prime}\left(t, \theta_{k, i}\right)\left(B(t)-B\left(\frac{i}{2^{k}}\right)\right)+B(t) \phi\left(t, \frac{1}{2^{k}}\right)+\phi\left(t, \frac{N_{k}}{2^{k}}\right)\left(B\left(\frac{N_{k}}{2^{k}}\right)-B(t)\right)=\widetilde{P}_{k, 2}(t)
$$

Finally, under condition (B), as $k \rightarrow \infty$ we have

$$
\widetilde{P}_{k, 2}(t) \rightarrow \widetilde{P}_{2}(t)=\int_{0}^{t} \phi^{\prime}(t, s)(B(t)-B(s)) d s+\phi(t, 0) B(t) \quad \text { a.s. }
$$

On the way we use that $B(t)$ is a realization of the Brownian motion, whose trajectories are almost surely Hölder continuous, i.e. for every $\varepsilon>0$ and compact set $K \subset \mathbb{R}$, there exists some $C>0$ such that

$$
\begin{equation*}
|B(t)-B(s)| \leq C|t-s|^{1 / 2-\varepsilon} \text { for all } t, s \in K \quad \text { a.s. } \tag{A.6}
\end{equation*}
$$

Now, since the integral sums $\widetilde{P}_{k, 2}(t)$ converge in probability (also in $\mathcal{L}^{2}$ ) to $\widetilde{P}_{2}(t)$, cf. (A.5), therefore $P_{2}(t)=\int_{0}^{t} \phi^{\prime}(t, s)(B(t)-B(s)) d s+\phi(t, 0) B(t)$ a.s., which completes the proof for the positive part.

Now let $T<0$ and set $M_{k}=\left|\left\lceil T 2^{k}\right\rceil\right|$. The integral sums for $P_{1}(t, T)$ in Eq. (A.3) can be transformed to

$$
\begin{aligned}
& \widetilde{P}_{k, 1}(t, T)=\sum_{i=1}^{M_{k}-1} \psi\left(t,-\frac{i}{2^{k}}\right)\left(B\left(-\frac{i}{2^{k}}\right)-B\left(-\frac{i+1}{2^{k}}\right)\right)=\psi\left(t,-\frac{1}{2^{k}}\right) B\left(-\frac{1}{2^{k}}\right) \\
& +\sum_{i=1}^{M_{k}-2}\left(\psi\left(t,-\frac{i+1}{2^{k}}\right)-\psi\left(t,-\frac{i}{2^{k}}\right)\right) B\left(-\frac{i}{2^{k}}\right)-\psi\left(t,-\frac{M_{k}-1}{2^{k}}\right) B\left(-\frac{M_{k}}{2^{k}}\right) \\
& =\psi\left(t,-\frac{1}{2^{k}}\right) B\left(-\frac{1}{2^{k}}\right)-\frac{1}{2^{k}} \sum_{i=1}^{M_{k}-2} \psi^{\prime}\left(t, \theta_{k, i}\right) B\left(-\frac{i}{2^{k}}\right)-\psi\left(t,-\frac{M_{k}-1}{2^{k}}\right) B\left(-\frac{M_{k}}{2^{k}}\right),
\end{aligned}
$$

for some $\theta_{k, i} \in\left(-\frac{i+1}{2^{k}},-\frac{i}{2^{k}}\right)$. Now, under condition (A), and due to the Hölder continuity of Brownian motion, cf. (A.6) we find that, as $k \rightarrow \infty$ :

- $\psi\left(t,-\frac{1}{2^{k}}\right) B\left(-\frac{1}{2^{k}}\right) \rightarrow 0$ a.s.,
- $\psi\left(t,-\frac{M_{k}-1}{2^{k}}\right) B\left(-\frac{M_{k}}{2^{k}}\right) \rightarrow \psi(t, T) B(T)$ a.s.,
- $\sum_{i=1}^{M_{k}-2} \frac{1}{2^{k}} \psi^{\prime}\left(t, \theta_{k, i}\right) B\left(-\frac{i}{2^{k}}\right) \rightarrow \int_{T}^{0} \psi^{\prime}(t, s) B(s) \mathrm{d} s$ a.s.

Therefore, as $k \rightarrow \infty$ we have

$$
P_{k, 1}(t, T) \rightarrow \widetilde{P}_{1}(t, T)=-\int_{T}^{0} \psi^{\prime}(t, s) B(s) \mathrm{d} s+\psi(t, T) B(T) \quad \text { a.s. }
$$

We emphasize that the condition (A) ensures the convergence of the integral around $s=0$. Finally, the condition (A) together with the property for Brownian motion (which is implied by the law of iterated logarithm) that for every $\varepsilon>0$ there exists $C>0$ such that $|B(T)| \leq C|T|^{1 / 2+\epsilon}$ for $T \downarrow-\infty$ almost surely, we have $\psi(t, T) B(T) \rightarrow 0$ a.s. for $T \rightarrow-\infty$. Similarly it ensures the convergence of the integral around $s=-\infty$. The proof is completed.

## B Proofs

Proof of Proposition 2.1. Notice that we may change the order of expectation and integration because the underlying functions are non-negative, and obtain

$$
\begin{aligned}
& \mathbb{E}\left(\left.\int_{0}^{t} \frac{Y(t)-Y(t-s)}{s} \mathrm{~d} s \right\rvert\, \tau(T)=t, \bar{Y}(T)=y\right) \\
& =\int_{0}^{t} \frac{\mathbb{E}(Y(t)-Y(t-s) \mid \tau(T)=t, \bar{Y}(T)=y)}{s} \mathrm{~d} s
\end{aligned}
$$

Furthermore, using the formula for the density of generalized Bessel bridge in (2.7) we find that

$$
\begin{aligned}
I(t, y) & =\int_{0}^{t} \int_{0}^{\infty} s^{-1} x \cdot g(x, s ; t, y) \mathrm{d} x \mathrm{~d} s \\
& =\frac{t^{3 / 2}}{y} \int_{0}^{t} \int_{0}^{\infty} \frac{x^{2}}{s^{5 / 2} \sqrt{2 \pi(t-s)}} e^{-\frac{x^{2}}{2 s}+\frac{y^{2}}{2 t}} \cdot\left[e^{-\frac{(x-y)^{2}}{2(t-s)}}-e^{-\frac{(x+y)^{2}}{2(t-s)}}\right] \mathrm{d} x \mathrm{~d} s \\
& =\frac{t^{3 / 2}}{\sqrt{2 \pi} y} \int_{0}^{t} \int_{0}^{\infty} \frac{x^{2}}{s^{5 / 2} \sqrt{t-s}} \cdot\left[e^{-\frac{1}{2}\left(\frac{x}{\sigma_{0}}-y \mu_{0}\right)^{2}}-e^{-\frac{1}{2}\left(\frac{x}{\sigma_{0}}+y \mu_{0}\right)^{2}}\right] \mathrm{d} x \mathrm{~d} s
\end{aligned}
$$

where $\mu_{0}^{2}:=\frac{s}{(t-s) t}$, and $\sigma_{0}^{2}:=\frac{(t-s) s}{t}$. Applying substitution $z:=\sigma_{0} x$, and $w:=s / t$ afterwards yields.

$$
I(t, y)=\frac{t}{\sqrt{2 \pi} y} \int_{0}^{1} \int_{0}^{\infty} \frac{(1-w) z^{2}}{w} \cdot\left[e^{-\frac{1}{2}(z-y q(w) / \sqrt{t})^{2}}-e^{-\frac{1}{2}(z+y q(w) / \sqrt{t})^{2}}\right] \mathrm{d} z \mathrm{~d} w
$$

where $q(w):=\sqrt{\frac{w}{1-w}}$. Substituting $q:=q(w)$ yields the result.
Next, we introduce a simple lemma. The intuition behind the proof is that, asymptotically, as $x \downarrow 0$ or $x \uparrow \infty$, the absolute value of $\log (x)$ grows to infinity slower than any power of $x^{-1}$ or $x$ respectively.
Lemma B.1. For every $n \in \mathbb{Z}_{+}$and $\varepsilon>0$ there exists $C>0$ such that $\left|\log ^{n}(x)\right| \leq$ $C\left(x^{\varepsilon}+x^{-\varepsilon}\right)$ for all $x>0$.

Proof of Proposition 4.2. The law of iterated logarithm for Brownian motion implies that

$$
\begin{equation*}
\limsup _{t \rightarrow \pm \infty} \frac{|B(t)|}{|t|^{1 / 2+\varepsilon}}=0 \quad \text { a.s. } \tag{B.1}
\end{equation*}
$$

for every $\varepsilon>0$. Moreover, it is also known that paths of Brownian motion are Hölder continuous almost surely on any bounded time interval, that is, for any $\varepsilon>0$ there exist finite random numbers $A_{k}^{\varepsilon}, k \in \mathbb{Z}_{+}$such that

$$
\begin{equation*}
|B(t)-B(s)| \leq A_{k}^{\varepsilon}|t-s|^{1 / 2-\varepsilon} \text { for all } t, s \in[-k, k] \text { a.s. } \tag{B.2}
\end{equation*}
$$

From now on, let us fix a single trajectory of Brownian motion, which satisfies (B.1) and (B.2) for all $k \in \mathbb{Z}_{+}$and for some $\varepsilon>0$ chosen at the end. Let us fix $n \in \mathbb{N}, H^{\prime} \in(0,1)$ and $t^{\prime} \in \mathbb{R}$. Given the representation (4.5), it suffices to show that

$$
\begin{align*}
& \lim _{(t, H) \rightarrow\left(t^{\prime}, H^{\prime}\right)} \int_{-\infty}^{0}\left[\frac{\partial}{\partial s} f_{H}^{(n)}(t, s)\right] B(s) \mathrm{d} s=\int_{-\infty}^{0}\left[\frac{\partial}{\partial s} f_{H^{\prime}}^{(n)}\left(t^{\prime}, s\right)\right] B(s) \mathrm{d} s,  \tag{B.3}\\
& \lim _{(t, H) \rightarrow\left(t^{\prime}, H^{\prime}\right)} g_{H}^{(n)}(t, 0) B(t)=g_{H^{\prime}}^{(n)}\left(t^{\prime}, 0\right) B\left(t^{\prime}\right),  \tag{B.4}\\
& \lim _{(t, H) \rightarrow\left(t^{\prime}, H^{\prime}\right)} \int_{0}^{t}\left[\frac{\partial}{\partial s} g_{H}^{(n)}(t, s)\right](B(t)-B(s)) \mathrm{d} s=\int_{0}^{t^{\prime}}\left[\frac{\partial}{\partial s} g_{H^{\prime}}^{(n)}\left(t^{\prime}, s\right)\right]\left(B\left(t^{\prime}\right)-B(s)\right) \mathrm{d} s . \tag{B.5}
\end{align*}
$$

Clearly, for each $s \in \mathbb{R}$ we have

$$
\lim _{(t, H) \rightarrow\left(t^{\prime}, H^{\prime}\right)} \frac{\partial}{\partial s} f_{H}^{(n)}(t, s)=\frac{\partial}{\partial s} f_{H^{\prime}}^{(n)}\left(t^{\prime}, s\right), \quad \text { and } \quad \lim _{(t, H) \rightarrow\left(t^{\prime}, H^{\prime}\right)} \frac{\partial}{\partial s} g_{H}^{(n)}(t, s)=\frac{\partial}{\partial s} g_{H^{\prime}}^{(n)}\left(t^{\prime}, s\right)
$$

so (B.4) is satisfied.
We will now show that an integrable majorant for (B.3) exists in two domains: (i) $s \in\left[-s^{*}, 0\right)$, and (ii) $s<-s^{*}$, where $s^{*}:=\max \left\{1, t^{\prime}\right\}$. Consider case (i) $s \in\left[-s^{*}, 0\right)$; we have

$$
\begin{aligned}
\left|\frac{\partial}{\partial s} f_{H}^{(n)}(t, s)\right| & \leq n\left|\log ^{n-1}(t-s)(t-s)^{H-\frac{3}{2}}\right|+\left|H-\frac{1}{2}\right| \cdot\left|\log ^{n}(t-s)(t-s)^{H-\frac{3}{2}}\right| \\
& +n\left|\log ^{n-1}(-s)(-s)^{H-\frac{3}{2}}\right|+\left|H-\frac{1}{2}\right| \cdot\left|\log ^{n}(-s)(-s)^{H-\frac{3}{2}}\right|
\end{aligned}
$$

According to Lemma B.1, for every $\delta>0$ there exists $C>0$ such that

$$
\left|\frac{\partial}{\partial s} f_{H}^{(n)}(t, s)\right| \leq C\left((t-s)^{H-3 / 2-\delta}+|s|^{H-3 / 2-\delta}\right)
$$

Therefore, there exists $C_{1}, C_{2}>0$ such that for all $(t, H)$ close enough to $\left(t^{\prime}, H^{\prime}\right)$ we have

$$
\left|\frac{\partial}{\partial s} f_{H}^{(n)}(t, s)\right| \leq C_{1}+C_{2}|s|^{H^{\prime}-3 / 2-\delta}
$$

for all $-s^{*}<s<0$. Moreover, using (B.2) with $k=\left\lceil s^{*}\right\rceil$ we find that

$$
\left|\frac{\partial}{\partial s} f_{H}^{(n)}(t, s) B(s)\right| \leq\left(C_{1}+C_{2}|s|^{H^{\prime}-3 / 2-\delta}\right) A_{k}^{\varepsilon}|s|^{1 / 2-\varepsilon}
$$

The upper bound above is integrable over $s \in\left[-s^{*}, 0\right)$ when $H^{\prime}-\delta-\varepsilon>0$, which completes the proof of case (i). We now consider the case (ii) $s<-s^{*}$; we have

$$
\left|\frac{\partial}{\partial s} f_{H}^{(n)}(t, s)\right| \leq n\left|F_{H}(t, s ; n-1)\right|+\left(H-\frac{1}{2}\right)\left|F_{H}(t, s ; n)\right|
$$

where $F_{H}(t, s ; n):=\log ^{n}(t-s)(t-s)^{H-\frac{3}{2}}-\log ^{n}(-s)(-s)^{H-\frac{3}{2}}$. Using the mean value theorem we find that for $s<-s^{*}$, where recall that $s^{*}=\max \left\{1, t^{\prime}\right\}$, we have

$$
\begin{aligned}
\left|F_{H}(t, s ; n)\right| & \leq t \sup _{\widetilde{s} \in[-s, t-s]}\left|n \log ^{n-1}(\widetilde{s})(\widetilde{s})^{H-5 / 2}+\left(H-\frac{3}{2}\right) \log ^{n}(\widetilde{s})(\widetilde{s})^{H-5 / 2}\right| \\
& \leq t \cdot\left(n \log ^{n-1}(t-s)(-s)^{H-5 / 2}+\left|H-\frac{3}{2}\right| \cdot \log ^{n}(t-s)(-s)^{H-5 / 2}\right)
\end{aligned}
$$

Since $s<-t^{\prime}$, then for all $t$ sufficiently close to $t^{\prime}$ we have $t<-2 s$. Hence,

$$
\left|F_{H}(t, s ; n)\right| \leq t \cdot\left(n \log ^{n-1}(-3 s)(-s)^{H-5 / 2}+\left|H-\frac{3}{2}\right| \cdot \log ^{n}(-3 s)(-s)^{H-5 / 2}\right)
$$

According to Lemma B.1, for any $\delta>0$ and $H$ close enough to $H^{\prime}$, the above is upper bounded by $C|s|^{H^{\prime}+\delta-5 / 2}$ for some $C>0$ and all $|s| \geq s^{*}$. Using (B.1) we thus find that there exists some $C^{\prime}>0$ such that

$$
\left|\frac{\partial}{\partial s} f_{H}^{(n)}(t, s) B(s)\right| \leq C^{\prime}|s|^{H^{\prime}+\delta-5 / 2}|s|^{1 / 2+\varepsilon}=C^{\prime}|s|^{H^{\prime}+\delta+\varepsilon-2}
$$

The upper bound above is an integrable function over $s<-s^{*}$ when $H^{\prime}+\delta+\varepsilon<1$, which completes the proof of case (ii).

It is left to show (B.5). Let $\eta>0$ and let $(t, H)$ be close enough to $\left(t^{\prime}, H^{\prime}\right)$ such that $\left|t^{\prime}-t\right|<\eta / 2$. We split the pre-limit integral (B.4) into two domains (i) $s \in\left[0, t^{\prime}-\eta\right]$, and (ii) $s \in\left[t^{\prime}-\eta, t\right]$. It is clear that

$$
\lim _{(t, H) \rightarrow\left(t^{\prime}, H^{\prime}\right)} \int_{0}^{t^{\prime}-\eta}\left[\frac{\partial}{\partial s} g_{H}^{(n)}(t, s)\right](B(t)-B(s)) \mathrm{d} s=\int_{0}^{t^{\prime}-\eta}\left[\frac{\partial}{\partial s} g_{H^{\prime}}^{(n)}\left(t^{\prime}, s\right)\right]\left(B\left(t^{\prime}\right)-B(s)\right) \mathrm{d} s
$$

for all $\eta$ small enough by the Lebesgue dominated convergence theorem with dominant being the supremum of $\left[\frac{\partial}{\partial s} g_{H}^{(n)}(t, s)\right](B(t)-B(s))$ over $s \in\left[0, t^{\prime}-\eta\right]$, which is uniformly bounded for all $(t, H)$ sufficiently close to $\left(t^{\prime}, H^{\prime}\right)$. It thus suffices to show that

$$
\lim _{\eta \rightarrow 0} \lim _{(t, H) \rightarrow\left(t^{\prime}, H^{\prime}\right)} \int_{t^{\prime}-\eta}^{t}\left[\frac{\partial}{\partial s} g_{H}^{(n)}(t, s)\right](B(t)-B(s)) \mathrm{d} s=0
$$

Since we have

$$
\left|\frac{\partial}{\partial s} g_{H}^{(n)}(t, s)\right| \leq n\left|\log ^{n-1}(t-s)(t-s)^{H-\frac{3}{2}}\right|+\left|H-\frac{1}{2}\right| \cdot\left|\log ^{n}(t-s)(t-s)^{H-\frac{3}{2}}\right|
$$

then according to Lemma B.1, for any $\varepsilon>0$ we can find a constant $C$ such that

$$
\left|\frac{\partial}{\partial s} g_{H}^{(n)}(t, s)\right| \leq C|t-s|^{H-\frac{3}{2}-\varepsilon}
$$

The above combined with the Hölder continuity of the Brownian trajectory, with $n=$ $\left\lceil t^{\prime}\right\rceil+1$, for any $\delta>0$ there exists some constant $C^{\prime}>0$ such that

$$
\left|\left[\frac{\partial}{\partial s} g_{H}^{(n)}(t, s)\right](B(t)-B(s))\right| \leq C^{\prime}|t-s|^{H-\frac{3}{2}-\varepsilon+1 / 2-\delta}
$$

Since $\left|t^{\prime}-t\right|<\eta / 2$, this gives us

$$
\left|\int_{t^{\prime}-\eta}^{t}\left[\frac{\partial}{\partial s} g_{H}^{(n)}(t, s)\right](B(t)-B(s)) \mathrm{d} s\right| \leq C^{\prime} \int_{0}^{\left|t-t^{\prime}\right|+\eta} s^{H-1-\varepsilon-\delta} \mathrm{d} s \leq \frac{C^{\prime}}{H-\varepsilon-\delta}(2 \eta)^{H-\varepsilon-\delta}
$$

so after taking $\varepsilon+\delta$ small enough, we find that the above converges to 0 , as $\eta \rightarrow 0$.
Proof of Proposition 4.3. Let us fix $n \in \mathbb{N}, H \in(0,1)$, and $t \in \mathbb{R}_{+}$. Given the representation (4.5), it suffices to show that
$\lim _{\Delta \rightarrow 0} \int_{-\infty}^{0} \frac{\frac{\partial}{\partial s} f_{H+\Delta}^{(n)}(t, s)-\frac{\partial}{\partial s} f_{H}^{(n)}(t, s)}{\Delta} \cdot B(s) \mathrm{d} s=\int_{-\infty}^{0}\left[\frac{\partial}{\partial s} f_{H}^{(n+1)}(t, s)\right] B(s) \mathrm{d} s$,
$\lim _{\Delta \rightarrow 0} \frac{\frac{\partial}{\partial s} g_{H+\Delta}^{(n)}(t, 0)-\frac{\partial}{\partial s} g_{H}^{(n)}(t, 0)}{\Delta}=g_{H}^{(n+1)}(t, 0)$,
$\lim _{\Delta \rightarrow 0} \int_{0}^{t} \frac{\frac{\partial}{\partial s} g_{H+\Delta}^{(n)}(t, s)-\frac{\partial}{\partial s} g_{H}^{(n)}(t, s)}{\Delta} \cdot(B(t)-B(s)) \mathrm{d} s=\int_{0}^{t}\left[\frac{\partial}{\partial s} g_{H}^{(n+1)}(t, s)\right](B(t)-B(s)) \mathrm{d} s$.

Clearly for each $s \in \mathbb{R}$ we have

$$
\begin{aligned}
& \lim _{\Delta \rightarrow 0} \frac{\frac{\partial}{\partial s} f_{H+\Delta}^{(n)}(t, s)-\frac{\partial}{\partial s} f_{H}^{(n)}(t, s)}{\Delta}=\frac{\partial}{\partial s} f_{H}^{(n+1)}(t, s), \\
& \lim _{\Delta \rightarrow 0} \frac{\frac{\partial}{\partial s} g_{H+\Delta}^{(n)}(t, s)-\frac{\partial}{\partial s} g_{H}^{(n)}(t, s)}{\Delta}=\frac{\partial}{\partial s} g_{H}^{(n+1)}(t, s),
\end{aligned}
$$

so (B.7) is satisfied. Using the mean value theorem we find that, for every $s<0$ there exists $\widetilde{H}_{s} \in[\min \{H, H+\Delta\}, \max \{H, H+\Delta\}]$ such that

$$
\frac{\frac{\partial}{\partial s} f_{H+\Delta}^{(n)}(t, s)-\frac{\partial}{\partial s} f_{H}^{(n)}(t, s)}{\Delta}=f_{\widetilde{H}_{s}}^{(n+1)}(t, s)
$$

Using the representation above, the proof of (B.6) is now analogous the the proof of (B.3) of Proposition 4.2. Similarly, for every $s>0$ there exists $\hat{H}_{s} \in[\min \{H, H+$ $\Delta\}, \max \{H, H+\Delta\}]$ such that

$$
\frac{\frac{\partial}{\partial s} g_{H+\Delta}^{(n)}(t, s)-\frac{\partial}{\partial s} g_{H}^{(n)}(t, s)}{\Delta}=g_{\hat{H}_{s}}^{(n+1)}(t, s)
$$

Using the above, the proof of (B.8) is now analogous to the proof of (B.5) of Proposition 4.2.

Proof of Lemma 4.5. Since $R_{H}(t ; N)$ has stationary increments; cf. Proposition 4.4, then $\mathbb{E}\left|R_{H}(t ; N)-R_{H}(s ; N)\right|^{2}=\operatorname{Var} R_{H}(|t-s| ; N)$, so it suffices to bound the variance
$\operatorname{Var} R_{H}(t ; N)$ for all $t>0$. We have

$$
\begin{aligned}
& R_{H}(t ; N)=\int_{-\infty}^{0}\left(f_{H}(t, s)-\sum_{n=0}^{N-1} \frac{f_{1 / 2}^{(n)}(t, s)}{n!}\left(H-\frac{1}{2}\right)^{n}\right) \mathrm{d} B(s) \\
& +\int_{0}^{t}\left(g_{H}(t, s)-\sum_{n=0}^{N-1} \frac{g_{1 / 2}^{(n)}(t, s)}{n!}\left(H-\frac{1}{2}\right)^{n}\right) \mathrm{d} B(s) .
\end{aligned}
$$

Thus, $\operatorname{Var}\left(R_{H}(t ; N)\right)=I_{1}+I_{2}$, where

$$
\begin{align*}
& I_{1}:=\int_{0}^{\infty}\left[f_{H}(t,-s)-\sum_{n=0}^{N-1} \frac{f_{1 / 2}^{(n)}(t,-s)}{n!}\left(H-\frac{1}{2}\right)^{n}\right]^{2} \mathrm{~d} s \\
& I_{2}:=\int_{0}^{t}\left[g_{H}(t, s)-\sum_{n=0}^{N-1} \frac{g_{1 / 2}^{(n)}(t, s)}{n!}\left(H-\frac{1}{2}\right)^{n}\right]^{2} \mathrm{~d} s . \tag{B.9}
\end{align*}
$$

Now, let $\underline{H}:=\min \left\{\frac{1}{2}, H\right\}, \bar{H}:=\max \left\{\frac{1}{2}, H\right\}$. Using Taylor's theorem with Lagrange remainder we find that

$$
\begin{aligned}
I_{1} & \leq\left(H-\frac{1}{2}\right)^{2 N} \frac{1}{(N!)^{2}} \int_{0}^{\infty} \sup _{\widetilde{H} \in[\underline{H}, \bar{H}]}\left(f_{\widetilde{H}_{s}}^{(N)}(t,-s)\right)^{2} \mathrm{~d} s \\
& I_{2} \leq\left(H-\frac{1}{2}\right)^{2 N} \frac{1}{(N!)^{2}} \int_{0}^{t} \sup _{\widetilde{H} \in[\underline{[ }, \bar{H}]}\left(g_{\widetilde{H}}^{(N)}(t, s)\right)^{2} \mathrm{~d} s .
\end{aligned}
$$

Consider the integral in the upper bound for $I_{1}$ above on two different domains: (i) $s \in(0, t]$, and (ii) $s \in(t, \infty)$. In case (i) we have

$$
\begin{aligned}
& \sup _{\widetilde{H} \in[\underline{H}, \bar{H}]}\left(f_{\widetilde{H}}^{(N)}(t,-s)\right)^{2}=\sup _{\widetilde{H} \in[\underline{H}, \bar{H}]}\left(\log ^{N}(t+s)(t+s)^{\widetilde{H}-\frac{1}{2}}-\log ^{N}(s) s^{\widetilde{H}-\frac{1}{2}}\right)^{2} \\
& \leq 2\left(\sup _{\widetilde{H} \in[\underline{H}, \bar{H}]} \log ^{2 N}(t+s)(t+s)^{2 \widetilde{H}-1}+\sup _{\widetilde{H} \in[\underline{H}, \bar{H}]} \log ^{2 N}(s) s^{2 \widetilde{H}-1}\right)
\end{aligned}
$$

Now, using Lemma B. 1 and the bound above, we find that for every $\varepsilon>0$ there exists $C>0$ such that

$$
\begin{aligned}
\sup _{\widetilde{H} \in[\underline{[ }, \bar{H}]}\left(f_{\widetilde{H}}^{(N)}(t,-s)\right)^{2} & \leq C\left(\sup _{\widetilde{H} \in[\underline{H}, \bar{H}]} s^{2 \widetilde{H}-1-\varepsilon}+\sup _{\widetilde{H} \in[\underline{H}, \bar{H}]} s^{2 \widetilde{H}-1+\varepsilon}\right) \\
& \leq 2 C\left(s^{2 \underline{H}-1-\varepsilon}+s^{2 \bar{H}-1+\varepsilon}\right) .
\end{aligned}
$$

This gives us that there exists some constant $C_{*}>0$ such that

$$
\int_{0}^{t} \sup _{\widetilde{H} \in[\underline{H}, \bar{H}]}\left(f_{\widetilde{H}}^{(N)}(t,-s)\right)^{2} \mathrm{~d} s \leq C_{*}\left(t^{2 \underline{H}-\varepsilon}+t^{2 \bar{H}+\varepsilon}\right),
$$

which completes case (i). For case (ii), using the mean value theorem we find that

$$
\begin{aligned}
& \left.f_{H}^{(N)}(t,-s) \leq t \cdot \sup _{\widetilde{s} \in[s, t+s]}\left|\frac{\partial}{\partial s} \log ^{N}(s) s^{H-\frac{1}{2}}\right|_{s=\widetilde{s}} \right\rvert\, \\
& \leq t \cdot \sup _{\widetilde{s} \in[s, t+s]}\left|N \log ^{N-1}(\widetilde{s})(\widetilde{s})^{H-\frac{3}{2}}\right|+t \cdot \sup _{\widetilde{s} \in[s, t+s]}\left|\left(H-\frac{1}{2}\right) \log ^{N}(\widetilde{s})(\widetilde{s})^{H-\frac{3}{2}}\right| .
\end{aligned}
$$

Again, using Lemma B. 1 and the bound above, we find that for every $\varepsilon>0$ there exists $C>0$ such that

$$
f_{H}^{(N)}(t,-s) \leq C t \cdot\left(\sup _{\widetilde{s} \in[s, t+s]}(\widetilde{s})^{H-\frac{3}{2}-\frac{\varepsilon}{2}}+\sup _{\widetilde{s} \in[s, t+s]}(\widetilde{s})^{H-\frac{3}{2}+\frac{\varepsilon}{2}}\right),
$$

so for sufficiently small $\varepsilon$, the suprema above are attained at $\widetilde{s}=s$ and thus

$$
f_{H}^{(N)}(t,-s) \leq C t\left(s^{H-\frac{3}{2}-\frac{\varepsilon}{2}}+s^{H-\frac{3}{2}+\frac{\varepsilon}{2}}\right) .
$$

Furthermore, this gives us that there exists some constant $C_{*}>0$ such that

$$
\begin{aligned}
& \int_{t}^{\infty} \sup _{\widetilde{H} \in[\underline{H}, \bar{H}]}\left(f_{\widetilde{H}}^{(N)}(t,-s)\right)^{2} \mathrm{~d} s \leq C_{*} \sup _{\widetilde{H} \in[\underline{H}, \bar{H}]} t^{2}\left(t^{2 \widetilde{H}-2-\varepsilon}+t^{2 \widetilde{H}-2+\varepsilon}\right) \\
& \leq 2 C_{*}\left(t^{2} \underline{H}-\varepsilon\right. \\
&\left.+t^{2 \bar{H}+\varepsilon}\right)
\end{aligned}
$$

which concludes the proof of case (ii) and shows that there exists $C>0$ such that

$$
I_{1} \leq\left(H-\frac{1}{2}\right)^{2 N} \frac{C}{(N!)^{2}}\left(t^{2 \underline{H}-\varepsilon}+t^{2 \bar{H}+\varepsilon}\right)
$$

for all $t>0, H \in[\underline{H}, \bar{H}]$. The upper bound for $I_{2}$ from (B.9) can be found using analogous method, as the upper bound for $I_{1}$ in case (i) above.

Proof of Proposition 4.4. Let $t, s \in \mathbb{R}_{+}$, then

$$
\mathbb{E}(\eta(t)-\eta(s))^{2}=\sum_{i=1}^{k} \sum_{j=1}^{k} a_{i} a_{j} \mathbb{E}\left(X_{H_{i}}^{\left(n_{i}\right)}(t)-X_{H_{i}}^{\left(n_{i}\right)}(s)\right)\left(X_{H_{j}}^{\left(n_{j}\right)}(t)-X_{H_{j}}^{\left(n_{j}\right)}(s)\right)
$$

Without the loss of generality, assume that $t>s \geq 0$. It suffices to show that each term of the double sum above depends only on $t-s$. Now, notice that for any $H \in(0,1), n \in \mathbb{Z}_{+}$ and $t>s>0$ we have

$$
X_{H}^{(n)}(t)-X_{H}^{(n)}(s)=\int_{-\infty}^{s}\left[g_{H}^{(n)}(t, w)-g_{H}^{(n)}(s, w)\right] \mathrm{d} B(w)+\int_{s}^{t} g_{H}^{(n)}(t, w) \mathrm{d} B(w)
$$

thus, for any $H, H^{\prime} \in(0,1), n, n^{\prime} \in \mathbb{Z}_{+}$, and $t>s$ we obtain

$$
\begin{aligned}
& \mathbb{E}\left(X_{H}^{(n)}(t)-X_{H}^{(n)}(s)\right)\left(X_{H^{\prime}}^{\left(n^{\prime}\right)}(t)-X_{H^{\prime}}^{\left(n^{\prime}\right)}(s)\right) \\
& =\int_{-\infty}^{s}\left(g_{H}^{(n)}(t, w)-g_{H}^{(n)}(s, w)\right)\left(g_{H^{\prime}}^{\left(n^{\prime}\right)}(t, w)-g_{H^{\prime}}^{\left(n^{\prime}\right)}(s, w)\right) \mathrm{d} w \\
& +\int_{s}^{t} g_{H}^{(n)}(t, w) g_{H^{\prime}}^{\left(n^{\prime}\right)}(t, w) \mathrm{d} w
\end{aligned}
$$

After applying the substitution $z:=w-s$ to both integrals above we find that

$$
\begin{aligned}
& \mathbb{E}\left(X_{H}^{(n)}(t)-X_{H}^{(n)}(s)\right)\left(X_{H^{\prime}}^{\left(n^{\prime}\right)}(t)-X_{H^{\prime}}^{\left(n^{\prime}\right)}(s)\right) \\
& =\int_{-\infty}^{0} f_{H}^{(n)}(t-s, z) f_{H^{\prime}}^{\left(n^{\prime}\right)}(0, z) \mathrm{d} z+\int_{0}^{t-s} g_{H}^{(n)}(t-s, z) g_{H^{\prime}}^{\left(n^{\prime}\right)}(t-s, z) \mathrm{d} z,
\end{aligned}
$$

whose value depends only on $t-s$.

## C Auxiliary results on extremes of Gaussian processes

We present some useful properties of extremes of Gaussian processes, which are also of independent interest.
Proposition C.1. Suppose that $0<H_{1}<H_{2} \leq 1$.
(i) For $a \in \mathbb{R}, T \in \mathbb{R}_{+}$and $\alpha \in \mathbb{R}$

$$
\begin{aligned}
& \mathbb{E}\left(\exp \left(\alpha\left(\sup _{t \in[0, T]} \sqrt{2} B_{H_{2}}(t)-a t^{2 H_{2}}\right)\right)\right) \\
& \leq \mathbb{E}\left(\exp \left(\alpha\left(\sup _{t \in\left[0, T^{\left.H_{2} / H_{1}\right]}\right.} \sqrt{2} B_{H_{1}}(t)-a t^{2 H_{1}}\right)\right)\right)<\infty .
\end{aligned}
$$

(ii) For $a>1$ and $\alpha<a$

$$
\begin{aligned}
& \mathbb{E}\left(\exp \left(\alpha\left(\sup _{t \in[0, \infty)} \sqrt{2} B_{H_{2}}(t)-a t^{2 H_{2}}\right)\right)\right) \\
& \leq \mathbb{E}\left(\exp \left(\alpha\left(\sup _{t \in[0, \infty)} \sqrt{2} B_{H_{1}}(t)-a t^{2 H_{1}}\right)\right)\right)<\infty .
\end{aligned}
$$

Proof. We note that $\operatorname{Var}\left(B_{H_{2}}(t)\right)=\operatorname{Var}\left(B_{H_{1}}\left(t^{H_{2} / H_{1}}\right)\right)$ and

$$
\begin{align*}
\operatorname{Cov}\left(B_{H_{2}}(s), B_{H_{2}}(t)\right) & =\frac{1}{2}\left(s^{2 H_{2}}+t^{2 H_{2}}-|s-t|^{2 H_{2}}\right) \\
& \geq \frac{1}{2}\left(s^{2 H_{2}}+t^{2 H_{2}}-\left|s^{H_{2} / H_{1}}-t^{H_{2} / H_{1}}\right|^{2 H_{1}}\right)  \tag{C.1}\\
& =\operatorname{Cov}\left(B_{H_{1}}\left(s^{H_{2} / H_{1}}\right), B_{H_{1}}\left(t^{H_{2} / H_{1}}\right)\right)
\end{align*}
$$

where inequality (C.1) is due to the superadditivity of the function $f(x)=x^{H_{2} / H_{1}}$ (recall that $H_{1}<H_{2}$ ).

Ad (i). From the Slepian inequality [1, Corollary 2.4], we get for each $u \in \mathbb{R}$

$$
\begin{aligned}
\mathbb{P}\left(\sup _{t \in[0, T]} \sqrt{2} B_{H_{2}}(t)-a t^{2 H_{2}}>u\right) & \leq \mathbb{P}\left(\sup _{t \in[0, T]} \sqrt{2} B_{H_{1}}\left(t^{H_{2} / H_{1}}\right)-a t^{2 H_{2}}>u\right) \\
& =\mathbb{P}\left(\sup _{t \in\left[0, T^{\left.H_{2} / H_{1}\right]}\right.} \sqrt{2} B_{H_{1}}(t)-a t^{2 H_{1}}>u\right) .
\end{aligned}
$$

Using that by (2.3) in [1], for $H=H_{1}, H_{2}$ and any $T \in \mathbb{R}_{+}$

$$
\lim _{u \rightarrow \infty} \frac{\log \left(\mathbb{P}\left(\sup _{t \in[0, T]} \sqrt{2} B_{H}(t)-a t^{2 H}>u\right)\right)}{u^{2}}=-\frac{1}{2 T^{2 H}}
$$

and hence for each $a \in \mathbb{R}$ and $\alpha \in \mathbb{R}$

$$
\mathbb{E}\left(\exp \left(\alpha\left(\sup _{t \in[0, T]} \sqrt{2} B_{H}(t)-a t^{2 H}\right)\right)\right)<\infty
$$

completes the proof of case (i).
Ad (ii). In view of Theorem 1 in [22], for any $H \in(0,1)$

$$
\lim _{u \rightarrow \infty} \frac{\log \left(\mathbb{P}\left(\sup _{t \in[0, \infty)} \sqrt{2} B_{H}(t)-a t^{2 H}>u\right)\right)}{u}=-a
$$

Thus, for all $\alpha<a$

$$
\mathbb{E}\left(\exp \left(\alpha\left(\sup _{t \in[0, \infty)} \sqrt{2} B_{H_{1}}(t)-a t^{2 H_{1}}\right)\right)\right)<\infty
$$

The rest of the proof follows by the same argument as the proof of (i).
Lemma C.2. Let $W(t)$ be a centered Gaussian process with stationary increments, bounded sample paths a.s. and $W(0)=0$ a.s. Then, for any $n \in \mathbb{N}$ there exists $C(n)>0$ such that

$$
\mathbb{E} \sup _{t \in[0, T]}|W(t)|^{n} \leq C(n) \cdot\left(T^{n+2}(\operatorname{Var} W(1))^{n / 2}+T^{2} \mathbb{E}\left(\sup _{t \in[0,1]} W(t)\right)^{n}\right)
$$

Proof of Lemma C.2. Without loss of generality, we suppose that $T \in \mathbb{N}$. We note that

$$
\mathbb{E} \sup _{t \in[0, T]}|W(t)|^{n} \leq T \sum_{i=0}^{T-1} \mathbb{E} \sup _{t \in[i, i+1]}|W(t)|^{n}
$$

Then, for $i=0, \ldots, T-1$

$$
\begin{aligned}
\mathbb{E} \sup _{t \in[i, i+1]}|W(t)|^{n} & \leq \mathbb{E}\left(|W(i)|+\sup _{t \in[i, i+1]}|W(t)-W(i)|\right)^{n} \\
& \leq \max \left\{1,2^{n-1}\right\}\left(\mathbb{E}|W(i)|^{n}+\mathbb{E} \sup _{t \in[0,1]}|W(t)|^{n}\right),
\end{aligned}
$$

where we used the stationarity of increments of $W(t)$. Moreover, for $i \geq 1$

$$
\begin{aligned}
\mathbb{E}\left(W^{2}(i)\right) & =\mathbb{E}\left(\sum_{k=0}^{i-1} W(k+1)-W(k)\right)^{2} \\
& \leq i \sum_{k=0}^{i-1} \mathbb{E}\left(W^{2}(1)\right)=\mathbb{E}\left(W^{2}(1)\right) i^{2}
\end{aligned}
$$

and since $\mathbb{E}|W(i)|^{n}=C_{n} \cdot\left(\mathbb{E}|W(i)|^{2}\right)^{n / 2}$, where $\left.C_{n}:=\mathbb{E} \mid \mathcal{N}(0,1)\right)\left.\right|^{n}=\frac{2^{n / 2} \Gamma((n+1) / 2)}{\sqrt{\pi}}$, then

$$
\mathbb{E}|W(i)|^{n} \leq C_{n}(\operatorname{Var} W(1))^{n / 2} i^{n}
$$

Finally, we have

$$
\mathbb{E} \sup _{t \in[0, T]}|W(t)|^{n} \leq 2 \max \left\{1,2^{n-1}\right\} T\left(\sum_{i=1}^{T-1}\left(C_{n}(\operatorname{Var} W(1))^{n / 2} i^{n}+\mathbb{E} \sup _{t \in[0,1]}|W(t)|^{n}\right)\right)
$$

which completes the proof.
Lemma C.3. Let $\{W(t) ; t \in[0,1]\}$ be a centered Gaussian process. If there exist $A>0$, $H \in(0,1)$ such that $\mathbb{E}|W(t)-W(s)|^{2} \leq A|t-s|^{2 H}$ for all $t, s \in[0,1]$, then for every $n \in \mathbb{Z}_{+}$ there exists a constant $C(n)>0$ such that

$$
\mathbb{E} \sup _{t \in[0,1]}|W(t)|^{n} \leq C(n) \cdot A^{n / 2}\left(1+\mu^{n-1}(H)+\mu^{n}(H)\right)
$$

where $\mu(H):=\mathbb{E s u p}_{t \in[0,1]} B_{H}(t)$, with $\left\{B_{H}(t): t \in \mathbb{R}_{+}\right\}$a fractional Brownian motion with Hurst parameter $H$.

Proof of Lemma C.3. When $n=1$ then the above holds by virtue of Sudakov-Fernique inequality [1, Theorem 2.9] with $C(1)=2$, that is we have

$$
\begin{equation*}
m:=\mathbb{E} \sup _{t \in[0,1]}|W(t)| \leq \mathbb{E} \sup _{t \in[0,1]}\{W(t)\}+\mathbb{E} \sup _{t \in[0,1]}\{-W(t)\} \leq 2 A^{1 / 2} \mu(H) \tag{C.2}
\end{equation*}
$$

Now let $n>1$. Using Borell inequality [1, Theorem 2.1], with $\sigma^{2}:=\sup _{t \in[0,1]} \mathbb{E} W^{2}(t)$ we have

$$
\begin{aligned}
\mathbb{E} \sup _{t \in[0,1]}|W(t)|^{n} & =\int_{0}^{\infty} \mathbb{P}\left(\sup _{t \in[0,1]}|W(t)|^{n}>u\right) \mathrm{d} u \\
& =\int_{0}^{m^{n}} \mathbb{P}\left(\sup _{t \in[0,1]}|W(t)|>u^{1 / n}\right) \mathrm{d} u+\int_{m^{n}}^{\infty} \mathbb{P}\left(\sup _{t \in[0,1]}|W(t)|>u^{1 / n}\right) \mathrm{d} u
\end{aligned}
$$

The first term above is bounded by $m^{n}$. Using Borell inequality, we can bound the second term above with

$$
\begin{aligned}
\int_{m^{n}}^{\infty} \exp \left\{-\frac{\left(u^{1 / n}-m\right)^{2}}{2 \sigma^{2}}\right\} \mathrm{d} u & =n \sigma \int_{0}^{\infty}(\sigma x+m)^{n-1} e^{-x^{2} / 2} \mathrm{~d} x \\
& \leq n \sigma \int_{0}^{\infty} \max \left\{1,2^{n-2}\right\}\left((\sigma x)^{n-1}+m^{n-1}\right) e^{-x^{2} / 2} \mathrm{~d} x \\
& \leq n\left(\sigma^{n}+\sigma m^{n-1}\right) \max \left\{1,2^{n-2}\right\} \int_{0}^{\infty}\left(x^{n-1}+1\right) e^{-x^{2} / 2} \mathrm{~d} x
\end{aligned}
$$

where in the first line we used the substitution $x:=\left(u^{1 / n}-m\right) / \sigma$. Now, let $C^{\prime}(n):=$ $n \max \left\{1,2^{n-2}\right\} \int_{0}^{\infty}\left(x^{n-1}+1\right) e^{-x^{2} / 2} \mathrm{~d} x$, then we have

$$
\mathbb{E} \sup _{t \in[0,1]}|W(t)|^{n} \leq m^{n}+C^{\prime}(n)\left(\sigma^{n}+\sigma m^{n-1}\right)
$$

By the assumption, we have $\sigma^{2} \leq A$ and also $m \leq 2 A^{1 / 2} \mu(H)$ due to Eq. (C.2) we obtain

$$
\begin{aligned}
\mathbb{E} \sup _{t \in[0,1]}|W(t)|^{n} & \leq\left(2 A^{1 / 2} \mu(H)\right)^{n}+C^{\prime}(n)\left(A^{n / 2}+A^{1 / 2}\left(2 A^{1 / 2} \mu(H)\right)^{n-1}\right) \\
& \left.=A^{n / 2}\left(C^{\prime}(n)+2 \mu^{n}(H)+C^{\prime}(n)(2 \mu(H))^{n-1}\right)\right)
\end{aligned}
$$

The claim follows with $C(n):=\max \left\{2,2^{n-1} C^{\prime}(n)\right\}$.
Lemma C.4. Let $W(t)$ be a centered Gaussian process with stationary increments and $\tau \geq 0$ be a random time. If there exist $\gamma, \beta, C>0$ satisfying $\mathbb{P}(\tau>T) \leq C e^{-\gamma T^{\beta}}$ for all $T \in \mathbb{R}_{+}$, then for each $n \in \mathbb{N}$ there exists $C^{\prime}>0$ depending only on $n, \gamma, \beta, C$ such that

$$
\mathbb{E}|W(\tau)|^{n} \leq C^{\prime}\left((\operatorname{Var} W(1))^{n / 2}+\mathbb{E}\left(\sup _{t \in[0,1]} W(t)\right)^{n}\right)
$$

Proof. See that

$$
\begin{aligned}
\mathbb{E}|W(\tau)|^{n} & \left.=\sum_{k=0}^{\infty} \mathbb{E}\left(|W(\tau)|^{n} ; \tau \in[k, k+1]\right)\right) \\
& \left.\leq \sum_{k=0}^{\infty} \mathbb{E}\left(\sup _{t \in[0, k+1]}|W(t)|^{n} ; \tau \geq k\right)\right) \\
& \left.\leq \sum_{k=0}^{\infty}\left[\mathbb{E}\left(\sup _{t \in[0, k+1]}|W(t)|^{2 n}\right) \mathbb{P}(\tau \geq k)\right)\right]^{1 / 2},
\end{aligned}
$$

where in the last line we used Hölder's inequality. Using Lemma C. 2 we further find that there exists a constant $C^{\prime}>0$ such that

$$
\mathbb{E}|W(\tau)|^{n} \leq C^{\prime}\left((\operatorname{Var} W(1))^{n / 2}+\mathbb{E}\left(\sup _{t \in[0,1]} W(t)\right)^{n}\right) \cdot \sum_{k=0}^{\infty}(k+1)^{n+1} e^{-\frac{\gamma}{2} k^{\beta}}
$$

which completes the proof because the sum above is finite.

## D Calculation of integrals in Corollary 3.3

In the proof of Corollary 3.3 one needs to handle complicated integrals. For convenience we present all these calculations here.

Proof of Corollary 3.3(i). Since

$$
p(t, y ; T, 0)=\frac{y \exp \left\{-y^{2} /(2 t)\right\}}{\pi t^{3 / 2} \sqrt{T-t}}
$$

cf. (2.5) (see also [6, 2.1.13.4]), then according to Theorem 3.1 we have $\mathscr{M}_{1 / 2}^{\prime}(T, 0)=$ $\mathcal{I}_{1}-\mathcal{I}_{2}$, where

$$
\mathcal{I}_{1}:=\int_{0}^{T} \int_{0}^{\infty} y(1+\log (t)) \frac{y \exp \left\{-\frac{y^{2}}{2 t}\right\}}{\pi t^{3 / 2} \sqrt{T-t}} \mathrm{~d} y \mathrm{~d} t, \quad \mathcal{I}_{2}:=\int_{0}^{T} \int_{0}^{\infty} I(t, y) \frac{y \exp \left\{-\frac{y^{2}}{2 t}\right\}}{\pi t^{3 / 2} \sqrt{T-t}} \mathrm{~d} y \mathrm{~d} t
$$

Till this end we show that

$$
\begin{equation*}
\mathcal{I}_{1}=\sqrt{\frac{2 T}{\pi}} \cdot(\log (T)+2 \log (2)-1), \quad \mathcal{I}_{2}=\sqrt{\frac{2 T}{\pi}}(1+2 \log (2)) \tag{D.1}
\end{equation*}
$$

which completes the proof. We first calculate $\mathcal{I}_{1}$. After substitution $y=\sqrt{t} z$ we have

$$
\mathcal{I}_{1}=\sqrt{\frac{2}{\pi}} \int_{0}^{T} \frac{1+\log (t)}{\sqrt{T-t}} \mathrm{~d} t \int_{0}^{\infty} \frac{z^{2}}{\sqrt{2 \pi}} e^{-\frac{z^{2}}{2}} \mathrm{~d} z=\sqrt{\frac{2}{\pi}} \cdot\left(2 \sqrt{T}+\int_{0}^{T} \frac{\log (t)}{\sqrt{T-t}} \mathrm{~d} t\right) \cdot \frac{1}{2} .
$$

Now, using substitution $t=T s$ we obtain

$$
\begin{aligned}
& \int_{0}^{T} \frac{\log (t)}{\sqrt{T-t}} \mathrm{~d} t \\
& =\sqrt{T} \int_{0}^{1} \frac{\log (s T)}{\sqrt{1-s}} \mathrm{~d} s=2 \sqrt{T} \log (T)+\sqrt{T} \int_{0}^{1} \frac{\log (s)}{\sqrt{1-s}} \mathrm{~d} s=\sqrt{T}(2 \log (T)+4 \log (2)-4)
\end{aligned}
$$

where we used that $\int_{0}^{1} \frac{x \log (x)}{\sqrt{1-x^{2}}} \mathrm{~d} x=\log (2)-1$; see formula 4.241-2 in [20]. This leads to the formula for $\mathcal{I}_{1}$ in (D.1).

We now calculate $\mathcal{I}_{2}$. Using e.g. [28, 5-(A6)] we find that for any $b \in \mathbb{R}$ we have

$$
\int_{0}^{\infty} x^{2} \mathrm{e}^{-\frac{(x+b)^{2}}{2}} \mathrm{~d} x=-b e^{-b^{2} / 2}+\left(1+b^{2}\right) \sqrt{\pi / 2} \cdot \operatorname{erfc}(b / \sqrt{2})
$$

Since for any $x \in \mathbb{R}$ we have $\operatorname{erfc}(-x)-\operatorname{erfc}(x)=2 \operatorname{erf}(x)$, then for any $b \in \mathbb{R}$

$$
\begin{equation*}
\int_{0}^{\infty} x^{2}\left(e^{-\frac{(x-b)^{2}}{2}}-e^{-\frac{(x+b)^{2}}{2}}\right) \mathrm{d} x=2 b e^{-b^{2} / 2}+\left(1+b^{2}\right) \sqrt{2 \pi} \cdot \operatorname{erf}(b / \sqrt{2}) \tag{D.2}
\end{equation*}
$$

Now, see that

$$
\mathcal{I}_{2}=\int_{0}^{T} \int_{0}^{\infty} \int_{0}^{\infty} \int_{0}^{\infty} \frac{\sqrt{2} t x^{2}}{y q\left(1+q^{2}\right)^{2} \sqrt{\pi}}\left(e^{-(x-y q / \sqrt{t})^{2} / 2}-e^{-(x+y q / \sqrt{t})^{2} / 2}\right) \frac{y \exp \left\{-\frac{y^{2}}{2 t}\right\}}{\pi t^{3 / 2} \sqrt{T-t}} \mathrm{~d} x \mathrm{~d} q \mathrm{~d} y \mathrm{~d} t .
$$

The substitution $y=\sqrt{t} z$ yields

$$
\mathcal{I}_{2}:=\frac{\sqrt{2}}{\pi^{3 / 2}} \int_{0}^{T} \frac{1}{\sqrt{T-t}} \mathrm{~d} t \cdot \int_{0}^{\infty} \int_{0}^{\infty} \int_{0}^{\infty} \frac{x^{2}}{q\left(1+q^{2}\right)^{2}}\left(e^{-\frac{(x-z q)^{2}}{2}}-e^{-\frac{(x+z q)^{2}}{2}}\right) e^{-\frac{z^{2}}{2}} \mathrm{~d} x \mathrm{~d} q \mathrm{~d} z
$$

Using (D.2) we obtain

$$
\int_{0}^{\infty} x^{2}\left(e^{-\frac{(x-z q)^{2}}{2}}-e^{-\frac{(x+z q)^{2}}{2}}\right) \mathrm{d} x=2 z q e^{-z^{2} q^{2} / 2}+\left(1+z^{2} q^{2}\right) \sqrt{2 \pi} \cdot \operatorname{erf}(z q / \sqrt{2})
$$

and thus

$$
\mathcal{I}_{2}:=\sqrt{\frac{2 T}{\pi}} \cdot \frac{2}{\pi} \int_{0}^{\infty} \frac{1}{q\left(1+q^{2}\right)^{2}}\left(J_{1}(q)+J_{2}(q)+J_{3}(q)\right) \mathrm{d} q
$$

where using the tables of integrals [28, 4.3.4 and 4.3.6] and the fact that $\pi-2 \arctan (1 / x)=$ $2 \arctan (x)$, we find that

$$
\begin{aligned}
& J_{1}(q):=2 q \int_{0}^{\infty} z e^{-z^{2}\left(1+q^{2}\right) / 2} \mathrm{~d} z=\frac{2 q}{1+q^{2}} \\
& J_{2}(q):=\sqrt{2 \pi} \int_{0}^{\infty} \operatorname{erf}(z q / \sqrt{2}) e^{-z^{2} / 2} \mathrm{~d} z=2 \arctan (q) \\
& J_{3}(q):=q^{2} \sqrt{2 \pi} \int_{0}^{\infty} z^{2} \operatorname{erf}(z q / \sqrt{2}) e^{-\frac{z^{2}}{2}} \mathrm{~d} z=2 q^{2}\left(\arctan (q)+\frac{q}{1+q^{2}}\right) .
\end{aligned}
$$

Finally, we have

$$
\begin{aligned}
& \int_{0}^{\infty} \frac{J_{1}(q)}{q\left(1+q^{2}\right)^{2}} \mathrm{~d} q=2 \int_{0}^{\infty} \frac{1}{\left(1+q^{2}\right)^{3}} \mathrm{~d} q \\
& \int_{0}^{\infty} \frac{J_{2}(q)}{q\left(1+q^{2}\right)^{2}} \mathrm{~d} q=2 \int_{0}^{\infty} \frac{\arctan (q)}{q\left(1+q^{2}\right)^{2}} \mathrm{~d} q \\
& \int_{0}^{\infty} \frac{J_{3}(q)}{q\left(1+q^{2}\right)^{2}} \mathrm{~d} q=2 \int_{0}^{\infty} \frac{q \arctan (q)}{\left(1+q^{2}\right)^{2}} \mathrm{~d} q+2 \int_{0}^{\infty} \frac{q^{2}}{\left(1+q^{2}\right)^{3}} \mathrm{~d} q .
\end{aligned}
$$

Using the identity

$$
\begin{equation*}
\frac{\arctan (q)}{q\left(1+q^{2}\right)^{2}}=\frac{\arctan (q)}{q\left(1+q^{2}\right)}-\frac{q \arctan (q)}{\left(1+q^{2}\right)^{2}} \tag{D.3}
\end{equation*}
$$

we obtain

$$
\begin{aligned}
\mathcal{I}_{2} & =\sqrt{\frac{2 T}{\pi}} \cdot \frac{4}{\pi}\left(\int_{0}^{\infty} \frac{1}{\left(1+q^{2}\right)^{2}} \mathrm{~d} q+\int_{0}^{\infty} \frac{\arctan (q)}{q\left(1+q^{2}\right)} \mathrm{d} q\right) \\
& =\sqrt{\frac{2 T}{\pi}} \cdot \frac{4}{\pi}\left(\frac{\pi}{4}+\frac{\pi}{2} \log (2)\right)
\end{aligned}
$$

where the value of second integral was found in 4.535-7 in [20]. We recognize that the result above matches the formula for $\mathcal{I}_{2}$ in (D.1). This completes the proof.

Proof of Corollary 3.3(ii). Since, for $a>0$ we have

$$
p(t, y, \infty, a)=\frac{\sqrt{2} a y \exp \left\{-\frac{(y+t a)^{2}}{2 t}\right\}}{t^{3 / 2} \sqrt{\pi}}
$$

see (2.6), then according to Theorem 3.1 we have $\mathscr{M}_{1 / 2}^{\prime}(\infty, a)=\mathcal{I}_{1}-\mathcal{I}_{2}$, where

$$
\begin{aligned}
& \mathcal{I}_{1}:=\int_{0}^{\infty} \int_{0}^{\infty}(y(1+\log (t))+a t \log (t)) \frac{\sqrt{2} a y \exp \left\{-\frac{(y+t a)^{2}}{2 t}\right\}}{t^{3 / 2} \sqrt{\pi}} \mathrm{~d} y \mathrm{~d} t \\
& \mathcal{I}_{2}:=\int_{0}^{\infty} \int_{0}^{\infty} I(t, y) \frac{\sqrt{2} a y \exp \left\{-\frac{(y+t a)^{2}}{2 t}\right\}}{t^{3 / 2} \sqrt{\pi}} \mathrm{~d} y \mathrm{~d} t
\end{aligned}
$$

At the end, we show that

$$
\begin{equation*}
\mathcal{I}_{1}=\frac{3}{2 a}-\frac{1}{a}\left(\gamma_{\mathrm{E}}-\log \left(2 a^{2}\right)\right), \quad \mathcal{I}_{2}=\frac{3}{2 a} \tag{D.4}
\end{equation*}
$$

which yields the result. First we calculate $\mathcal{I}_{1}$. After substitution $y=t z$, we obtain

$$
\begin{aligned}
\mathcal{I}_{1} & =\sqrt{\frac{2}{\pi}} \cdot a \int_{0}^{\infty} \int_{0}^{\infty}(z+(z+a) \log (t)) \cdot z t^{3 / 2} \exp \left\{-\frac{(z+a)^{2}}{2} \cdot t\right\} \mathrm{d} t \mathrm{~d} z \\
& =\sqrt{\frac{2}{\pi}} \cdot a \int_{0}^{\infty}\left(z^{2} J_{1}(z)+z(z+a) J_{2}(z)\right) \mathrm{d} z
\end{aligned}
$$

where

$$
J_{1}(z):=\int_{0}^{\infty} t^{3 / 2} \exp \left\{-t \cdot \frac{(z+a)^{2}}{2}\right\} \mathrm{d} t, \quad J_{2}(z):=\int_{0}^{\infty} \log (t) \cdot t^{3 / 2} \exp \left\{-t \cdot \frac{(z+a)^{2}}{2}\right\} \mathrm{d} t
$$

The first integral can easily be calculated using the formula $\int_{0}^{\infty} t^{3 / 2} e^{-\mu t} \mathrm{~d} t=\frac{\Gamma(5 / 2)}{\mu^{5 / 2}}$, while the second integral can be calculated using the formula 4.352-1 in [20] (with $\nu=5 / 2$, $\mu>0$ ) which for convenience we recall below

$$
\begin{equation*}
\int_{0}^{\infty} \log (t) t^{3 / 2} e^{-\mu t} \mathrm{~d} t=\frac{\Gamma(5 / 2)}{\mu^{5 / 2}}\left(\frac{8}{3}-\gamma_{\mathrm{E}}-2 \log (2)-\log (\mu)\right) \tag{D.5}
\end{equation*}
$$

After straightforward algebraic manipulations we arrive at

$$
\mathcal{I}_{1}=\sqrt{\frac{2}{\pi}} \cdot a \int_{0}^{\infty}\left(\frac{3 z^{2} \sqrt{2 \pi}}{(a+z)^{5}}-\frac{\sqrt{2 \pi} z}{(a+z)^{4}} \cdot\left(-8+3 \gamma_{\mathrm{E}}+\log (8)+6 \log (a+z)\right)\right) \mathrm{d} z
$$

The calculation of the integrals of rational functions above is standard. It is left to find the integral of the product of rational function and logarithm, i.e.

$$
\int_{0}^{\infty} \frac{z}{(a+z)^{4}} \log (a+z) \mathrm{d} z=a^{-2} \int_{0}^{\infty} \frac{y}{(1+y)^{4}}(\log (a)+\log (1+y)) \mathrm{d} y
$$

where we used the substitution $z=a y$. Again, the integral of the rational function can be found, in particular, we have $\int_{0}^{\infty} \frac{y}{(1+y)^{4}} \mathrm{~d} y=\frac{1}{6}$, so it is left to find

$$
\int_{0}^{\infty} \frac{y \log (1+y)}{(1+y)^{4}} \mathrm{~d} y=\int_{0}^{1} w(1-w) \log (w) \mathrm{d} w=-\frac{5}{36}
$$

where we substituted $w=(1+y)^{-1}$ and used formula 4.253-1 in [20]. Finally, simple algebraic manipulations yield the value of $\mathcal{I}_{1}$ in (D.4).

We now calculate the value of $\mathcal{I}_{2}$. After the substitutions $y=t z$ and $x=\sqrt{t} w$ we obtain:

$$
\mathcal{I}_{2}=\frac{2 a}{\pi} \int_{0}^{\infty} \int_{0}^{\infty} \int_{0}^{\infty} \frac{w^{2}\left(J_{1}(w, z, q)-J_{2}(w, z, q)\right)}{q\left(1+q^{2}\right)^{2}} \mathrm{~d} w \mathrm{~d} z \mathrm{~d} q
$$

where

$$
J_{1}(w, z, q)=\int_{0}^{\infty} t^{2} \cdot e^{-t \cdot \frac{(a+z)^{2}+(w-q z)^{2}}{2}} \mathrm{~d} t, \quad J_{2}(w, z, q)=\int_{0}^{\infty} t^{2} \cdot e^{-t \cdot \frac{(a+z)^{2}+(w+q z)^{2}}{2}} \mathrm{~d} t .
$$

Since $\int_{0}^{\infty} t^{2} e^{-\mu t} d t=\mu^{-3} \Gamma(3)$, then we obtain

$$
\mathcal{I}_{2}=\frac{32 a}{\pi} \int_{0}^{\infty} \int_{0}^{\infty} \frac{J_{1}^{\prime}(z, q)-J_{2}^{\prime}(z, q)}{q\left(1+q^{2}\right)^{2}} \mathrm{~d} z \mathrm{~d} q
$$

where

$$
J_{1}^{\prime}(z, q)=\int_{0}^{\infty} \frac{w^{2}}{\left((a+z)^{2}+(w-q z)^{2}\right)^{3}} \mathrm{~d} w, \quad J_{2}^{\prime}(z, q)=\int_{0}^{\infty} \frac{w^{2}}{\left((a+z)^{2}+(w+q z)^{2}\right)^{3}} \mathrm{~d} w
$$

Functions $J_{1}^{\prime}(z, q)$ and $J_{2}^{\prime}(z, q)$ are definite integrals of rational functions, so it is wellknown how to find them. Below, we write the explicit formula for $\mathcal{I}_{2}$

$$
\mathcal{I}_{2}=\frac{8 a}{\pi} \int_{0}^{\infty} \int_{0}^{\infty} \frac{3 q z(a+z)+\left((a+z)^{2}+3 q^{2} z^{2}\right) \arctan \left(\frac{q z}{a+z}\right)}{q\left(1+q^{2}\right)^{2}(a+z)^{5}} \mathrm{~d} q \mathrm{~d} z
$$

After the substitution $z=a y$ we obtain

$$
\mathcal{I}_{2}=\frac{8}{\pi a} \int_{0}^{\infty} \int_{0}^{\infty} \frac{3 q y(1+y)+\left((1+y)^{2}+3 q^{2} y^{2}\right) \arctan \left(\frac{q y}{1+y}\right)}{q\left(1+q^{2}\right)^{2}(1+y)^{5}} \mathrm{~d} q \mathrm{~d} y
$$

and further substituting $z=\frac{y}{1+y}$ we find that

$$
\mathcal{I}_{2}=\frac{8}{\pi a} \int_{0}^{1}(1-z) \cdot\left(I_{1}(z)+I_{2}(z)+I_{3}(z)\right) \mathrm{d} z
$$

where

$$
\begin{aligned}
& I_{1}(z):=3 z \int_{0}^{\infty} \frac{1}{\left(1+q^{2}\right)^{2}} \mathrm{~d} q=\frac{\pi}{4} \cdot 3 z \\
& I_{2}(z):=\int_{0}^{\infty} \frac{\arctan (q z)}{q\left(1+q^{2}\right)^{2}} \mathrm{~d} q=\frac{\pi}{4} \cdot\left(2 \log (1+z)-\frac{z}{1+z}\right) \\
& I_{3}(z):=3 z^{2} \int_{0}^{\infty} \frac{q \arctan (q z)}{\left(1+q^{2}\right)^{2}} \mathrm{~d} q=\frac{\pi}{4} \cdot \frac{3 z^{3}}{1+z}
\end{aligned}
$$

where the formula for $I_{2}(z)$ was found using (D.3) and formulas 4.535-7 and 4.535-11 in [20], while $I_{3}(z)$ follows from 4.535-11 in [20]. This gives us

$$
\begin{aligned}
\mathcal{I}_{2} & =\frac{2}{a} \int_{0}^{1}\left(3 z(1-z)+2(1-z) \log (1+z)+\frac{z(1-z)\left(3 z^{2}-1\right)}{1+z}\right) \mathrm{d} z \\
& =\frac{2}{a} \cdot\left(\frac{1}{2}+\log (16)-\frac{5}{2}+\frac{11}{4}-\log (16)\right)=\frac{3}{2 a}
\end{aligned}
$$

which completes the proof.
Proof of Corollary 3.3(iii). The following calculations are very similar to the ones did in the proof of Corollary 3.3(ii). For $a>1$ we have

$$
p(t, y, \infty, a / \sqrt{2})=\frac{a y \exp \left\{-\frac{(y+t a / \sqrt{2})^{2}}{2 t}\right\}}{t^{3 / 2} \sqrt{\pi}}
$$

see (2.6). Therefore, according to Theorem 3.2 we have $\mathscr{P}_{1 / 2}^{\prime}(\infty, a)=\mathcal{I}_{1}-\mathcal{I}_{2}$, where

$$
\begin{aligned}
& \mathcal{I}_{1}:=\sqrt{\frac{2}{\pi}} \cdot a \int_{0}^{\infty} \int_{0}^{\infty}\left(y(1+\log (t))-\frac{a}{\sqrt{2}} t \log (t)\right) e^{\sqrt{2} y} \cdot \frac{y \exp \left\{-\frac{(y+t a / \sqrt{2})^{2}}{2 t}\right\}}{t^{3 / 2}} \mathrm{~d} y \mathrm{~d} t \\
& \mathcal{I}_{2}:=\sqrt{\frac{2}{\pi}} \cdot a \int_{0}^{\infty} \int_{0}^{\infty} I(t, y) e^{\sqrt{2} y} \cdot \frac{y \exp \left\{-\frac{(y+t a / \sqrt{2})^{2}}{2 t}\right\}}{t^{3 / 2}} \mathrm{~d} y \mathrm{~d} t .
\end{aligned}
$$

In the following $\operatorname{acoth}(\cdot)$ is the inverse hyperbolic cotangent function, i.e. for $|z|>1$

$$
\operatorname{acoth}(z)=\frac{1}{2}\left(\log \left(1+\frac{1}{z}\right)-\log \left(1-\frac{1}{z}\right)\right)=\frac{1}{2} \log \left(\frac{z+1}{z-1}\right) .
$$

Till this end, we show that

$$
\begin{equation*}
\mathcal{I}_{1}=\frac{a\left(3-2 a-4(a-1)^{2} \operatorname{acoth}(1-2 a)\right)}{(a-1)^{2}}, \quad \mathcal{I}_{2}=\frac{a(1-4(a-1) \operatorname{acoth}(1-2 a))}{(a-1)^{2}} \tag{D.6}
\end{equation*}
$$

which yields the result. First we calculate $\mathcal{I}_{1}$. After substitution $y=t z$, we obtain

$$
\mathcal{I}_{1}:=\sqrt{\frac{2}{\pi}} \cdot a \int_{0}^{\infty} \int_{0}^{\infty}\left(z+\left(z-\frac{a}{\sqrt{2}}\right) \log (t)\right) \cdot z t^{3 / 2} \exp \left\{-\frac{(z+a / \sqrt{2})^{2}-2 \sqrt{2} z}{2} \cdot t\right\} \mathrm{d} t \mathrm{~d} z
$$

Define $z_{1}^{2}:=\left(z+\frac{a}{\sqrt{2}}\right)^{2}-2 \sqrt{2} z=\left(z+\frac{a-2}{\sqrt{2}}\right)^{2}+2(a-1)$. We have

$$
\mathcal{I}_{1}=\sqrt{\frac{2}{\pi}} \cdot a \int_{0}^{\infty}\left(z^{2} J_{1}(z)+z\left(z-\frac{a}{\sqrt{2}}\right) J_{2}(z)\right) \mathrm{d} z
$$

where

$$
J_{1}(z):=\int_{0}^{\infty} t^{3 / 2} e^{-t \cdot z_{1}^{2} / 2} \mathrm{~d} t, \quad J_{2}(z):=\int_{0}^{\infty} \log (t) \cdot t^{3 / 2} e^{-t \cdot z_{1}^{2} / 2} \mathrm{~d} t
$$

The integrals $J_{1}$ and $J_{2}$ can be calculated explicitly in the same way as in the proof of Corollary 3.3(ii); see Eq. (D.5). After some algebraic manipulations this gives us that $\mathcal{I}_{1}$ equals to

$$
2 a \int_{0}^{\infty}\left(\frac{z^{2}\left(11-\left(3 \gamma_{\mathrm{E}}+3 \log (2)+6 \log \left(z_{1}\right)\right)\right)}{z_{1}^{5}}-\frac{a z\left(8-\left(3 \gamma_{\mathrm{E}}+3 \log (2)+6 \log \left(z_{1}\right)\right)\right)}{\sqrt{2} z_{1}^{5}}\right) \mathrm{d} z
$$

Finally, we calculate the quantity above using the definite integrals below

$$
\begin{aligned}
\int_{0}^{\infty} \frac{z}{z_{1}^{5}} \mathrm{~d} z & =\frac{1}{3 \sqrt{2}(a-1)^{2} a} \\
\int_{0}^{\infty} \frac{z^{2}}{z_{1}^{5}} \mathrm{~d} z & =\frac{1}{6(a-1)^{2}} \\
\int_{0}^{\infty} \frac{z \log \left(z_{1}\right)}{z_{1}^{5}} \mathrm{~d} z & =\frac{8-3 \log (2)-3 a+3(a-2) a \log \left(\frac{a}{a-1}\right)+6 \log (a)}{18 \sqrt{2}(a-1)^{2} a} \\
\int_{0}^{\infty} \frac{z^{2} \log \left(z_{1}\right)}{z_{1}^{5}} \mathrm{~d} z & =\frac{4+6 a-6 \log (2)+6(2+(a-2) a) \log (a-1)-6(a-2) a \log (a)}{72(a-1)^{2}}
\end{aligned}
$$

which, after algebraic manipulations yields the formula for $\mathcal{I}_{1}$ in (D.6).

We now proceed to calculate the value of $\mathcal{I}_{2}$. After substitution $y=t z$ and $x=\sqrt{t} w$ we obtain:

$$
\mathcal{I}_{2}=\frac{2 a}{\pi} \int_{0}^{\infty} \int_{0}^{\infty} \int_{0}^{\infty} \frac{w^{2}\left(J_{1}(w, z, q)-J_{2}(w, z, q)\right)}{q\left(1+q^{2}\right)^{2}} \mathrm{~d} w \mathrm{~d} z \mathrm{~d} q
$$

where

$$
\begin{aligned}
& J_{1}(w, z, q)=\int_{0}^{\infty} t^{2} \cdot e^{-t \cdot \frac{\left(z+\frac{a}{\sqrt{2}}\right)^{2}-2 \sqrt{2} z+(w-q z)^{2}}{2}} \mathrm{~d} t \\
& J_{2}(w, z, q)=\int_{0}^{\infty} t^{2} \cdot e^{-t \cdot \frac{\left(z+\frac{a}{\sqrt{2}}\right)^{2}-2 \sqrt{2} z+(w+q z)^{2}}{2}} \mathrm{~d} t .
\end{aligned}
$$

Since $\int_{0}^{\infty} t^{2} e^{-\mu t}=\mu^{-3} \Gamma(3)$, then

$$
\mathcal{I}_{2}=\frac{32 a}{\pi} \int_{0}^{\infty} \int_{0}^{\infty} \int_{0}^{\infty} \frac{J_{1}^{\prime}(z, q)-J_{2}^{\prime}(z, q)}{q\left(1+q^{2}\right)^{2}} \mathrm{~d} z \mathrm{~d} q
$$

where

$$
J_{1}^{\prime}(z, q)=\int_{0}^{\infty} \frac{w^{2}}{\left(z_{1}^{2}+(w-q z)^{2}\right)^{3}} \mathrm{~d} w, \quad J_{2}^{\prime}(z, q)=\int_{0}^{\infty} \frac{w^{2}}{\left(z_{1}^{2}+(w+q z)^{2}\right)^{3}} \mathrm{~d} w
$$

with $z_{1}^{2}:=\left(z+\frac{a}{\sqrt{2}}\right)^{2}-2 \sqrt{2} z=\left(z+\frac{a-2}{\sqrt{2}}\right)^{2}+2(a-1)$. Functions $J_{1}^{\prime}(z, q)$ and $J_{2}^{\prime}(z, q)$ are definite integrals of rational functions, so it is well-known how to find them. Now, we write

$$
\mathcal{I}_{2}=\frac{8 a}{\pi} \int_{0}^{\infty} \int_{0}^{\infty} \frac{3 q z z_{1}+\left(z_{1}^{2}+3 q^{2} z^{2}\right) \arctan \left(\frac{q z}{z_{1}}\right)}{q\left(1+q^{2}\right)^{2} z_{1}^{5}} \mathrm{~d} q \mathrm{~d} z
$$

We thus have

$$
\mathcal{I}_{2}=\frac{8 a}{\pi} \int_{0}^{\infty} \frac{I_{1}(z)+I_{2}(z)+I_{3}(z)}{z_{1}^{5}} \mathrm{~d} z
$$

where

$$
\begin{aligned}
& I_{1}(z):=3 z z_{1} \int_{0}^{\infty} \frac{1}{\left(1+q^{2}\right)^{2}} \mathrm{~d} q=\frac{\pi}{4} \cdot 3 z z_{1}, \\
& I_{2}(z):=z_{1}^{2} \int_{0}^{\infty} \frac{\arctan \left(\frac{q z}{z_{1}}\right)}{q\left(1+q^{2}\right)^{2}} \mathrm{~d} q=\frac{\pi}{4} \cdot z_{1}^{2}\left(2 \log \left(1+\frac{z}{z_{1}}\right)-\frac{z}{z+z_{1}}\right) \\
& I_{3}(z):=3 z^{2} \int_{0}^{\infty} \frac{q \arctan \left(\frac{q z}{z_{1}}\right)}{\left(1+q^{2}\right)^{2}} \mathrm{~d} q=\frac{\pi}{4} \cdot \frac{3 z^{3}}{z+z_{1}}
\end{aligned}
$$

where the formula from $I_{2}(z)$ follows from (D.3) and formulas 4.535-7 and 4.535-11 in [20], while $I_{3}(z)$ was found using formula 4.535-11. This gives us

$$
\mathcal{I}_{2}=2 a \cdot\left(3 \int_{0}^{\infty} \frac{z}{z_{1}^{4}} \mathrm{~d} z+2 \int_{0}^{\infty} \frac{\log \left(1+\frac{z}{z_{1}}\right)}{z_{1}^{3}} \mathrm{~d} z-\int_{0}^{\infty} \frac{z}{z_{1}^{3}\left(z+z_{1}\right)} \mathrm{d} z+3 \int_{0}^{\infty} \frac{z^{3}}{z_{1}^{5}\left(z+z_{1}\right)} \mathrm{d} z\right) .
$$

Finally, we calculate the quantity above using the definite integrals below

$$
\begin{aligned}
& \int_{0}^{\infty} \frac{z}{z_{1}^{4}} \mathrm{~d} z= \frac{2 \sqrt{a-1}+(a-2)\left(\operatorname{atan}\left(\frac{a-2}{2 \sqrt{a-1}}\right)-\frac{\pi}{2}\right)}{8(a-1)^{3 / 2}}, \\
& \int_{0}^{\infty} \frac{\log \left(1+\frac{z}{z_{1}}\right)}{z_{1}^{3}} \mathrm{~d} z= \frac{-a+4(a-1) \log (2)+(a-2) \sqrt{a-1}\left(\operatorname{atan}\left(\frac{a-2}{2 \sqrt{a-1}}\right)-\frac{\pi}{2}\right)}{(a-1) a^{2}} \\
&-\frac{(a-2)^{2} \operatorname{acoth}(1-2 a)}{(a-1) a^{2}}, \\
& \int_{0}^{\infty} \frac{z}{z_{1}^{3}\left(z+z_{1}\right)} \mathrm{d} z= \frac{a-(a-2) \sqrt{a-1}\left(\operatorname{atan}\left(\frac{a-2}{2 \sqrt{a-1}}\right)-\frac{\pi}{2}\right)}{(a-1) a^{2}} \\
& \int_{0}^{\infty} \frac{-\frac{4(a-1)(\operatorname{acoth}(1-2 a)+\log (2))}{a^{2}},}{z_{1}^{5}\left(z+z_{1}\right)} \mathrm{d} z= \frac{2 a\left(-3 a^{2}+17 a-12\right)-3(a-2) \sqrt{a-1}\left(a^{2}+8 a-8\right)\left(\operatorname { a t a n } \left(\frac{a-2}{\left.2 \sqrt{a-1})-\frac{\pi}{2}\right)}\right.\right.}{24(a-1)^{2} a^{2}} \\
& a^{2}
\end{aligned}
$$

which, after some algebraic manipulations yields (D.6) for $\mathcal{I}_{2}$. This completes the proof.

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